



# The KATRIN experiment and next-generation tritium $\beta$ -decay experiments for keV-scale sterile neutrinos

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**Introduction**

**The Karlsruhe Tritium Neutrino experiment  
KATRIN**

**Beyond KATRIN and TOF@KATRIN**

**Sterile neutrino searches with KATRIN**

**Conclusions**

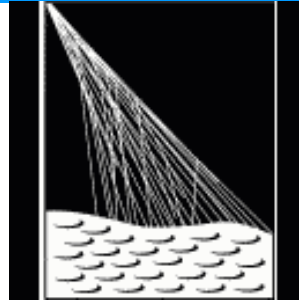




# Positive results from $\nu$ oscillation experiments

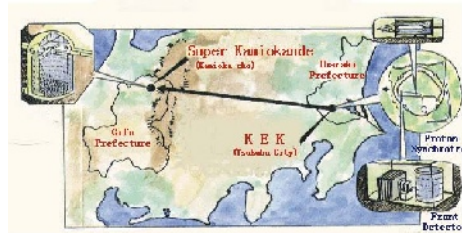
## atmospheric neutrinos

(Kamiokande, Super-Kamiokande, ...)



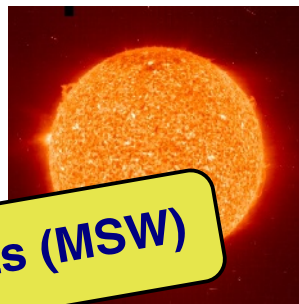
## accelerator neutrinos

(K2K, T2K, MINOS, OPERA, MiniBoone)



## solar neutrinos

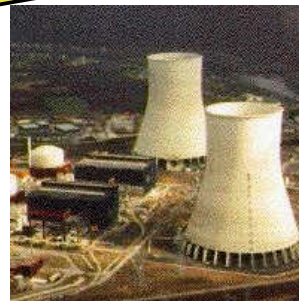
(Homestake, Gallex, Sage, Super-Kamiokande, SNO, Borexino)



**Matter effects (MSW)**

## reactor neutrinos

(KamLAND, CHOOZ, Daya Bay, DoubleCHOOZ, RENO, ...)



$\Rightarrow$  **non-trivial  $\nu$ -mixing**

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{pmatrix} \cdot \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

**with:**

$$0.37 < \sin^2(\theta_{23}) < 0.63 \quad \text{maximal!}$$

$$0.26 < \sin^2(\theta_{12}) < 0.36 \quad \text{large!}$$

$$0.018 < \sin^2(\theta_{13}) < 0.030 \quad 8.9^\circ$$

$$7.0 \cdot 10^{-5} \text{ eV}^2 < \Delta m_{12}^2 < 8.2 \cdot 10^{-5} \text{ eV}^2$$

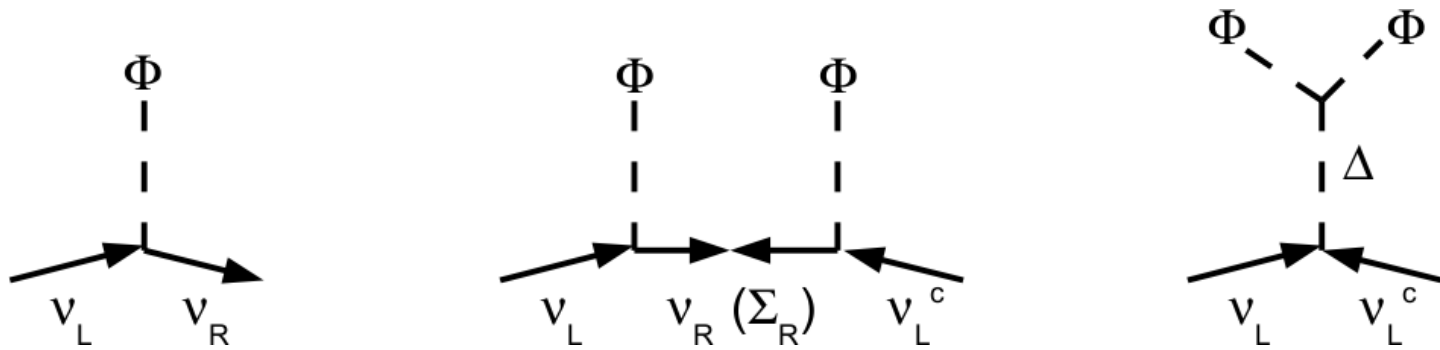
$$2.2 \cdot 10^{-3} \text{ eV}^2 < |\Delta m_{13}^2| < 2.6 \cdot 10^{-3} \text{ eV}^2$$

$\Rightarrow m(\nu_j) \neq 0$ , but unknown!

$\rightarrow$  **direct  $m(\nu)$  &  $0\nu\beta\beta$  searches, cosmology**

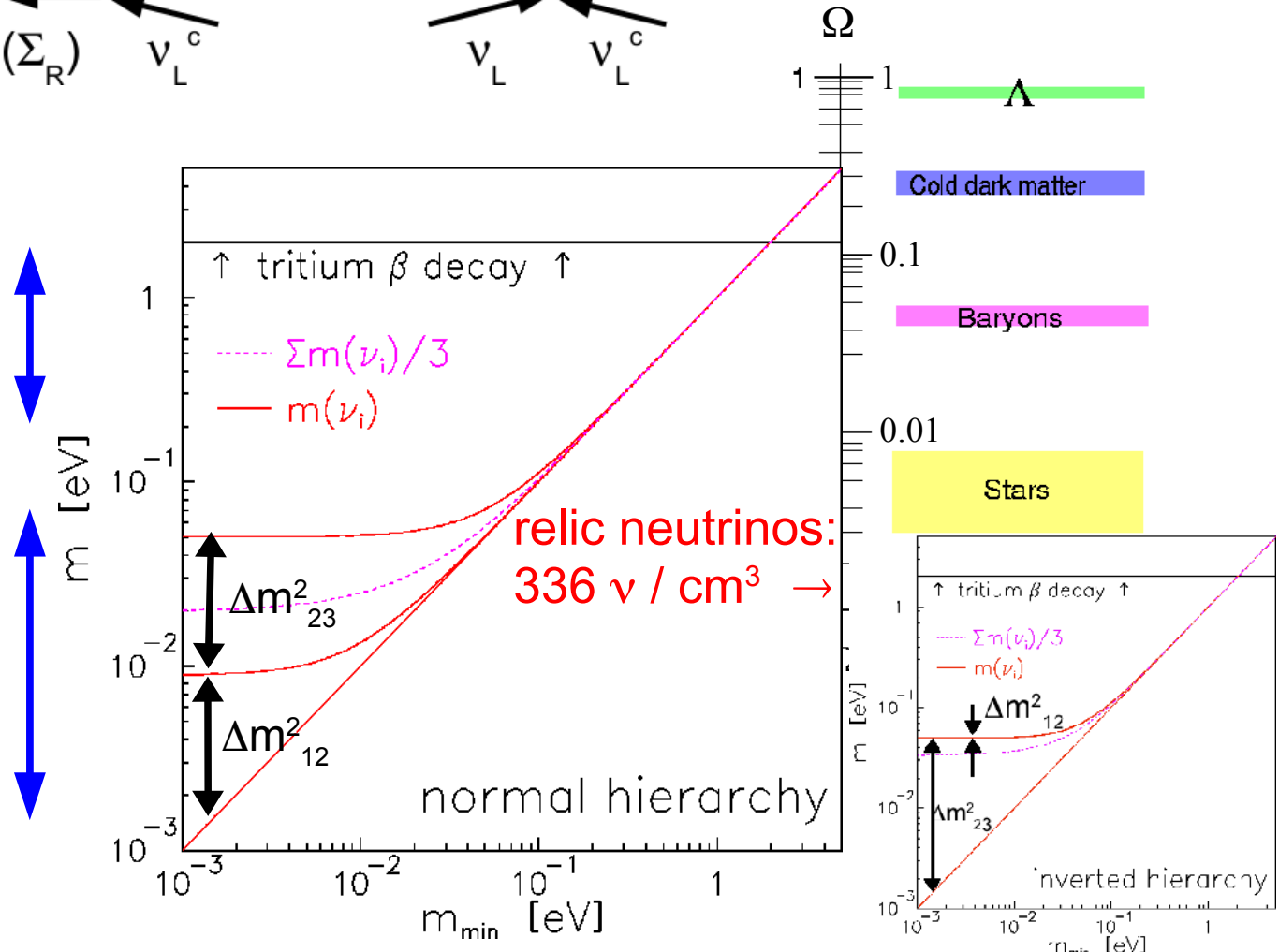
# Need for the absolute $\nu$ mass determination

**Non-zero neutrino masses go beyond the usual Yukawa coupling to the Higgs → Beyond the SM physics**



degenerated masses  
cosmological relevant  
e.g. seesaw mechanism type 2

hierarchical masses  
e.g. seesaw mechanism type 1  
explains smallness of masses, but not large (maximal) mixing

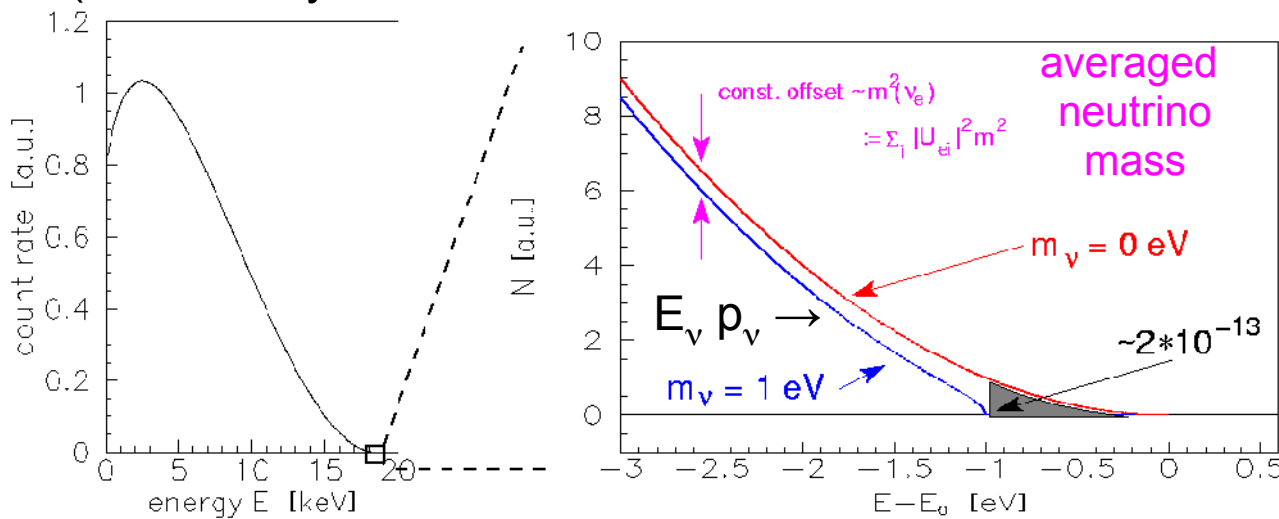


# Direct determination of $m(\nu_e)$ from $\beta$ decay

$$\beta: dN/dE = K \underbrace{F(E,Z)}_{\text{phase space}} \underbrace{p}_{p_e} \underbrace{E_{\text{tot}}}_{E_e} \underbrace{(E_0 - E_e)}_{E_\nu} \underbrace{\sum |U_{ei}|^2 \sqrt{(E_0 - E_e)^2 - m(\nu_i)^2}}_{p_\nu}$$

with “electron neutrino mass”:  $m(\nu_e)^2 := \sum |U_{ei}|^2 m(\nu_i)^2$

(modified by electronic final states, recoil corrections, radiative corrections)



**$m(\nu) < 2$  eV (Mainz, Troitsk)**

*Recent review:*  
**G. Drexlin, V. Hannen, S. Mertens,  
C. Weinheimer, Adv. High Energy  
Phys., 2013 (2013) 293986**

**Need:** low endpoint energy  $\Rightarrow$  Tritium  $^3\text{H}$  ( $^{187}\text{Re}$ ,  $^{163}\text{Ho}$ )  
 very high energy resolution &  
 very high luminosity &  
 very low background  $\Rightarrow$  MAC-E-Filter  
 (or bolometer for  $^{187}\text{Re}$ ,  $^{163}\text{Ho}$ )

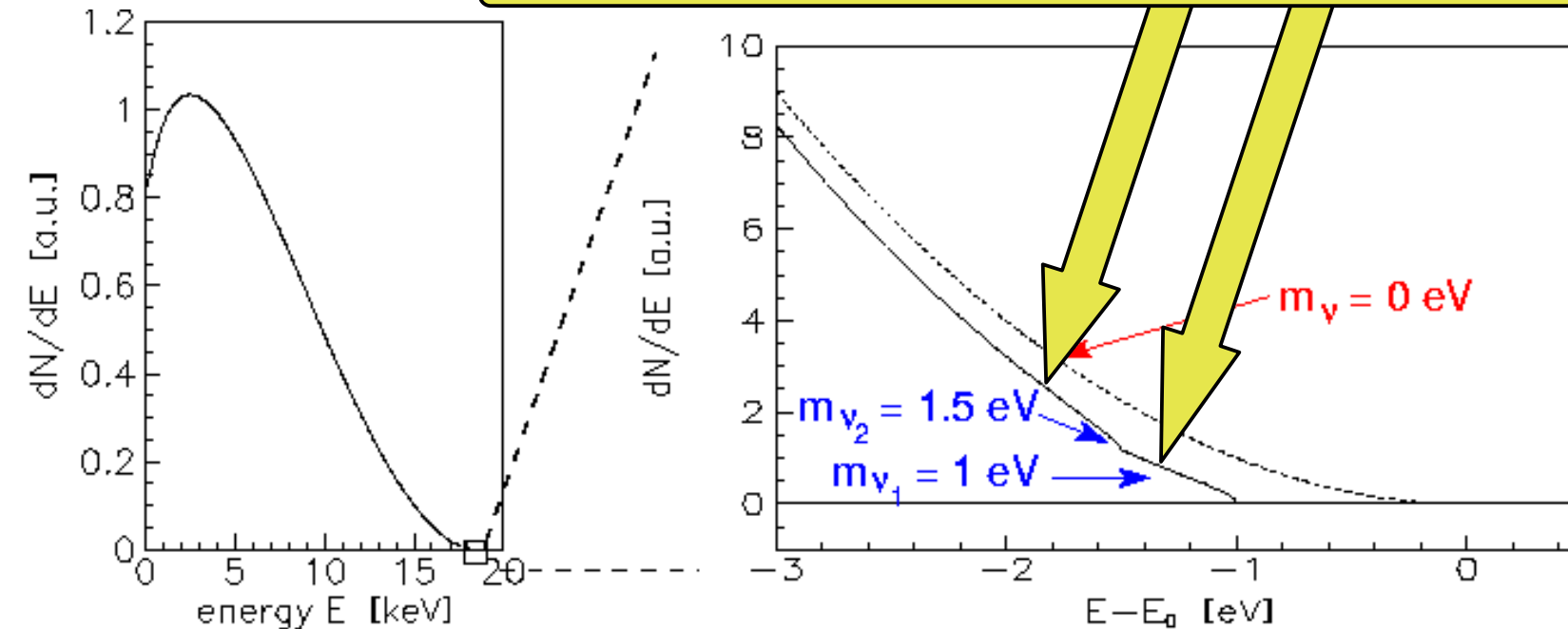
# Direct determination of $m(\nu_e)$ from $\beta$ decay

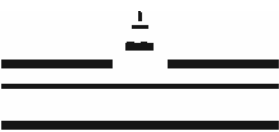
The neutrino mass eigenstates  
are observable as kinks in principle  
→ amplitudes add incoherently

If the resolution is better than the mass splitting  
then the various mass eigenstate appear

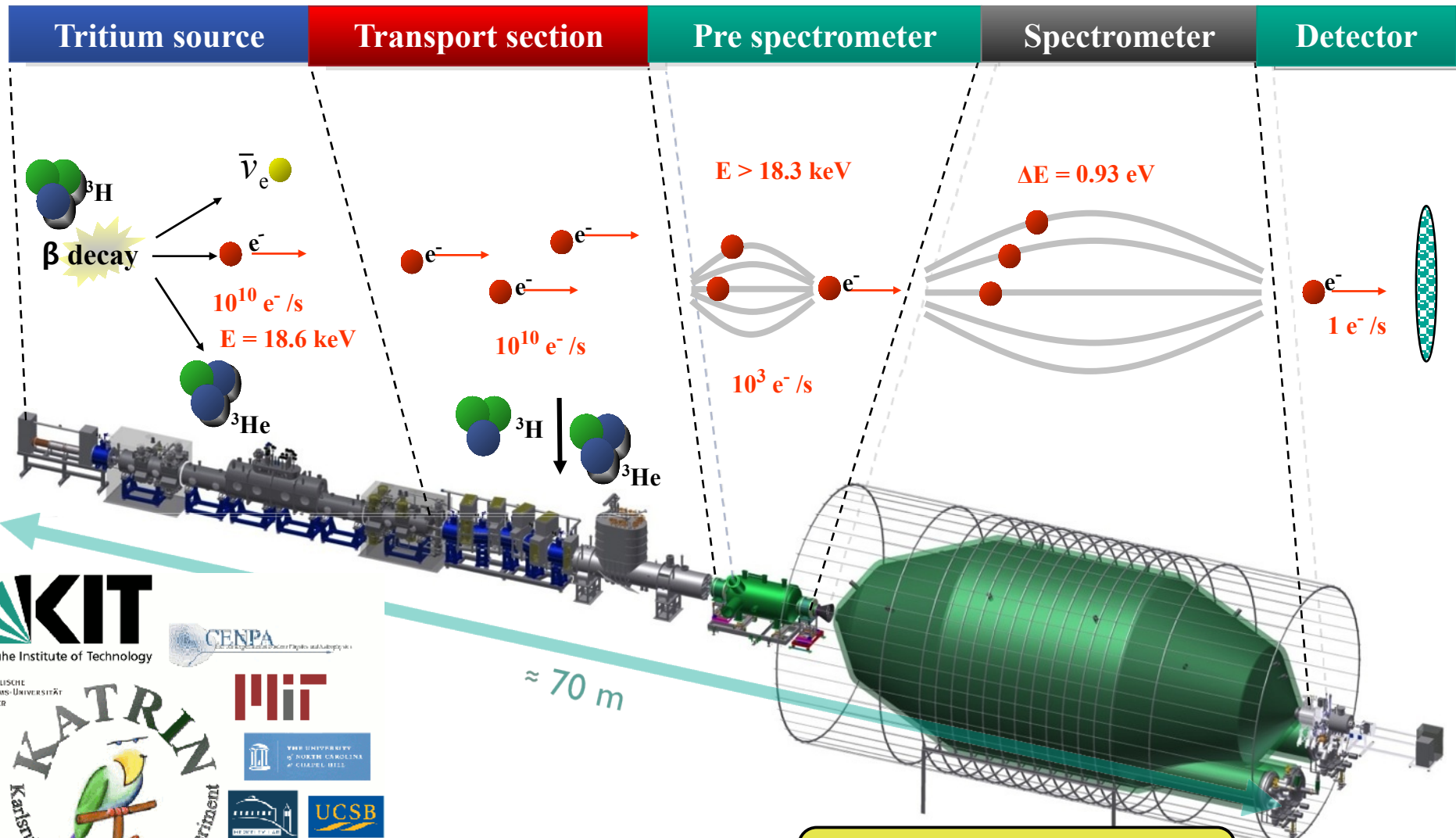
If mass eigenstates are not resolved:

$$m^2(\nu_e) = \sum |U_{ei}|^2 m^2(\nu_i)$$





# The Karlsruhe Tritium Neutrino Experiment KATRIN - overview



**Sensitivity on  $m(\nu_e)$ :**  
 $2 \text{ eV} \rightarrow 200 \text{ meV}$

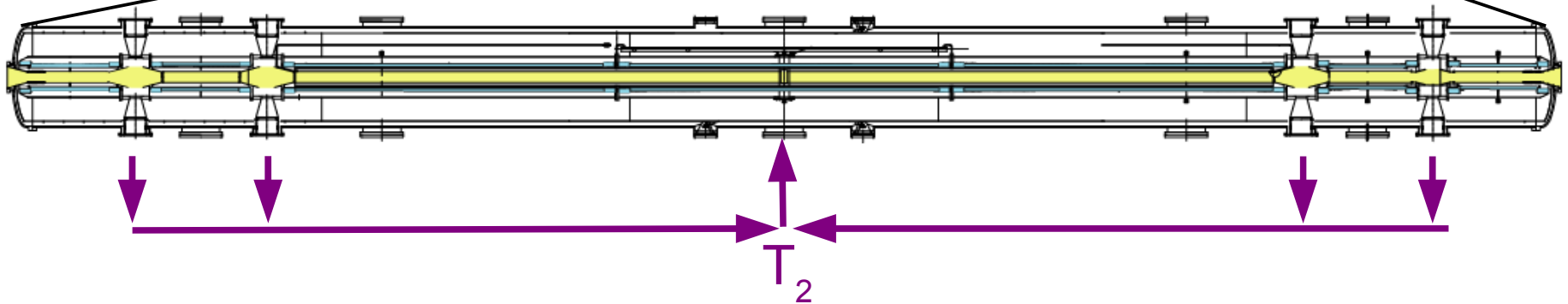
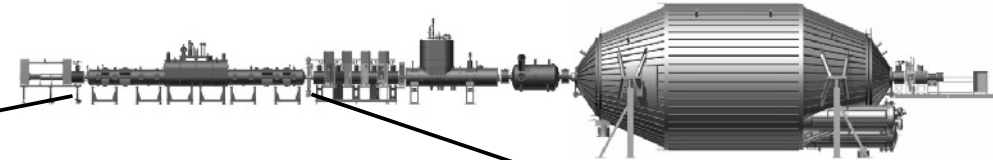


# Molecular Windowless Gaseous Tritium Source WGTS

per mill stability source strength request:

$$dN/dt \sim f_T \cdot N / \tau \sim n = f_T \cdot p \cdot V / R T$$

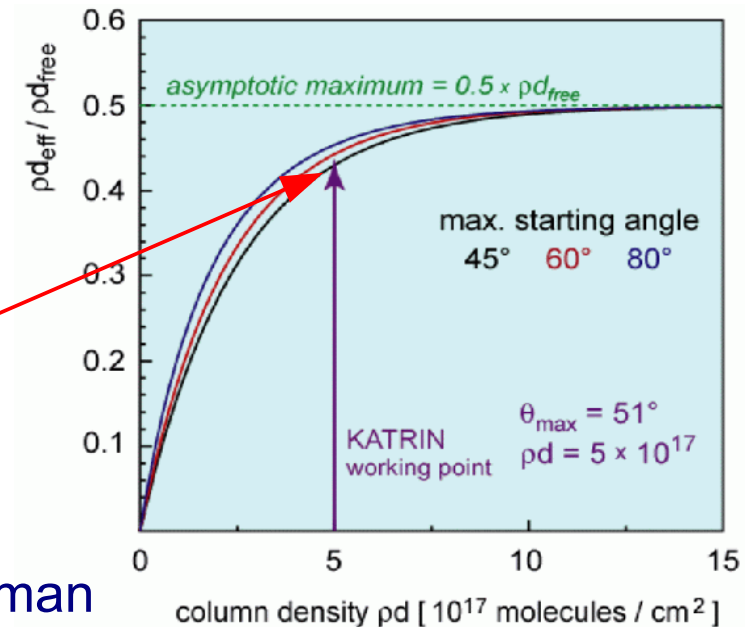
tritium fraction  $f_T$  & ideal gas law



WGTS: tube in long superconducting solenoids  
 $\varnothing$  9cm, length: 10m, T = 30 K

Tritium recirculation (and purification)  
 $p_{inj} = 0.003$  mbar,  $q_{inj} = 4.7$  Ci/s

allows to measure with near to maximum count rate using  
 $\rho d = 5 \cdot 10^{17}/\text{cm}^2$   
 with small systematics



check column density by e-gun,  $T_2$  purity by laser Raman



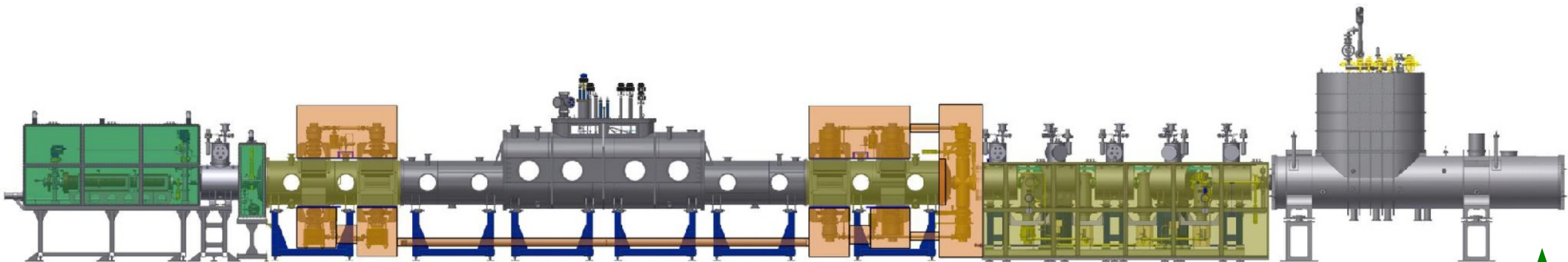
# Transport and differential & cryo pumping sections

Monitoring & calibration system

Molecular windowless gaseous tritium source

Differential pumping

Cryogenic pumping with Argon snow at LHe temperatures



$T_2$ -injection 1.8 mbar l/s (STP)  
=  $1.7 \cdot 10^{11}$  Bq/s = 40 g/d

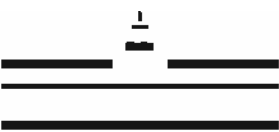
FT-ICR Penning traps to measure ions from WGTS

$\approx 10^{-7}$  mbar l/s

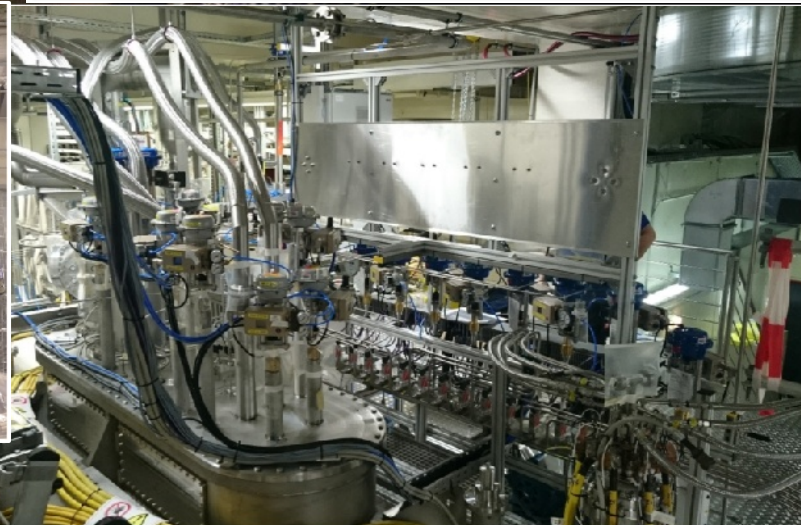
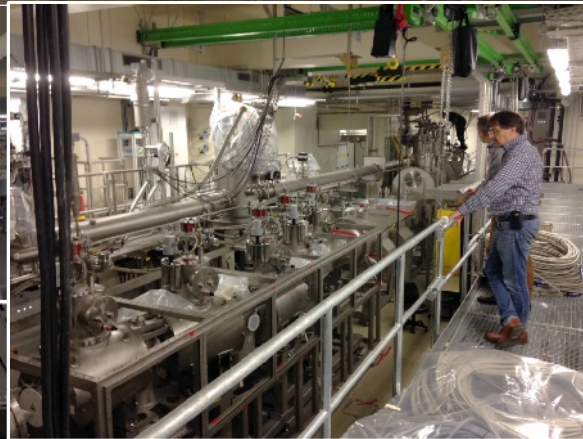
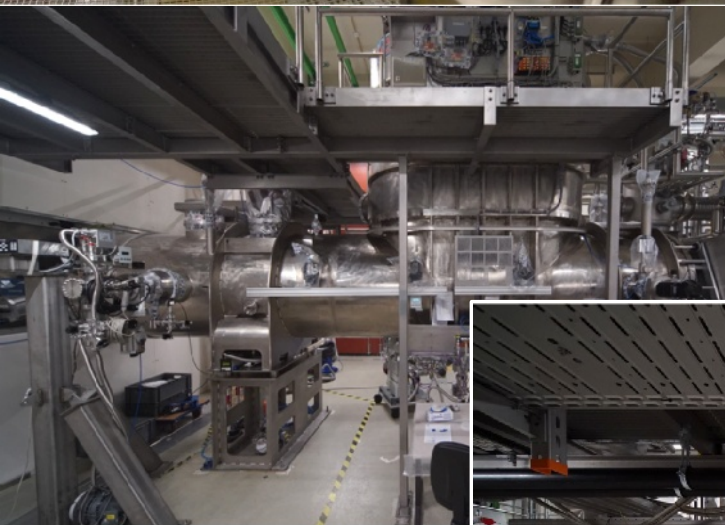
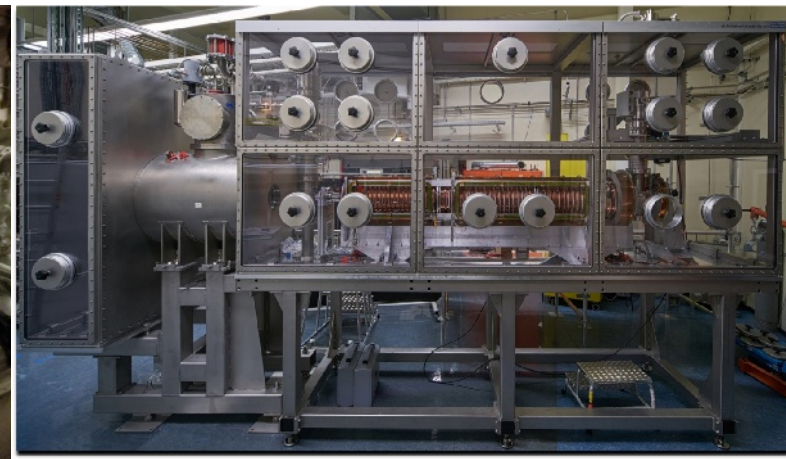
$< 2.5 \cdot 10^{-14}$  mbar l/s

$\Rightarrow$  adiabatic electron guiding &  $T_2$  reduction factor of  $\sim 10^{14}$

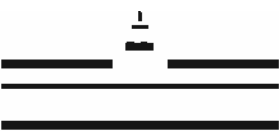




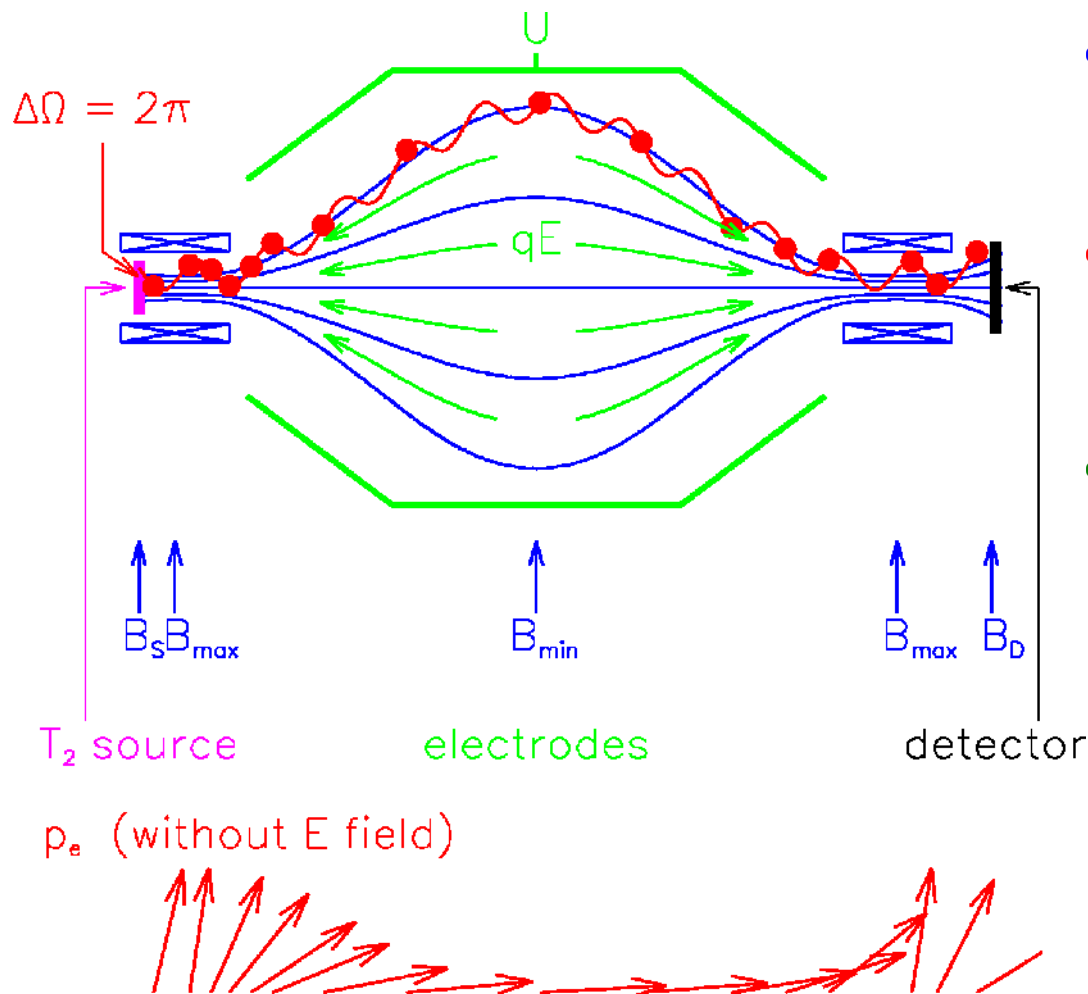
# The tritium source and transport section is really there !



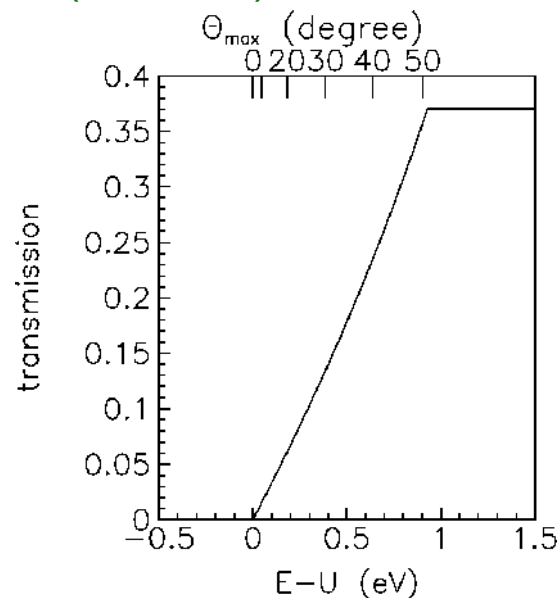
**Full commissioning under way  
start data taking with tritium  
in 2017**



# The classical way: Tritium $\beta$ -spectroscopy with a MAC-E-Filter



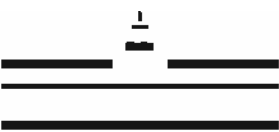
- Two supercond. solenoids compose magnetic guiding field
- adiabatic transformation:  
 $\mu = E_{\perp} / B = \text{const.}$   
 $\Rightarrow$  parallel  $e^-$  beam
- Energy analysis by electrostat. retarding field  
 $\Delta E = E \cdot B_{min} / B_{max}$   
 $= 0.93 \text{ eV (KATRIN)}$



$\Rightarrow$  sharp integrating transmission function without tails  $\rightarrow$

Magnetic Adiabatic Collimation + Electrostatic Filter  
(A. Picard et al., Nucl. Instr. Meth. 63 (1992) 345)

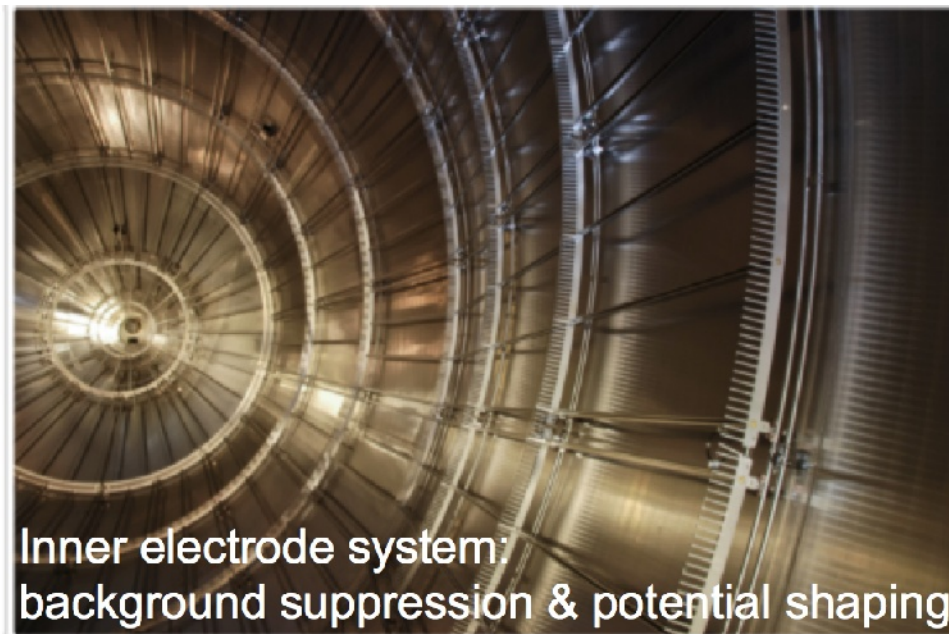




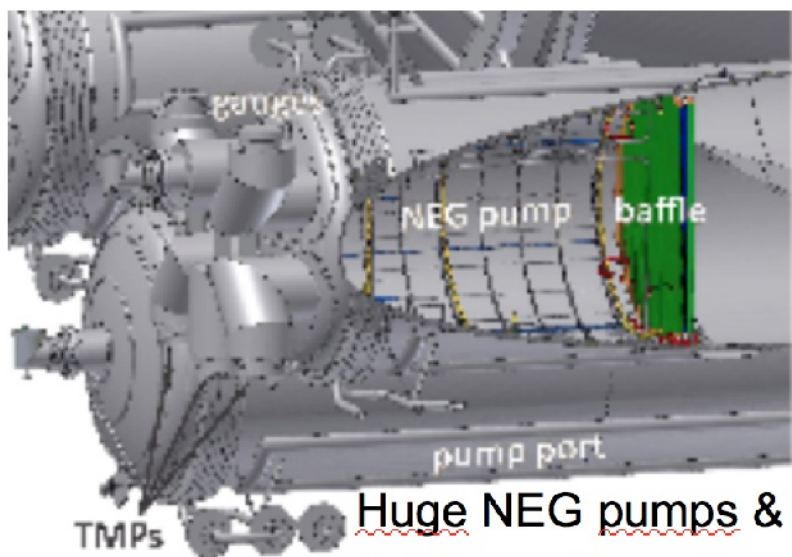
# KATRIN main spectrometer



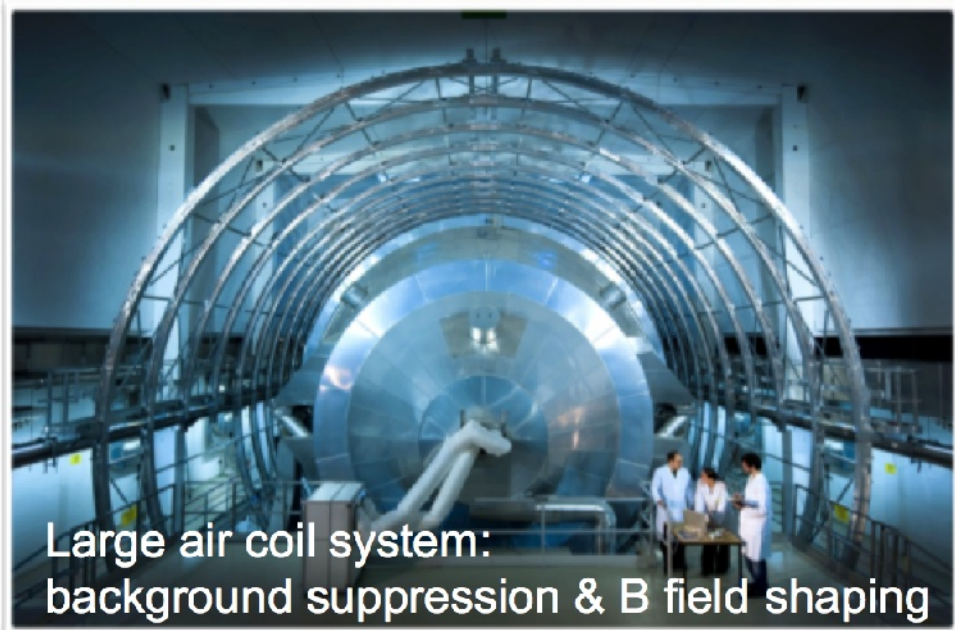
Huge spectrometer:  
high energy resolution & huge acceptance



Inner electrode system:  
background suppression & potential shaping

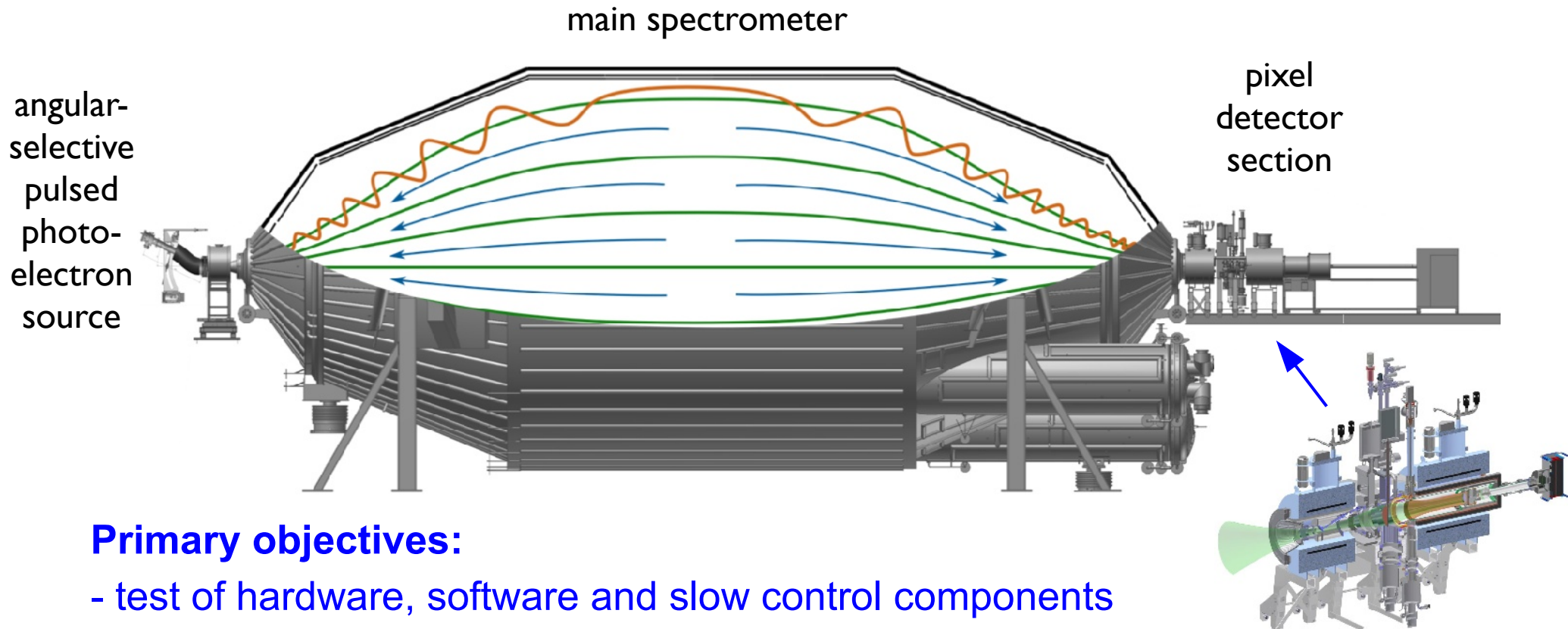


Huge NEG pumps & LN<sub>2</sub> baffles:  
UHV conditions & radon trapping



Large air coil system:  
background suppression & B field shaping

# Main spectrometer and detector commissioning – objectives

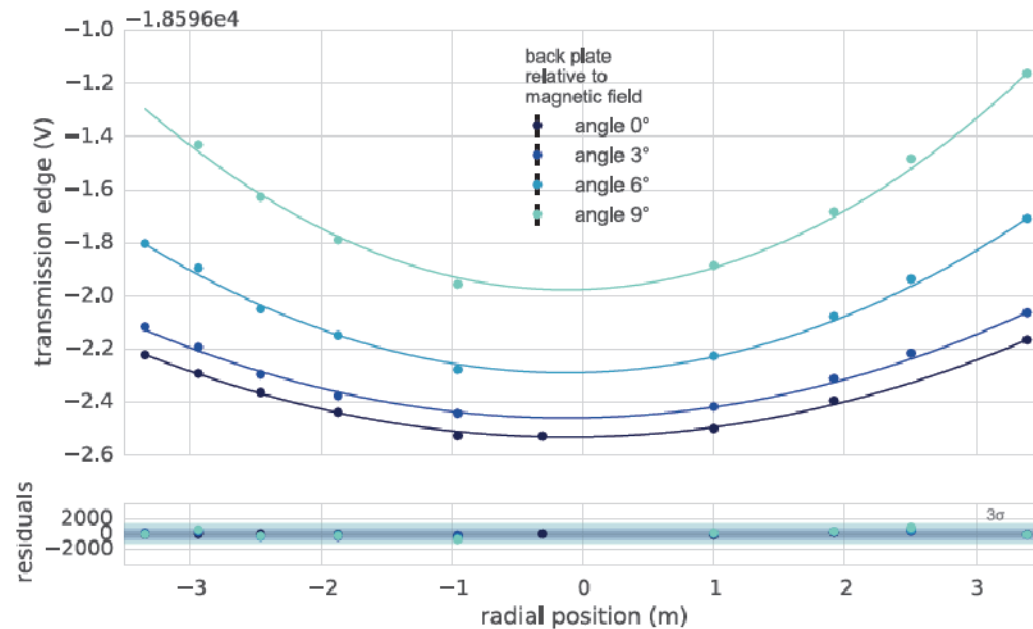
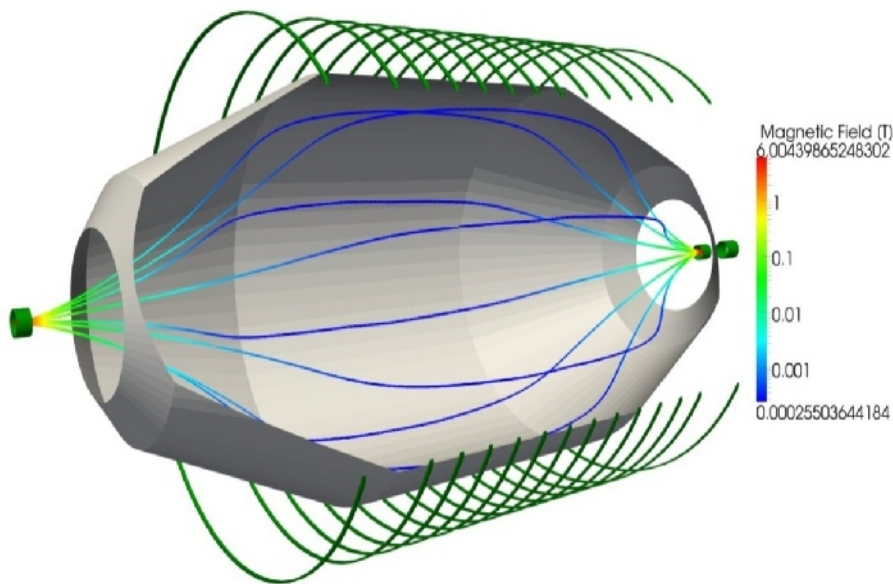
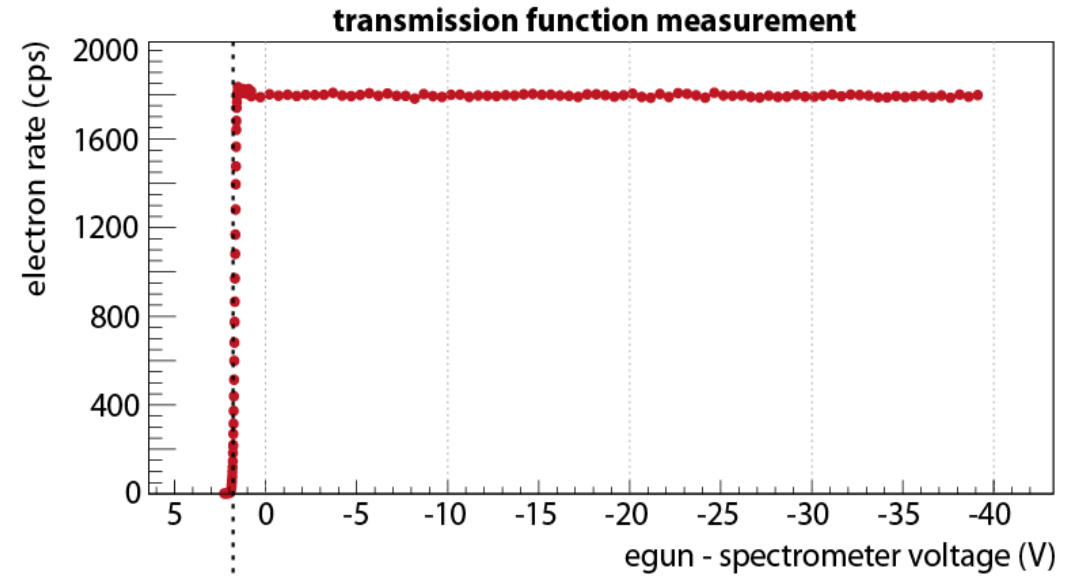
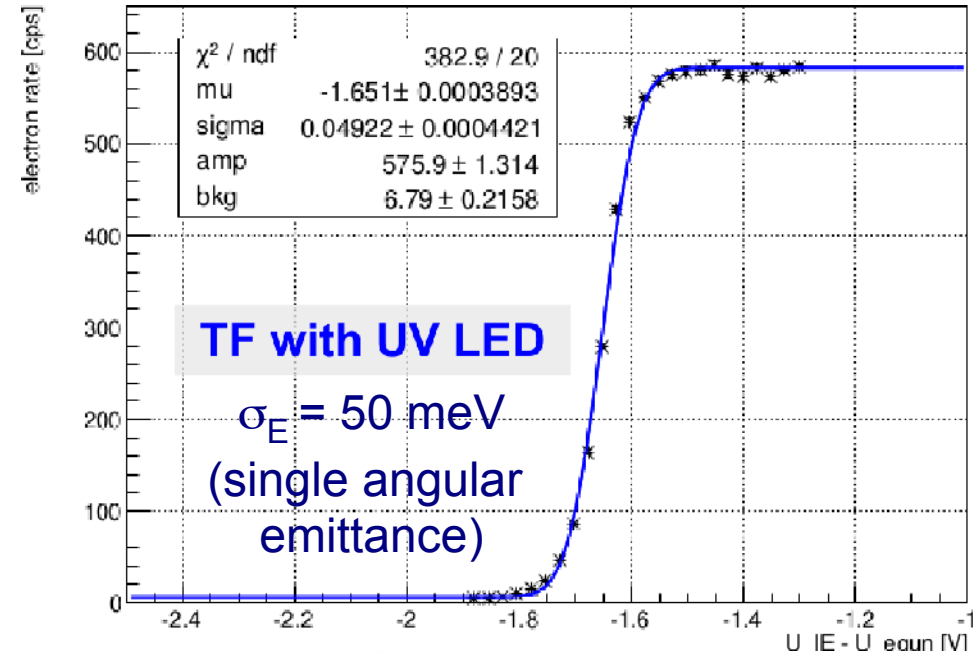


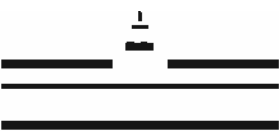
## Primary objectives:

- test of hardware, software and slow control components
- provide ultra high vacuum conditions at the  $p \approx 10^{-11}$  mbar level
- detailed understanding of the transmission properties of this MAC-E-Filter ( $E = 18.6$  keV with  $\Delta E = 0.93$  eV resolution) and compare to simulation with Kasseiopeia
- detailed understanding and passive & active control of background processes

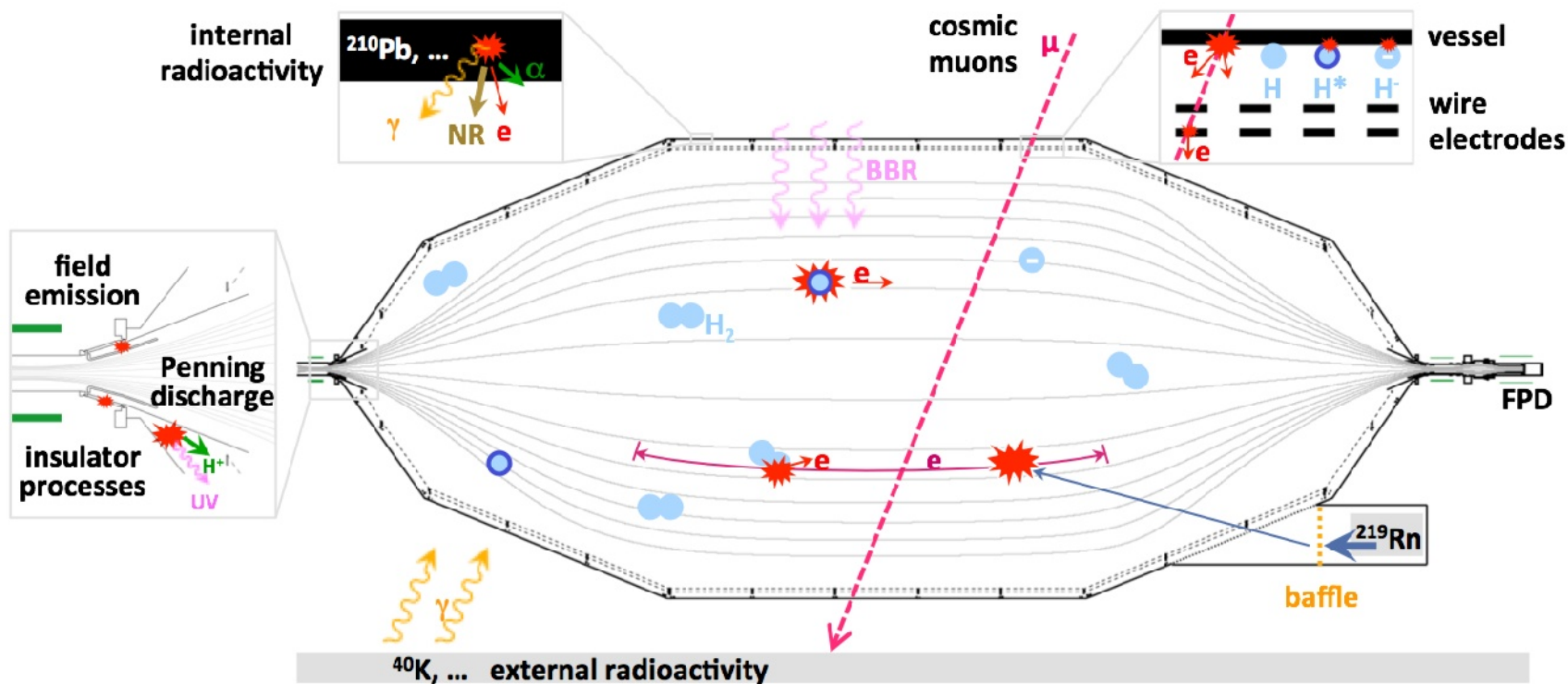


# Commissioning of main spectrometer ( $\Delta E = 0.93$ eV) and detector

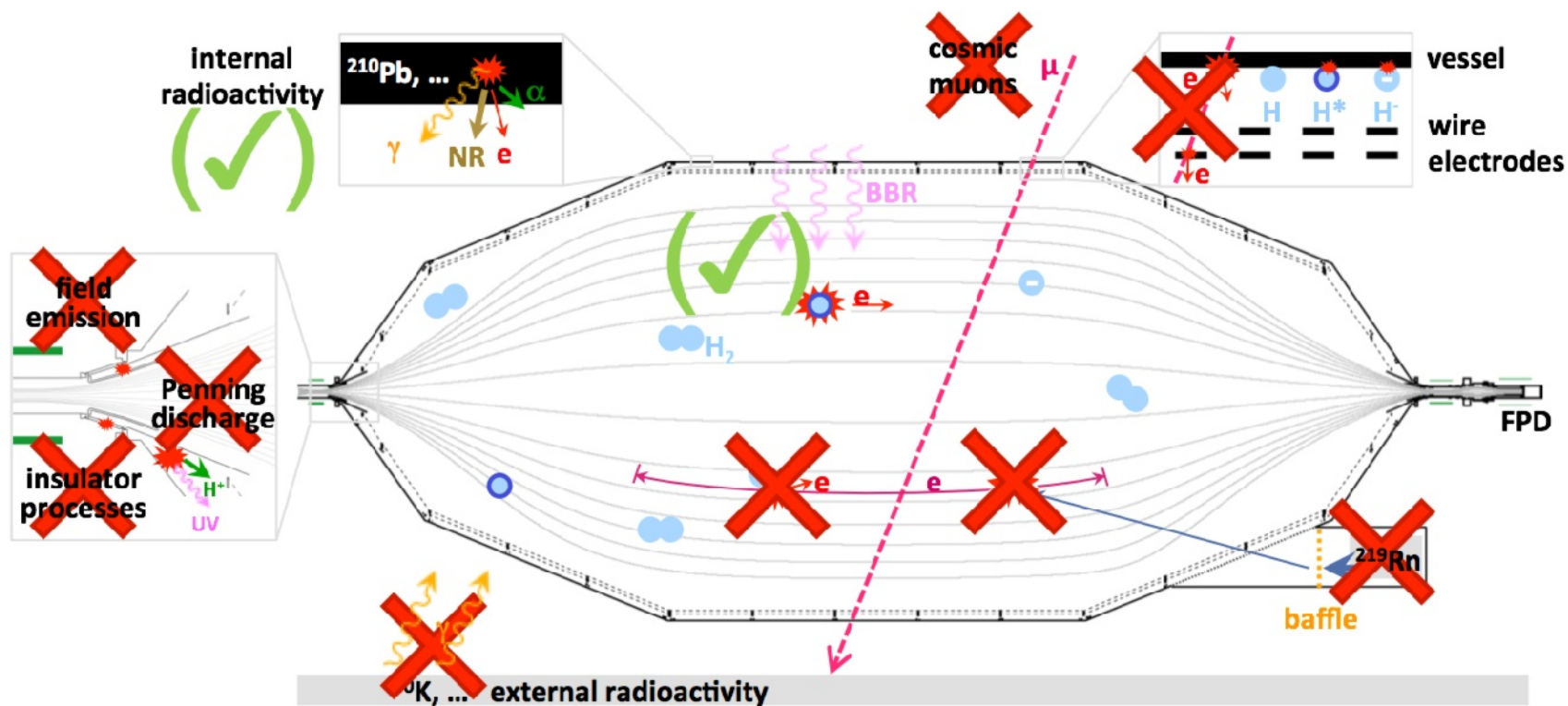




# Possible background sources at KATRIN



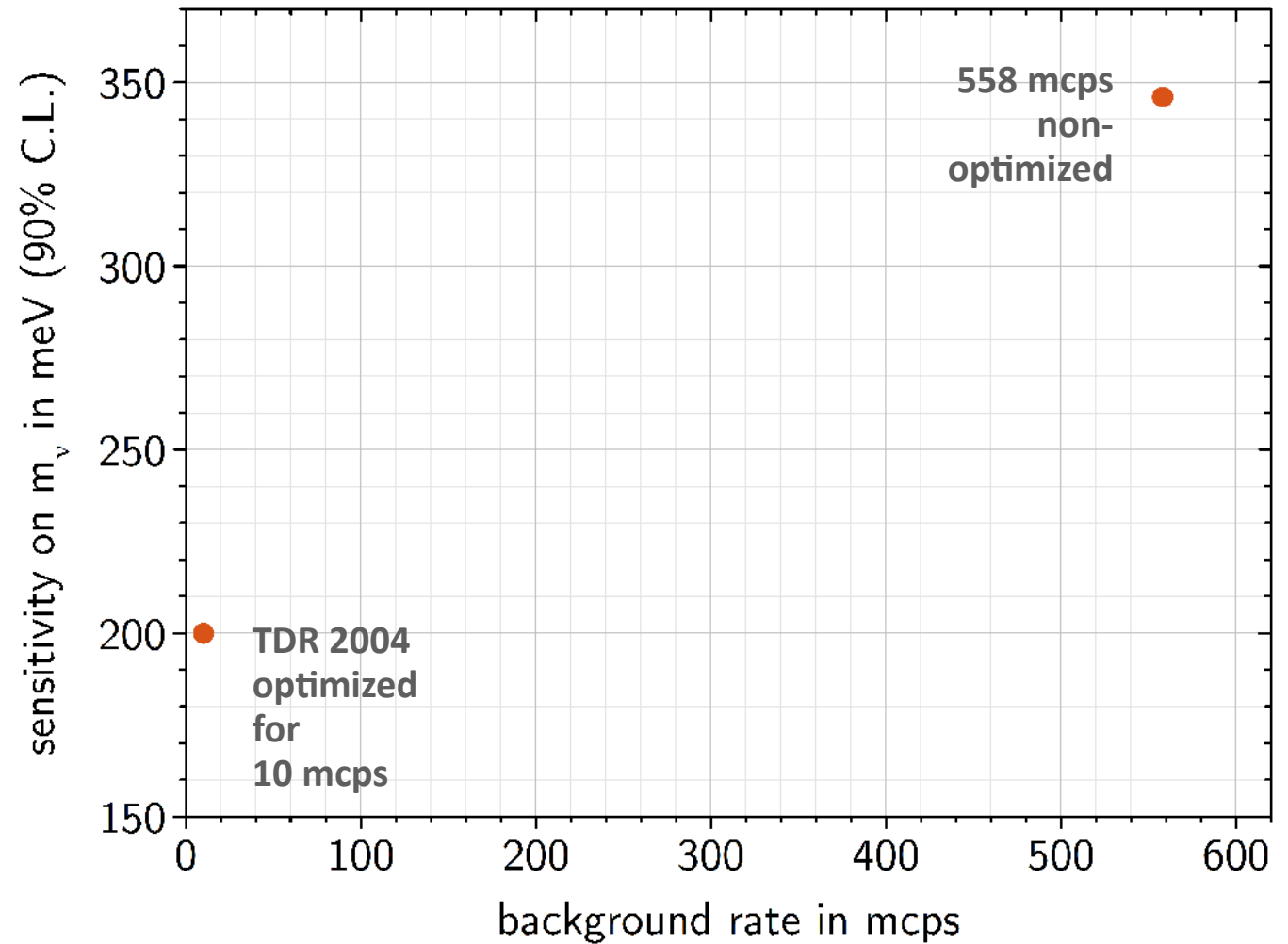
# Background sources at KATRIN: detailed understanding, but ...



- 8 sources of background investigated and understood
- 7 out of 8 avoided or actively eliminated by
  - fine-shaping of special electrodes
  - symmetric magnetic fields
  - LN<sub>2</sub>-cooled baffles (cold traps)
  - wire electrode grids

- 1 out of 8 remaining:  
caused by  $^{210}\text{Pb}$  on spectrometer walls (neutral H\* atoms ionised by black-body radiation in spectrometer)

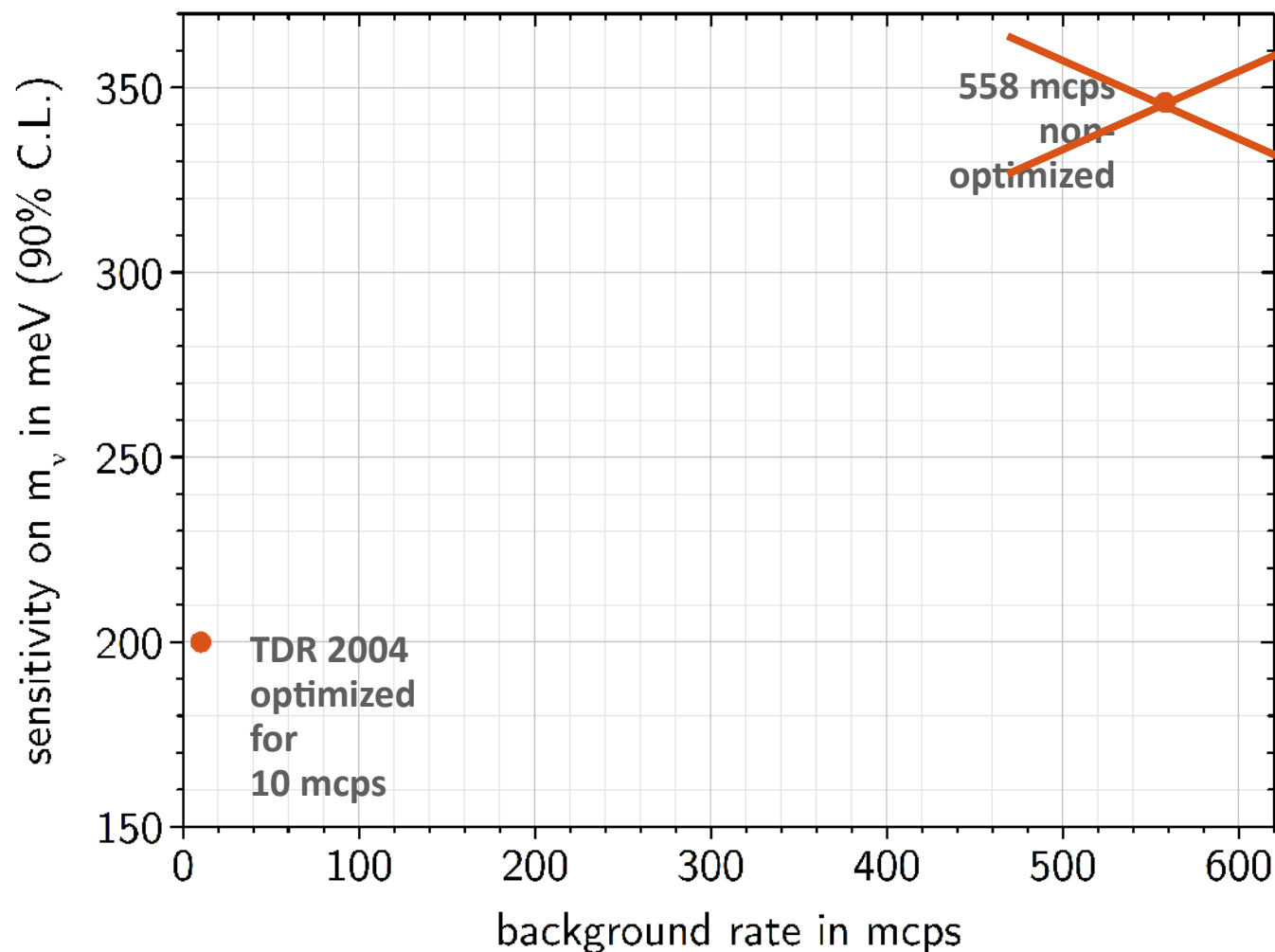
# Sensitivity & background rate





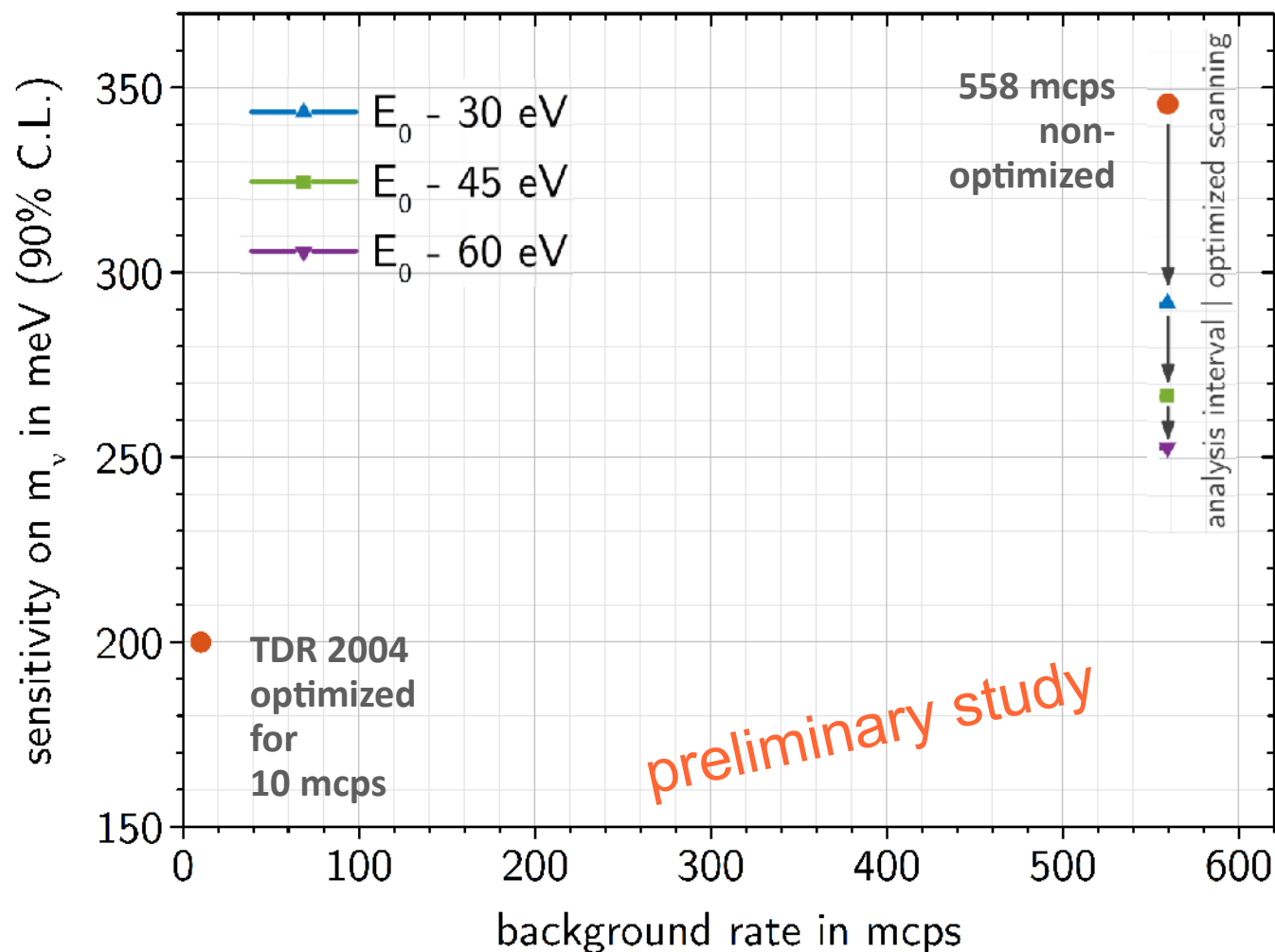
# Sensitivity & background rate

- Further **background reduction measures** being studied
- In addition: **several mitigation strategies** currently under investigation:



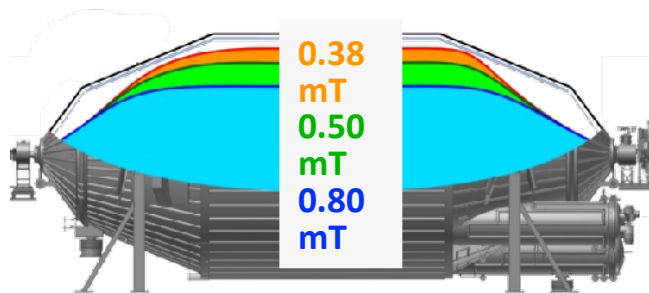
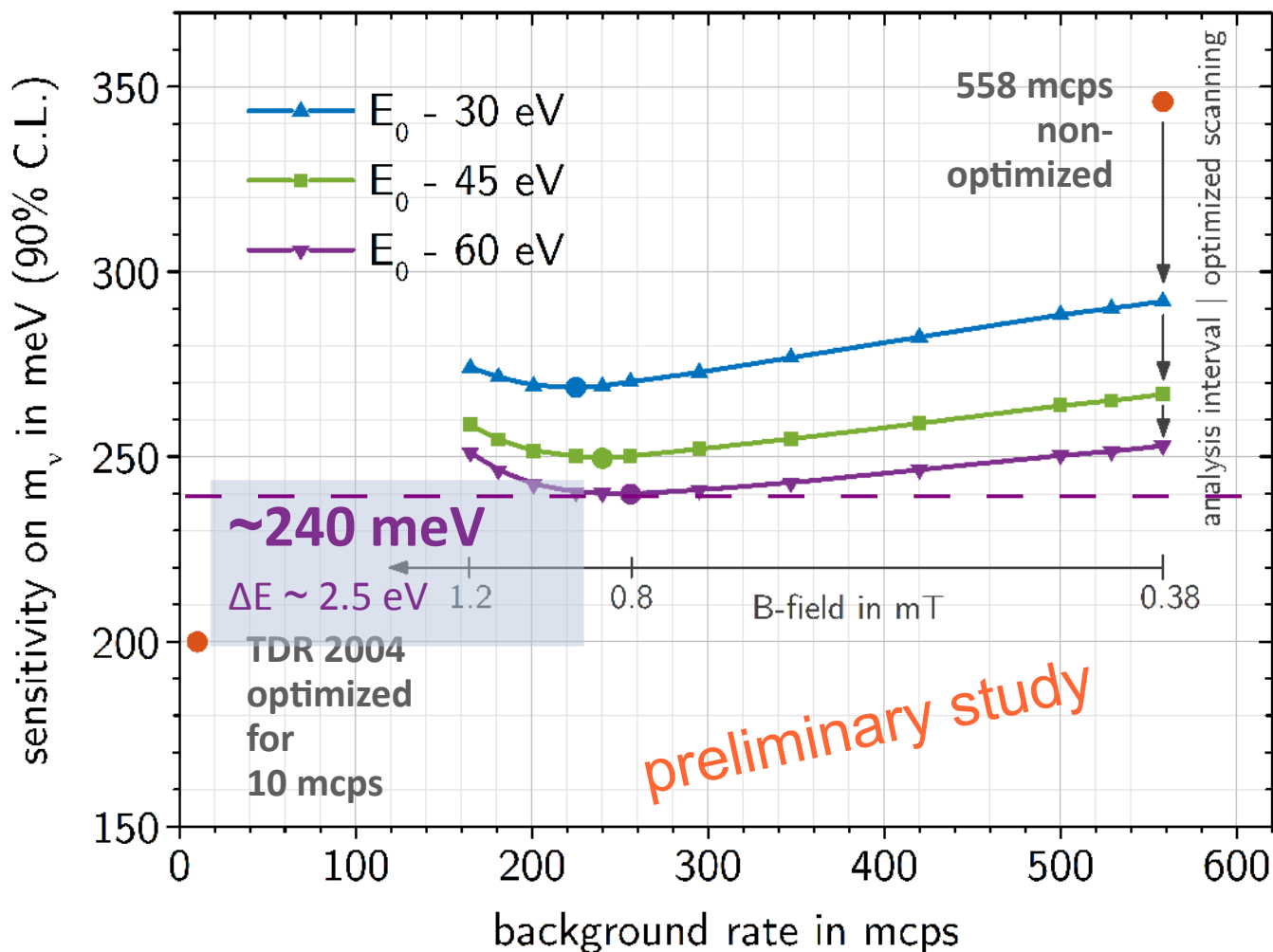
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  - optimized scanning
  - range of spectral analysis



# Sensitivity & background rate

- Further **background reduction measures** being studied
- In addition: **several mitigation strategies** currently under investigation:
  - optimized scanning
  - range of spectral analysis
  - flux tube compression by increasing B





2016: Commissioning of pre and main spectrometer with detector SDS 3  
Continue commissioning of KATRIN Source and Transport Section  
Sending electrons through the 70m long beamline

2017: Ramping up Windowless Gaseous Tritium Source:  $D_2$ ,  $D_2(T_2)$ ,  $T_2$   
Commissioning of complete KATRIN system  
First tritium data  
First chance to look for keV sterile neutrinos  
Regular neutrino mass measurements



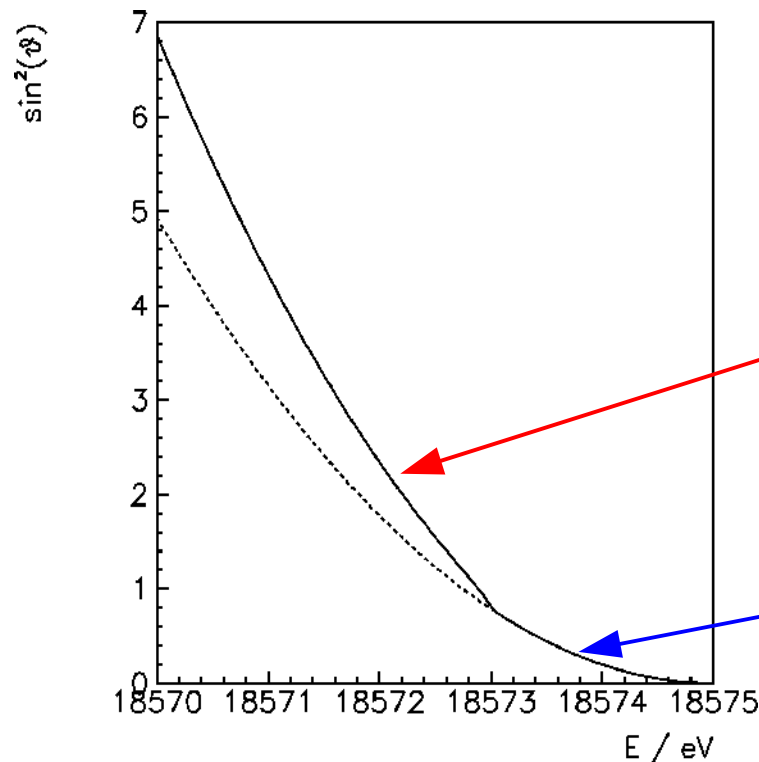
# Influence of a 4<sup>th</sup> sterile neutrino near the endpoint $E_0$

Weak mixing of a  
4<sup>th</sup> sterile neutrino state:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_{\text{sterile}} \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \end{pmatrix}$$

If  $m(\nu_1)$ ,  $m(\nu_2)$ ,  $m(\nu_3)$  are similarly light:

$$dN/dE = K F(E,Z) p E_{\text{tot}} (E_0 - E_e) \left( \cos^2(\theta) \sqrt{(E_0 - E_e)^2 - m(\nu_{1,2,3})^2} + \sin^2(\theta) \sqrt{(E_0 - E_e)^2 - m(\nu_4)^2} \right)$$



e.g.

$$m(\nu_4) = 2 \text{ eV}$$

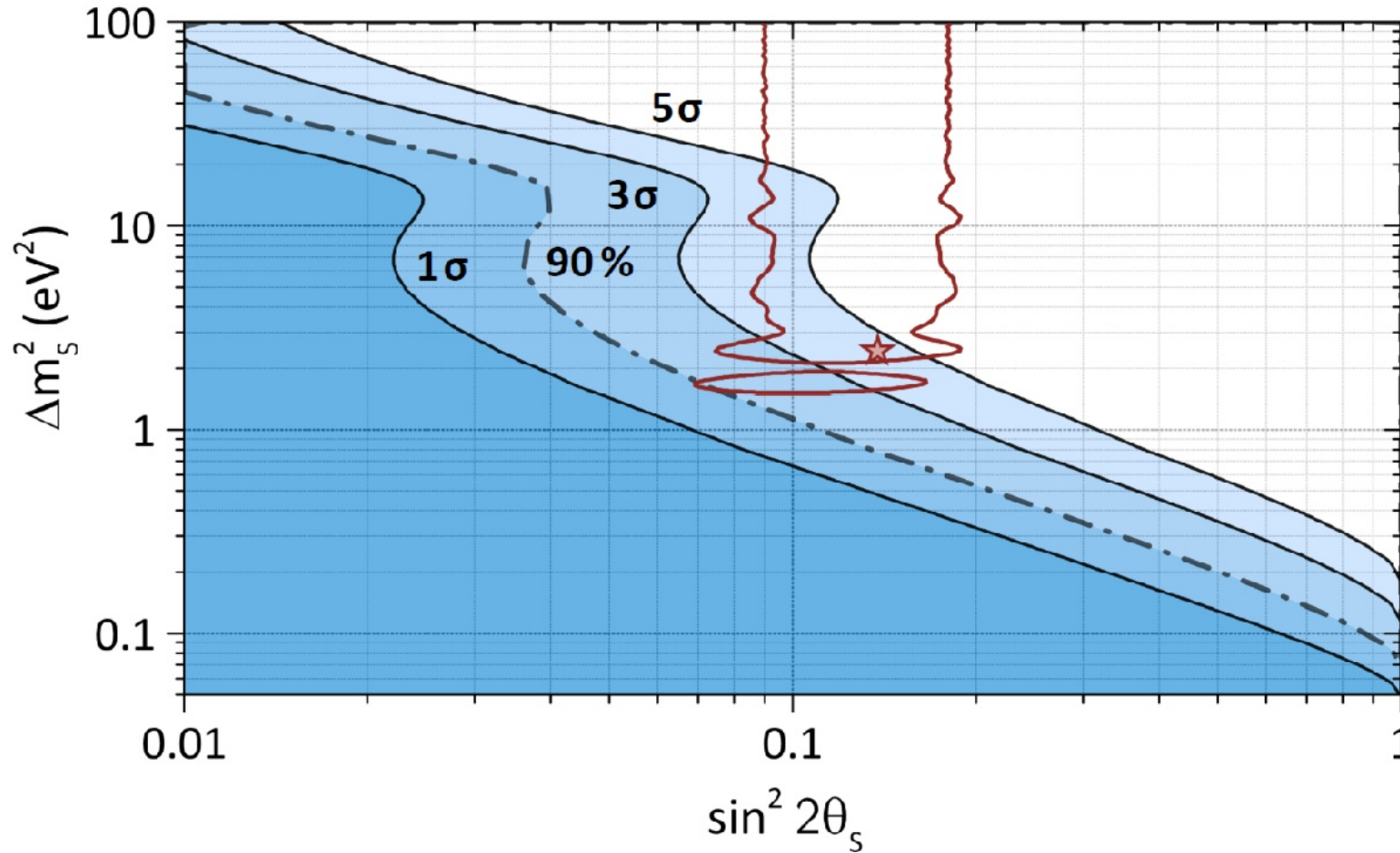
$$\sin^2(\theta) = 0.3$$

e.g.

$$m(\nu_{1,2,3}) \gg 0 \text{ eV}$$

$$\cos^2(\theta) = 0.7$$

# Sensitivity on sterile eV neutrinos



M.Kleesiek,  
PhD thesis,  
KIT (2014)

see also:

J. A. Formaggio, J. Barret, PLB 706 (2011) 68

A. Sejersen Riis, S. Hannestad, JCAP02 (2011) 011

A. Esmaili, O.L.G. Peres, arXiv:1203.2632

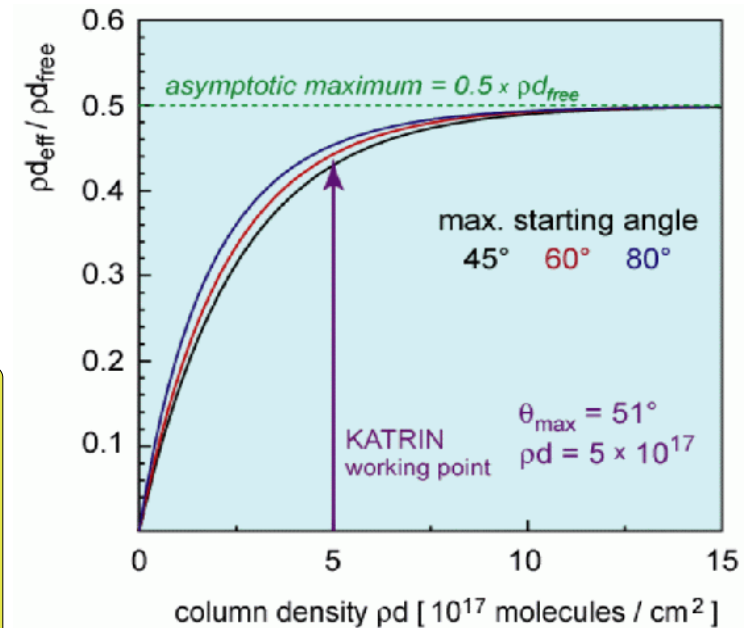
# Can KATRIN be largely improved ?

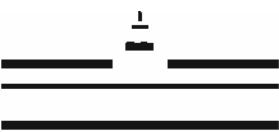
## Problems to be solved

- 1) The source is already opaque  
 → need to increase size transversally  
 magnetic flux tube conservation  
 requests larger spectrometer too  
 but a  $\varnothing 100\text{m}$  spectrometer is not feasible

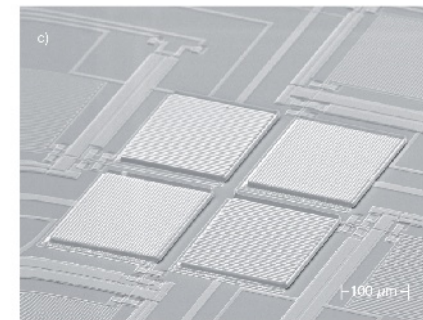
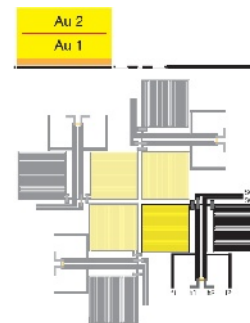
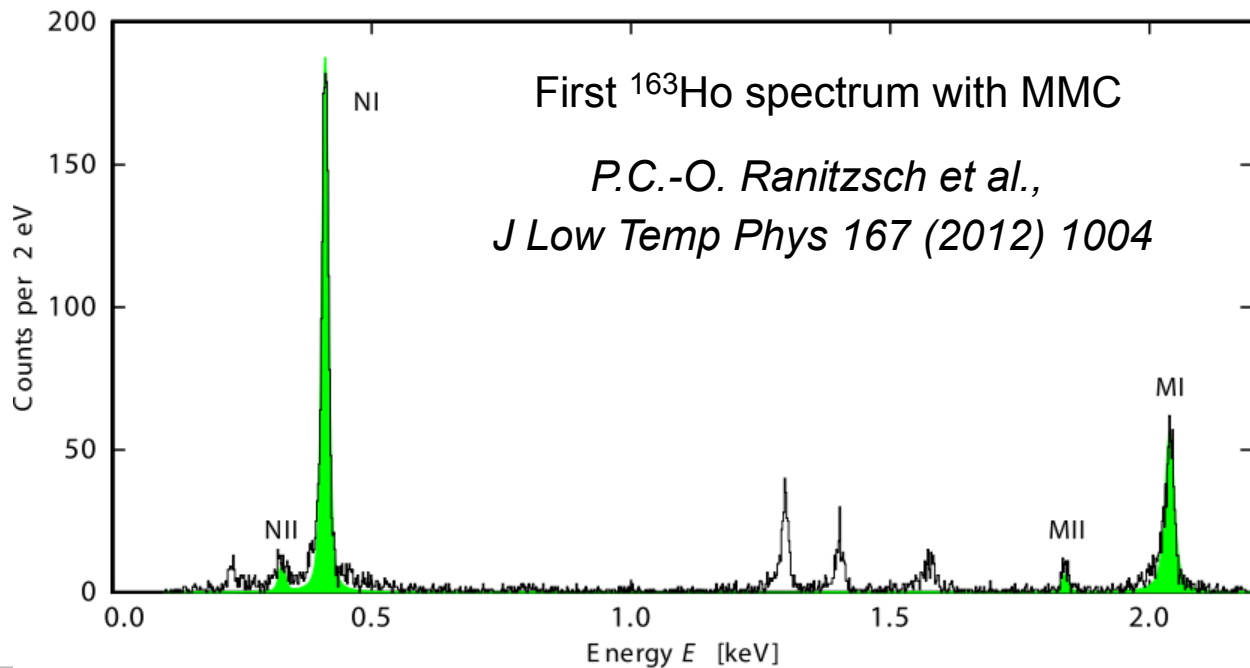
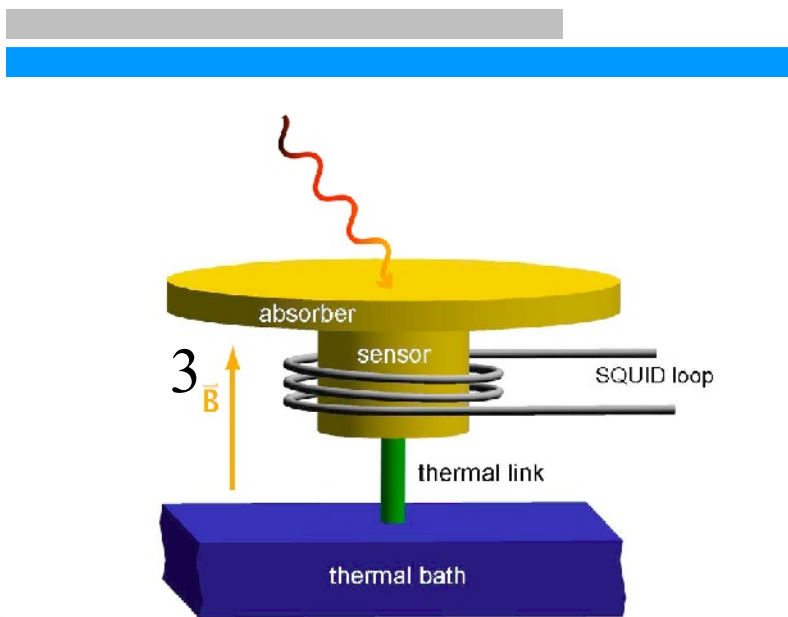
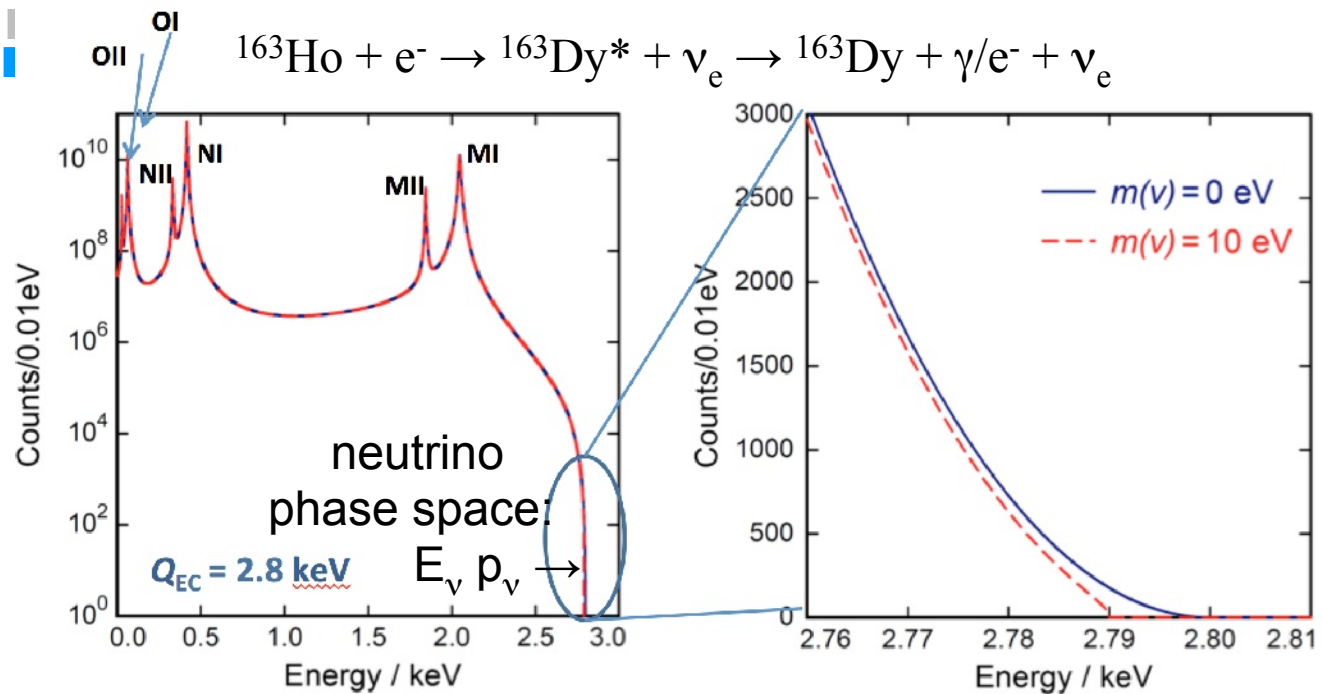
### Possible ways out:

- a) source inside detector (compare to  $0\nu\beta\beta$ )  
 using cryogenic bolometers (ECHO, HOLMES, ..)





# ECHO neutrino mass project: $^{163}\text{Ho}$ electron capture with metallic magnetic calorimeters



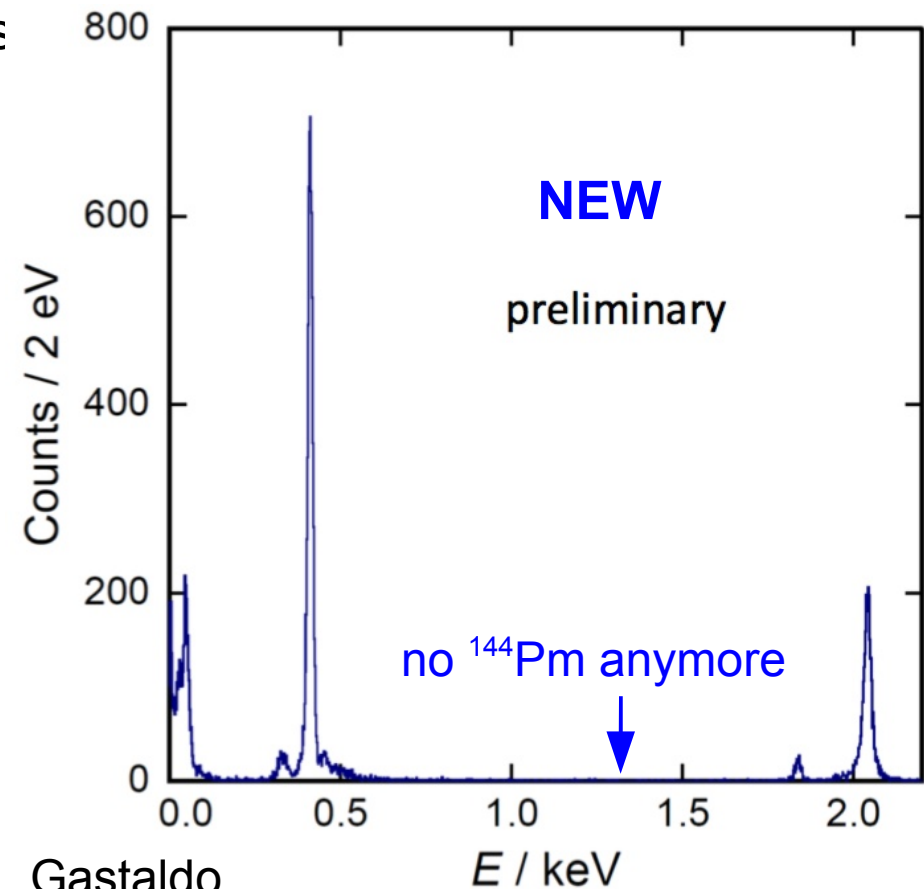
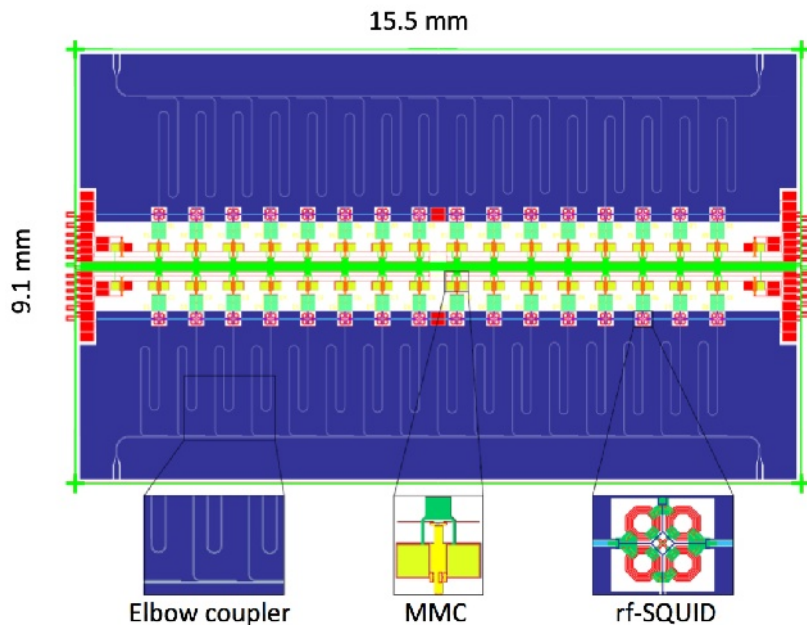
courtesy L. Gastaldo



# ECHO neutrino mass project: $^{163}\text{Ho}$ electron capture with metallic magnetic calorimeters

Recent achievements by ECHO:

- new Q-value: 2.8 keV (independently by MMC & Penning trap, was 2.5 keV before!)
- new source production: chemical purification + mass separation  $\rightarrow$  no  $^{144}\text{Pm}$  or  $^{166\text{m}}\text{Ho}$
- very good energy resolution of this technology ( $\Delta E_{\text{FWHM}} = 1.6 \text{ eV}$  at 6 keV)
- ultra-short response (pile-up!): risetime 90 ns
- 128 pixels: microwave SQUID multiplexing
- funding for ECHO-1k



courtesy L. Gastaldo

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**$^{163}\text{Ho}$  EC is investigated by 3 collaborations:  
ECHO, HOLMES, NuMECS**

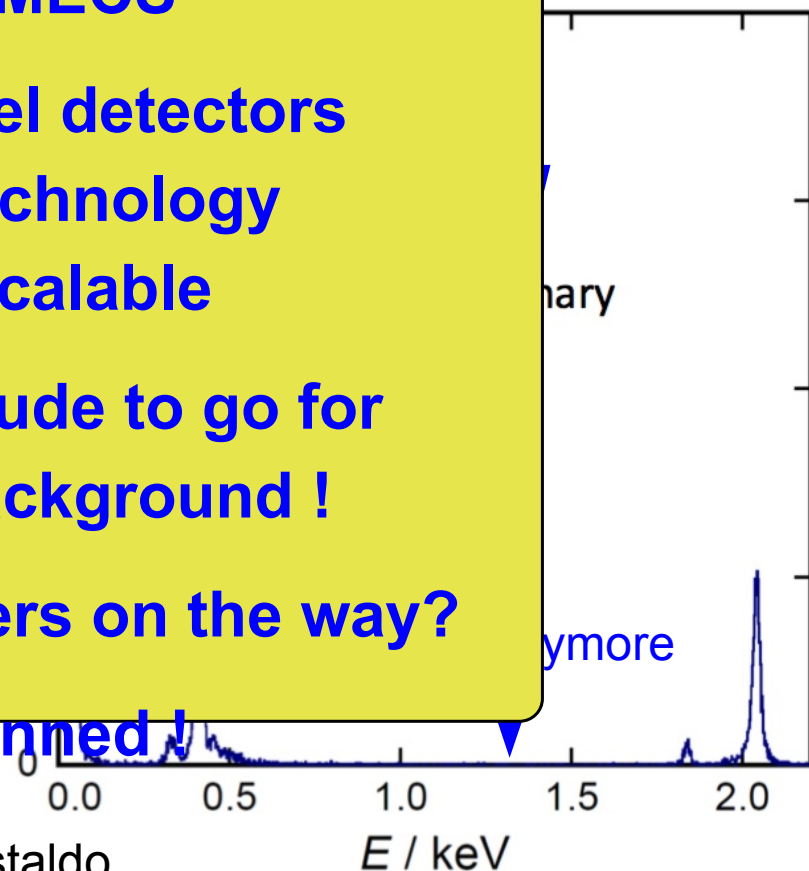
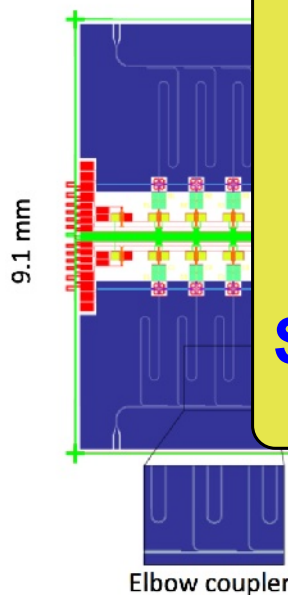
**Cryo-calorimetric multipixel detectors  
are a very interesting technology**

**$\rightarrow$  starts to become scalable**

**But many orders of magnitude to go for  
required statistics and background !**

**Systematics and show stoppers on the way?**

**$\rightarrow$  We should stay tuned !**



courtesy L. Gastaldo

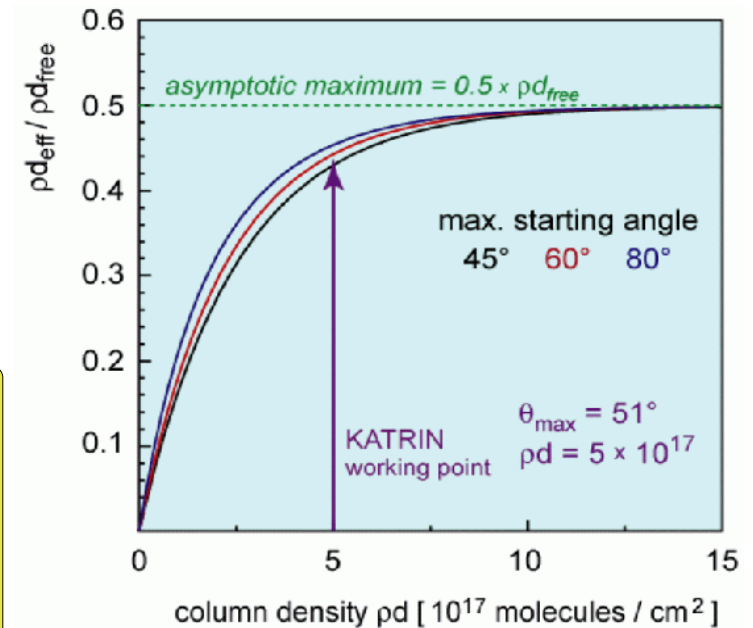
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### Possible ways out:

- a) source inside detector (compare to  $0\nu\beta\beta$ )  
 using cryogenic bolometers (ECHO, HOLMES, ..)
- b) hand-over energy information of  $\beta$  electron  
 to other particle (radio photon),  
 which can escape tritium source (Project 8)



# Project 8's goal: Measure coherent cyclotron radiation of tritium $\beta$ electrons

## General idea:

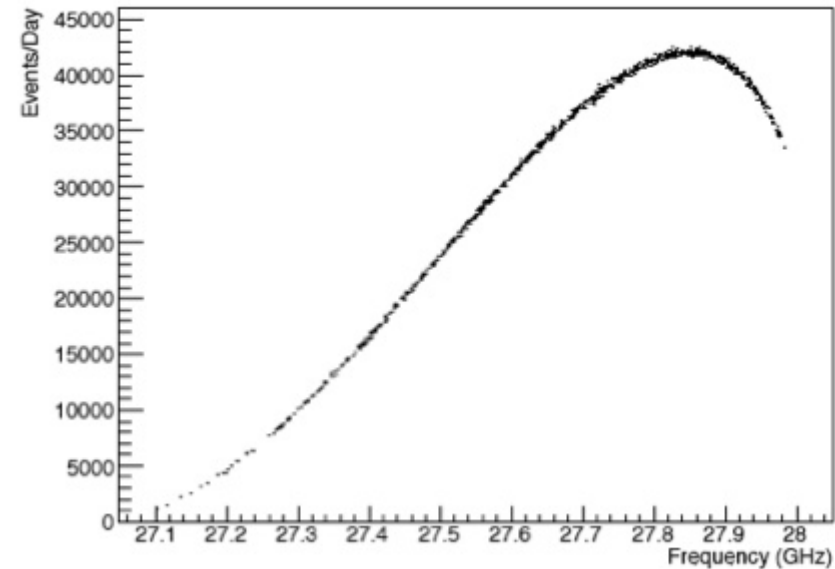
*B. Monreal and J. Formaggio, PRD 80 (2009) 051301*

- Source = KATRIN tritium source technology :

uniform B field + low pressure  $T_2$  gas

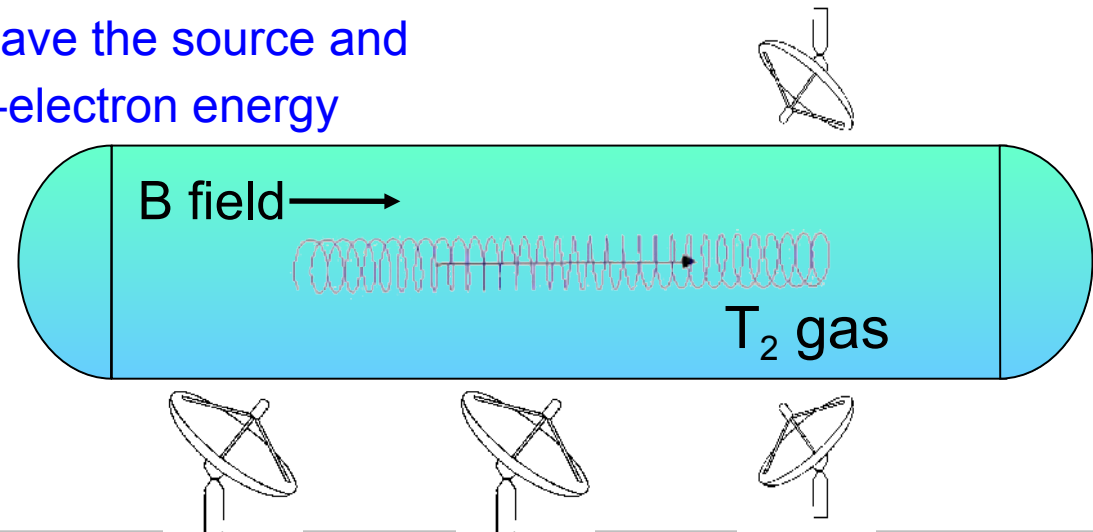
$\beta$  electron radiates coherent cyclotron radiation

$$\omega(\gamma) = \frac{\omega_0}{\gamma} = \frac{eB}{K + m_e}$$



- Antenna array (interferometry) for cyclotron radiation detection

since cyclotron radiation can leave the source and carries the information of the  $\beta$ -electron energy

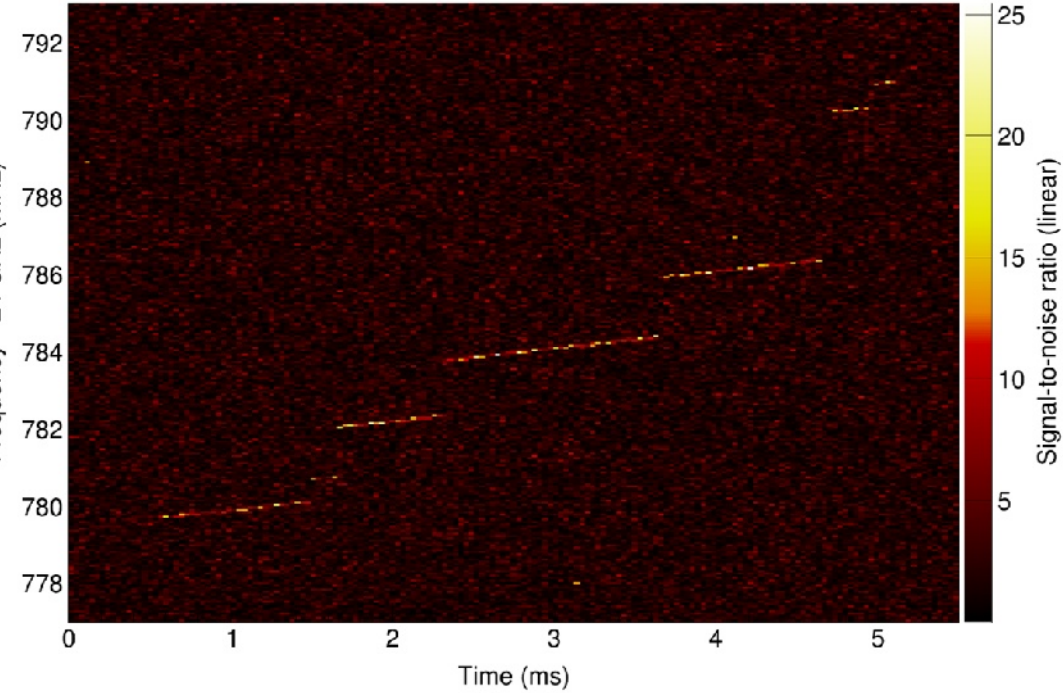
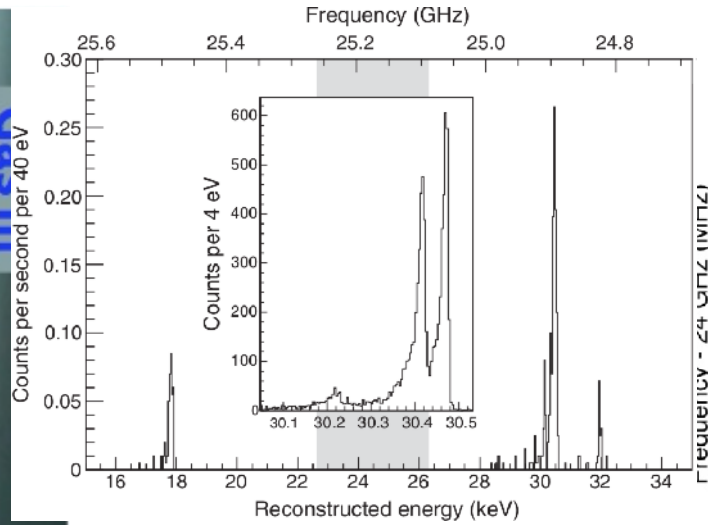
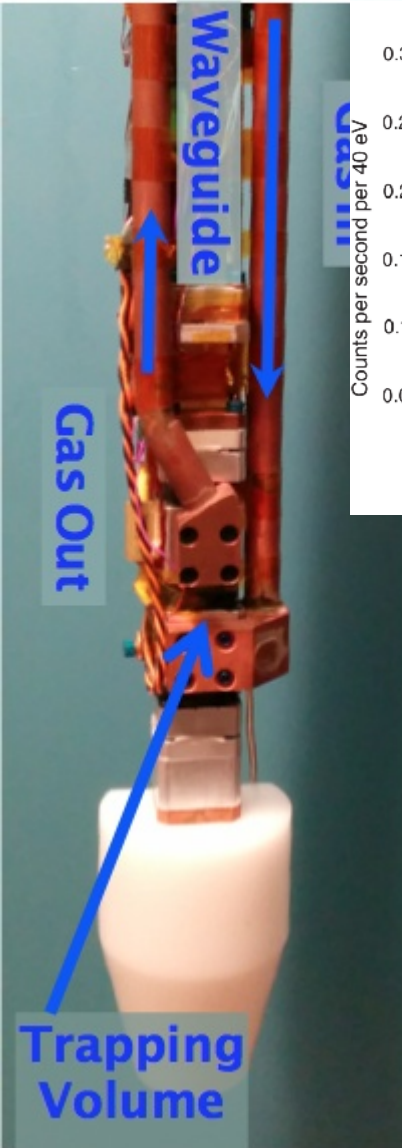






# Project 8's phase 1: detection single electrons from $^{83m}\text{Kr}$

D. M. Asner et al., Phys. Rev. Lett. 114, 162501

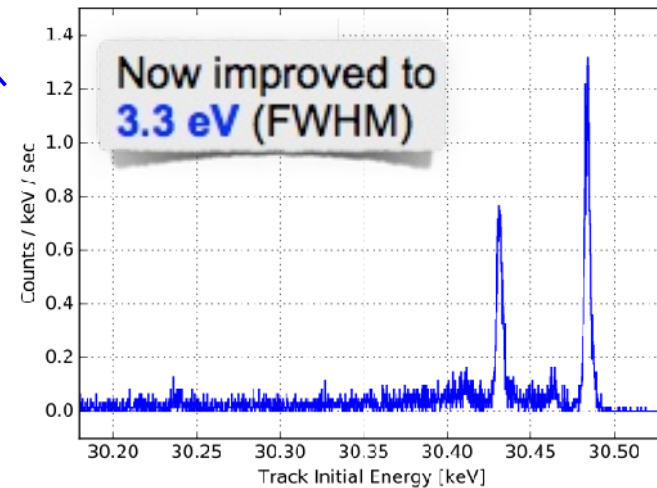
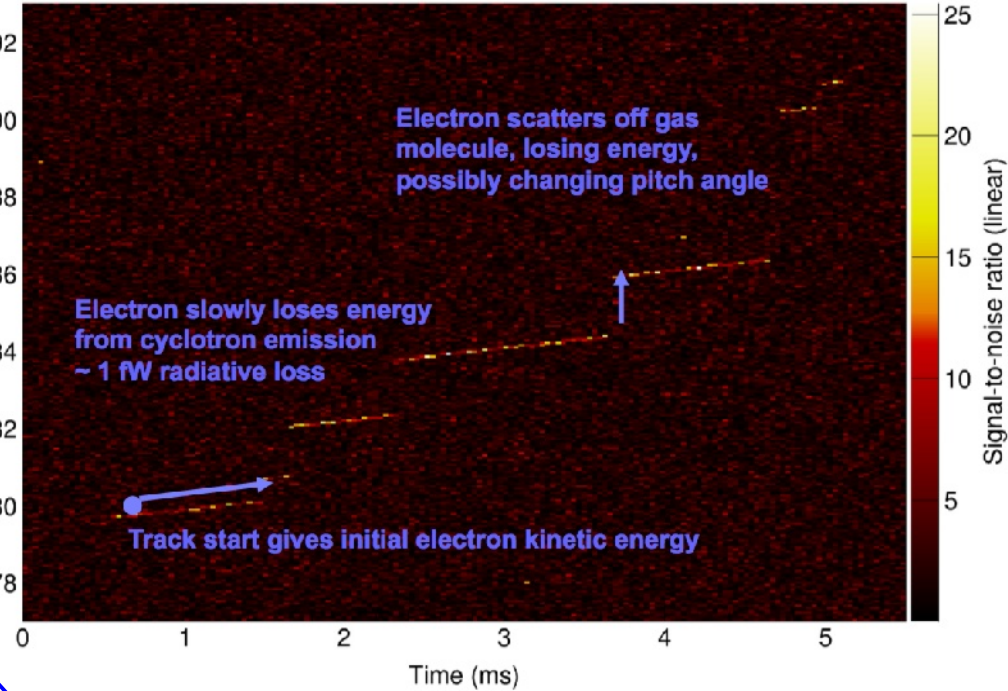
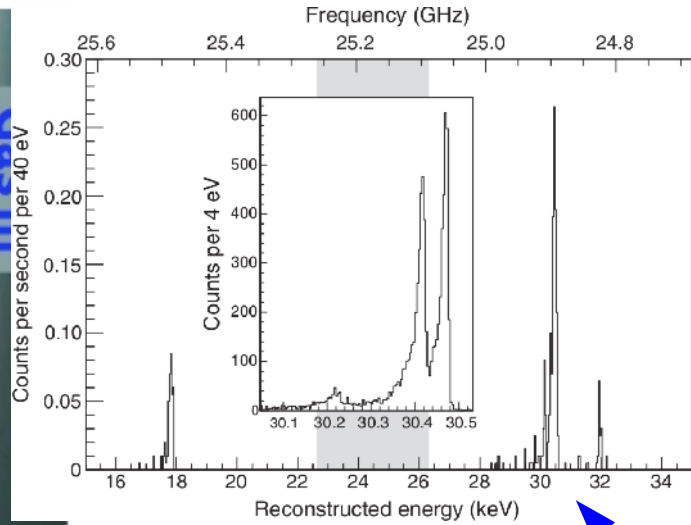
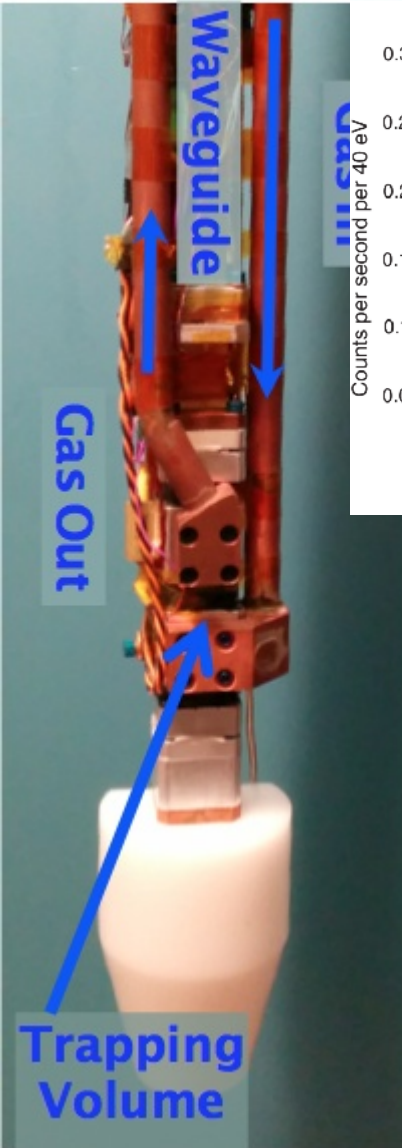


courtesy J. Formaggio, RGH Robertson



# Project 8's phase 1: Detection single electrons from $^{83m}\text{Kr}$

D. M. Asner et al., Phys. Rev. Lett. 114, 162501



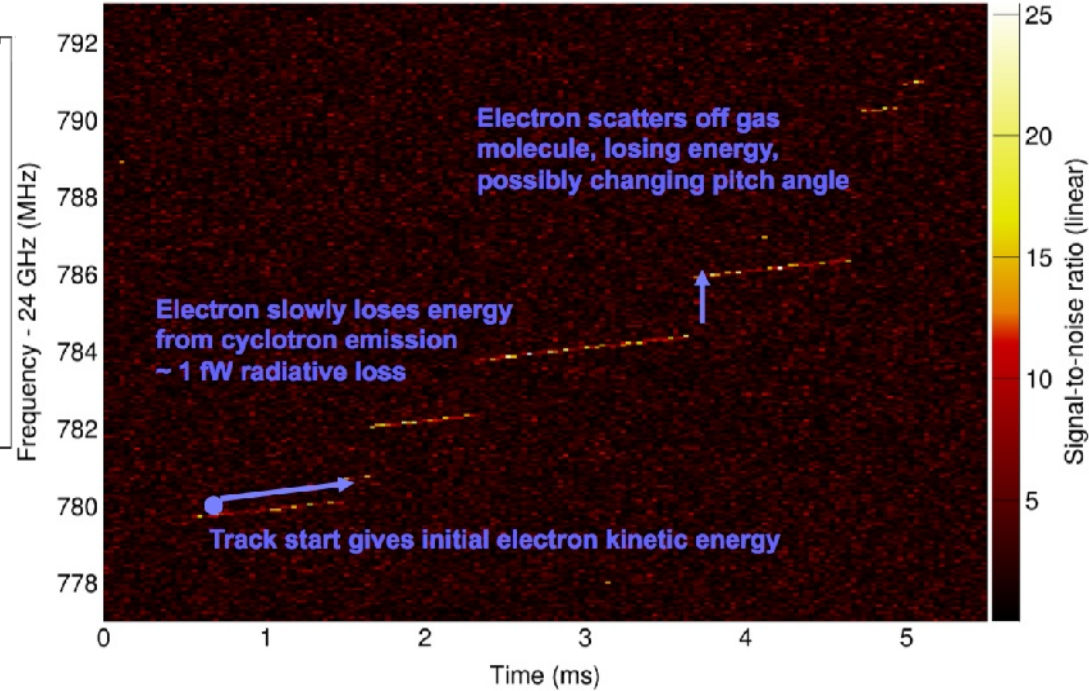
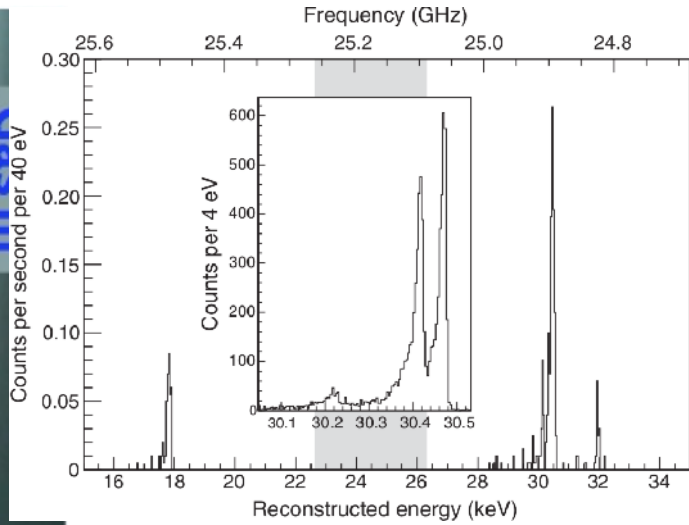
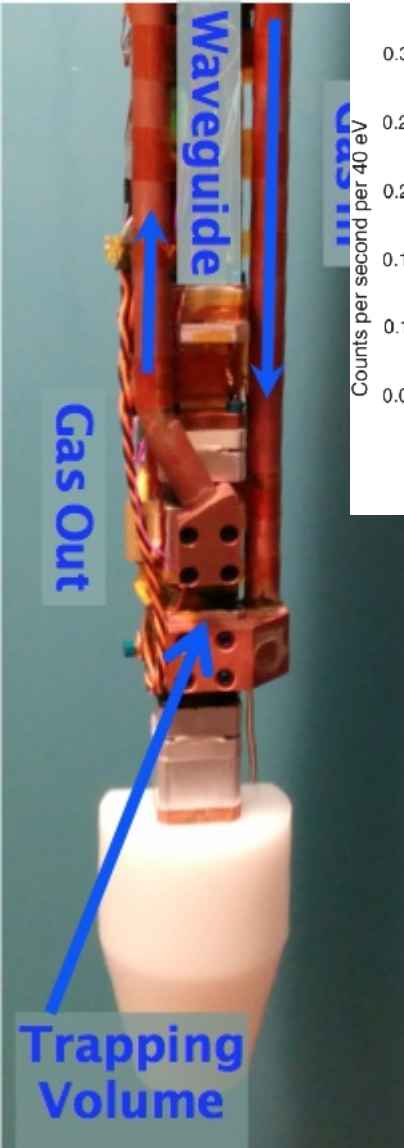
courtesy J. Formaggio, RGH Robertson





# Project 8's phase 1: Detection single electrons from $^{83m}\text{Kr}$

D. M. Asner et al., Phys. Rev. Lett. 114, 162501



**First detection of single electrons successful but still a lot of R&D necessary**

- Is a large scale experiment possible ?
- What are the systematic uncertainties & other limitations?

courtesy J. Formaggio, RGH Robertson

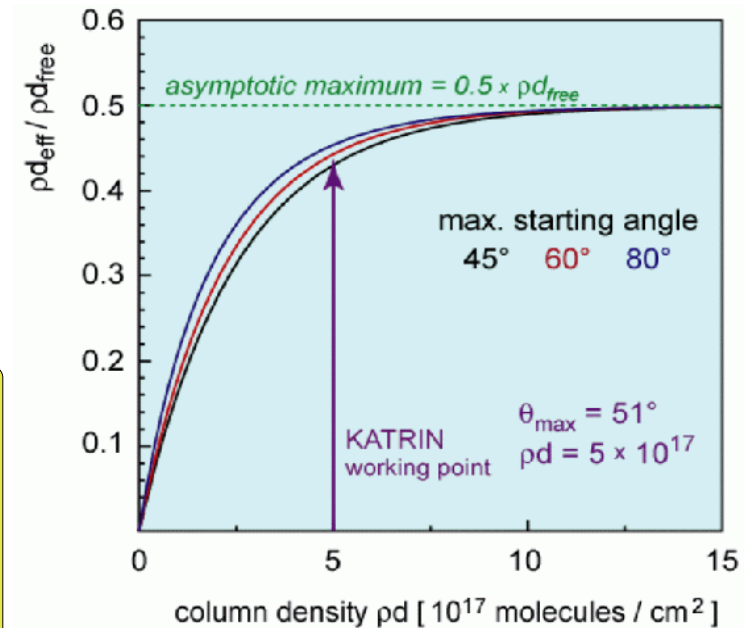
# Can KATRIN be largely improved ?

## Problems to be solved

- 1) The source is already opaque
  - need to increase size transversally
  - magnetic flux tube conservation requests larger spectrometer too
  - but a  $\varnothing 100\text{m}$  spectrometer is not feasible

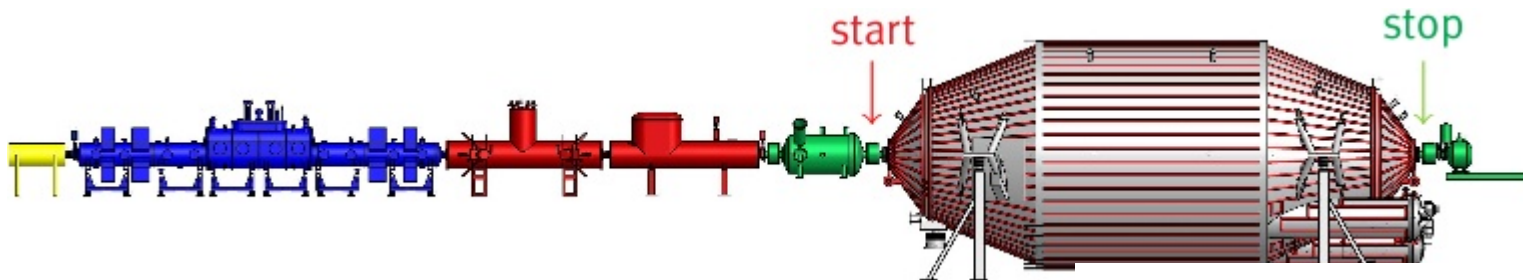
### Possible ways out:

- a) source inside detector (compare to  $0\nu\beta\beta$ )  
using cryogenic bolometers (ECHO, HOLMES, ..)
- b) hand-over energy information of  $\beta$  electron to other particle (radio photon),  
which can escape tritium source (Project 8)
- c) make better use of the electrons  
→ time-of-flight spectroscopy

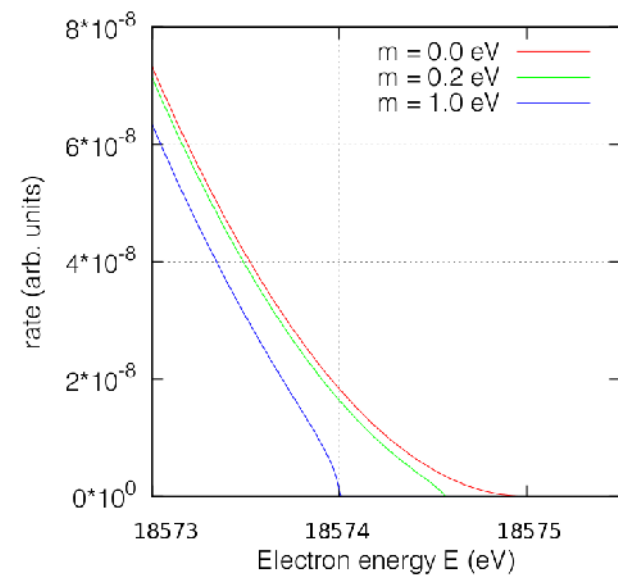
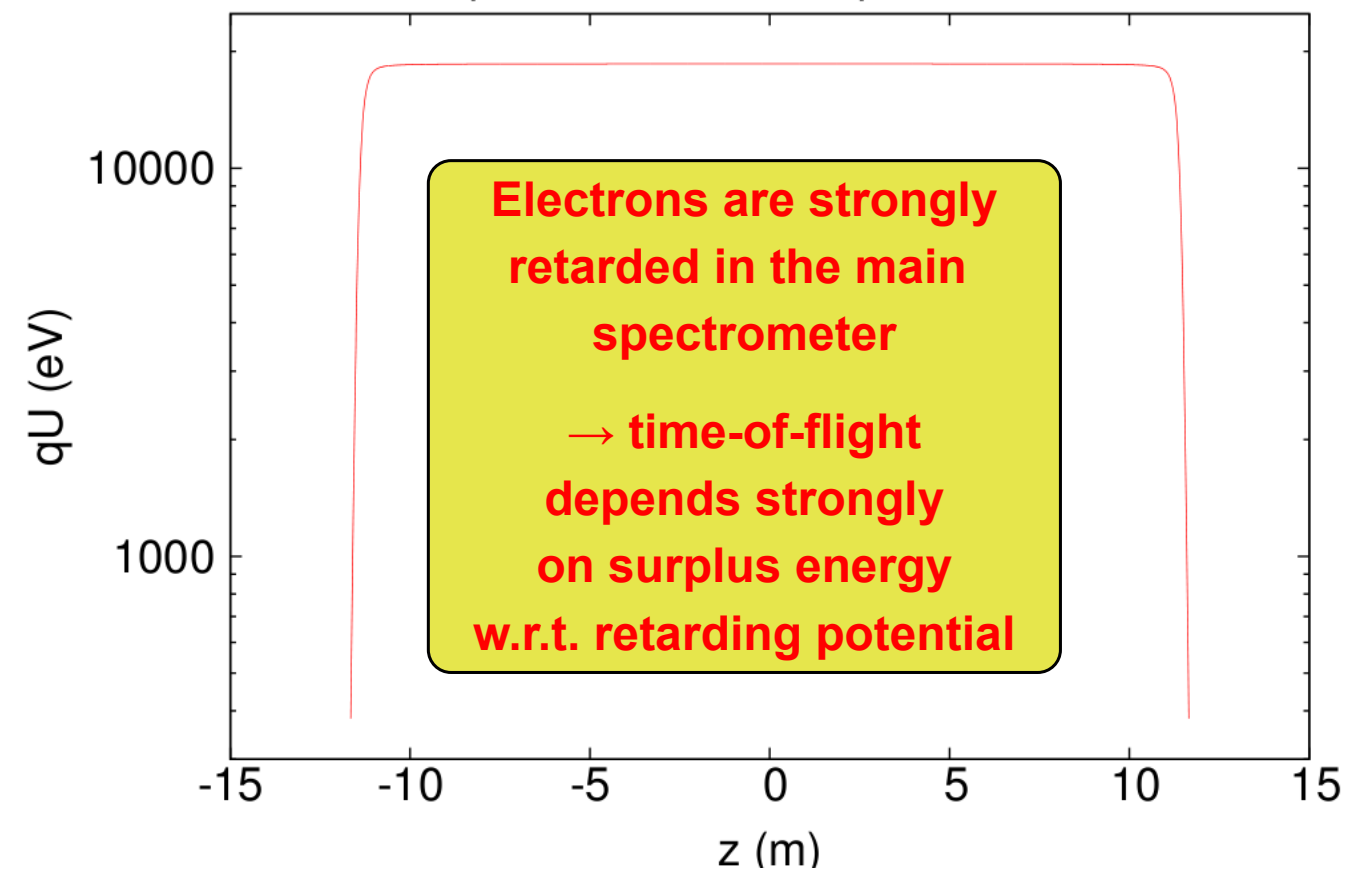




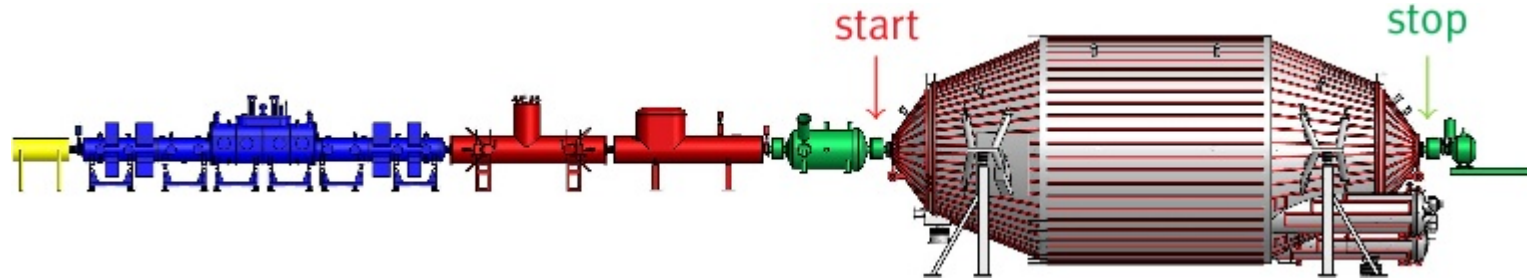
# Alternative spectroscopy: measure time-of-flight TOF through KATRIN spectrometer



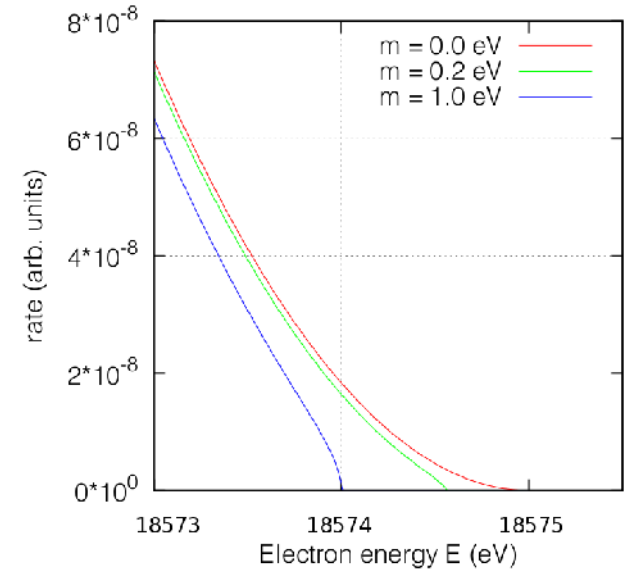
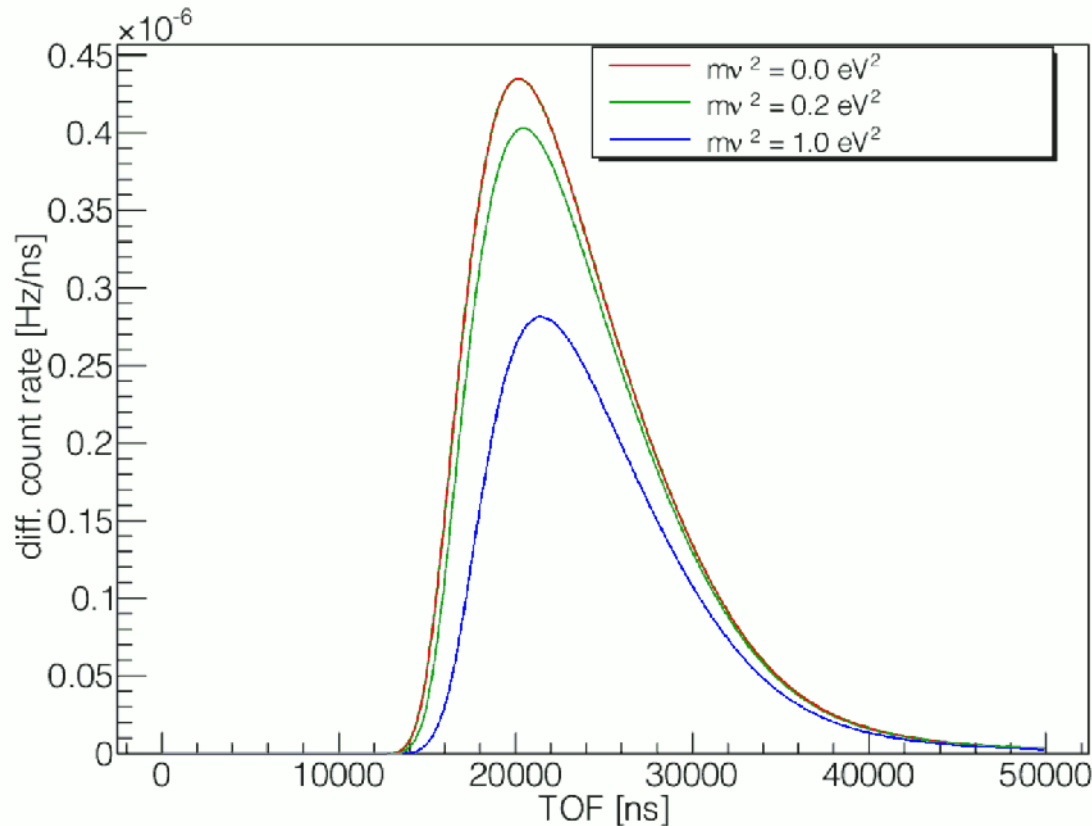
Electric potential on main spectrometer z axis



# Alternative spectroscopy: measure time-of-flight TOF through KATRIN spectrometer



Comparison of TOF spectra for different neutrino masses for  $E_0 = 18574.0$  eV,  $U_{ret} = -18570.0$  eV



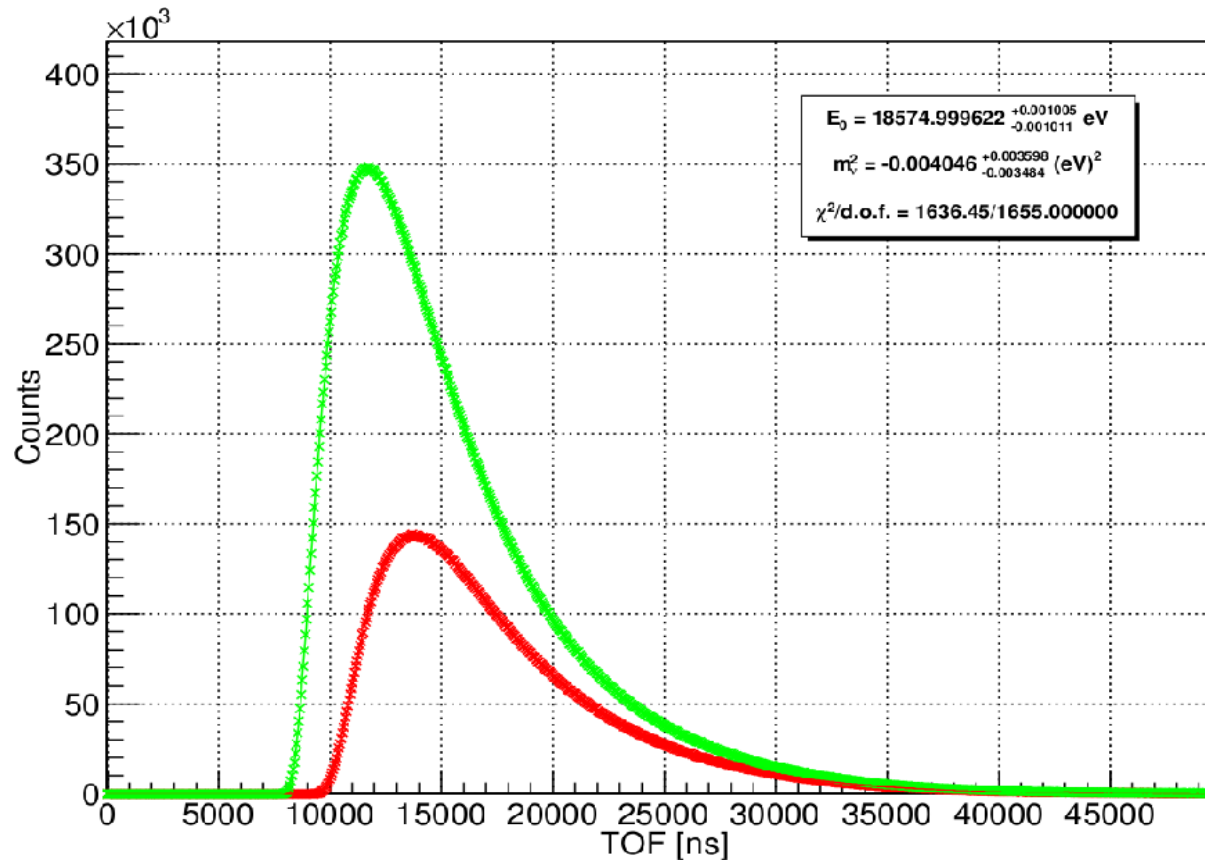
**Time-of-flight spectrum corresponds to a full energy spectrum**

**→ sensitive to the neutrino mass**

# Sensitivity improvement on $m^2(\nu_e)$ by ideal TOF determination

Measure at 2 (instead of  $\approx 30$ ) different retarding potentials  
since TOF spectra contain already all the information

→ Factor 5 improvement in  $m_\nu^2$  w.r.t. standard KATRIN in ideal case !

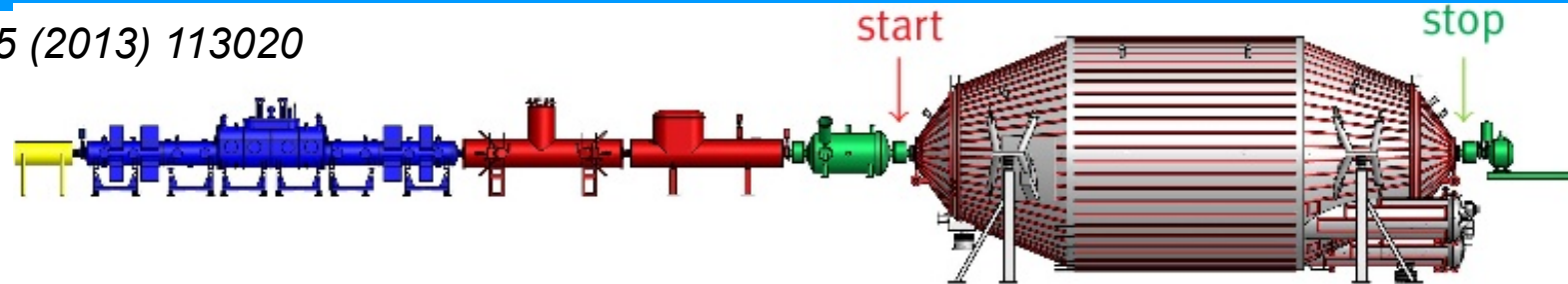


*N. Steinbrink et al.*  
*NJP 15 (2013) 113020*

Coincidence request between start and stop signal → nice background suppression

# How to realize time-of-flight spectroscopy @KATRIN

*N. Steinbrink et al., NJP 15 (2013) 113020*



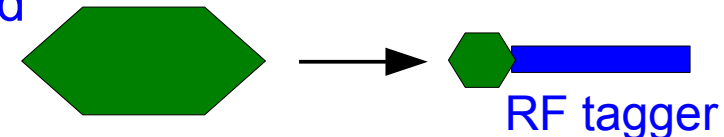
**Advantage: measure  $\beta$ -spectrum by TOF at one (a few) retarding potential(s)  
i.e. as measuring differential  $\beta$ -spectrum**

Stop: Can measure time-of-arrival with KATRIN detector with  $\Delta t = 50$  ns  $\rightarrow$  ok

Start:  **$e^-$ -tagger**: Need to determine time-of-passing-by of  $e^-$  before main spectrometer without disturbing energy and momentum by more than 10 meV:

**$\rightarrow$  factor 5 in  $\Delta m(\nu)_{\text{stat}}^2$  under ideal conditions  
added value: significant background reduction !**

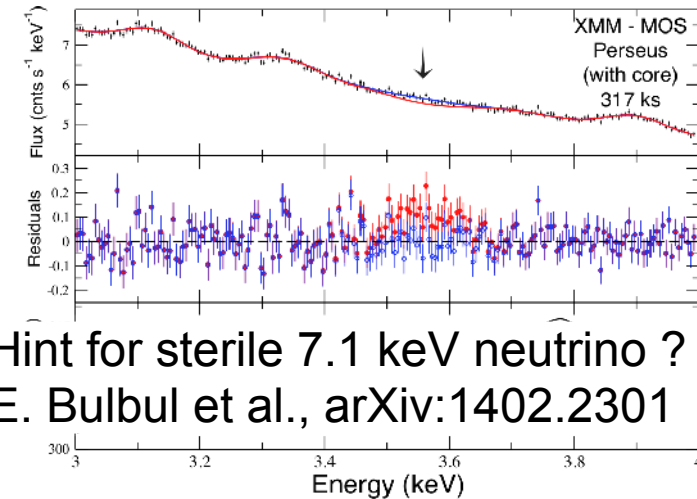
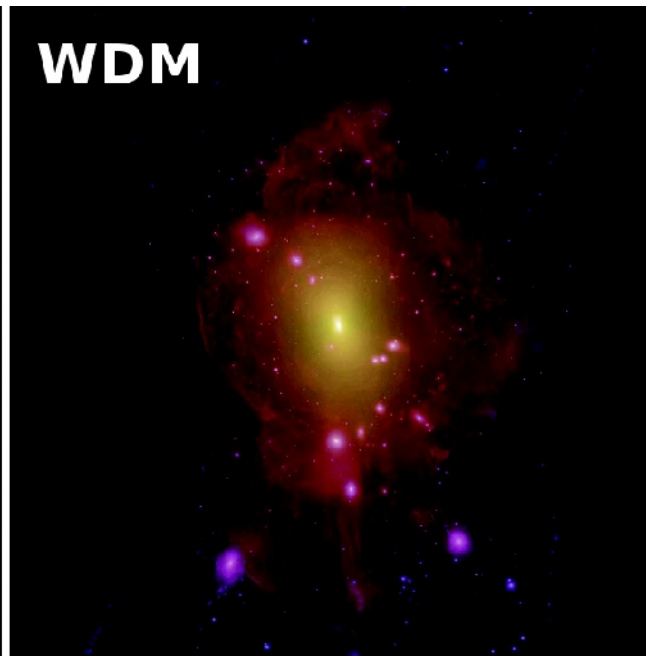
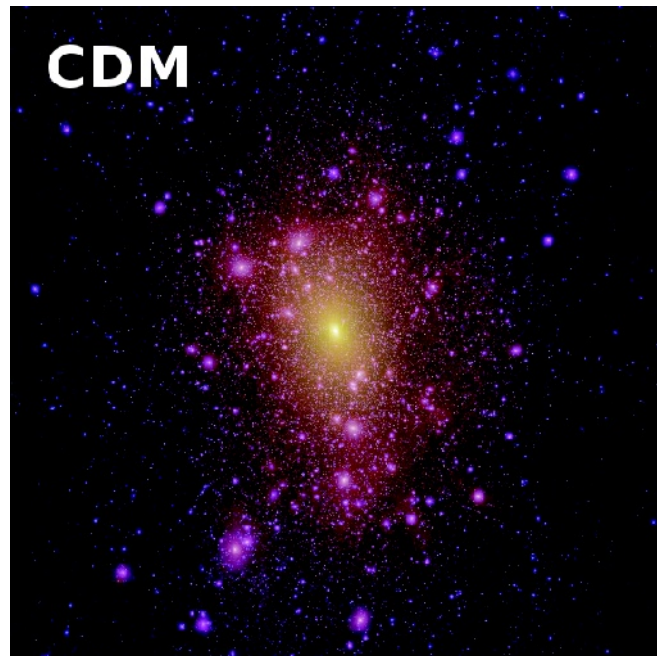
One implementation: reduce **pre spectrometer** length  
& add a **Project 8-type tagger** within a long solenoid  
or another type of electron tagger



or: Use **use tagger-less methods**: „gated filter“ or „time-focusing time-of-flight“

# Hints for a 2<sup>nd</sup> sterile neutrino: Warm Dark Matter in the universe

$\Lambda$ CDM (Cold Dark Matter with cosmological constant) models (masses of about 100 GeV) predict too much structure at galactic scales (too many satellite galaxies)



(e.g. Lovell et al. at Meudon Workshop 2012)

In contrast to observations ! (here only artist view on the right)

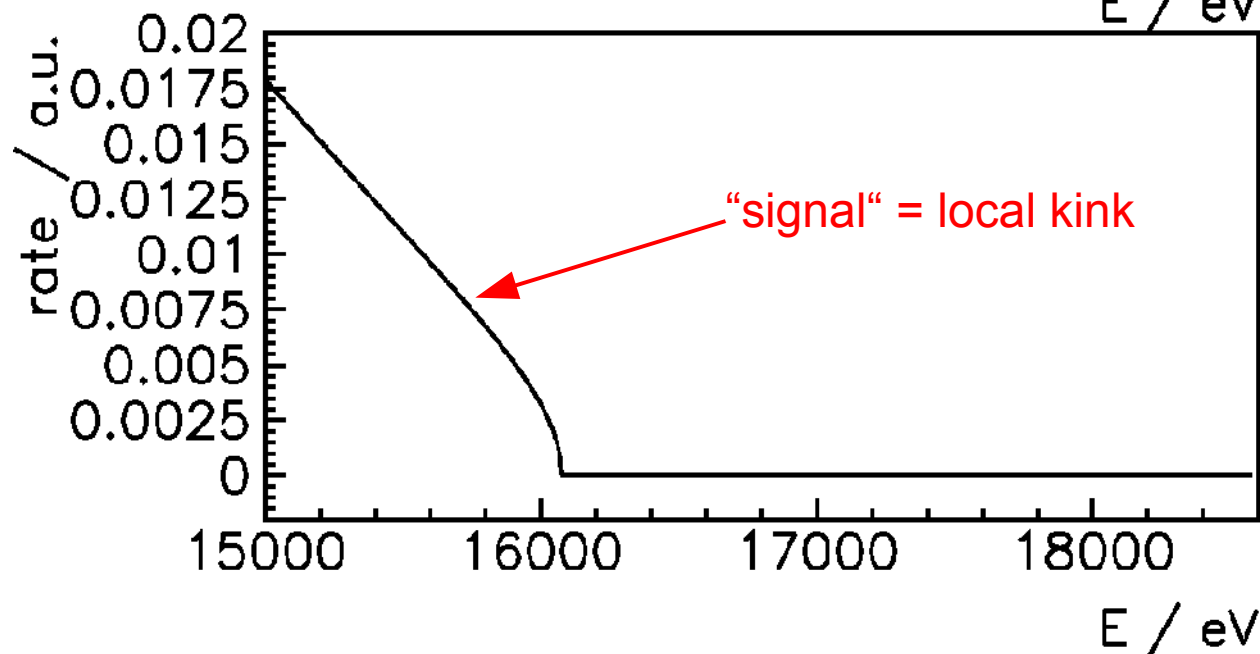
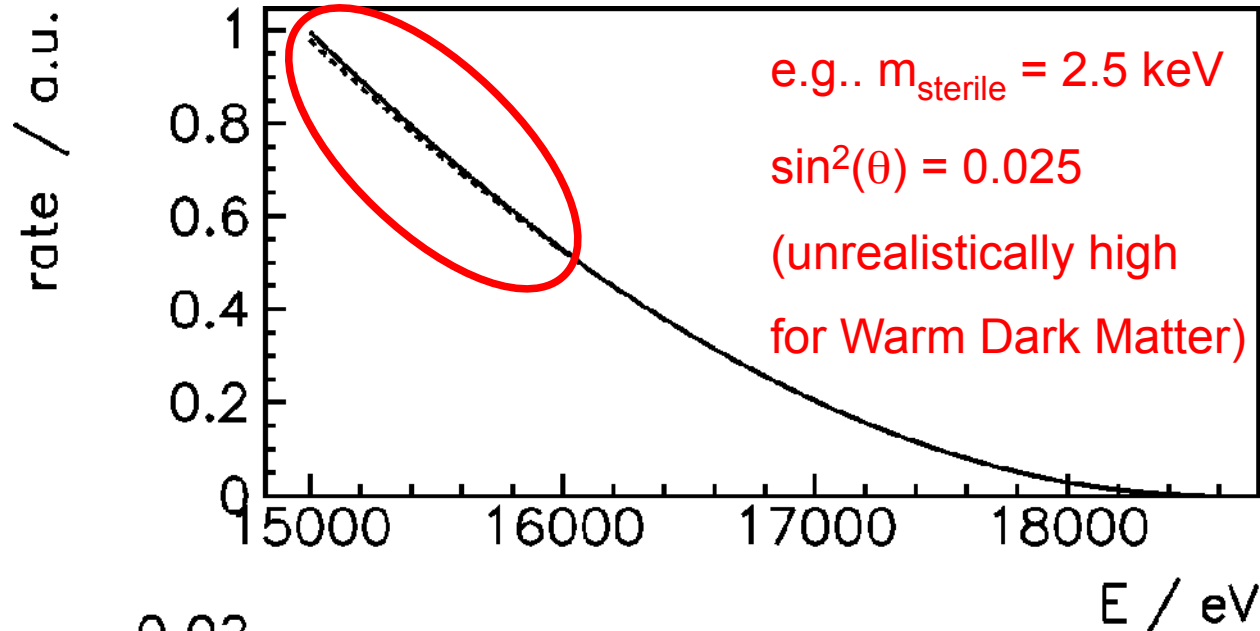
Warm Dark Matter (masses of a few keV, e.g. sterile neutrinos) would smear out these structures



[http://chandra.harvard.edu/graphics/resources/illustrations/milkyWay/milkyway\\_magellanic\\_clouds.jpg](http://chandra.harvard.edu/graphics/resources/illustrations/milkyWay/milkyway_magellanic_clouds.jpg)



# Search for a tiny kink of a keV neutrino

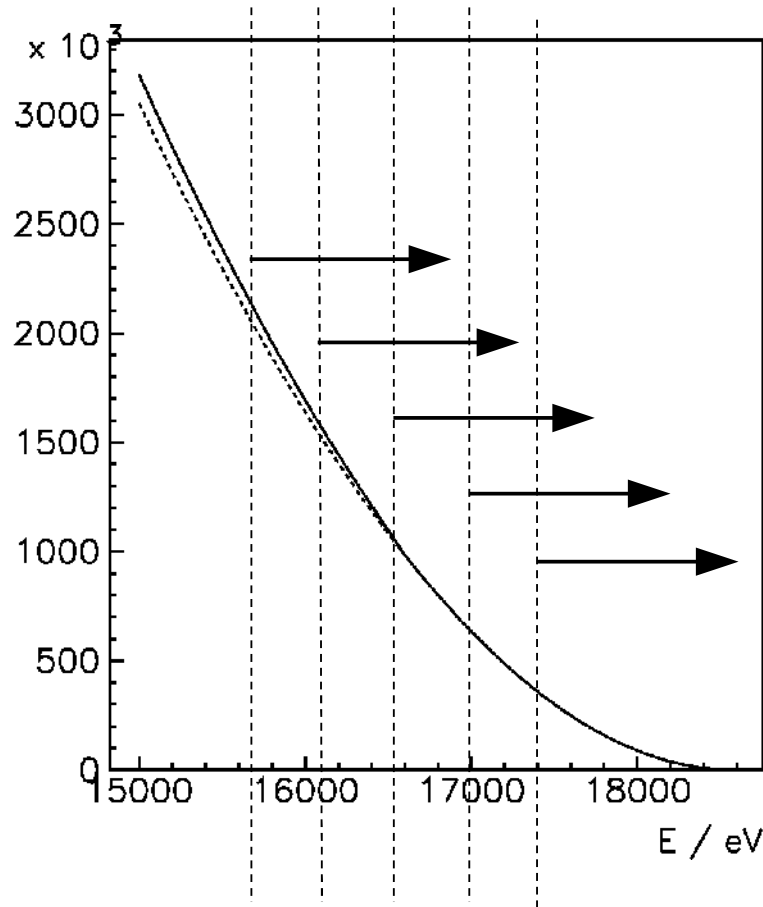


## Main questions:

- 1) How to measure this tiny kink a few keV below the endpoint ?  
 In parallel, in addition or after KATRIN's  $m(\nu_e)$  mission ?
- 2) How to get enough statistics ?
- 3) How to fight against the systematics ?

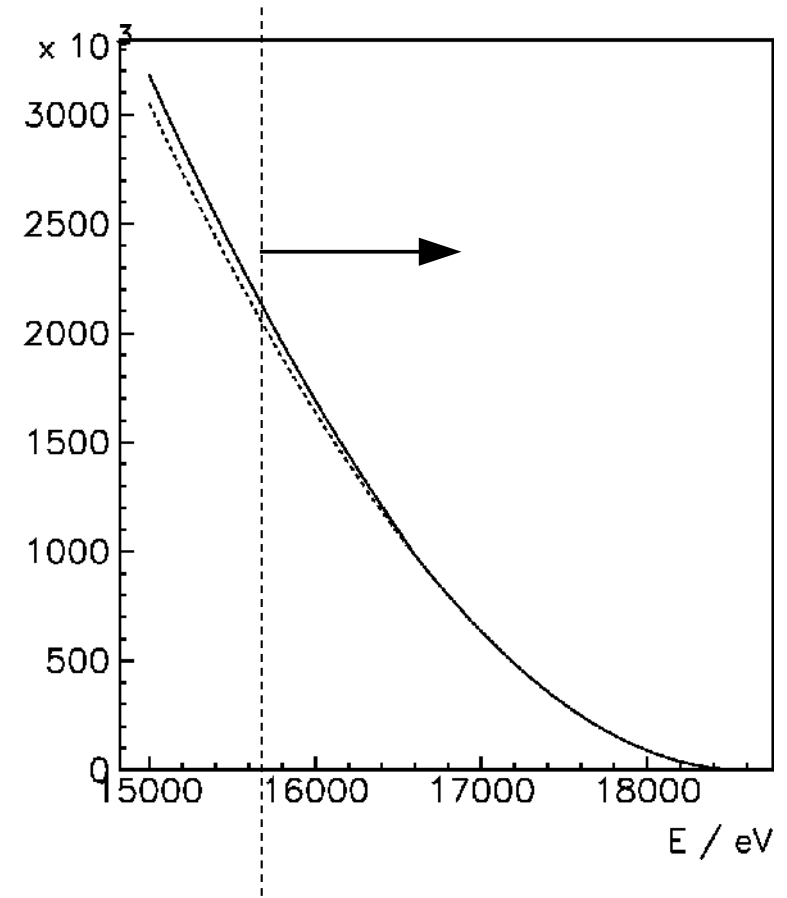
# Second gain of integral versus differential: Avoid many steps in MAC-E-Filter mode

## Integral – MAC-E-Filter method



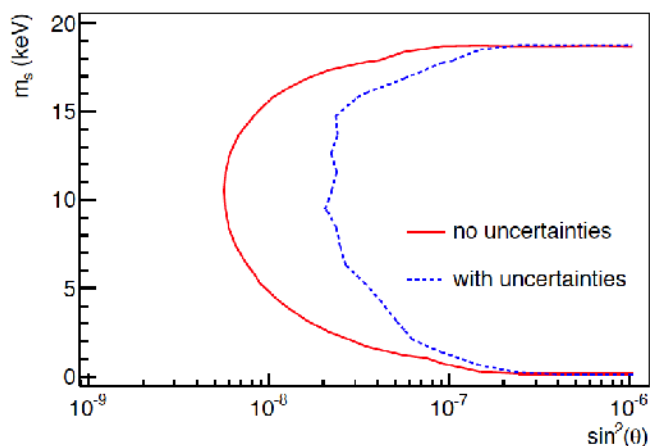
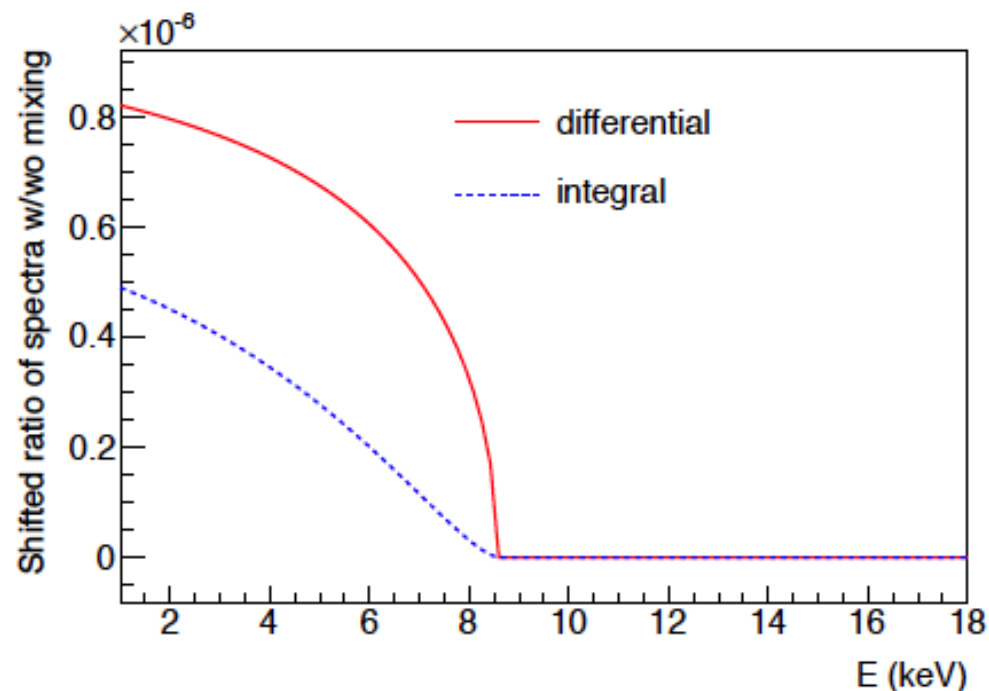
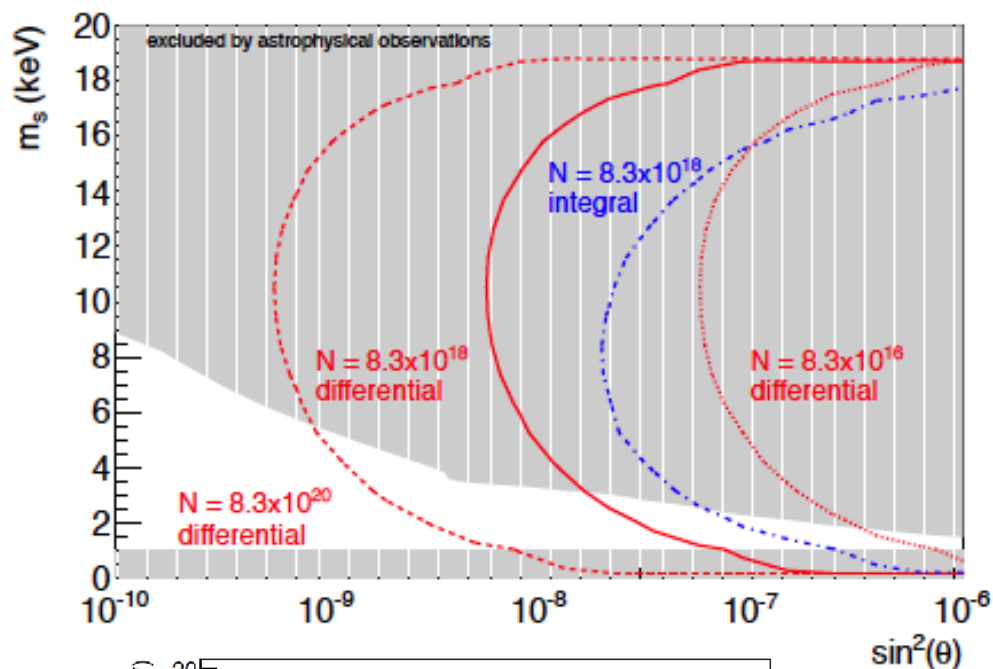
need many retardation voltages  
to obtain spectral information

## Differential measurement



need one retardation voltage  
and other means (TRISTAN-detector, TOF)  
to obtain spectral information

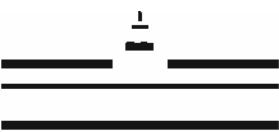
# Statistical sensitivity for integral and differential measurement



→ **Potential statistical uncertainty is not a problem for  $10^{-7}$  even including systematics (1<sup>st</sup> investigation)! but would require different measurement strategy**

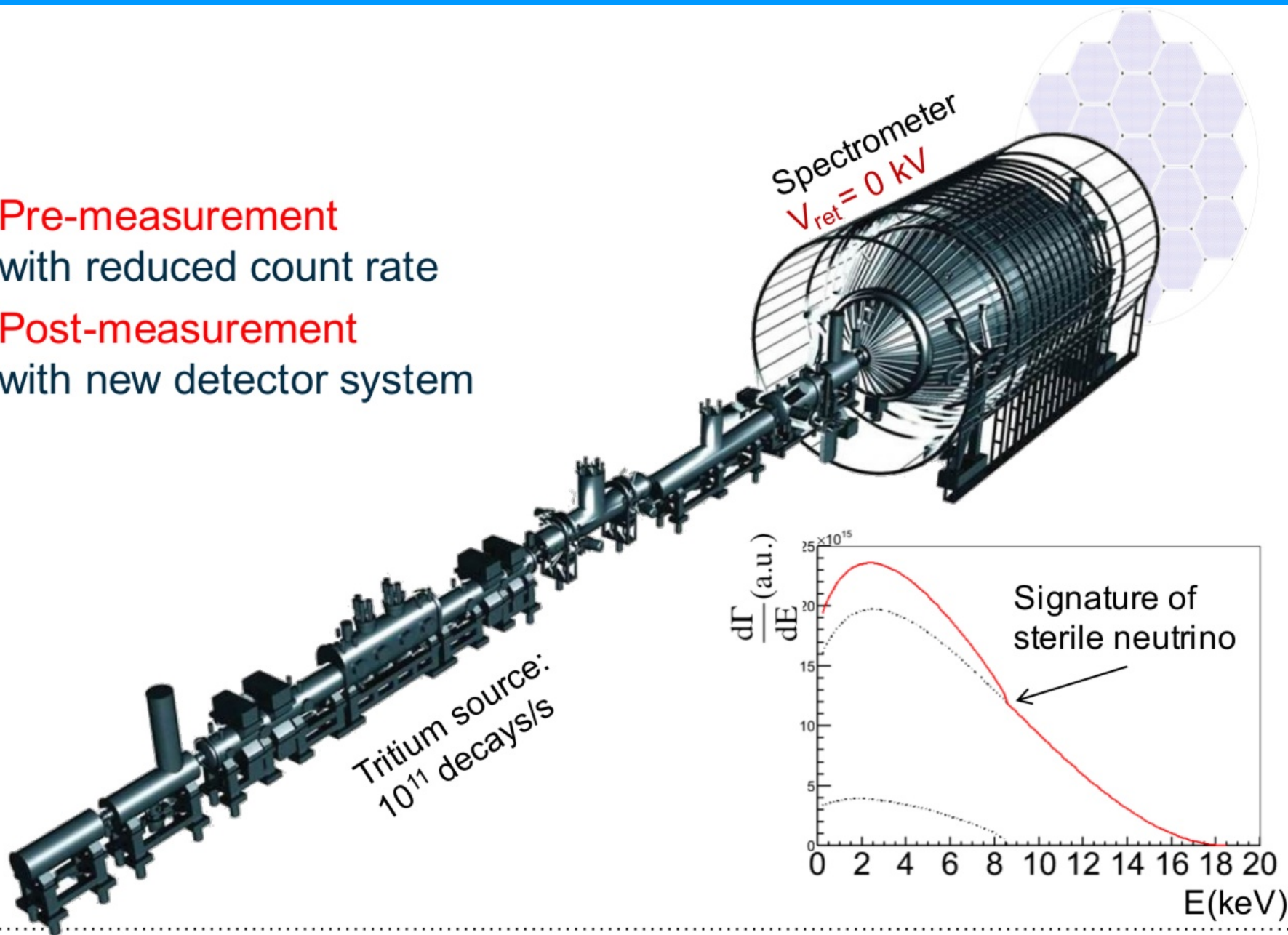
S. Mertens et al., JCAP 02 (2015) 020

„Sensitivity of Next Generation Tritium  $\beta$ -Decay Experiments for keV-Scale Sterile Neutrinos“



# How to search for keV neutrinos with KATRIN: several options

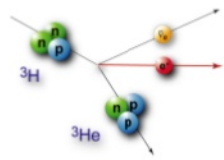
- **Pre-measurement** with reduced count rate
- **Post-measurement** with new detector system





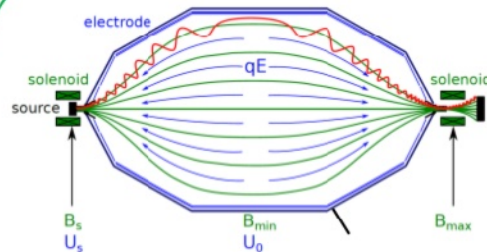
# How to search for keV neutrinos with KATRIN: could be done soon (2017)

## Pre-Measurement



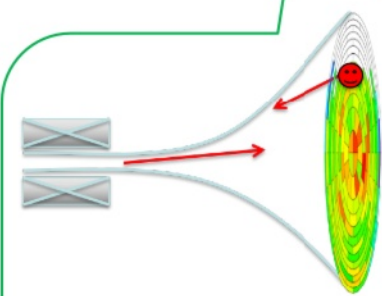
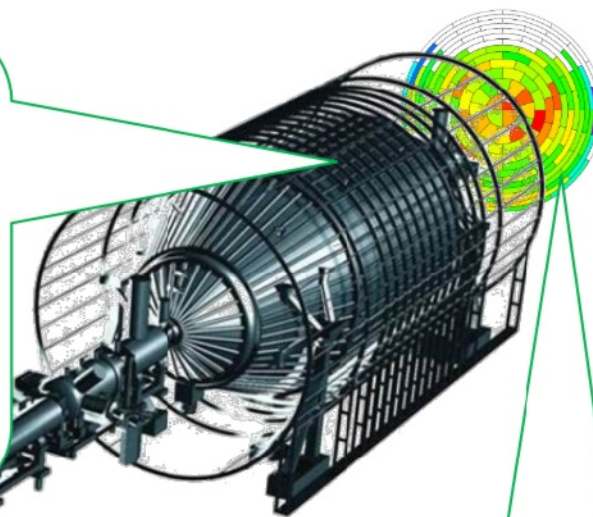
### Reduce count rate

- Less  $T_2$
- Small  $B_{\text{source}}$



### Adiabatic transport

- Large  $B_{\text{aircoil}}$



### Backscattering

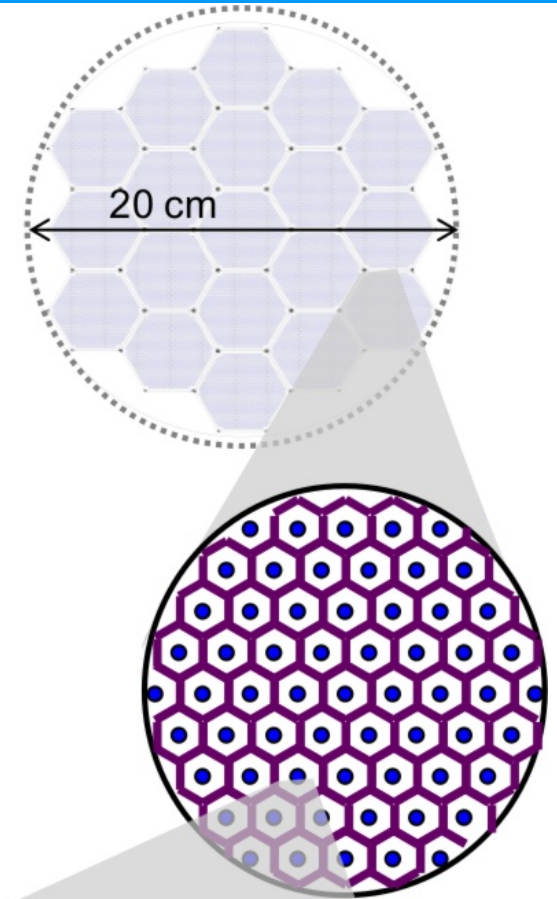
- Small  $B_{\text{detector}}$

Optimized electro-magnetic design seems to make sterile neutrino search with KATRIN possible

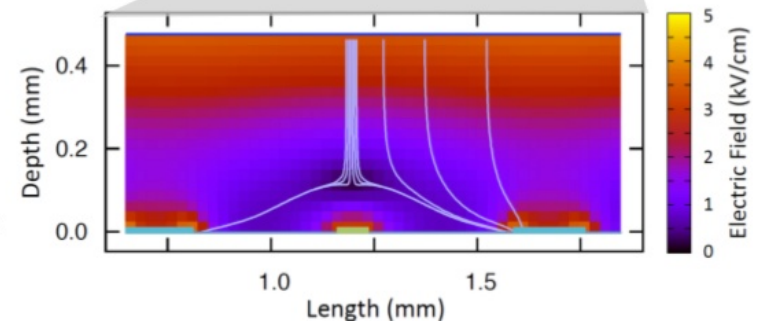
## Post-Measurement

### Novel detector design:

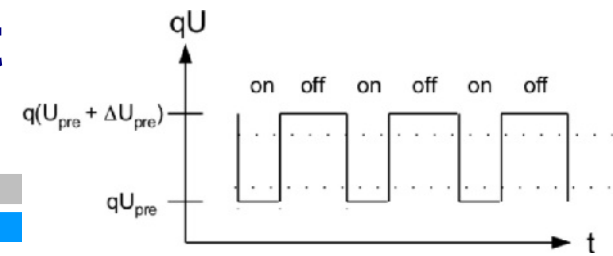
- Capability of handling high rates ( $>10^9$  cnts/s)
  - $>10\ 000$  pixel
- High energy resolution (300 eV @ 20 keV)
  - Thin deadlayer ( $\sim 10$  nm)
- Large pixels ( $\sim 1$  mm) with small capacity ( $<0.2$  pF)
  - Multi-drift-ring design (SDD)
- Minimize systematics (ppm-level)
  - Sophisticated read-out\*



TRISTAN detector  
design simulations



# Smearing of time-of-flight in gated-filter mode

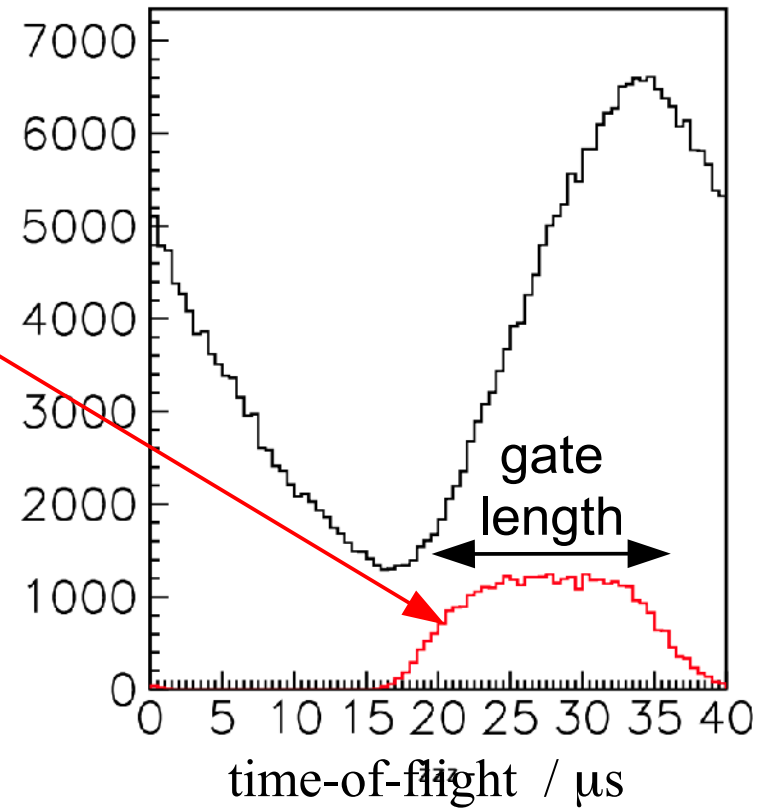
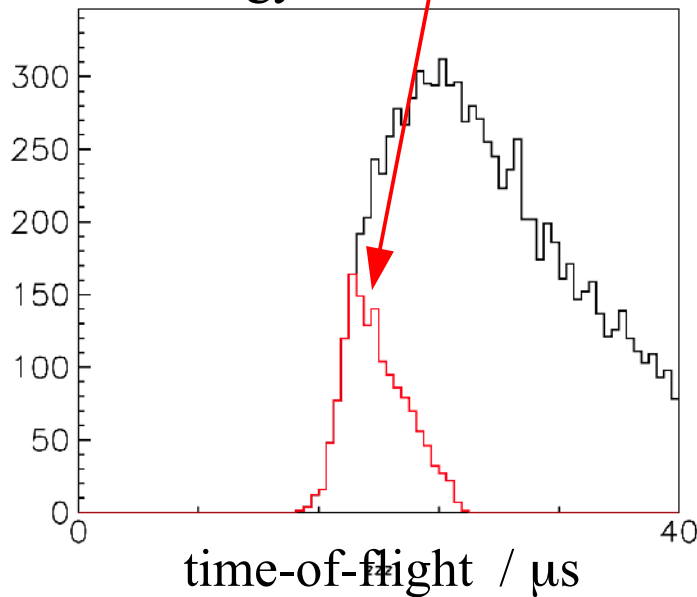
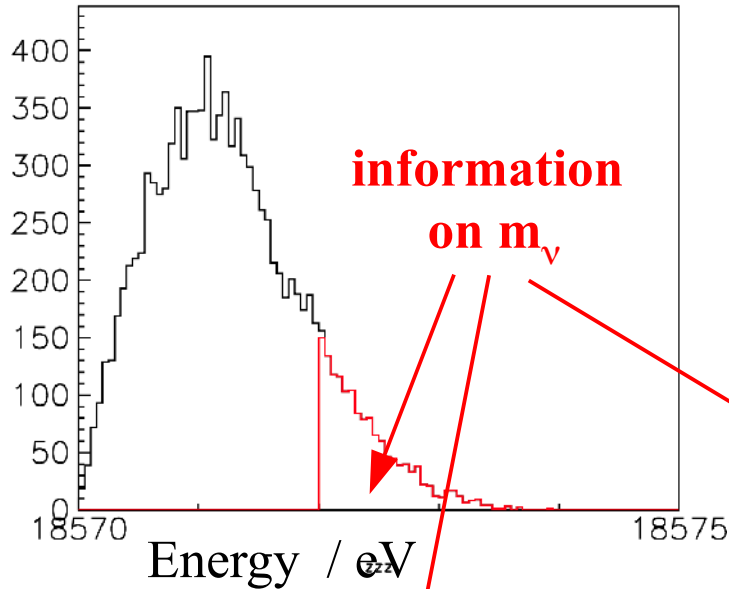


Gated filter:

$$T = t_{\text{on}} + t_{\text{off}} = 40 \mu\text{s}, \quad t_{\text{on}} = 16 \mu\text{s},$$

$$E_0 = 18.575 \text{ eV}, \quad qU = 18.570 \text{ eV}$$

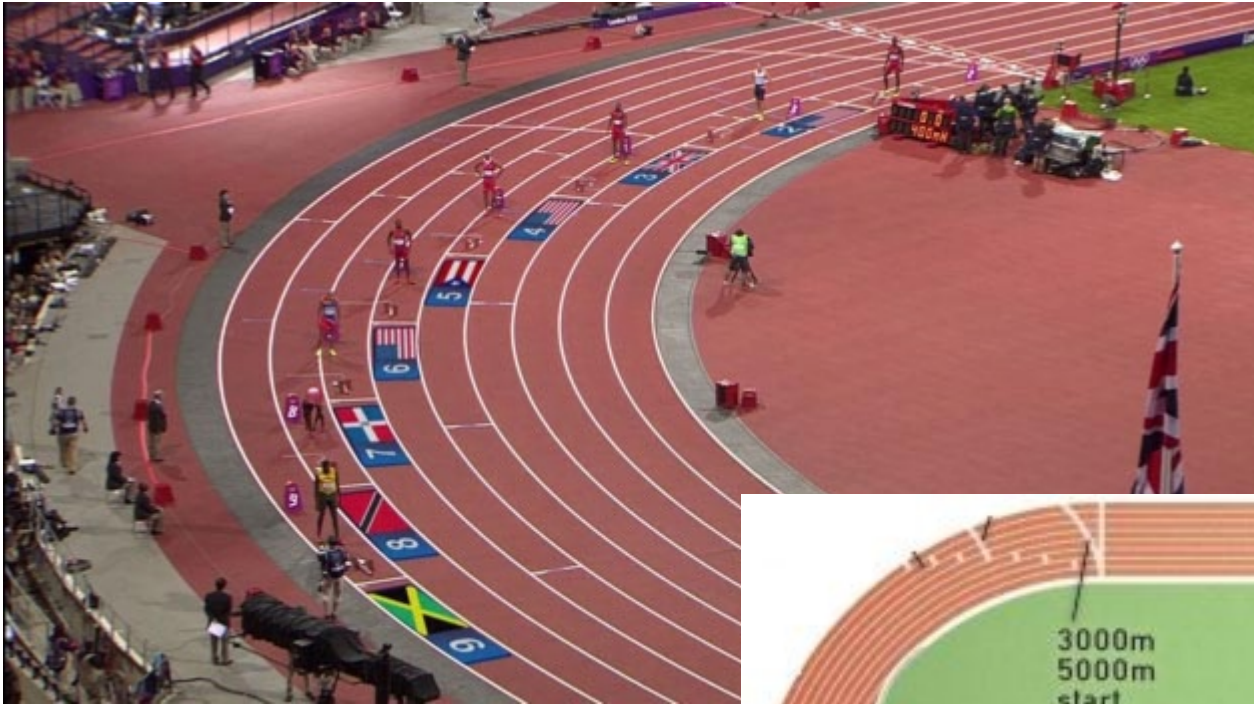
$E_0 = 18.575 \text{ eV}$ ,  $U_0 = -18570\text{V}$ , last 3 eV



→  $m_v$ -sensitive part of tof spectrum  
is smeared out → loss of sensitivity!



# Avoid smearing by gated-filter: “Time-focusing time-of-flight”



Idea: “Electrons starting behind/late  
get a shortcut/inner track  
or get accelerated (shortcut in time)  
in order to arrive in time”

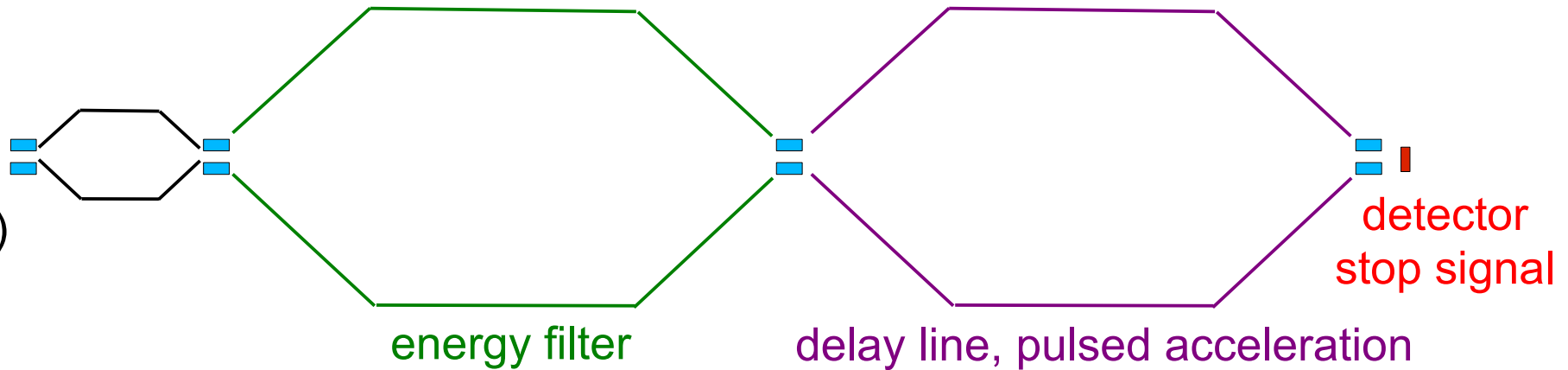


Pictures:  
<http://olympic.org>  
<http://100milesisnotthatfar.com>



# Time-focusing at KATRIN

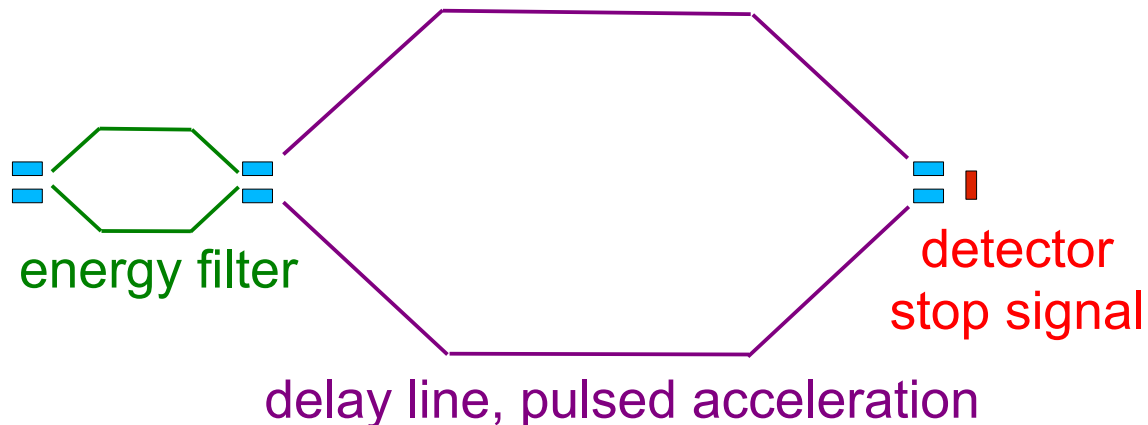
**ideal**  
(but  
not  
realistic)



„inner track“: give electron higher energies by periodic acceleration:

$$U = U_0 + U_1(t) \quad \text{e.g. } U_1(t) = U_{10} * (1 + \sin(\omega t)) / 2$$

**for keV sterile  $\nu$  search:** hardware exists in principle

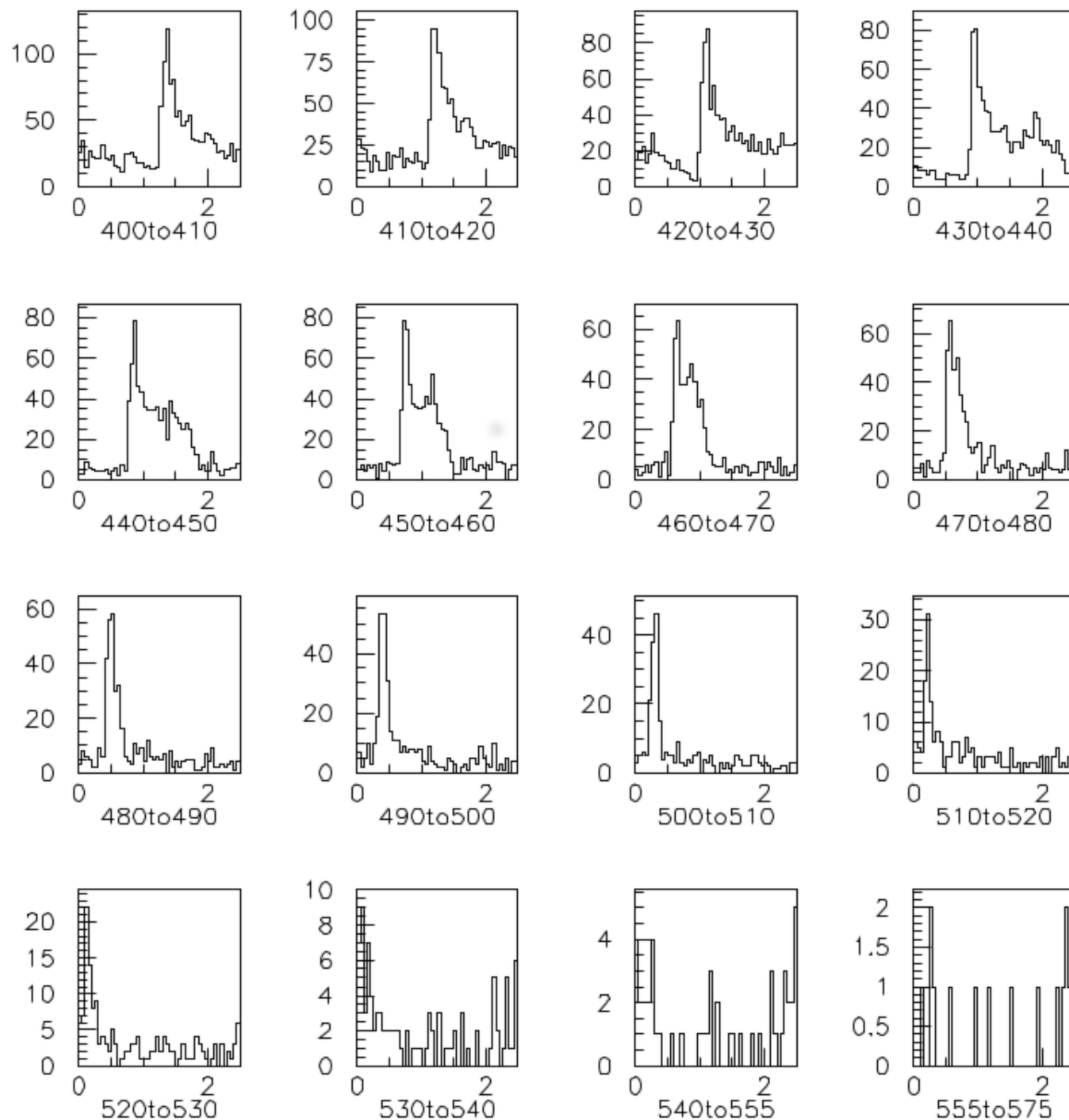
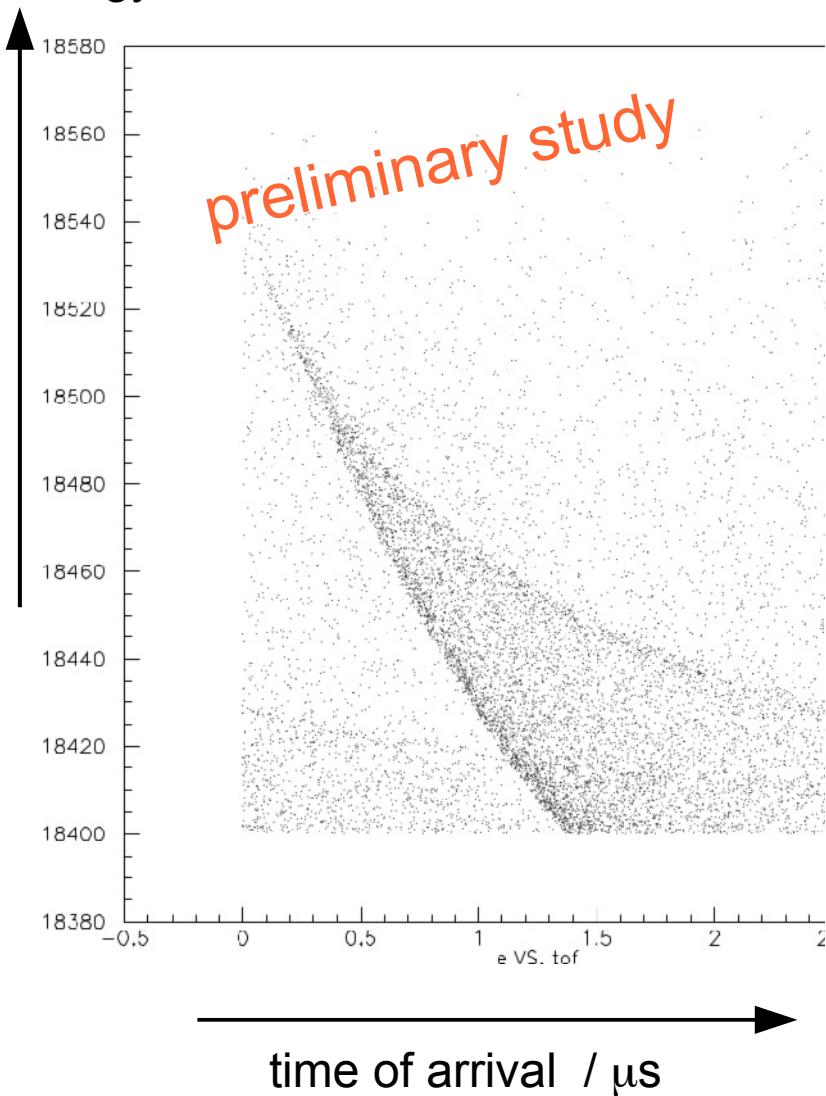


# Time-focusing at main spectrometer

sine wave 400 kHz, 100 Vpp, U = -18400 V

Ideal scenario, toy simulations

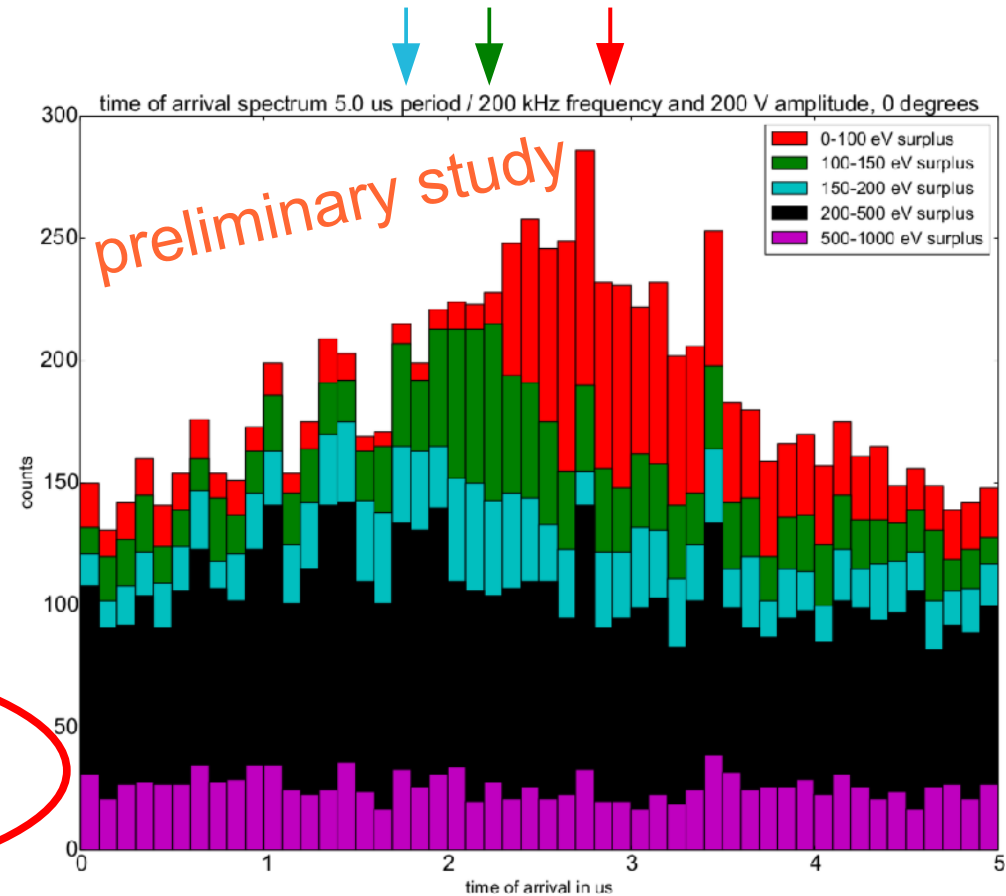
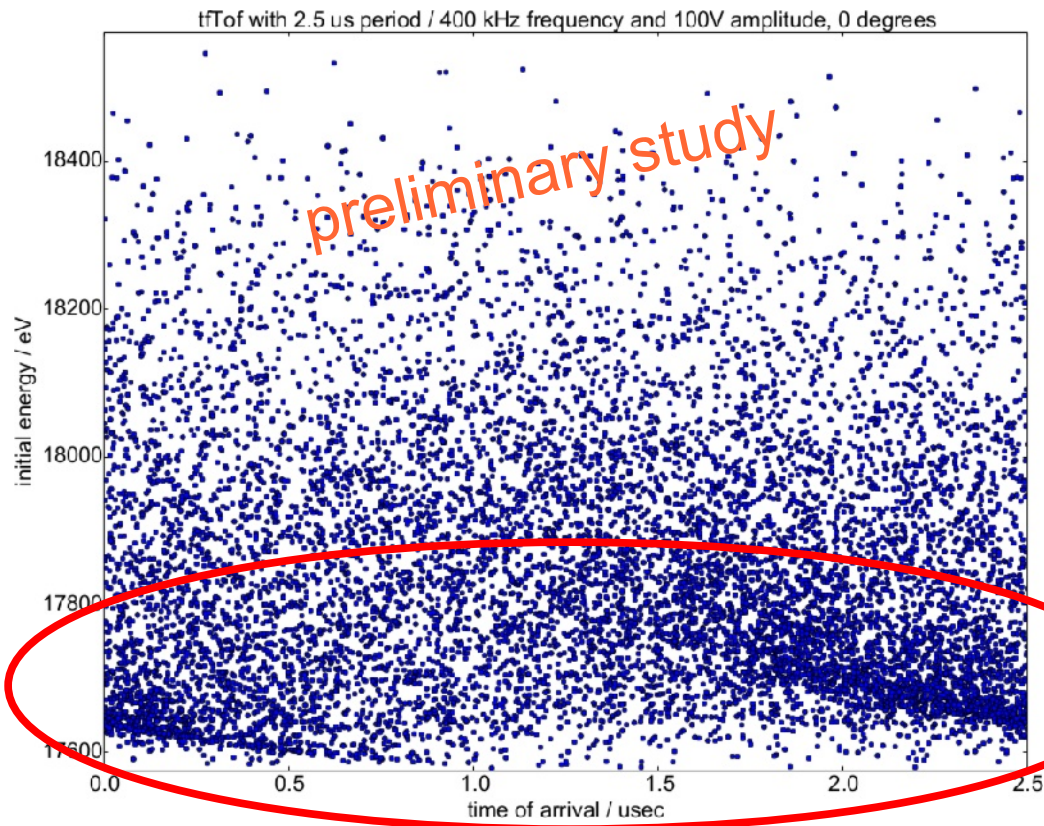
energy / eV



# Time-focusing at main spectrometer

## sine wave 400 kHz, 100 Vpp, $U = -17575$ V

Realistic scenario, modulating HV on central wire electrode system



**time-focusing energy band  
just above energy threshold**

# Statistical sensitivity of a $^{163}\text{Ho}$ EC experiment (ECHO, HOLMES, NuMECS)

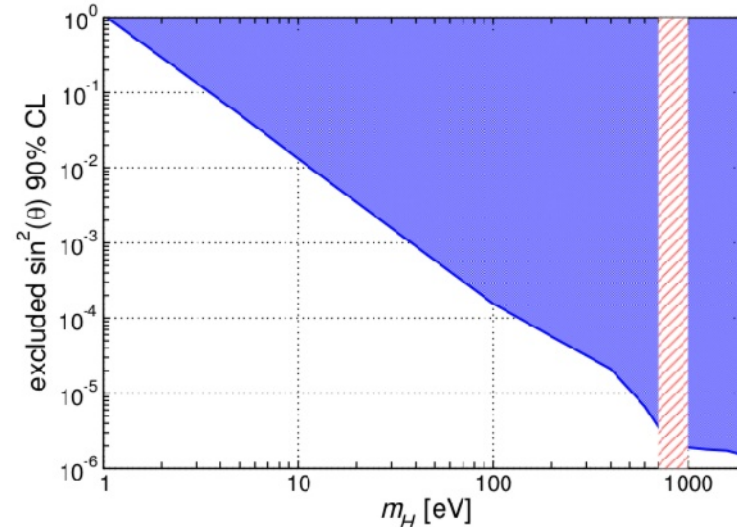
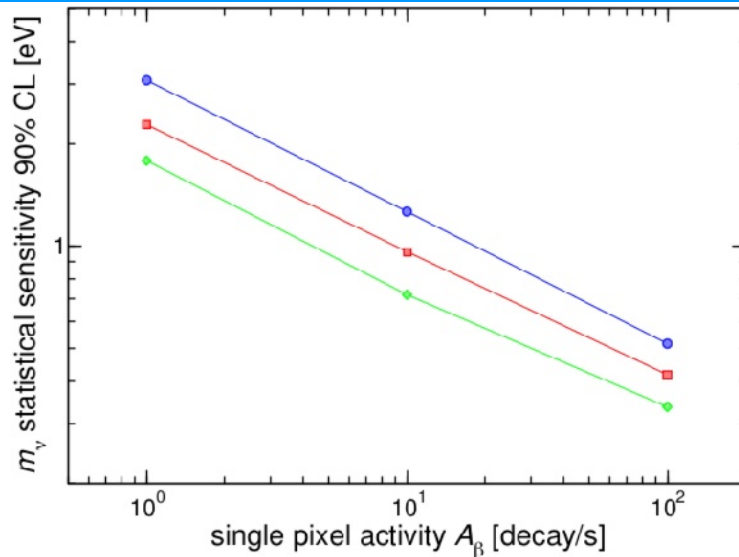
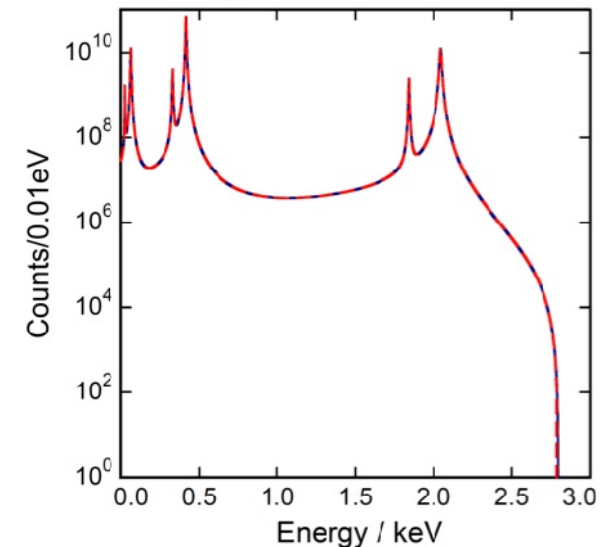


Figure 28: Calculated experimental  $^{163}\text{Ho}$  EC calorimetric spectrum for  $Q = 2.8 \text{ keV}$ ,  $\Delta E = 1 \text{ eV}$ , for a constant exposure of  $10^5 \text{ detector} \times \text{year}$ , and for (top to bottom)  $\tau_R = 10 \mu\text{s}$ ,  $1 \mu\text{s}$ , and  $0.1 \mu\text{s}$  (left). Sensitivity to heavy sterile neutrinos detected from kinks in a  $^{163}\text{Ho}$  calorimetric spectrum with  $Q = 2.8 \text{ keV}$ ,  $N_{ev} = 3 \times 10^{13}$ ,  $\Delta E = 1 \text{ eV}$ , and  $f_{pp} = 3 \times 10^{-4}$ . (right).

A. Nucciotti, arXiv:1511.00968

**limited by statistics**

**shape is non-trivial below the endpoint**



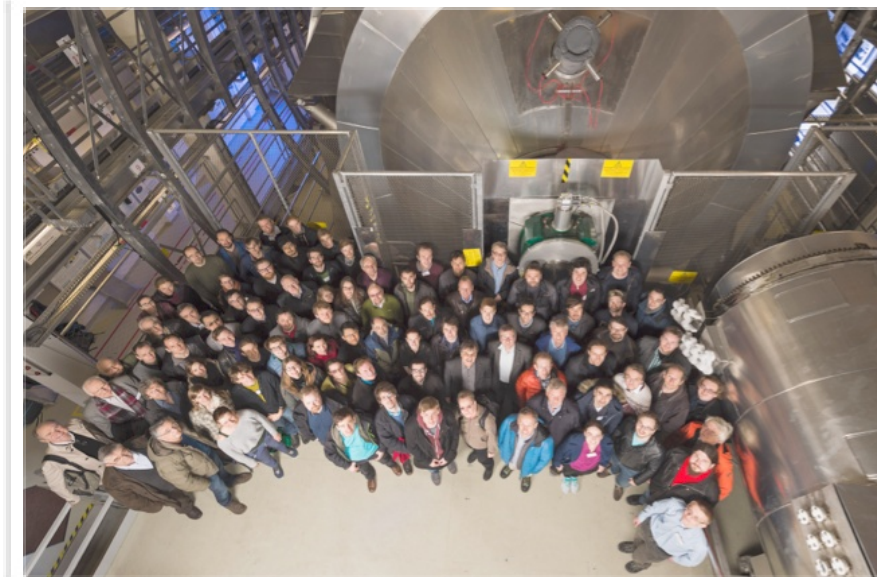


KATRIN is the direct neutrino mass experiment complementary to cosmological analyses and  $0\nu\beta\beta$  searches

KATRIN will start direct neutrino mass measurements in 2017

KATRIN's sensitivity: 200 meV

KATRIN can also look for sterile neutrinos (eV, keV)



Is 200 meV the end of direct neutrino mass searches? No!

- significant developments on  $^{163}\text{Ho}$  micro calorimeters (ECHO, HOLMES, NuMECS)
- new ideas like Project 8, ...
- addition differential methods to KATRIN by TOF, new detectors, ..

keV neutrino search possible with KATRIN, first search will be done in 2017 !

**THANK YOU FOR YOUR ATTENTION !**