

**Baryon asymmetry of the Universe**

**from**

**sterile neutrinos**

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## Outline

baryon asymmetry of the Universe from neutrinos: leptogenesis

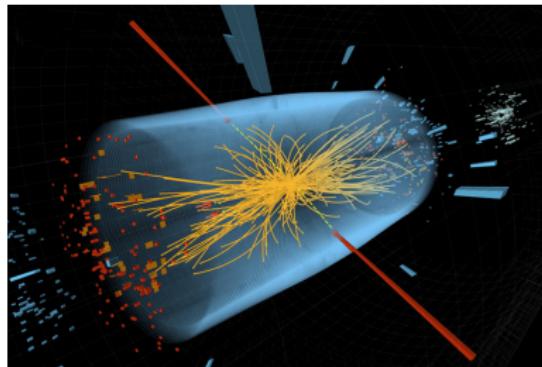
first principle approach: effective equations of motion

coefficients from thermal Green's functions: washout

production of ultrarelativistic sterile neutrinos

## Physics beyond the Standard Model

all particles of the SM have  
been found



physics beyond the SM must account for

dark matter particles

neutrino masses

baryon asymmetry

# Baryon asymmetry of the Universe (BAU)

net baryon density

$$n_B \equiv n_b - n_{\bar{b}}$$

baryon to photon ratio

$$\eta \equiv n_B / n_\gamma$$

Big Bang Nucleosynthesis

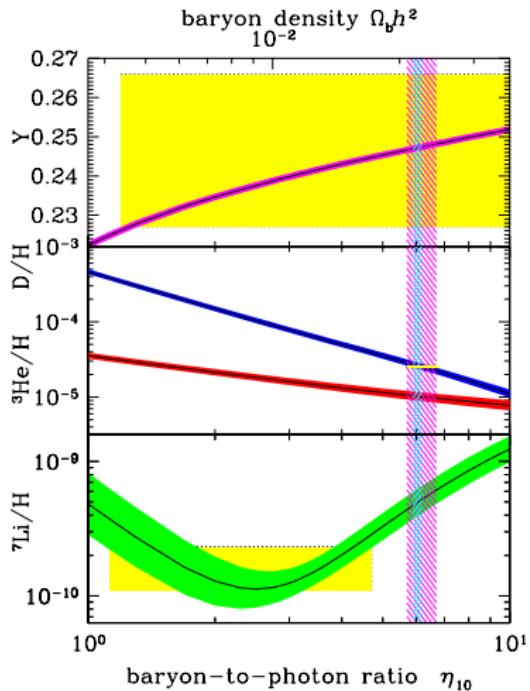
( $1 \text{ MeV} \gtrsim T \gtrsim 10 \text{ keV}$ )  $\rightsquigarrow$

$$5.7 < (\eta \times 10^{10}) < 6.7$$

Cosmic Microwave Background

( $T \sim 0.25 \text{ eV}$ )

$$\eta \times 10^{10} = 6.04 \pm 0.08 \text{ [Planck]}$$



[Particle Data Group]

## Baryon asymmetry of the Universe (BAU)

why is  $\eta \neq 0$ ?

initial condition?

not if there was inflation!

Sakharov: asymmetry can be  
dynamically generated if there is

1. baryon number violation
2.  $C$  and  $CP$  violation
3. non-equilibrium



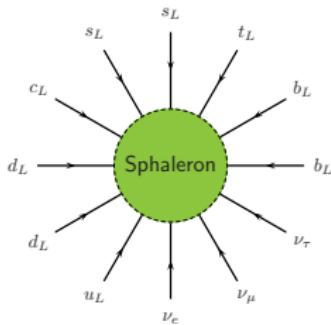
## Baryon + lepton number violation

$B - L$  is conserved in the Standard Model

$B + L$  is not [t'Hooft]

$B + L$  violation unsuppressed for  $T \gtrsim 160$  GeV

'sphaleron' processes



Lepton asymmetry  $\leftrightarrow$  Baryon asymmetry

## Neutrino masses

SM: massless neutrinos, but from neutrino oscillations:

$$\Delta m_{\text{solar}}^2 \simeq 7.6 \times 10^{-5} \text{ eV}, \quad \Delta m_{\text{atmospheric}}^2 \simeq 2.4 \times 10^{-3} \text{ eV}$$

add right-handed (sterile) neutrinos  $N_I = N_I^c$ :

$$\mathcal{L}_N = \frac{i}{2} \overline{N} \not{\partial} N - \frac{1}{2} \overline{N^c} M N - \left( \overline{N} h \widetilde{\varphi}^\dagger \ell + \text{h.c.} \right)$$

$M \gg h v \Rightarrow$  see-saw formula for light  $\nu$  mass matrix

$$m_\nu = h M^{-1} h^T v^2$$

$$m_\nu \sim 0.1 \text{ eV},$$

$$m_e/v < h < 1 \Leftrightarrow$$

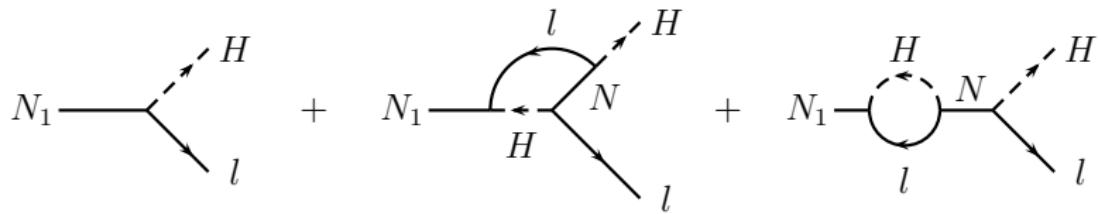
$$\text{TeV} \lesssim M \lesssim M_{\text{GUT}}$$



## Baryogenesis through leptogenesis [Fukugita, Yanagida]

Majorana masses  $M_{ij} \rightarrow$  lepton number violation

complex Yukawa couplings  $h_{ij} \rightarrow$  CP-violation



decay rates

$$\Gamma(N \rightarrow \ell\varphi) \neq \Gamma(N \rightarrow \bar{\ell}\bar{\varphi})$$

$\Gamma \lesssim H \rightarrow$  non-equilibrium

## Baryogenesis through leptogenesis: scenarios

thermal leptogenesis (non-resonant)

- asymmetry from sterile neutrino **decay**
- no fine tuning
- close to thermal equilibrium, non-relativistic [DB, Wörmann]
- lightest sterile neutrino  $M_1 \gtrsim 10^9$  GeV [Davidson, Ibarra; di Bari]
- GUT scale physics

resonant leptogenesis:  $M_2 - M_1 \sim$  thermal width [Pilaftsis, Underwood]

- no mass bound

## Baryogenesis through leptogenesis: scenarios

asymmetry from **production** of sterile  $N_i$  [Akhmedov, Smirnov, Rubakov]

- far from equilibrium
- $M_2, M_3 \gtrsim \text{MeV}$  [Canetti, Drewes, Shaposhnikov '13]
- thermal effects  $\rightarrow$  no fine tuned mass degeneracy needed

[Drewes, Garbrecht]

## Decay and non-equilibrium

$$K \equiv \left. \frac{\Gamma_0}{H} \right|_{T=M_1} \text{ 'washout factor'}$$

$K \gg 1$ : close to equilibrium when  $T \sim M_1$  'strong washout'

$K \ll 1$ : far from equilibrium when  $T \sim M_1$  'weak washout'

$$K = \frac{\tilde{m}_1}{m_*}, \quad \tilde{m}_1 = \frac{(m_D m_D^\dagger)_{11}}{M_1} \quad m_* \simeq 10^{-3} \text{ eV}$$

$\tilde{m}_1 >$  smallest light neutrino mass [Fujii, Hamaguchi, Yanagida]

$$(\Delta m_{\text{solar}}^2)^{1/2} < \tilde{m}_1 < (\Delta m_{\text{atmospheric}}^2)^{1/2} \Leftrightarrow 7.4 < K < 46$$

## Traditional approach to leptogenesis

Boltzmann equations for phase space densities  $f_a(t, |\mathbf{p}|)$

$$D_t f_a = \text{Coll}_a[f]$$

collision term (for leptons)

$$\begin{aligned} \text{Coll}_\ell[f] &= \int_{\mathbf{p}_i} (2\pi)^4 \delta(p_\ell + p_{\bar{\varphi}} - p_N) \\ &\times \left[ |\mathcal{M}|_{N \rightarrow \ell \bar{\varphi}}^2 f_N (1 - f_\ell) (1 + f_{\bar{\varphi}}) - |\mathcal{M}|_{\ell \bar{\varphi} \rightarrow N}^2 f_\ell f_{\bar{\varphi}} (1 - f_N) \right] + \dots \end{aligned}$$

problems:

double counting of resonant intermediate states

unclear how to include medium effects

theoretical error  $\leftrightarrow$  radiative corrections ???

## First principles approaches to leptogenesis

identify slow and fast variables  $X$

$\gamma_X$  = relaxation rate

$\gamma_X \gg H$  fast, in thermal equilibrium *spectator processes*

$\gamma_X \sim H$ , slow, interesting non-equilibrium dynamics

$\gamma_X \ll H$  practically conserved

write effective equations of motion for slow ones

computation = ‘two step procedure’

1. short time/distance physics  $\rightsquigarrow$  coefficients
2. solve effective equations of motion

## First principles approaches to leptogenesis

effective equations of motion for slow variables

$$D_t X_a = -\gamma_{ab} X_b$$

valid

on time scales  $\gtrsim \gamma^{-1}$

to all orders in Standard Model couplings

coefficients

$$\gamma_{ab} = \gamma_{ab}(T)$$

determined by short time physics

only depend on temperature

radiative corrections can be systematically computed

## Non-relativistic unflavored leptogenesis

non-relativistic approximation: neglect motion of  $N_i$

Simplest case:

$$M_1 \ll M_2, M_3$$

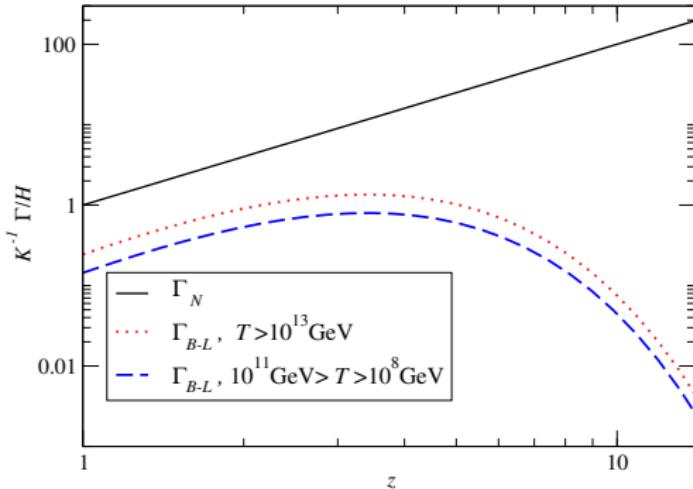
only one charged lepton flavor relevant

$\rightsquigarrow$  only  $n_N \equiv n_{N_1}$  and  $n_{B-L}$  need to be considered

$$\left( \frac{d}{dt} + 3H \right) n_N = -\gamma_N (n_N - n_{N,eq}) + \gamma_{N,B-L} n_{B-L}$$

$$\left( \frac{d}{dt} + 3H \right) n_{B-L} = \gamma_{B-L,N} (n_N - n_{N,eq}) - \gamma_{B-L} n_{B-L}$$

## Rates vs $z \equiv M_1/T$

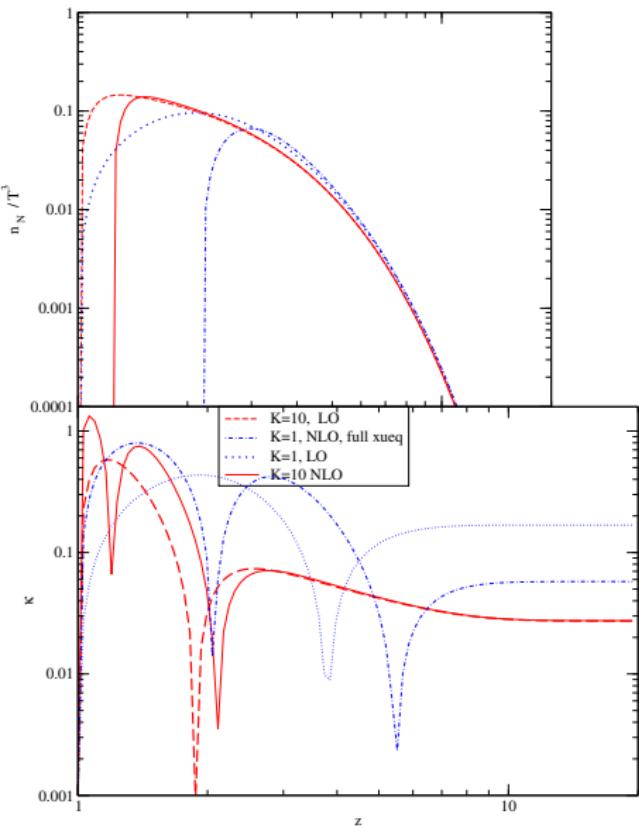


$n_N$  approaches equilibrium exponentially small for  $T \ll M_1$   
I'genesis must happen before  $B - L$  washout rate  $\gamma_{B-L}$   
maximal for  $z \sim 4$   
 $\rightarrow 0$  for  $z \rightarrow \infty$   
asymmetry freezes in

[DB, M. Wörmann]

$$K \equiv \left. \frac{\Gamma_0}{H} \right|_{T=M_1} \text{'washout factor'}$$

## Non-relativistic, relativistic corrections



relativistic correction:  
include also kinetic energy  
density of sterile  $N$

## $\gamma_{ab}$ from correlation of thermal fluctuations

[DB, M. Laine]

effective eqs. of motion for thermal fluctuations

$$\dot{X}_a = -\gamma_{ab}X_b + \xi_a$$



Langevin equation

random force  $\xi$

$$\langle \xi_a(t)\xi_b(t') \rangle \propto \delta(t-t')$$

real time correlation function

$$\langle X_a(t)X_c(0) \rangle = (e^{-\gamma t})_{ab} \langle X_b X_c \rangle$$

## Correlations from finite temperature QFT

$$C_{ab}(t) \equiv \frac{1}{2} \langle \{X_a(t), X_b(0)\} \rangle$$

match results at time/frequency scales

$$t_{\text{UV}} \ll t \ll \gamma^{-1}, \quad \omega_{\text{UV}} \gg \omega \gg \gamma$$

at leading order in  $\hbar$ :

$$\gamma_{ab} = \frac{1}{2V} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \int dt e^{i\omega t} \left\langle \left[ \dot{X}_a(t), \dot{X}_b(0) \right] \right\rangle_0 (\Xi^{-1})_{cb}$$

similar to Kubo relation for transport coefficients:  
viscosity, in particular diffusion constants

matrix of susceptibilities  $\Xi_{ab} \equiv \frac{1}{TV} \langle X_a X_b \rangle$

## Washout rate

$$\left( \frac{d}{dt} + 3H \right) n_N = -\gamma_N (n_N - n_{N,eq}) + \gamma_{N,B-L} n_{B-L}$$

$$\left( \frac{d}{dt} + 3H \right) n_{B-L} = \gamma_{B-L,N} (n_N - n_{N,eq}) - \gamma_{B-L} n_{B-L}$$

## Washout rate

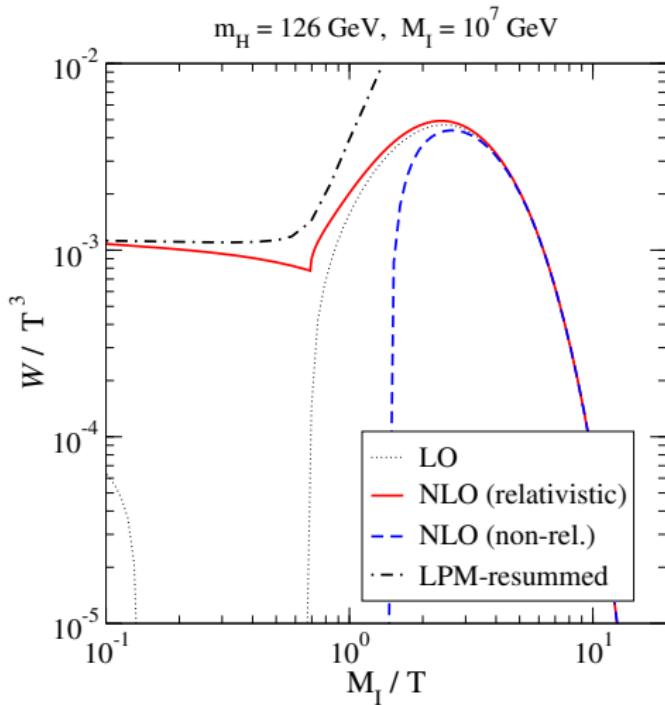
integrate out  $N_I$  at leading order  $\Rightarrow$

$$\gamma_{ab} = -\frac{1}{2} \sum_I \int_{\mathbf{k}} \frac{f'_F(E_I)}{2E_I} h_{Ii} \operatorname{tr} \left[ k \left( T_a^\ell [\tilde{\rho}(k) + \tilde{\rho}(-k)] T_c^\ell \right. \right. \\ \left. \left. + T_c^\ell [\tilde{\rho}(k) + \tilde{\rho}(-k)] T_a^\ell \right)_{ij} \right] h_{Ij}^* (\Xi^{-1})_{cb}$$

with spectral function

$$\tilde{\rho}_{ij\alpha\beta}(k) \equiv \int_x e^{ik \cdot x} \left\langle \left\{ (\tilde{\varphi}^\dagger \ell_{i\alpha})(x), (\bar{\ell}_{j\beta} \tilde{\varphi})(0) \right\} \right\rangle_0$$

## Washout rate



integral over spectral function

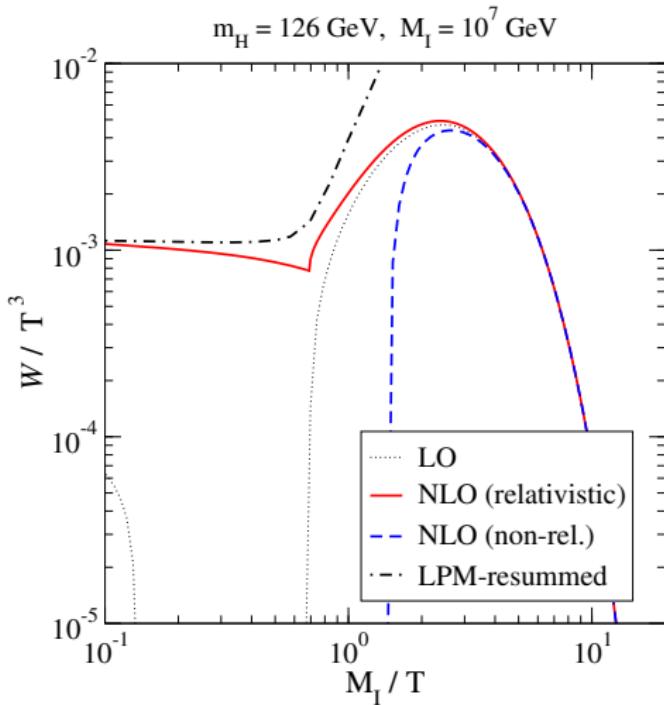
## spectral function results

ultra-relativistic ( $M_I \lesssim g^2 T$ ):  
complete LO [D Besak, DB]

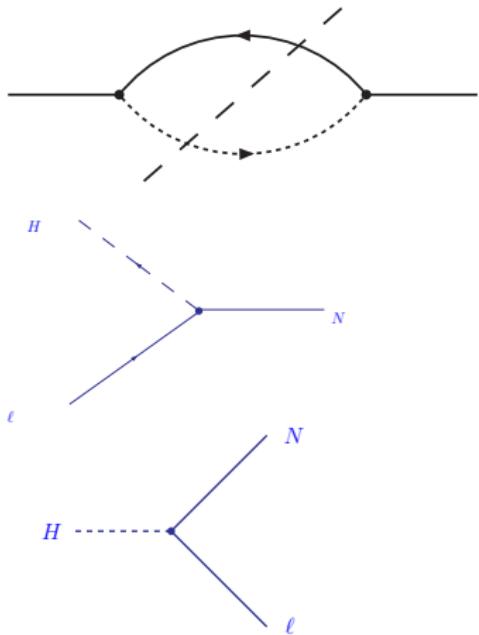
relativistic ( $M_I \sim T$ ):  
NLO [M Laine]

non-relativistic ( $M_I \gg T$ ):  
NLO  
[A Salvio et al., Laine, Y Schröder]

## Washout rate

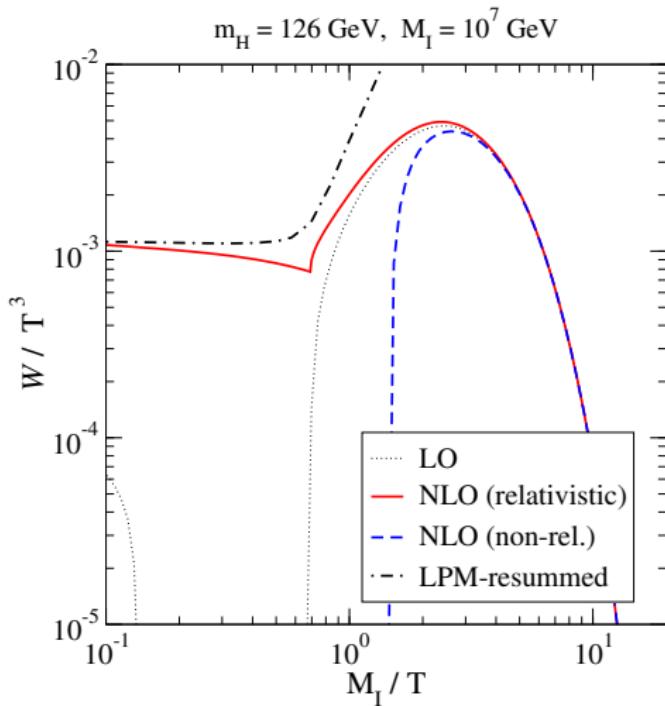


(naive) LO

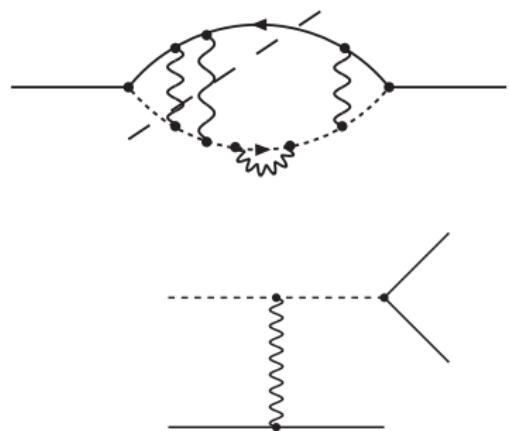


integral over spectral function

## Washout rate



complete LO



integral over spectral function

## Washout rate

corrections to spectral functions (non-relativistic and relativistic )  
 $= O(g^2)$

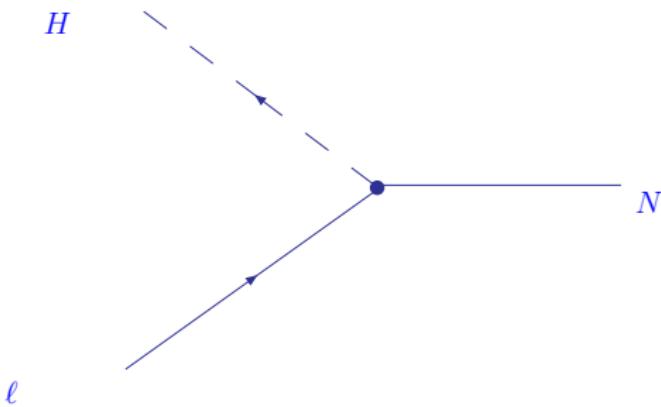
corrections to susceptibilities  $= O(g)$ , infrared effect

leading corrections from 'simple' thermodynamics

complete  $O(g^2)$  computed [DB, M. Sangel]  
corrections  $\leq 4\%$ , mostly from QCD

## *N*-production: inverse decay

momenta =  $O(T)$



at high temperature  $T \gtrsim M_N$ :  
momenta are

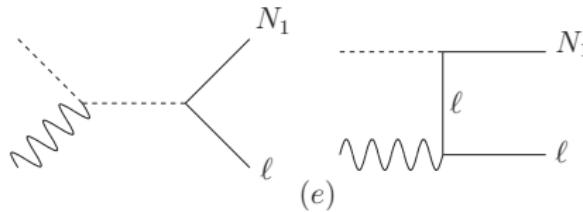
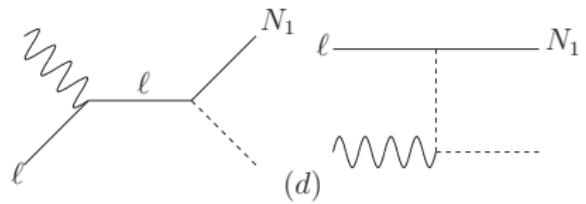
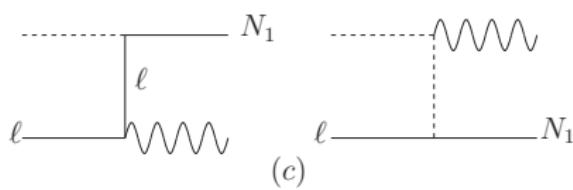
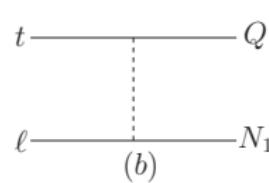
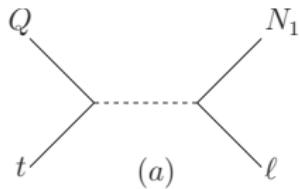
close to light cone  
nearly collinear

when  $M_N \lesssim gT$ :  
- opening angle =  $O(g)$   
-  $O(g^2)$  phase space suppression

$H$  = Higgs

$\ell$  = Standard Model lepton

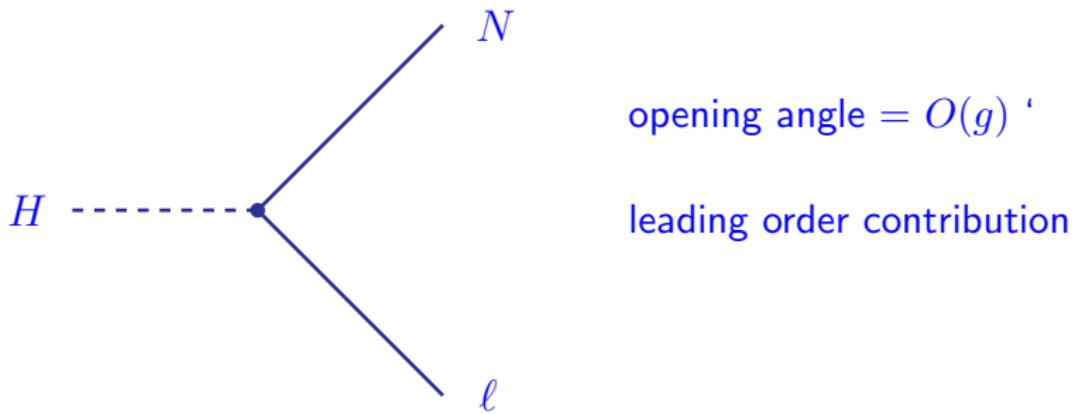
## **N-production: $2 \rightarrow 2$ scattering**



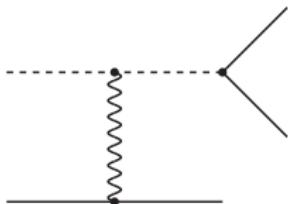
Besak, DB

## Thermal mass effects

medium effects  $\rightsquigarrow$  thermal masses  
 $\rightsquigarrow$  new channel



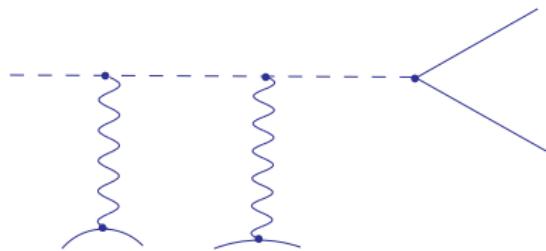
## Soft gauge interactions



collinear enhancement

compensates additional vertices

~ $\rightarrow$  leading order contribution



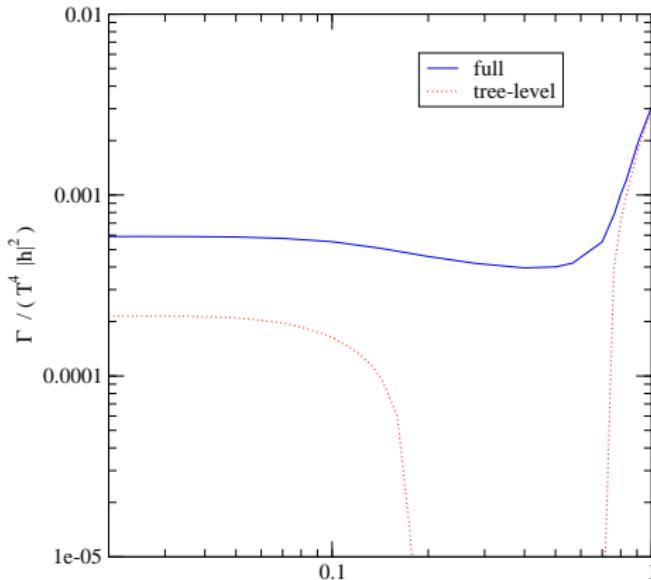
multiple soft scattering unsuppressed

leading order contribution

Landau-Pomeranchuk-Migdal effect

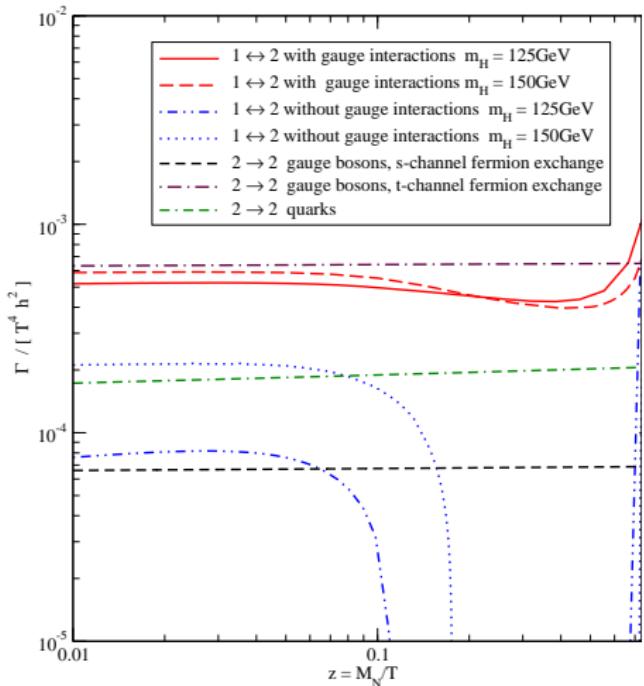
[Anisimov, Besak, DB]

## Thermal masses + soft gauge interactions



soft gauge interactions  
~~~  
factor 3 enhancement  
at high  $T$

# Complete LO production rate



gauge interactions  
dominate

## Summary and outlook

- lot of theoretical progress in leptogenesis
- systematic approach by identifying fast, slow and quasi-static quantities
- effective equations of motion for slow quantities
- coefficients in effective equations of motion related to real time correlation functions at finite temperature
- Kubo-type relations, valid to all orders in Standard Model couplings
- NLO and NNLO corrections computed for washout rate
- thermal production of ultrarelativistic  $N$ : thermal mass effects, soft gauge interactions very important

## Outlook

CP-asymmetry ?

error bars for leptogenesis

CP asymmetry in ultrarelativistic regime