

Dark Matter Halos: with and without Baryons

Yehuda Hoffman
Racah Institute of Physics
Hebrew University, Jerusalem, Israel

Hebrew University: Isaac Rodrigez

Emilio Romano-Diaz, Isaac Shlosman (Kentucky), Clayton Heller
(Statesboro, GA)

CLUES collaboration: Stefan Gottloeber (Potsdam), Gustavo Yepes,
Luis Martinez-Vaquero, Alexander Knebe, Steffen Knollmann

◆ General remarks: with & without baryons

◆ PDM halos: Λ CDM phenomenology

◆ Theoretical Considerations:

1. Adiabatic contraction

2. Dynamical friction

◆ Analysis

1. Dissecting the Romano-Diaz et al halo

I. Evolution in phase space

II. Radial profiles

III. Subhalos

2. CLUES

I. Radial profiles

II. Subhalos

◆ Final statements

General remarks: DM halos with & without baryons in the Λ CDM cosmology

Notations:

- Pure DM simulations (PDM)
- **Baryons + DM simulations (BDM)**

- A PDM HALO is a well defined object. Almost a general consensus on radial structure, substructures, shape, angular momentum, ...
- The structure of DM halos is well known (from simulations) but is hardly understood (analytically)
- **No consensus on BDM halos**
- **No numerical convergence**
- **No general consensus on the subgrid processes**
- **Results depends on the numerical implementations of subgrid processes**

Theoretical Considerations: A. Adiabatic Contraction

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CONTRACTION OF DARK MATTER GALACTIC HALOS DUE TO BARYONIC INFALL¹

GEORGE R. BLUMENTHAL AND S. M. FABER

Lick Observatory, Board of Studies in Astronomy and Astrophysics, University of California, Santa Cruz

RICARDO FLORES

Department of Physics, Brandeis University, Waltham, Massachusetts

AND

JOEL R. PRIMACK

Board of Studies in Physics, University of California, Santa Cruz

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$$M(r)r \approx \text{const}$$

$$[M_{\text{dm}}(r) + M_b(r)]r = [M_{\text{dm}}(r) + M_b(r_f)]r_f$$

initial

final

Theoretical Considerations: A. Adiabatic Contraction

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THE ASTROPHYSICAL JOURNAL, 616:16–26, 2004 November 20
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RESPONSE OF DARK MATTER HALOS TO CONDENSATION OF BARYONS: COSMOLOGICAL SIMULATIONS AND IMPROVED ADIABATIC CONTRACTION MODEL

OLEG Y. GNEDIN,¹ ANDREY V. KRAVTSOV,² ANATOLY A. KLYPIN,³ AND DAISUKE NAGAI²

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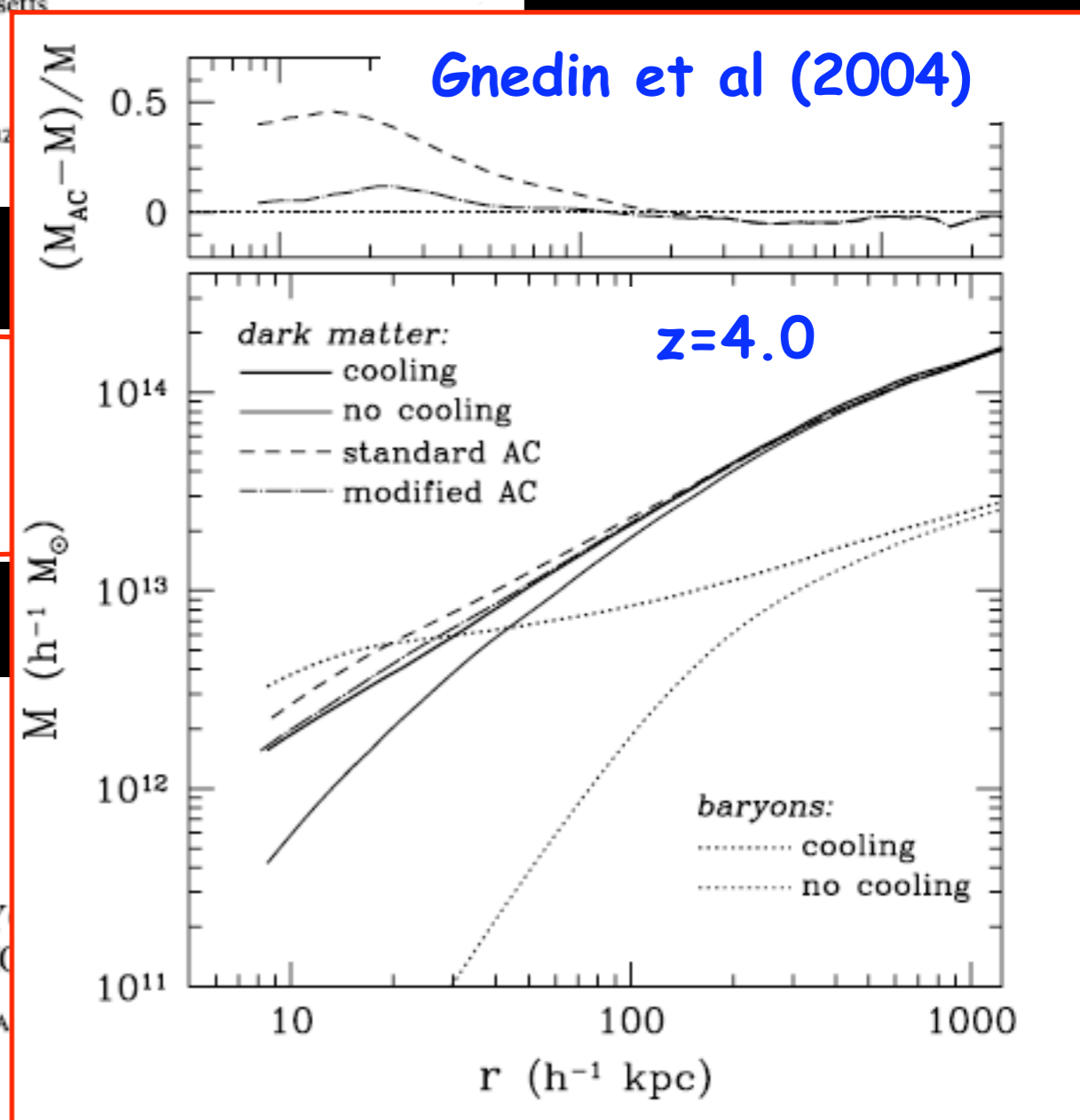
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Received 1985 June 14; accepted 1985 July 9

$$M(r)r \approx \text{const}$$

$$[M_{\text{dm}}(r) + M_b(r)]r = [M_{\text{dm}}(r) +$$

initial



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RESPONSE OF DARK MATTER HALOS TO CONDENSATION OF BARYONS IN SIMULATIONS AND IMPROVED ADIABATIC CONTRACTION

OLEG Y. GNEDIN,¹ ANDREY V. KRAVTSOV,² ANATOLY A. KLYPIN,³ AND DAVID N. SPOFFORD

Received 2004 June 9; accepted 2004 August 3

Theoretical Considerations: B. Dynamical Friction

THE ASTROPHYSICAL JOURNAL, 560:636–643, 2001 October 20
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DARK HALOS: THE FLATTENING OF THE DENSITY CUSP BY DYNAMICAL FRICTION

AMR EL-ZANT AND ISAAC SHLOSMAN

Department of Physics and Astronomy, 177 Chemistry/Physics Building, University of Kentucky, Lexington, KY 40506-0055;
elzant@pa.uky.edu, shlosman@pa.uky.edu

AND

YEHUDA HOFFMAN

Racah Institute of Physics, Hebrew University, Jerusalem, Israel 91904; hoffman@vms.huji.ac.il

THE ASTROPHYSICAL JOURNAL, 607:L75–L78, 2004 June 1
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FLAT-CORED DARK MATTER IN CUSPY CLUSTERS OF GALAXIES

AMR A. EL-ZANT,¹ YEHUDA HOFFMAN,² JOEL PRIMACK,³ FRANCOISE COMBES,⁴ AND ISAAC SHLOSMAN⁵

- Clumpy mixture of DM and baryons - clumps loose energy to the ambient DM -> heating and expansion of the DM
- A key element - clumps need to be baryon rich or otherwise there is no effect!

Theoretical Considerations: B. Dynamical Friction

THE ASTROPHYSICAL JOURNAL, 560:636–643, 2001 October 20

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DARK HALOS: THE

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FLAT

AMR A. EL-ZANT,¹ YEHUDA HOFFMAN,² JOEL PRIMACK,³ FRANCOISE COMBES,⁴ AND ISAAC SHLOSMAN⁵

Simplified Dynamical Model

- Substructures modeled as point-like particles
- Dynamical friction is modeled by Chandrasekhar (1943) formula
- No evolution of substructures - no attempt to account for star formation, feedback, ...
- Start from NFW DM halos

FRICION

0506-0055;

- Clumpy mixture of DM and baryons - clumps loose energy to the ambient DM -> heating and expansion of the DM
- A key element - clumps need to be baryon rich or otherwise there is no effect!

Note on simulations:

- Initial conditions set by constrained realizations of Gaussian fields
- Romano-Diaz et al simulations - a first step in a project to 'design' a halo 'on demand'
- Constrained Local Universe Simulations - use observational data to constrain the 'local universe' - Local Volume, Local Group

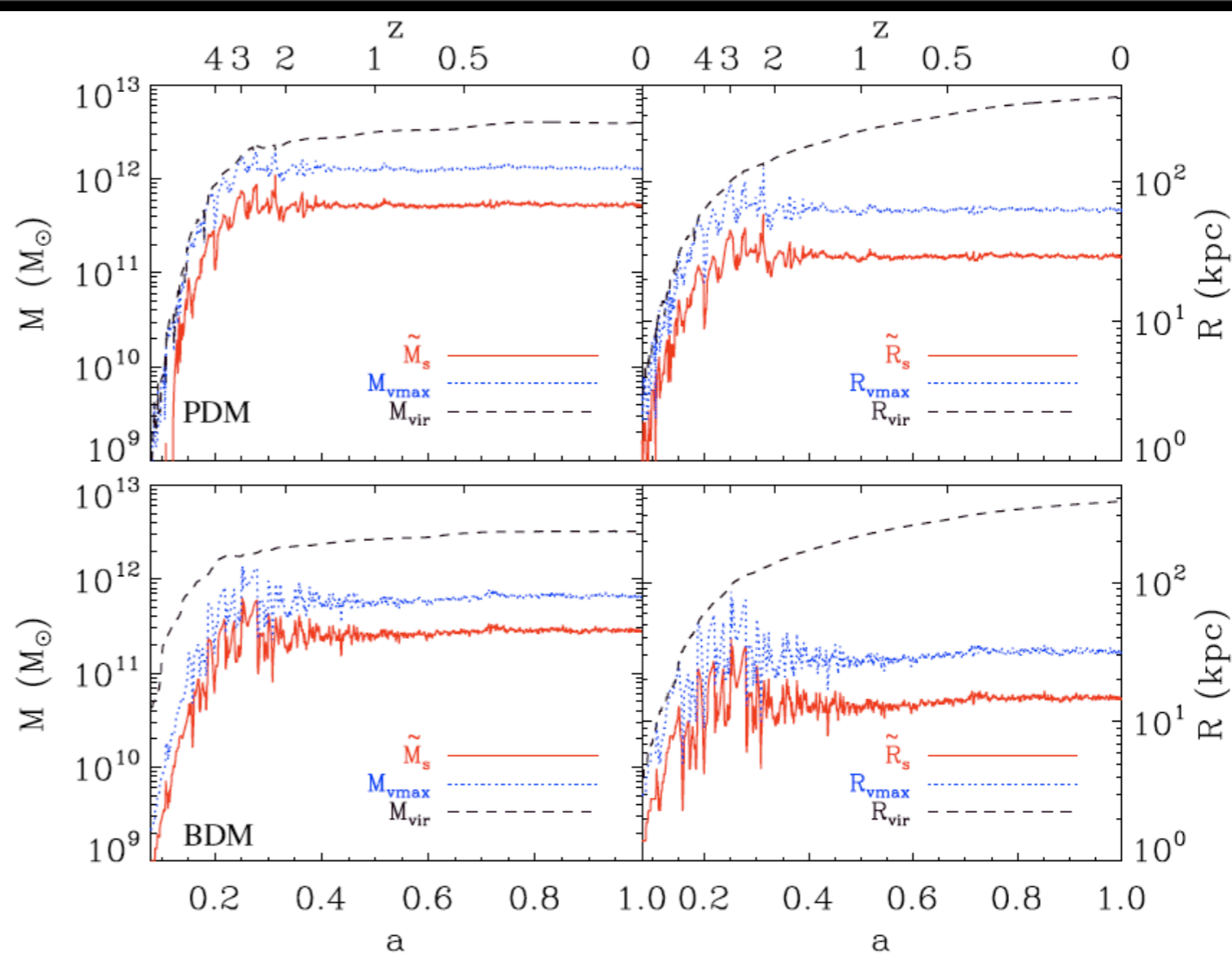
Dissecting Galaxy Formation: Comparison of the DM in PDM and BDM simulations

Romano-Diaz, Shlosman, Heller & YH (2008 - 2010)

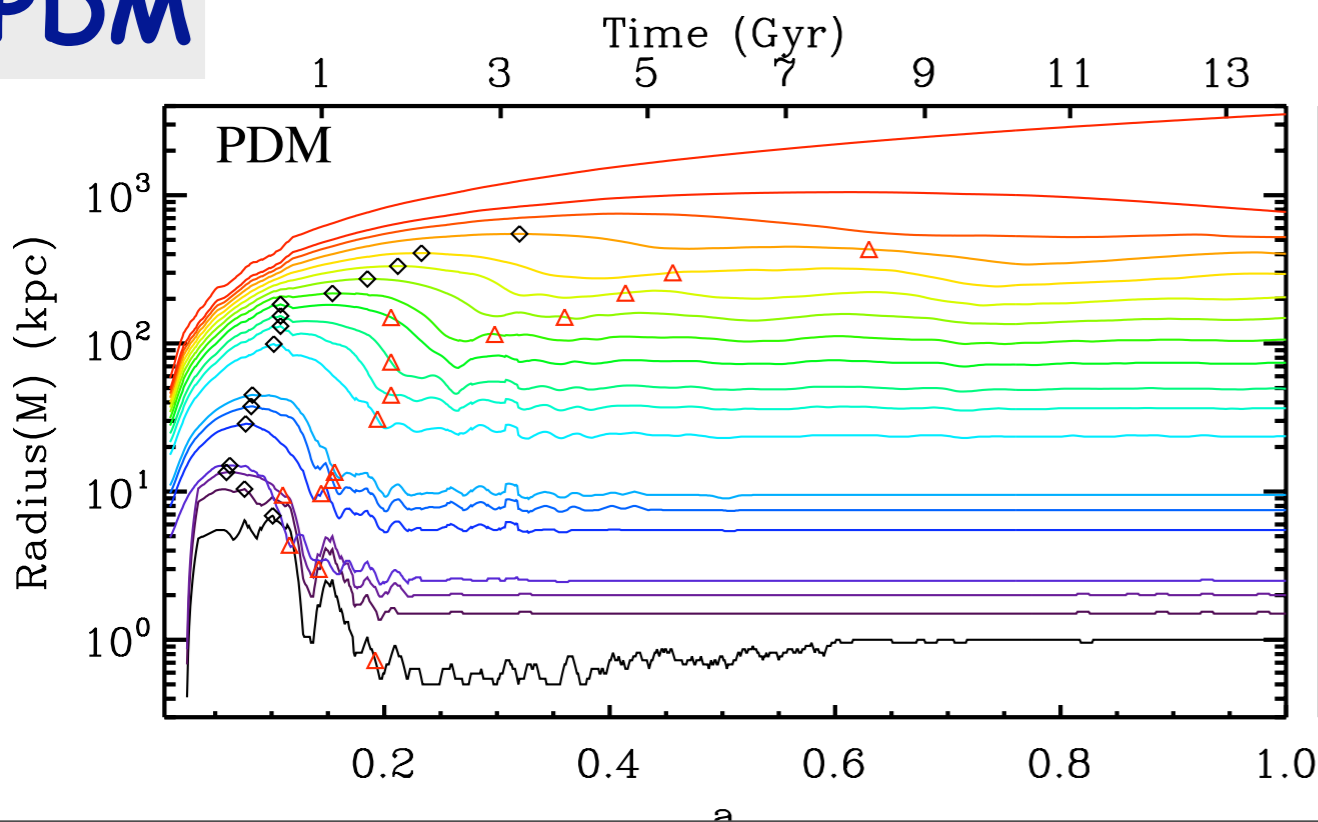
- Code SPH (FTM - Heller & Shlosman 1994, Heller et al 2007)
- Physics: feedback - stellar winds & SN -> delayed cooling
- Physical coordinates, vacuum boundary conditions
- Computational sphere 6 Mpc/h
- $N_{\text{tot}} \sim 6 \cdot 10^6$, $m(\text{DM particle}) \sim 2 \cdot 10^6 M_{\text{sun}}$
- Single halo sets as a constrained realization

General overview of the PDM & BDM halos:

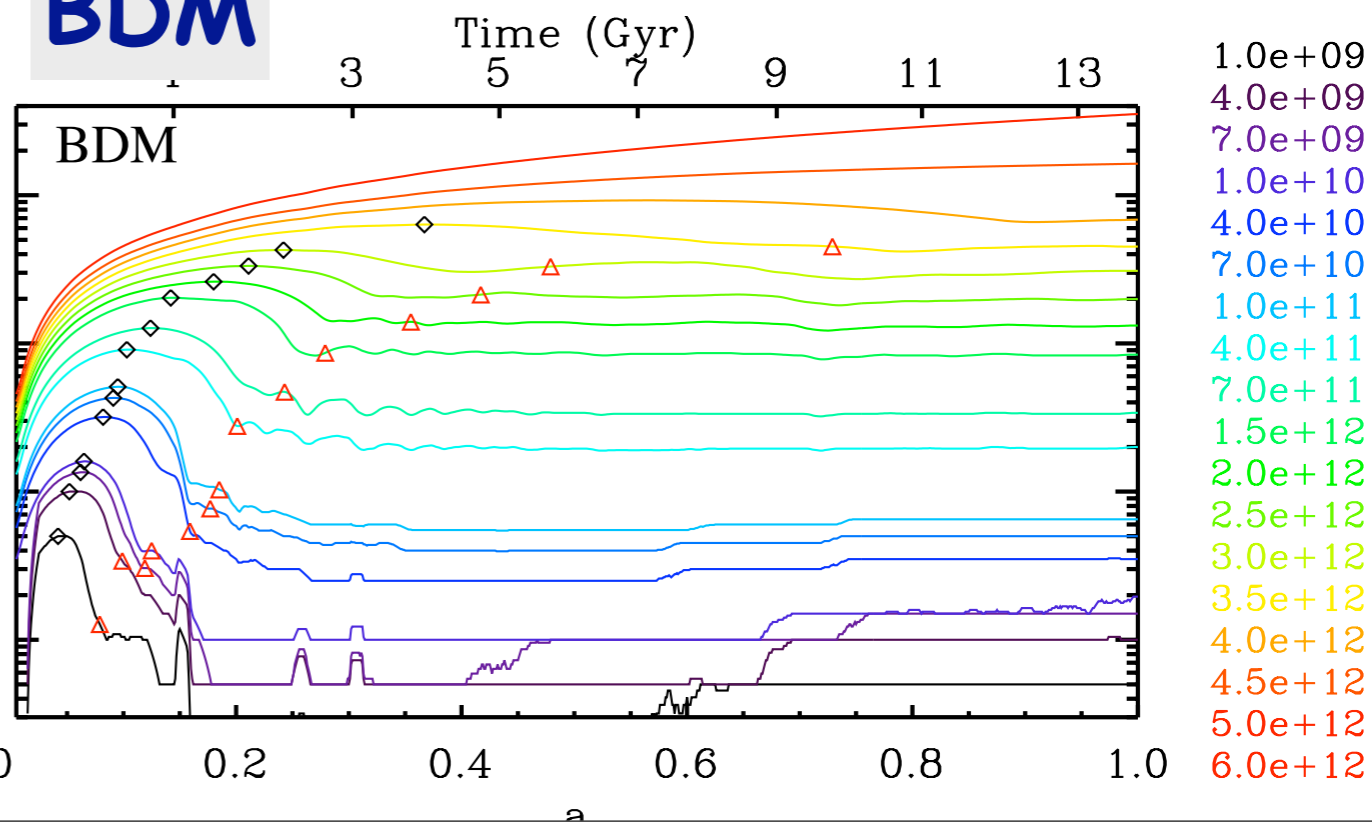
$M_{\text{vir}} \sim 3.5e12 M_{\text{sun}}$
 $R_{\text{vir}} \sim 400 \text{ kpc}$



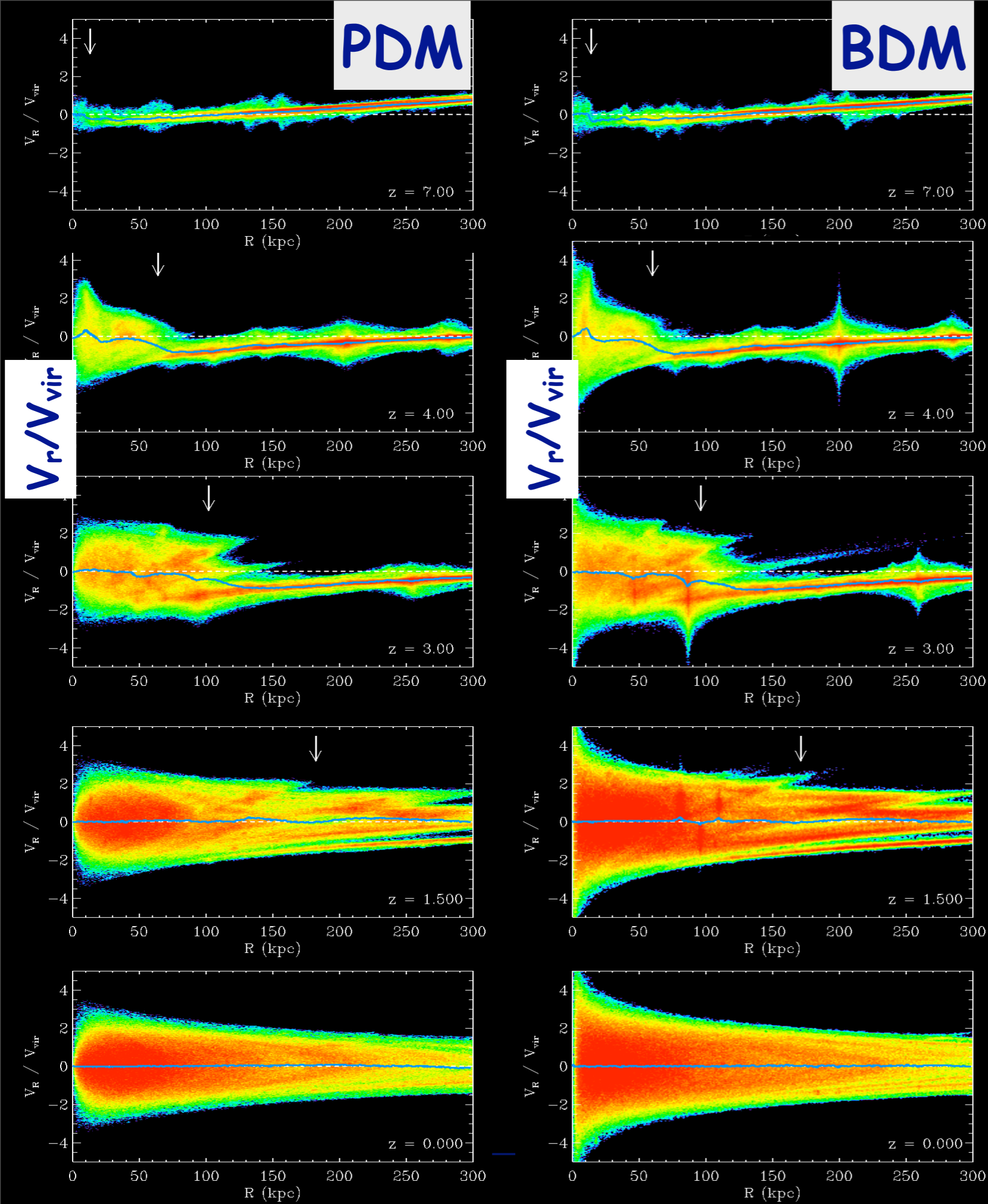
PDM



BDM

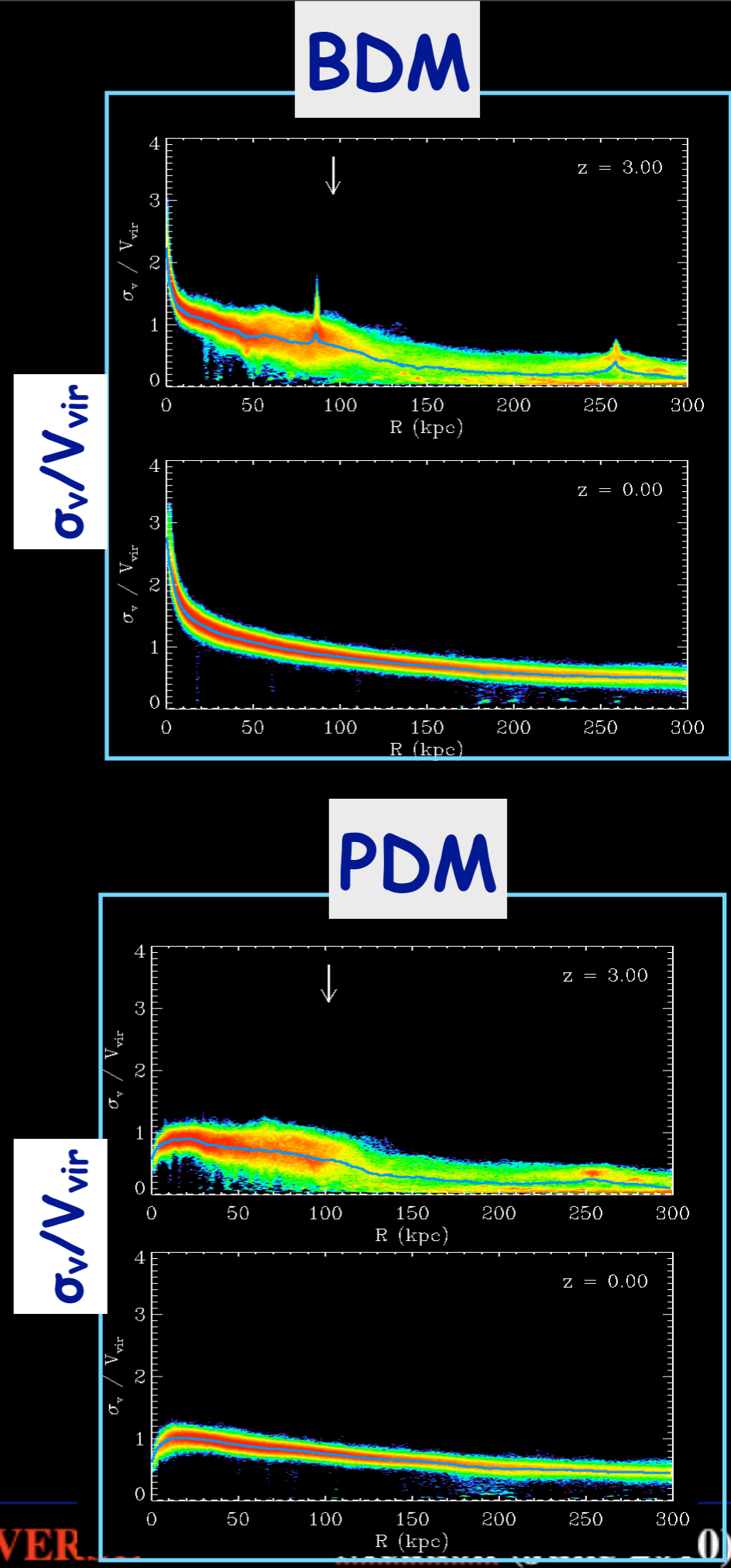
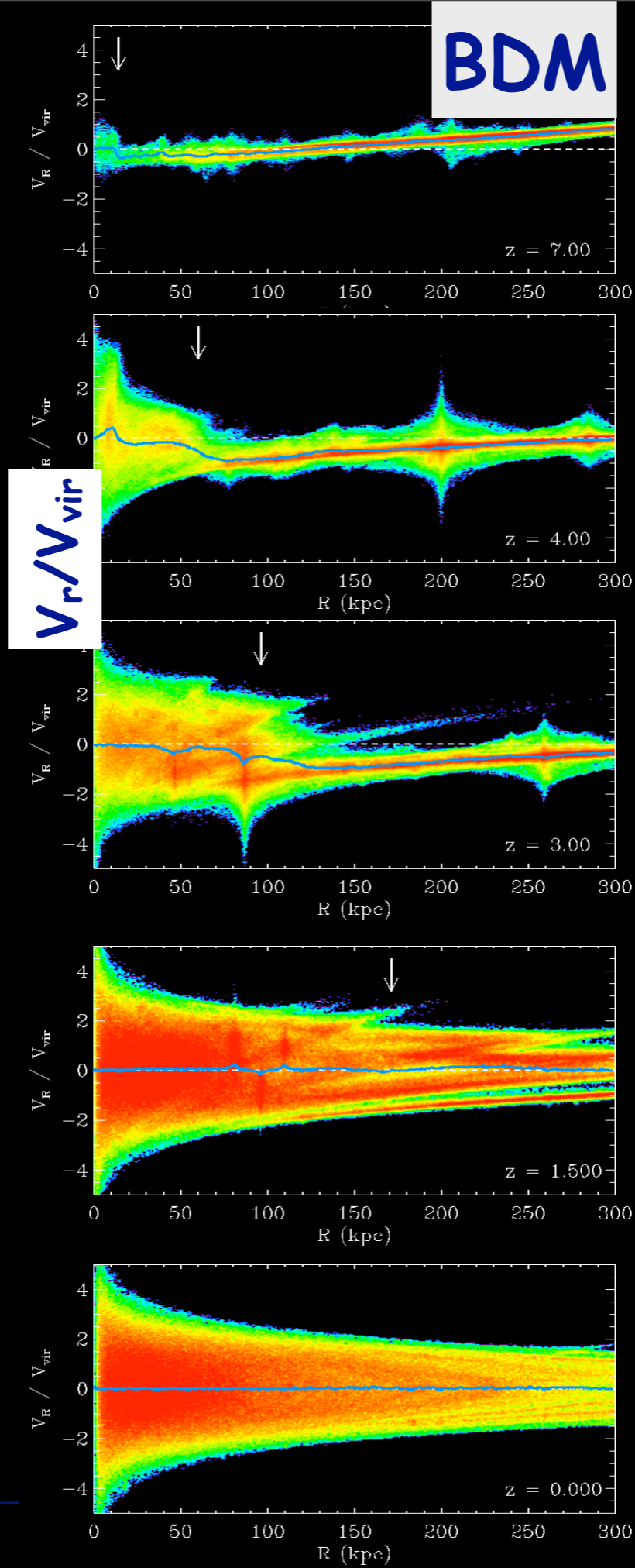
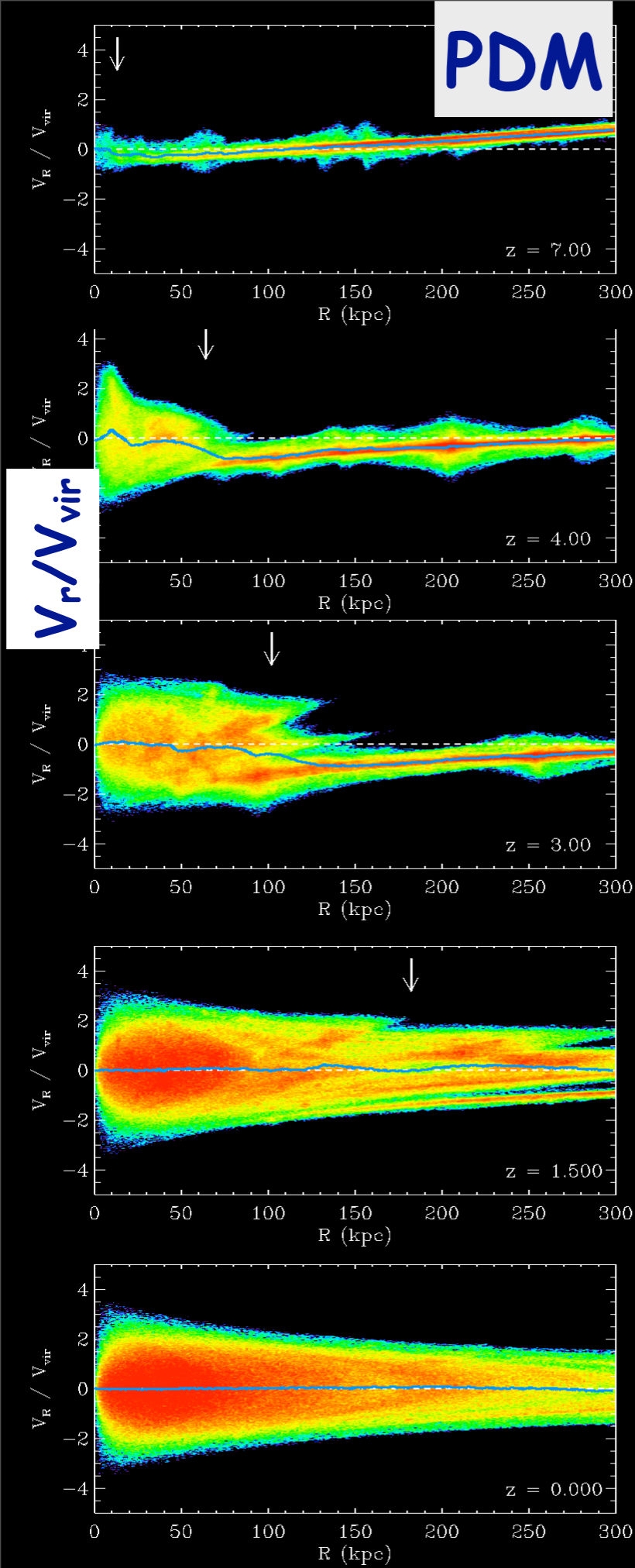


- 1.0e+09
- 4.0e+09
- 7.0e+09
- 1.0e+10
- 4.0e+10
- 7.0e+10
- 1.0e+11
- 4.0e+11
- 7.0e+11
- 1.5e+12
- 2.0e+12
- 2.5e+12
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- 4.0e+12
- 4.5e+12
- 5.0e+12
- 6.0e+12



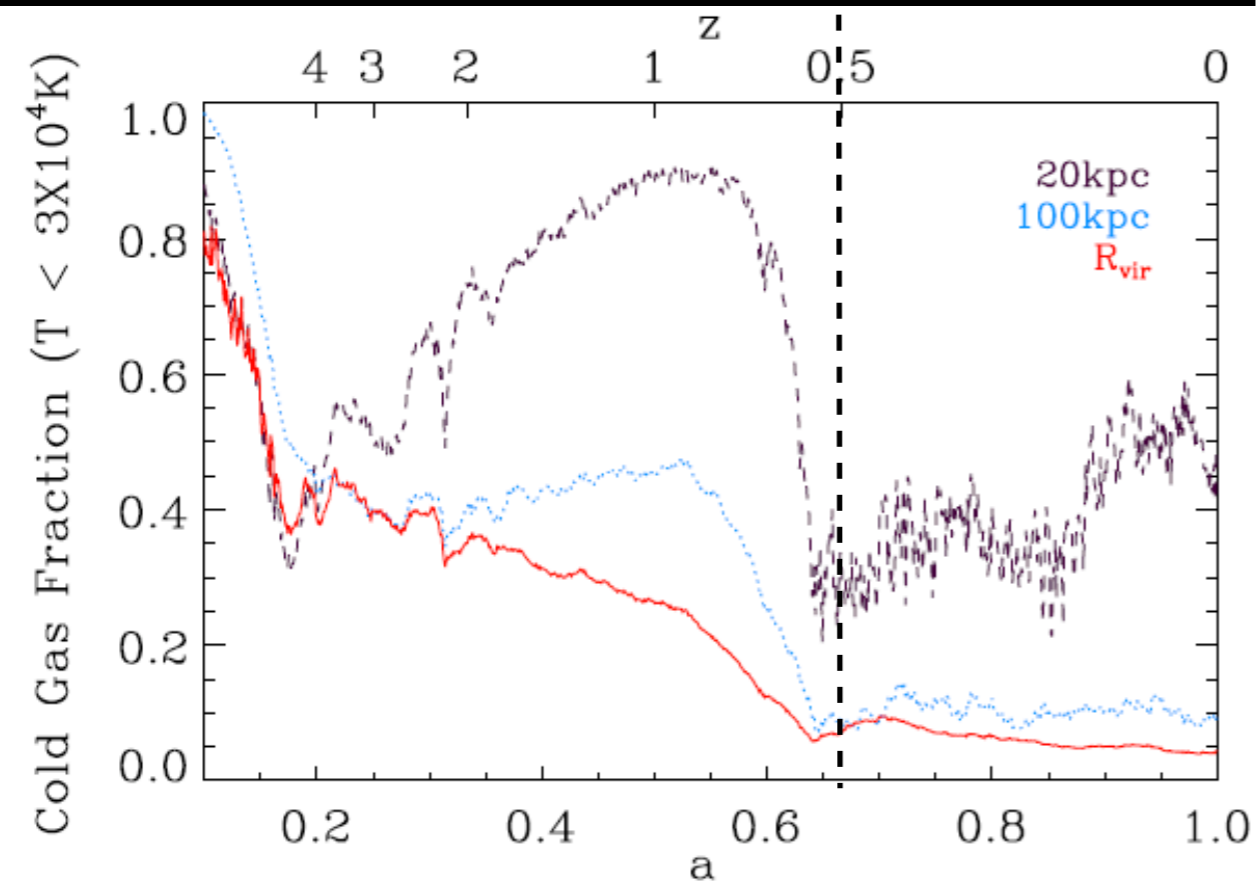
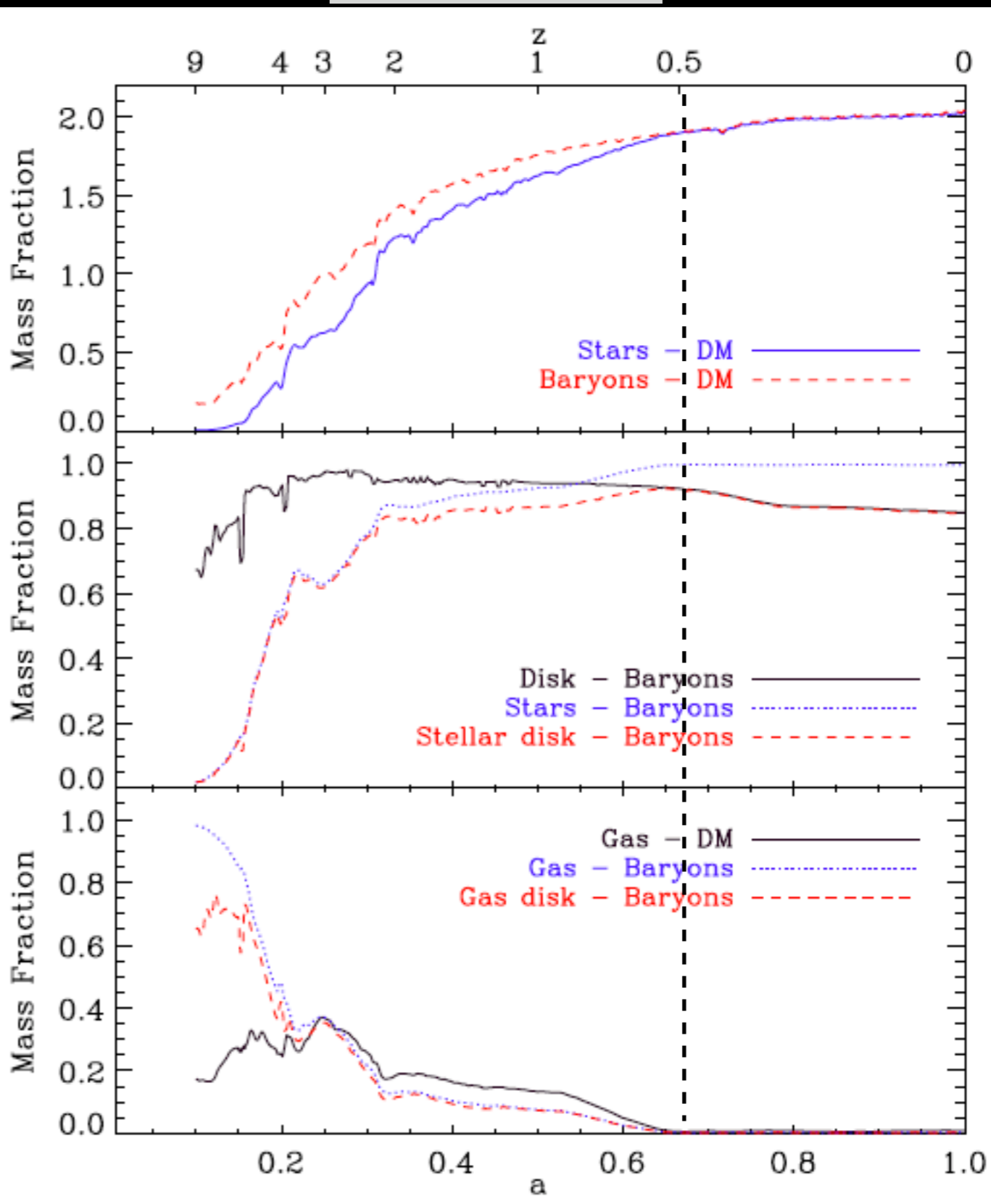
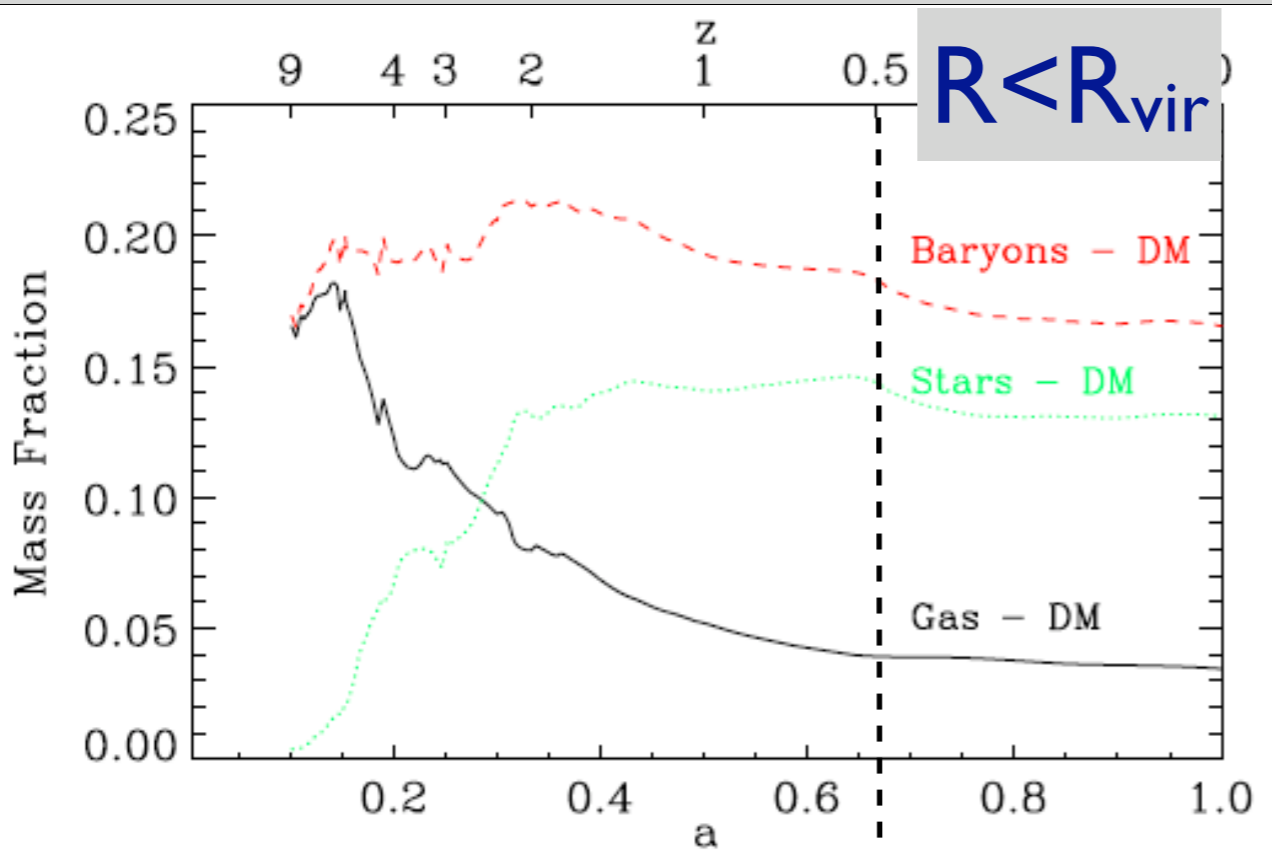
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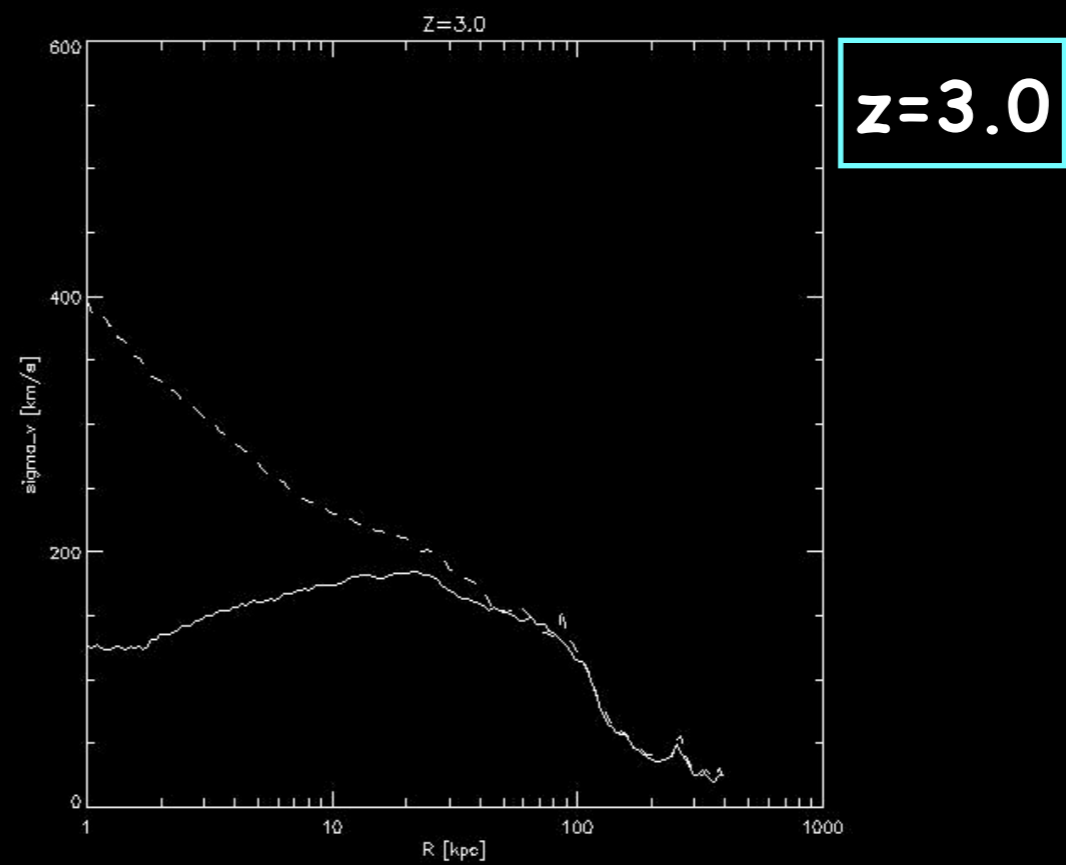
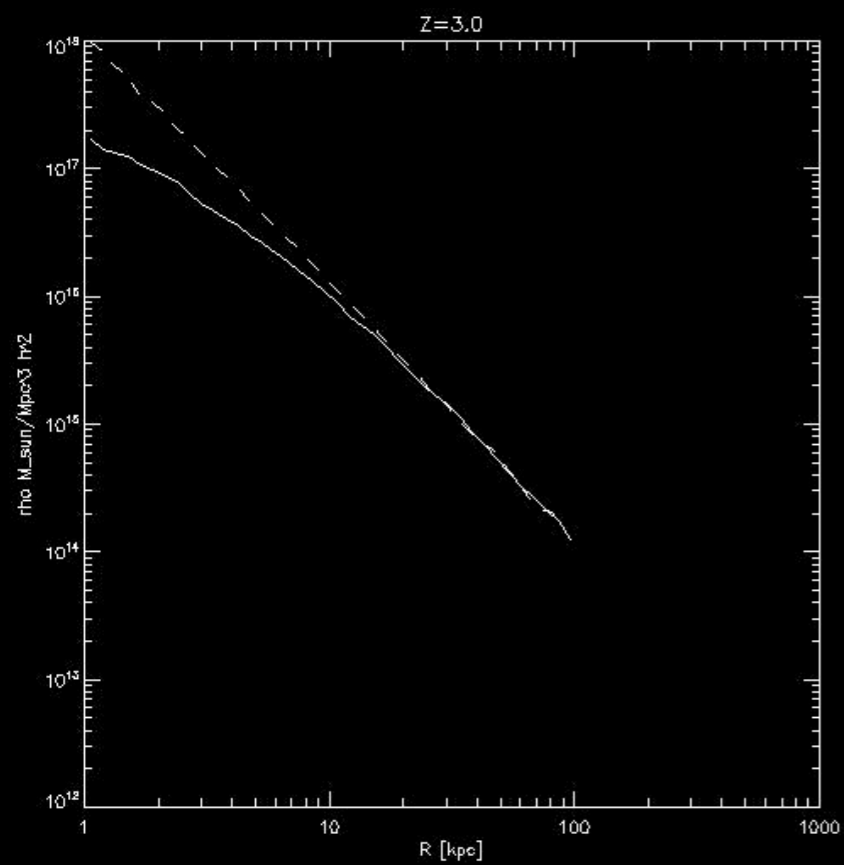
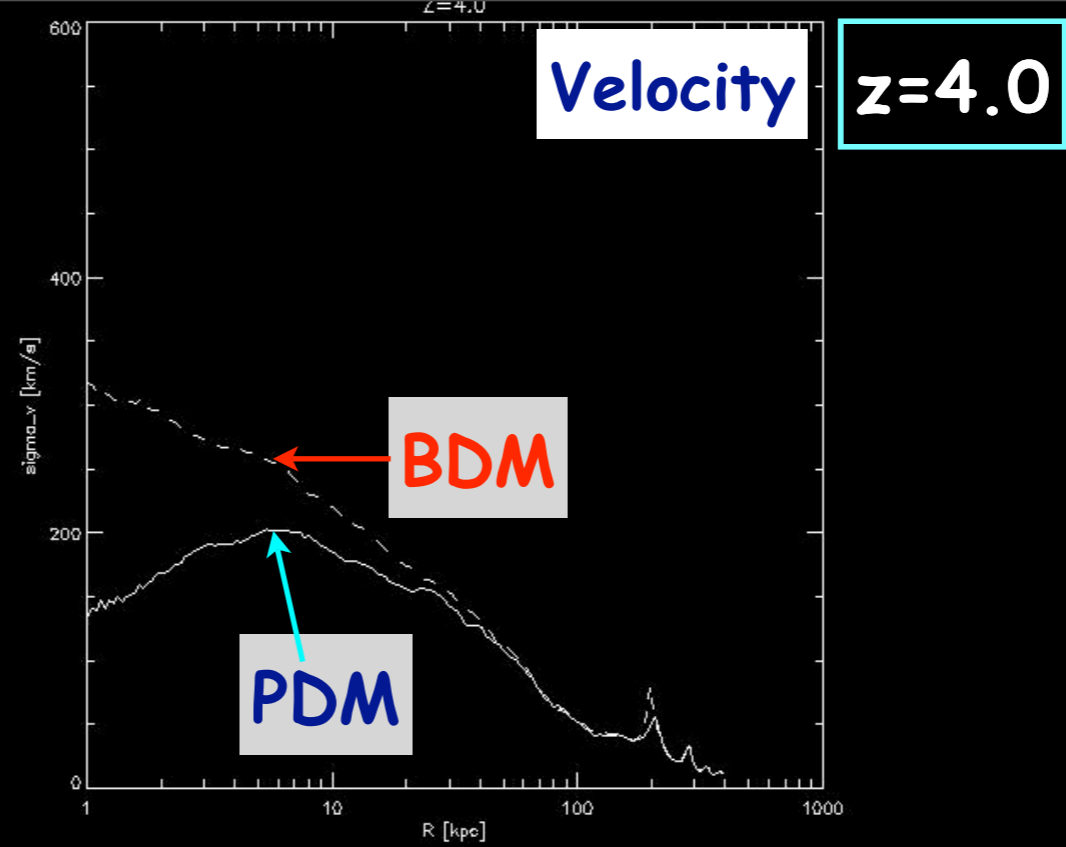
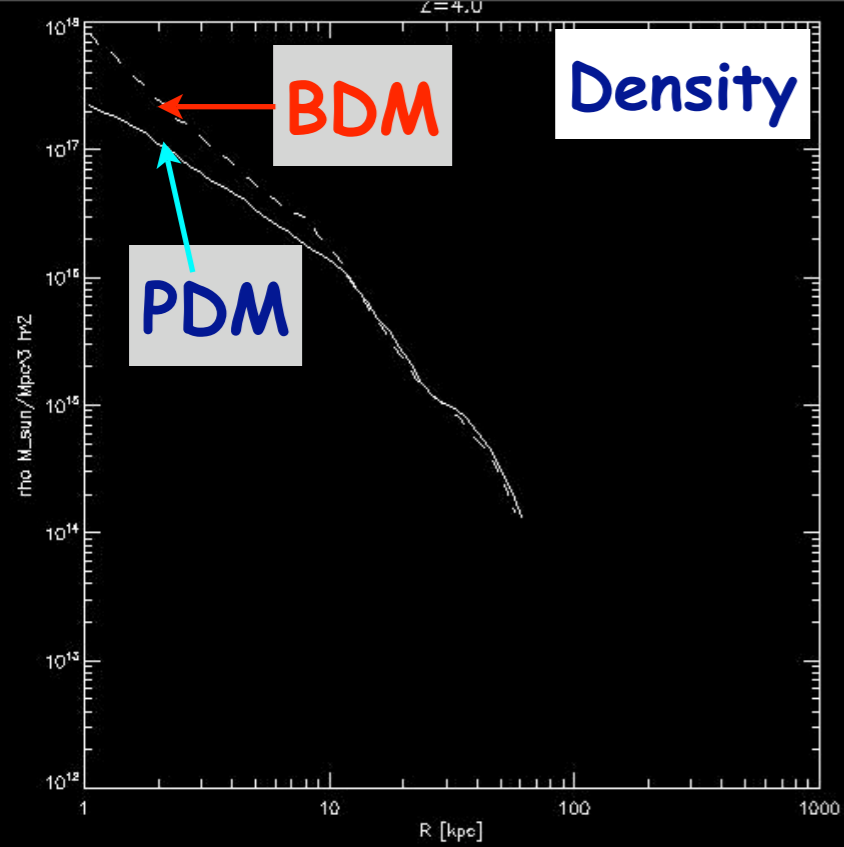
Meudon (June 2010)

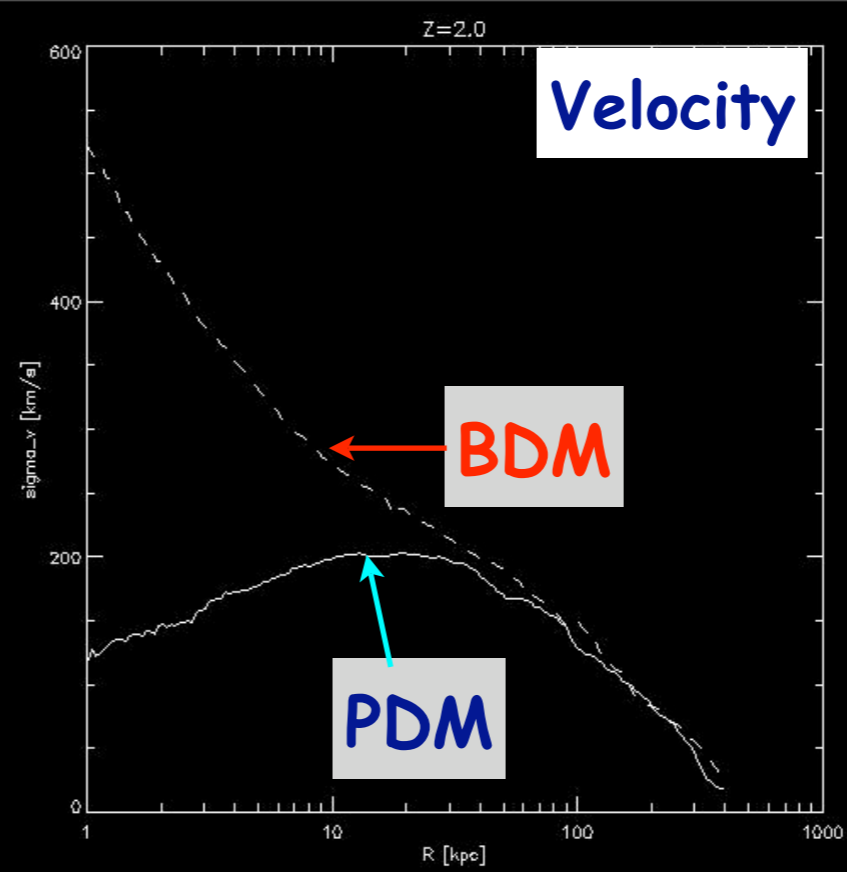
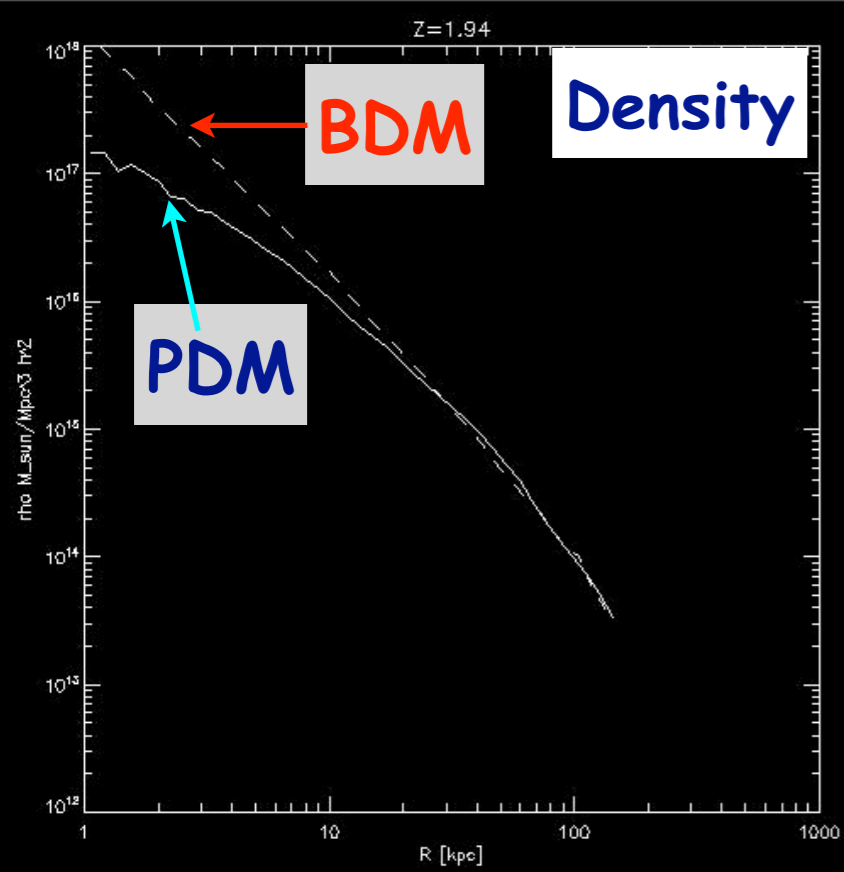


The state of the baryons:

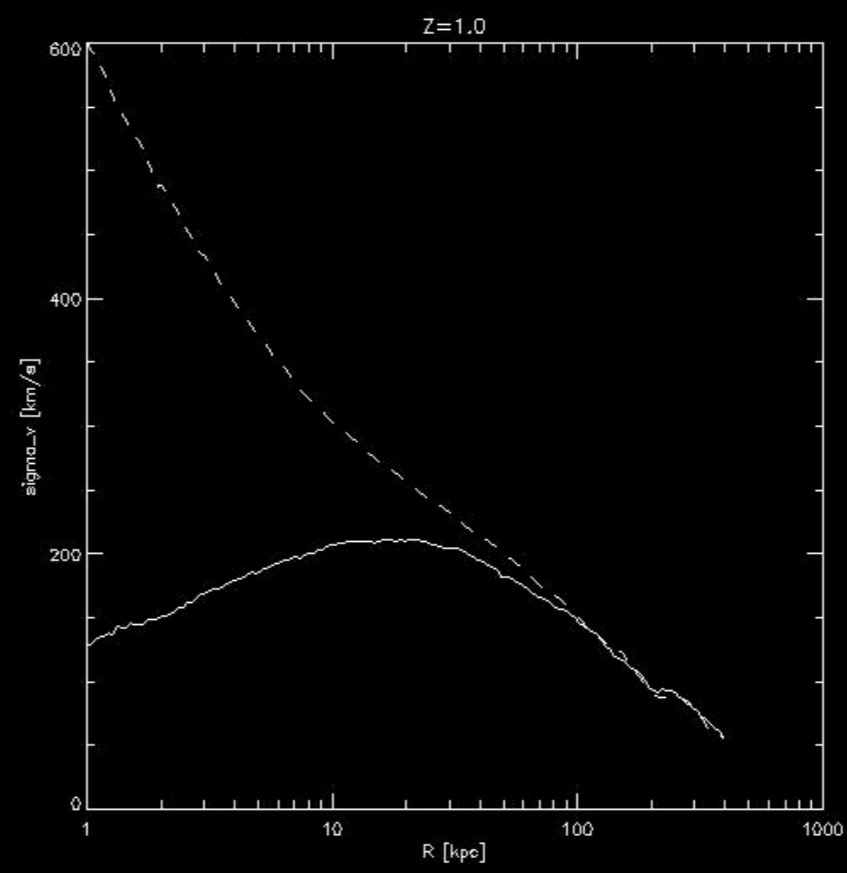
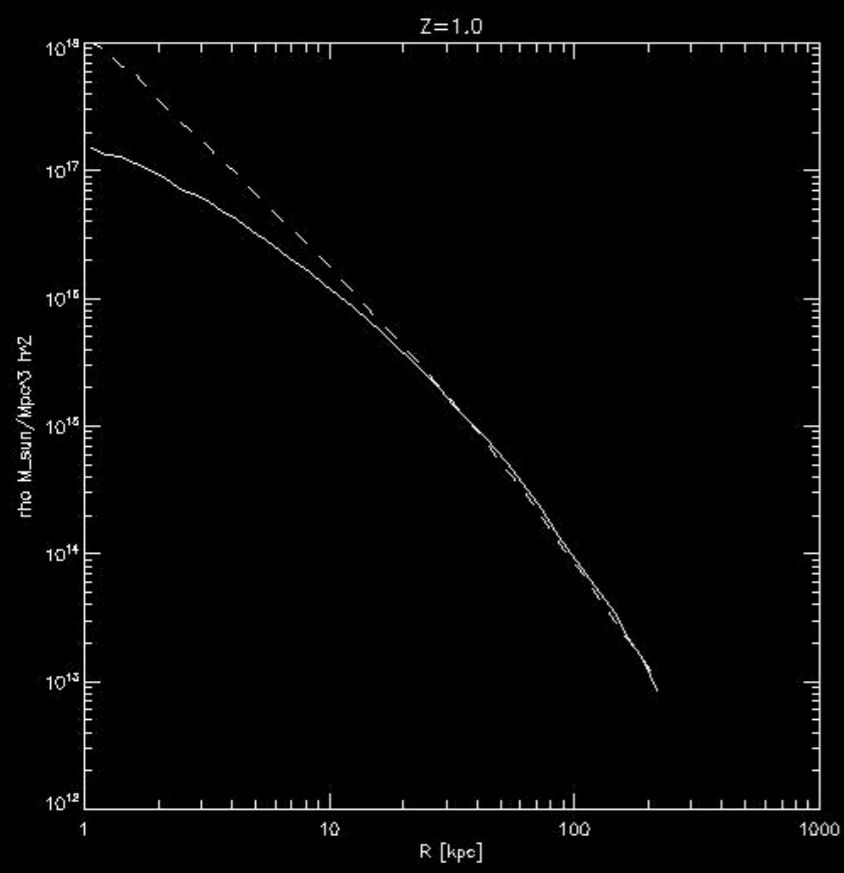
$R < 10 \text{ kpc}$



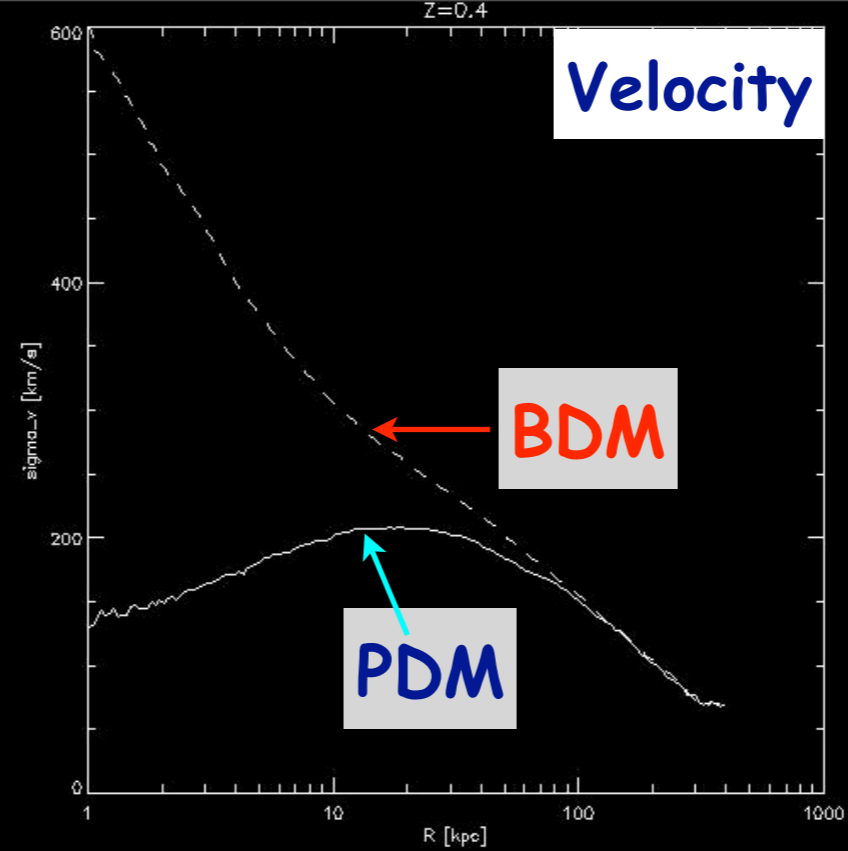
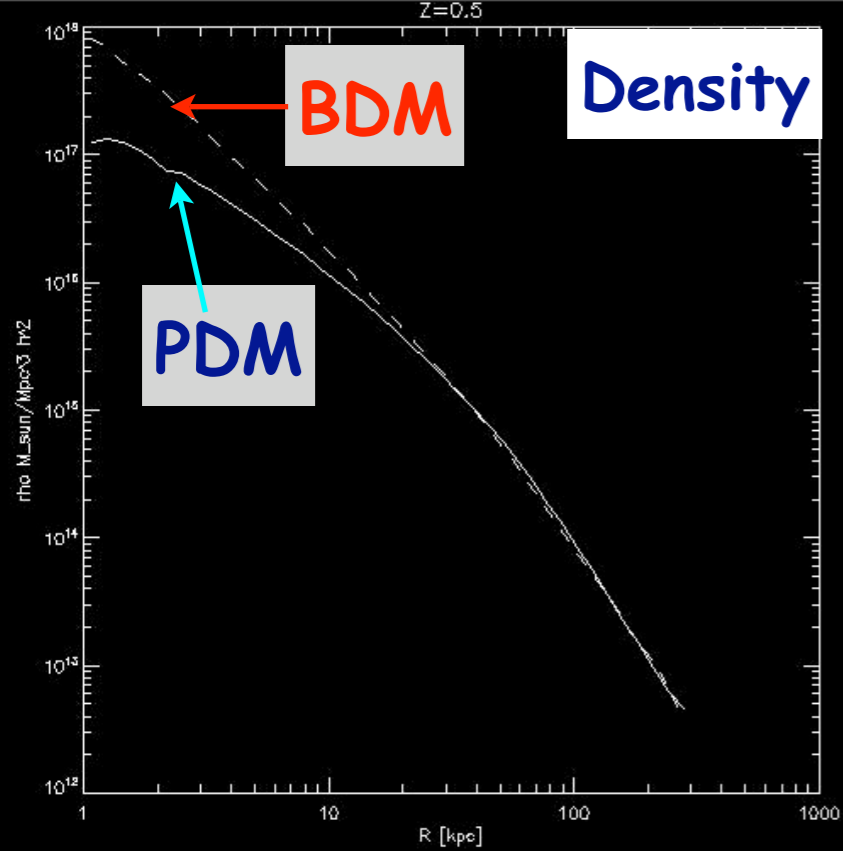




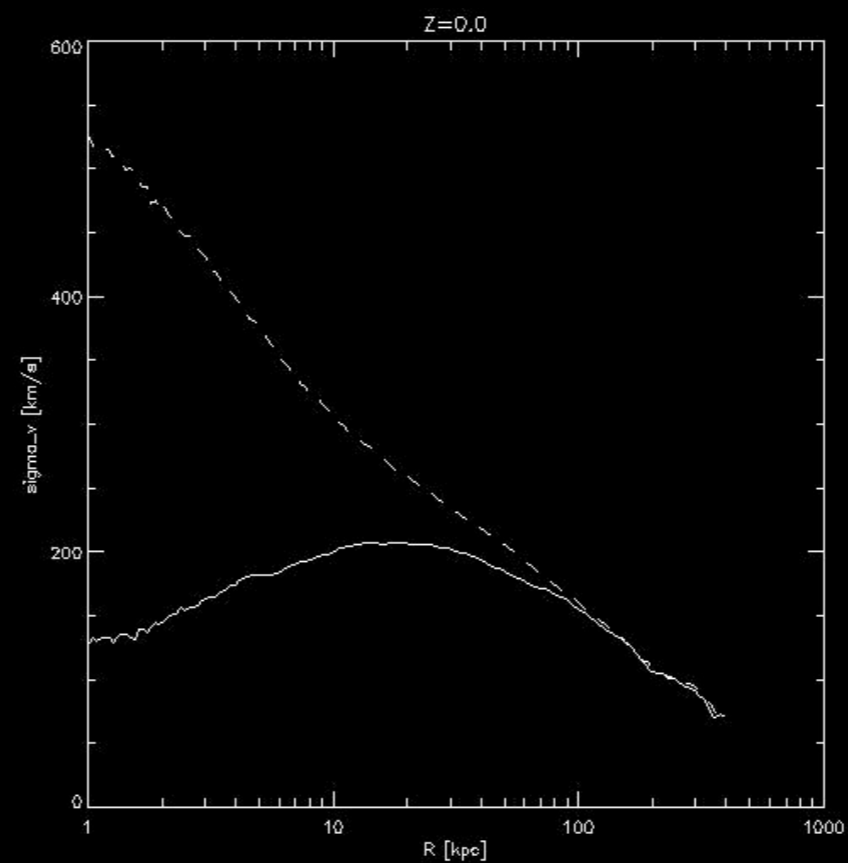
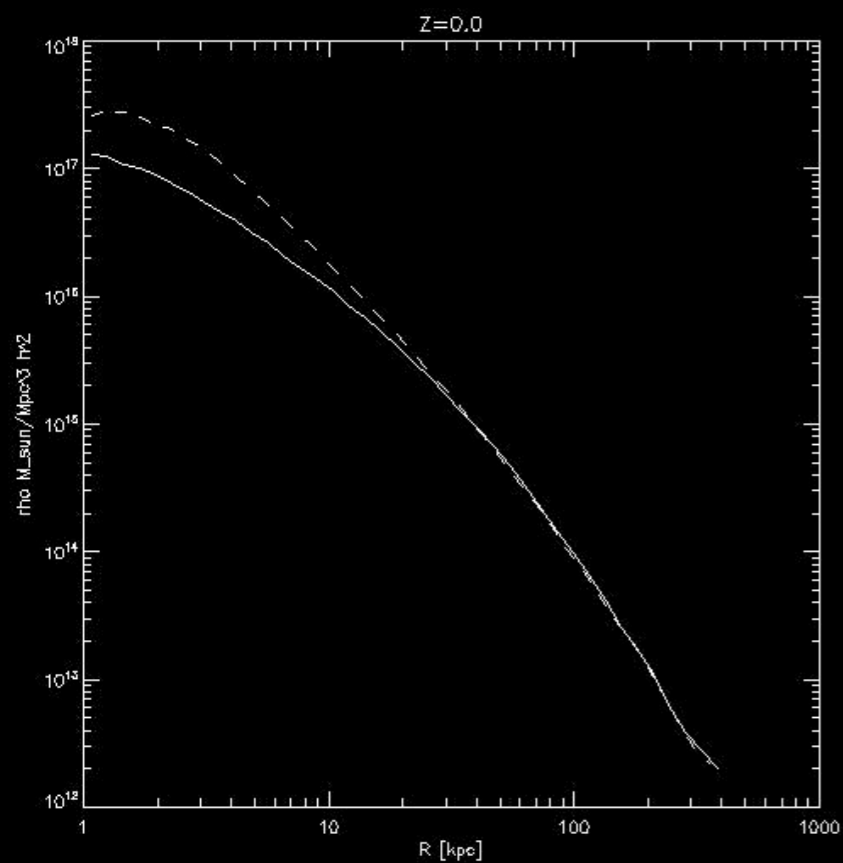
z=2.0



z=1.0

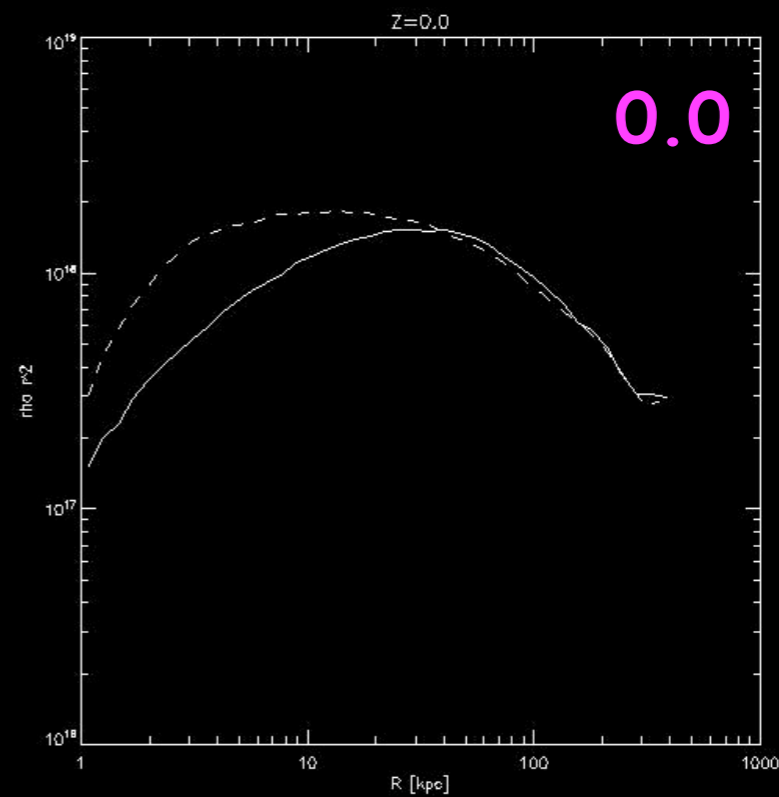
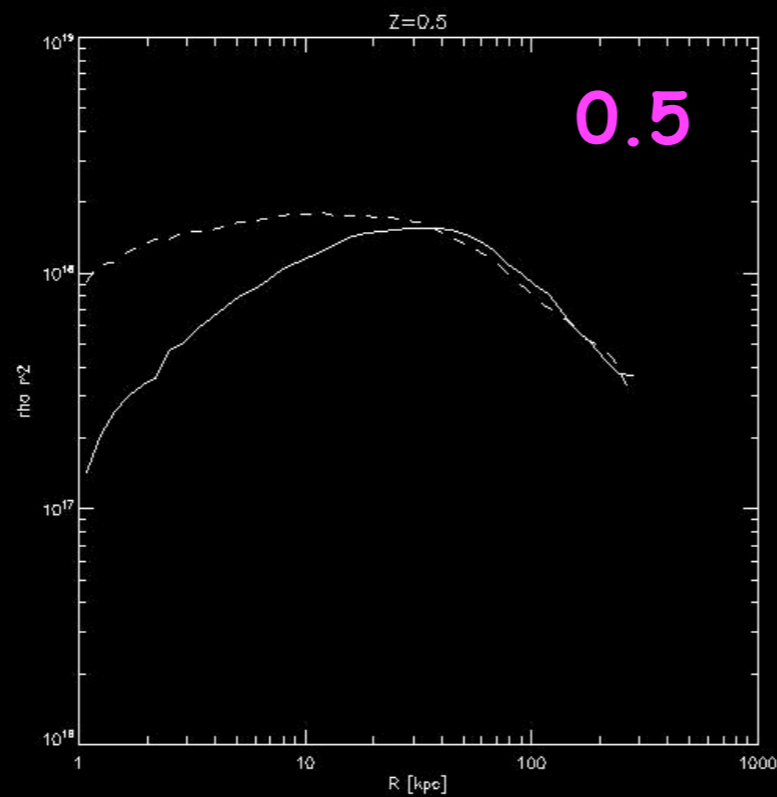
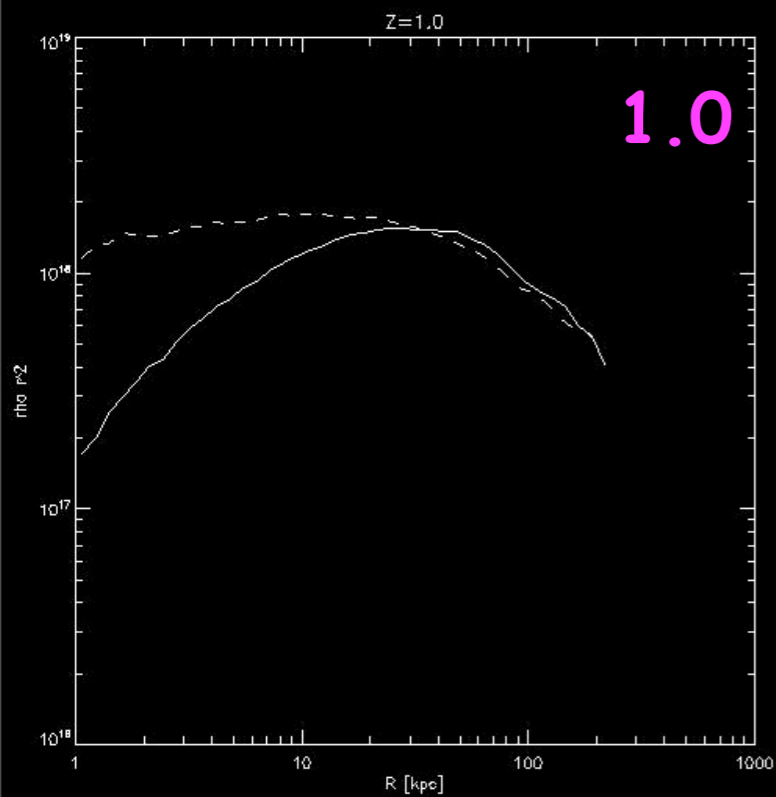
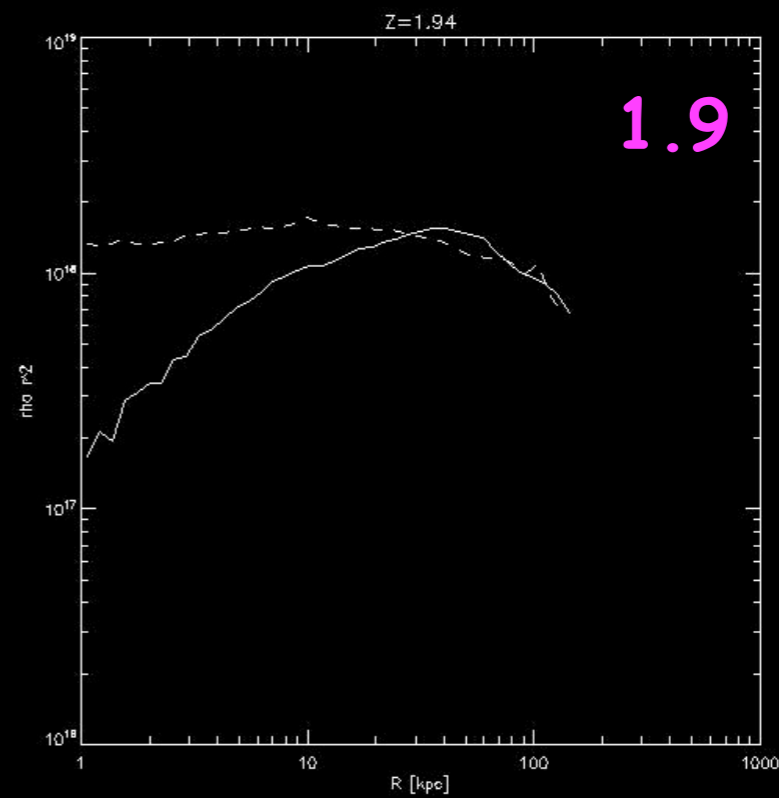
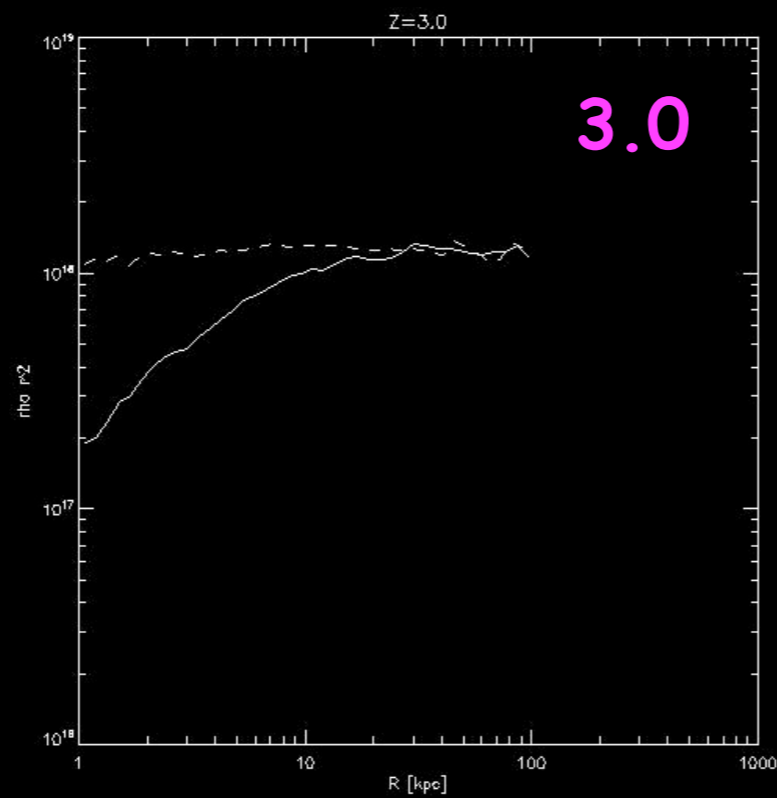
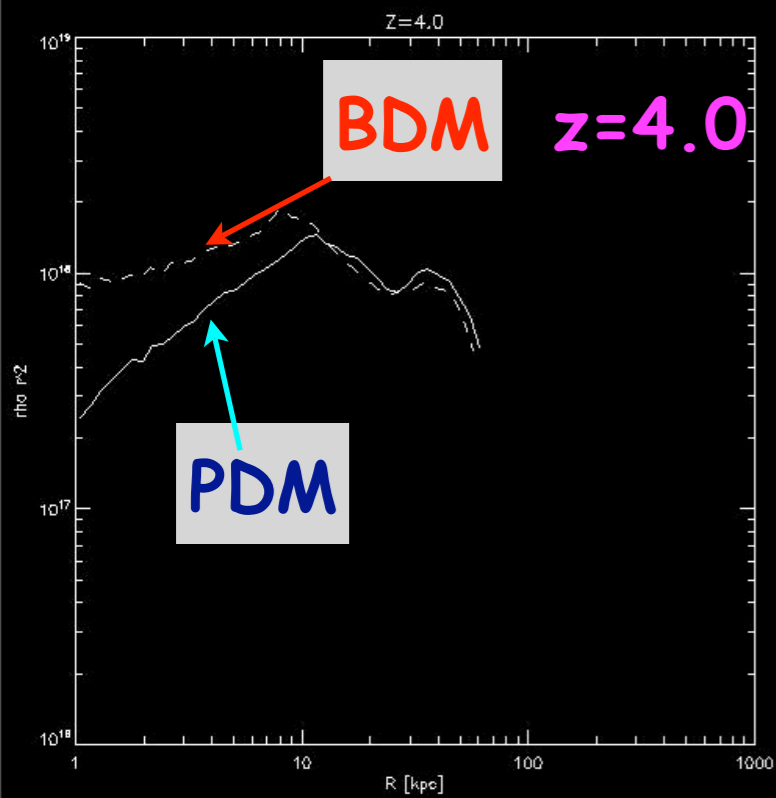


z=0.5

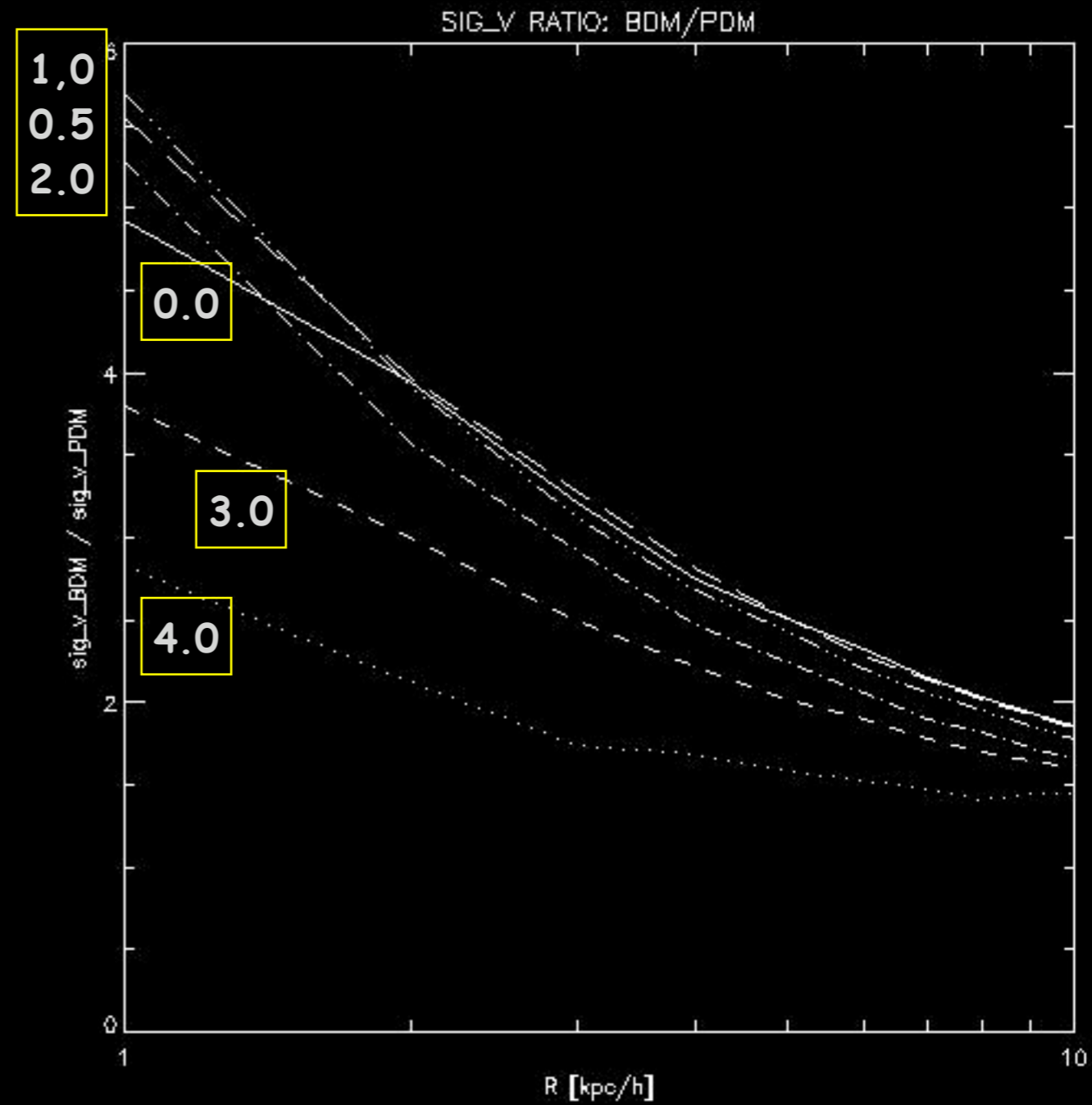
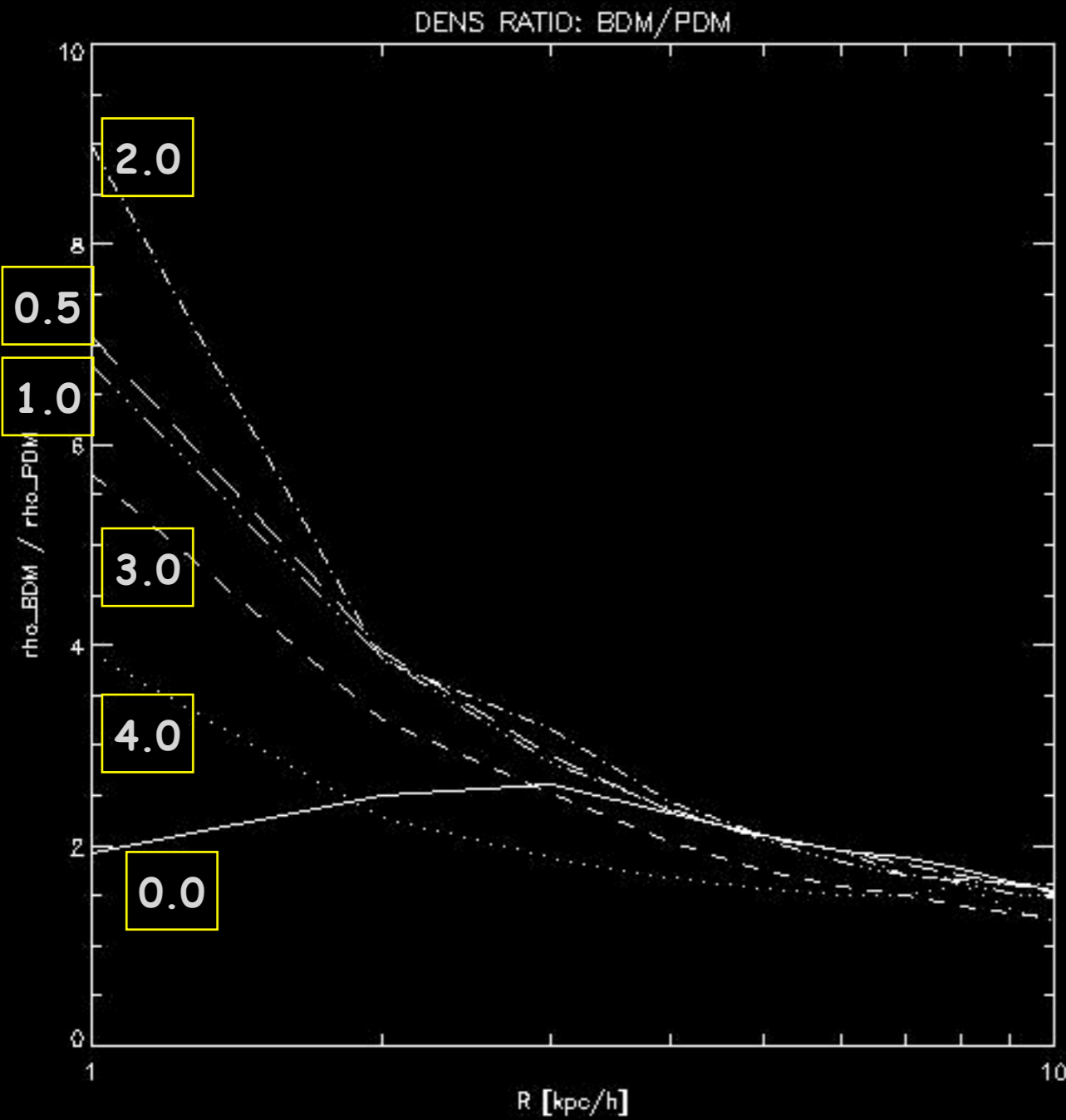


z=0.0

$$\rho r^2$$



Z



1st Interim Report

- Adiabatic contraction works:
 1. At $z \sim 4$ the DM density profile is almost isothermal
 2. At all times BDM density profile exceeds the PDM (at small r 's).
- Dynamical friction works:
 1. At $z < 1$ the BDM density profile flattens
 2. The excess of DM (in the BDM vs PDM) at the center decreases in time
- But the total effect depends on the details: mostly on the feedback - controls the amount of baryons in the substructures

SUBSTRUCTURES

$z = 2.33$

$z = 2.33$

PDM

$z = 2.33$

$z = 2.33$

BDM

600 kpc

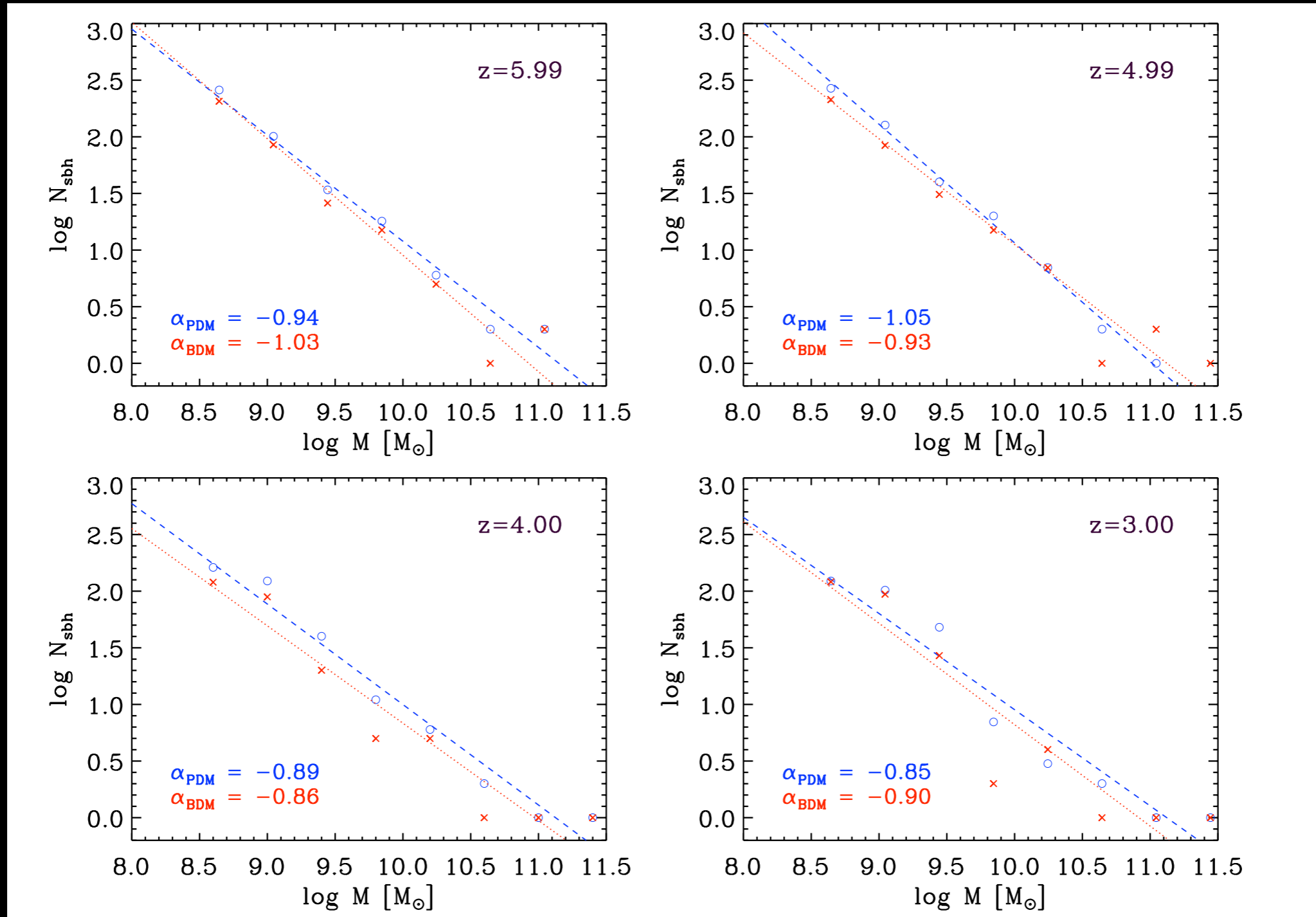
Two competing effects:

- The dissipative gas makes the DM substructures more resilient to tidal disruption.
- The host halo has a deeper potential well - hence stronger tidal field.

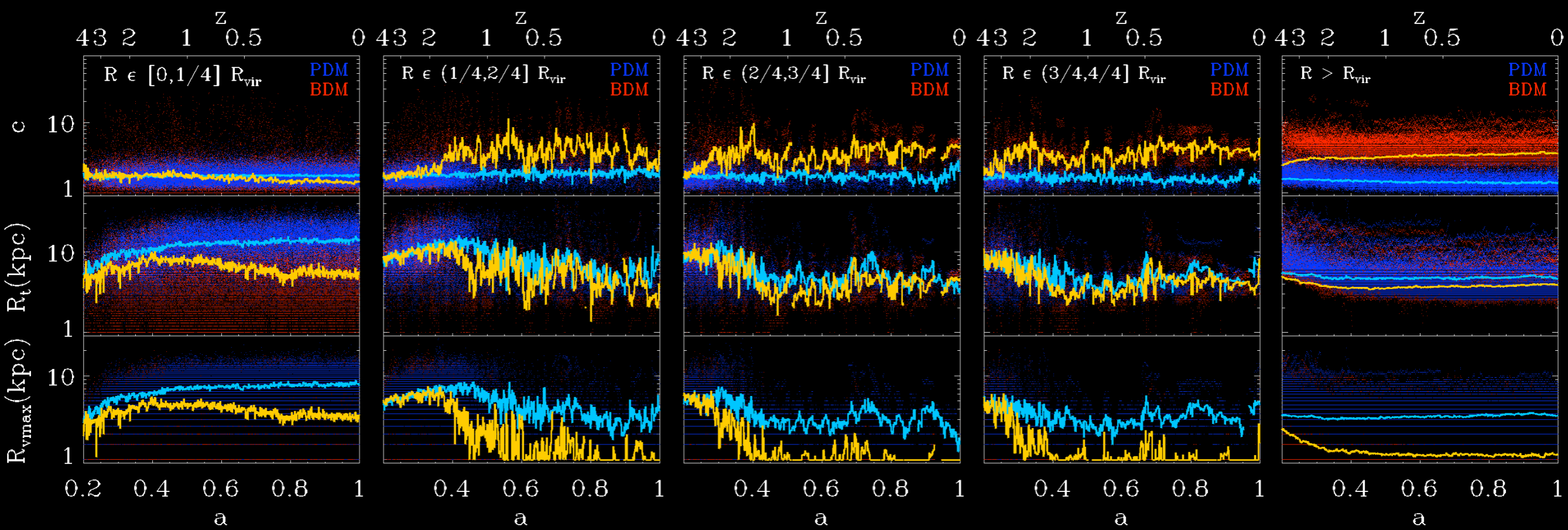
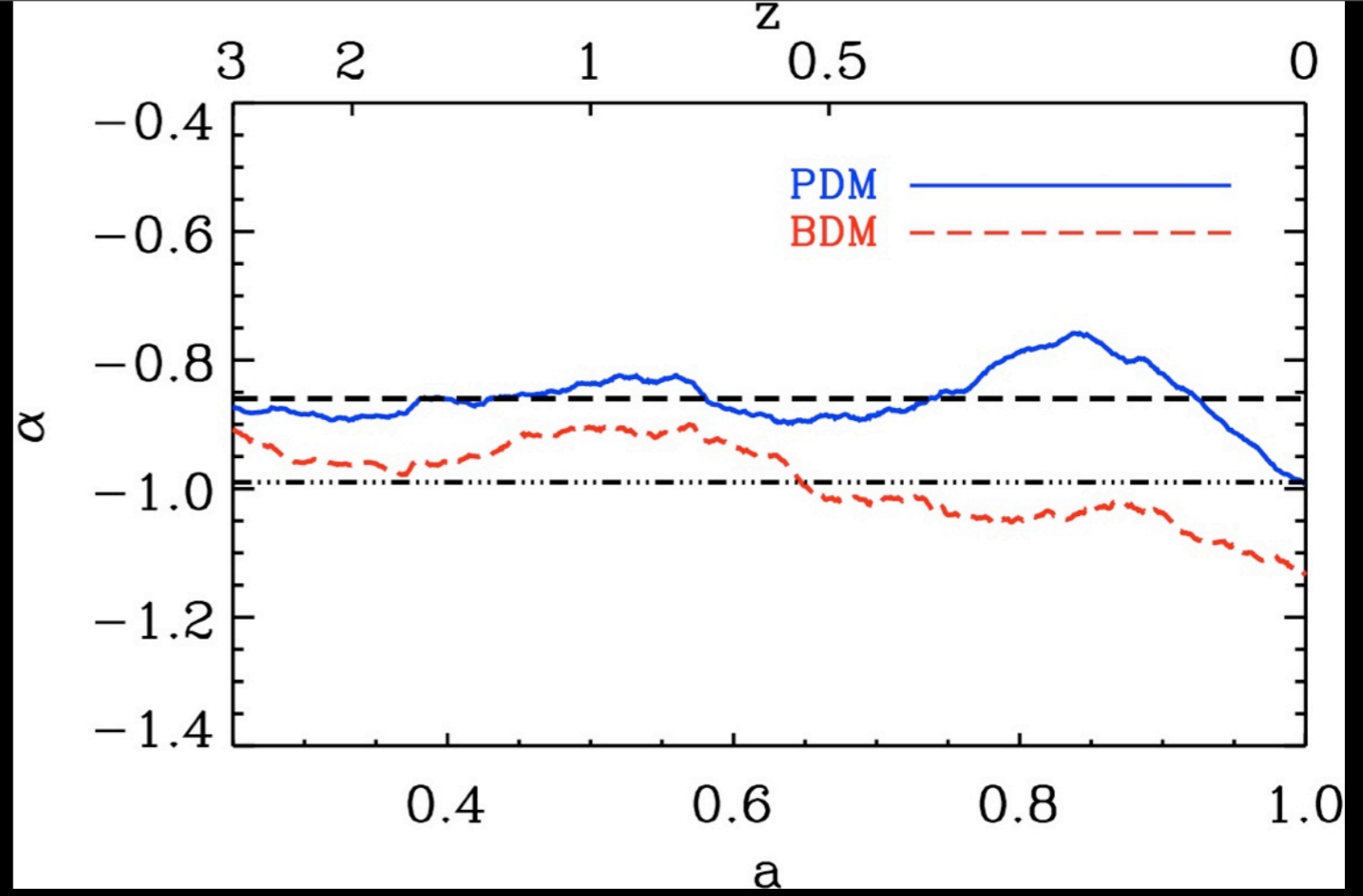
And the winner is ...?

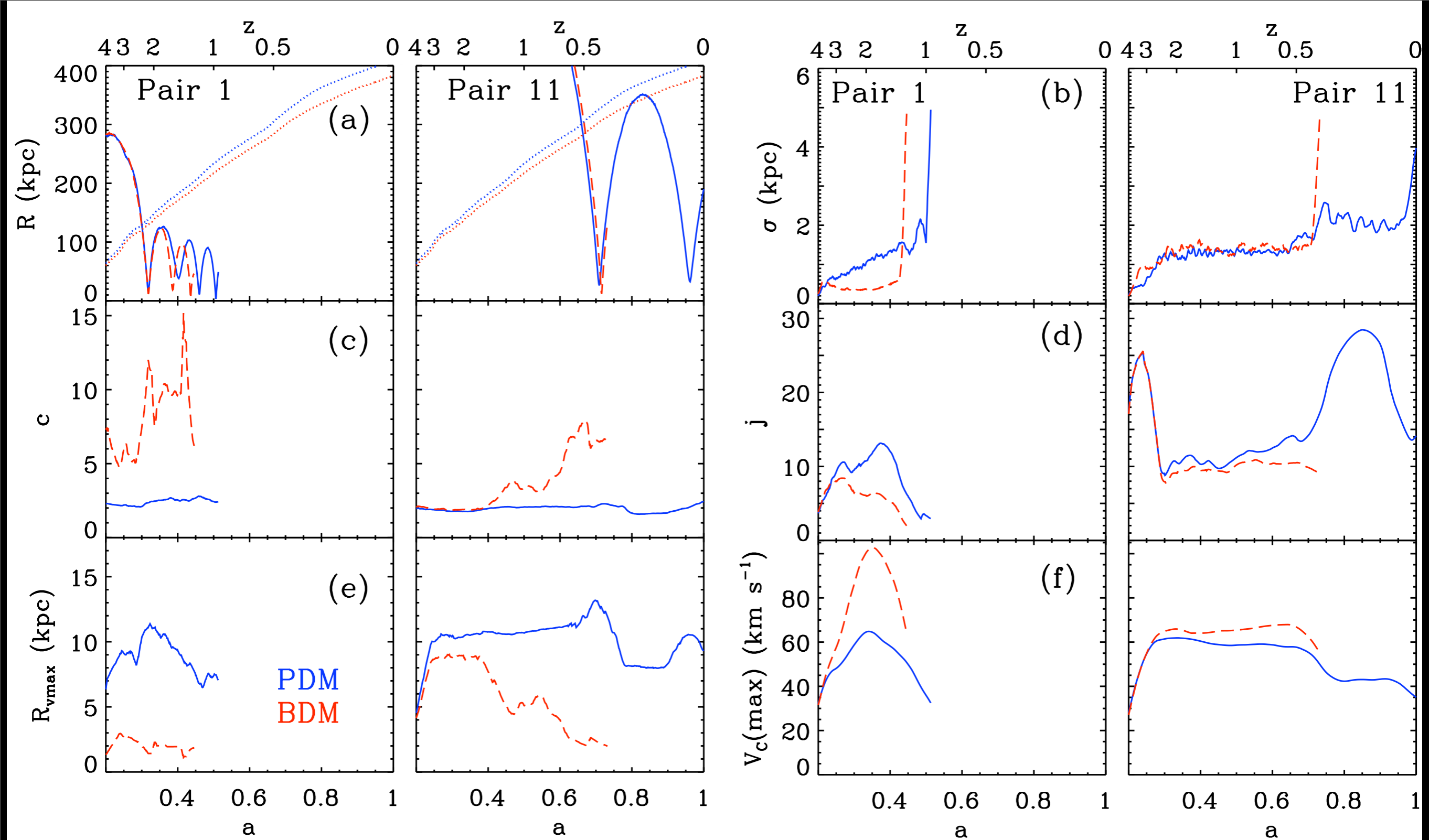
Meudon (June 2010)

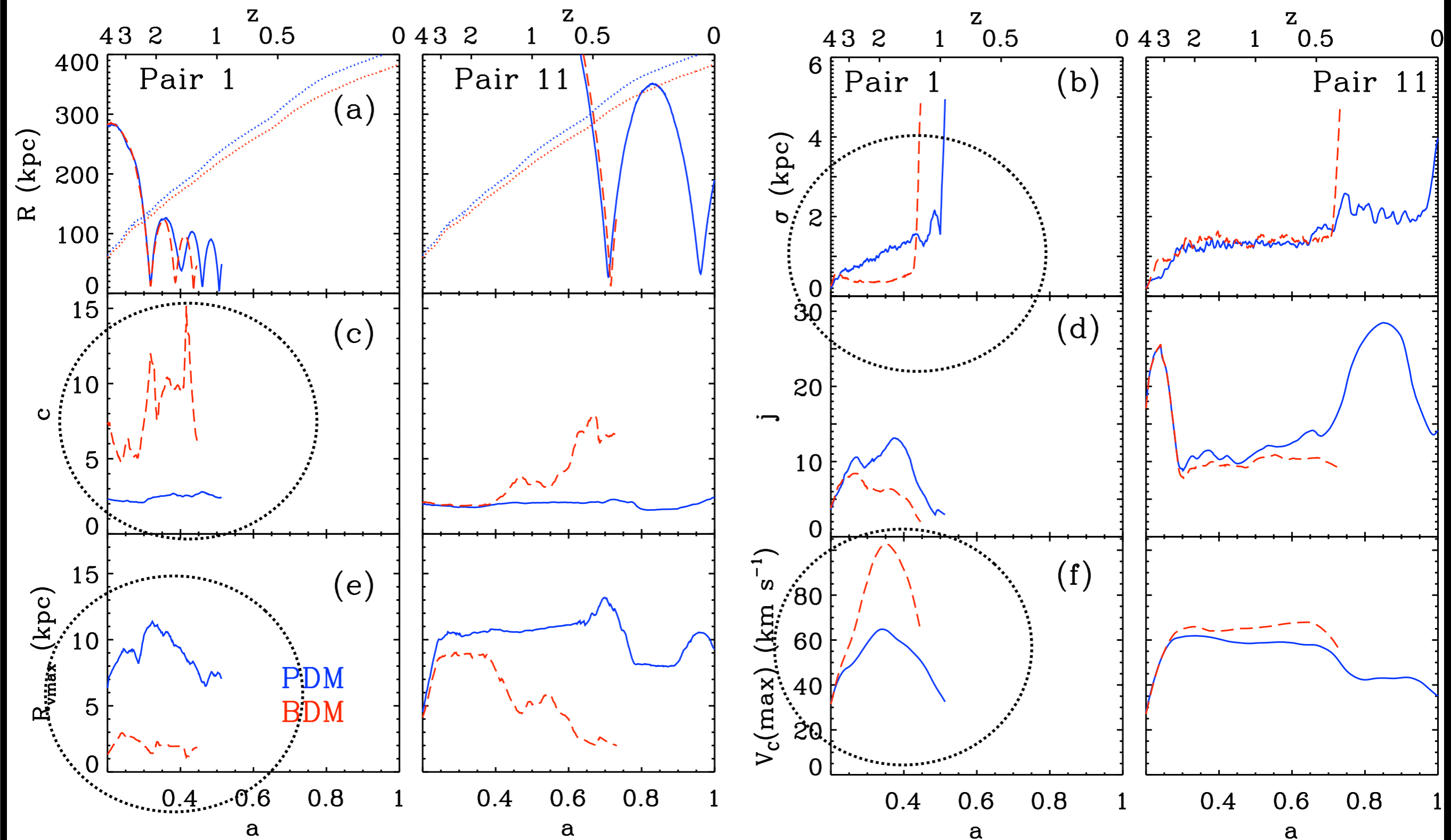
Substructure (DM) mass function:



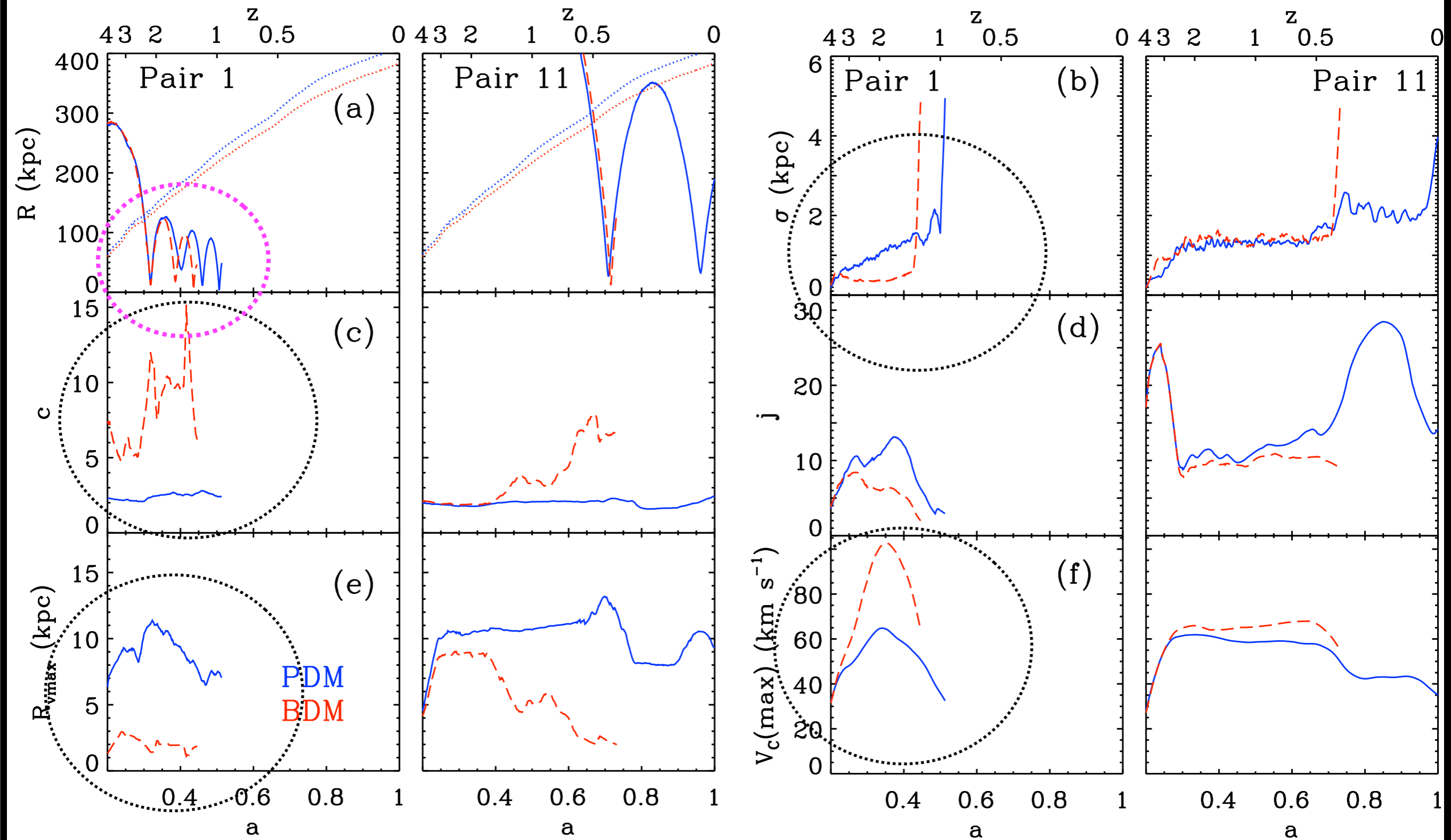
Time evolution of substructures parameters:





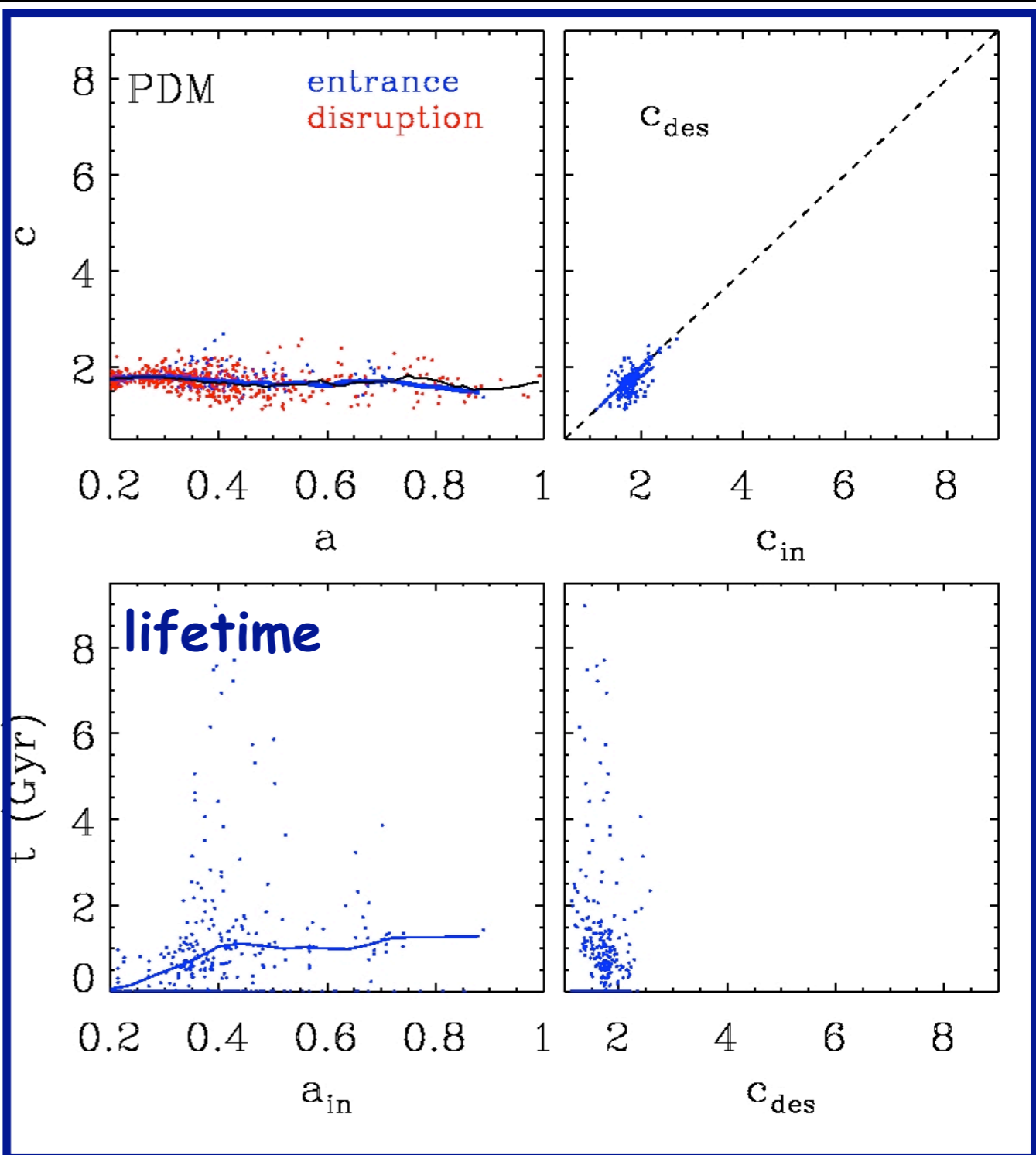
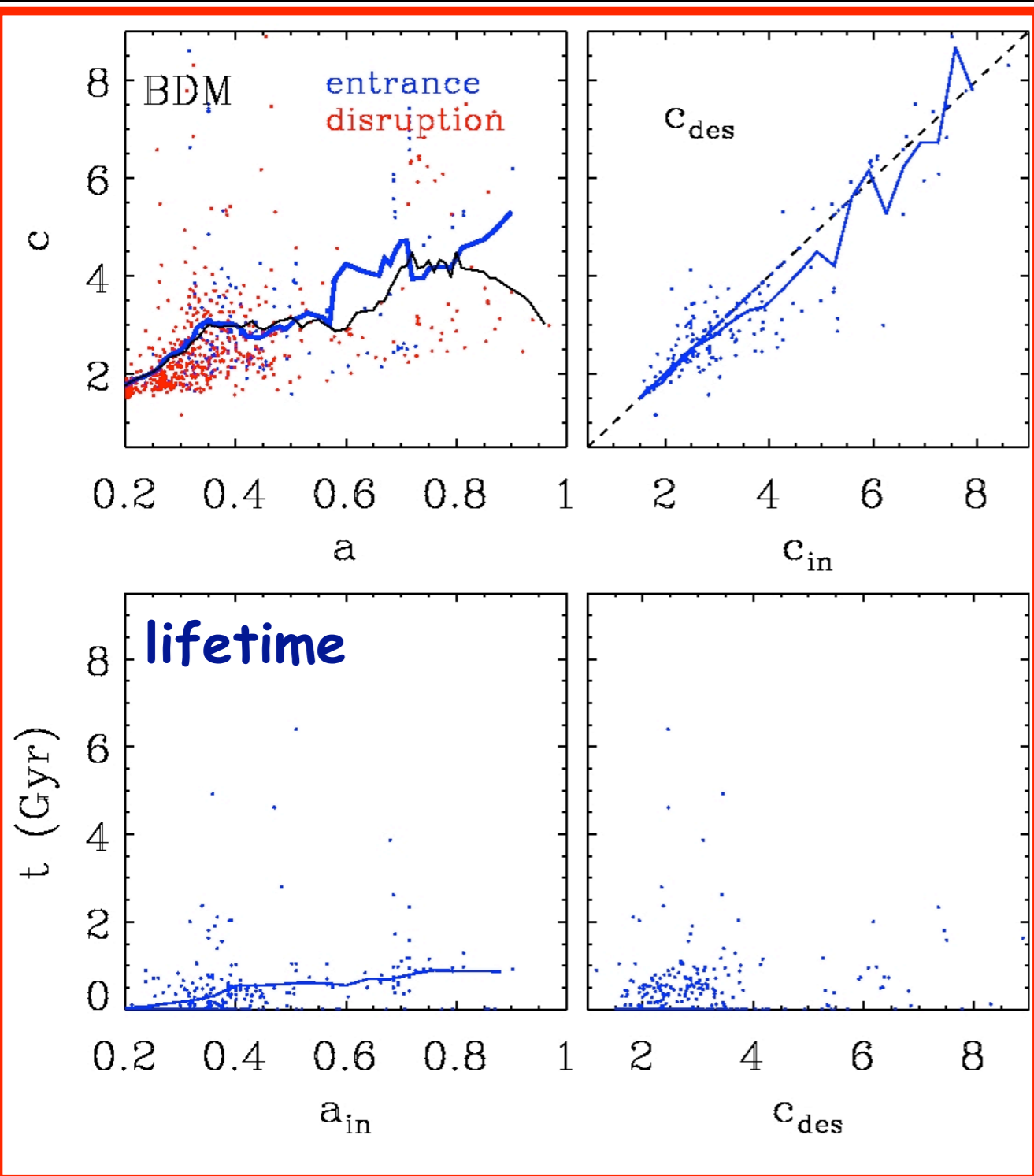


BDM subhalos are more tightly bound



BDM subhalos are more tightly bound

BDM subhalos loose orbital energy faster



2nd Interim Report: Substructures

- Subhalos mass function: $n_{\text{BDM}}(M) \propto n_{\text{PDM}}(M) \propto M^{-\alpha}$ ($\alpha \sim 0.9$)
- BDM subhalos are more tightly bound, but so is the host halo:
Compared with the PDM the BDM subhalos
 1. 'die younger'
 2. loose more of their mass
 3. loose more of their orbital energy
 4. population is depleted faster
- But the total effect depends on the details: mostly on the feedback - controls the amount of baryons in the substructures



AIP



CLUES

People

Simulations

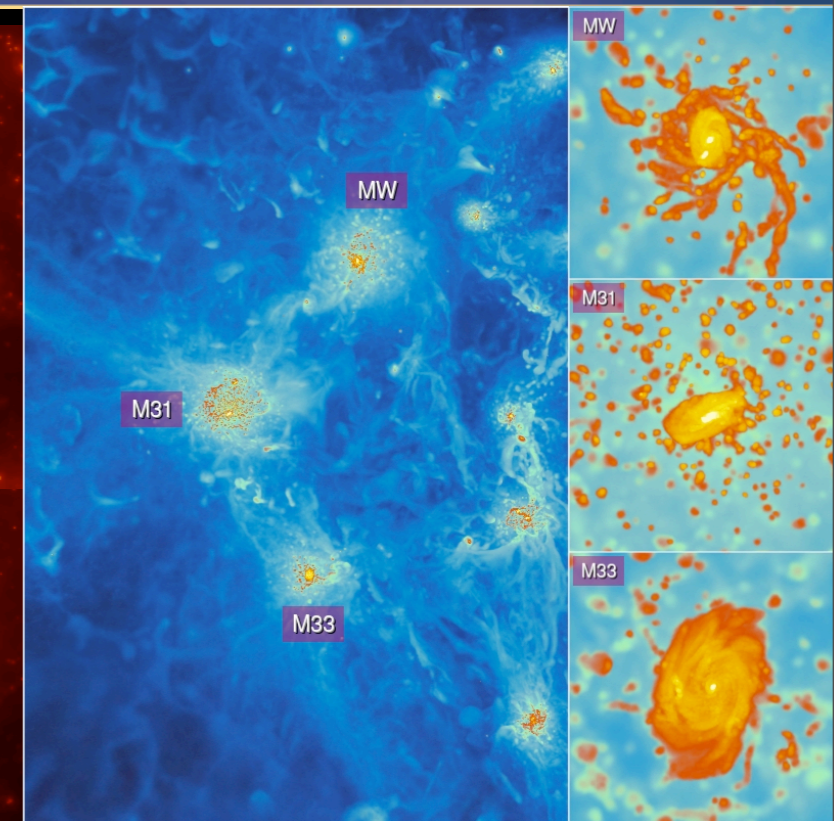
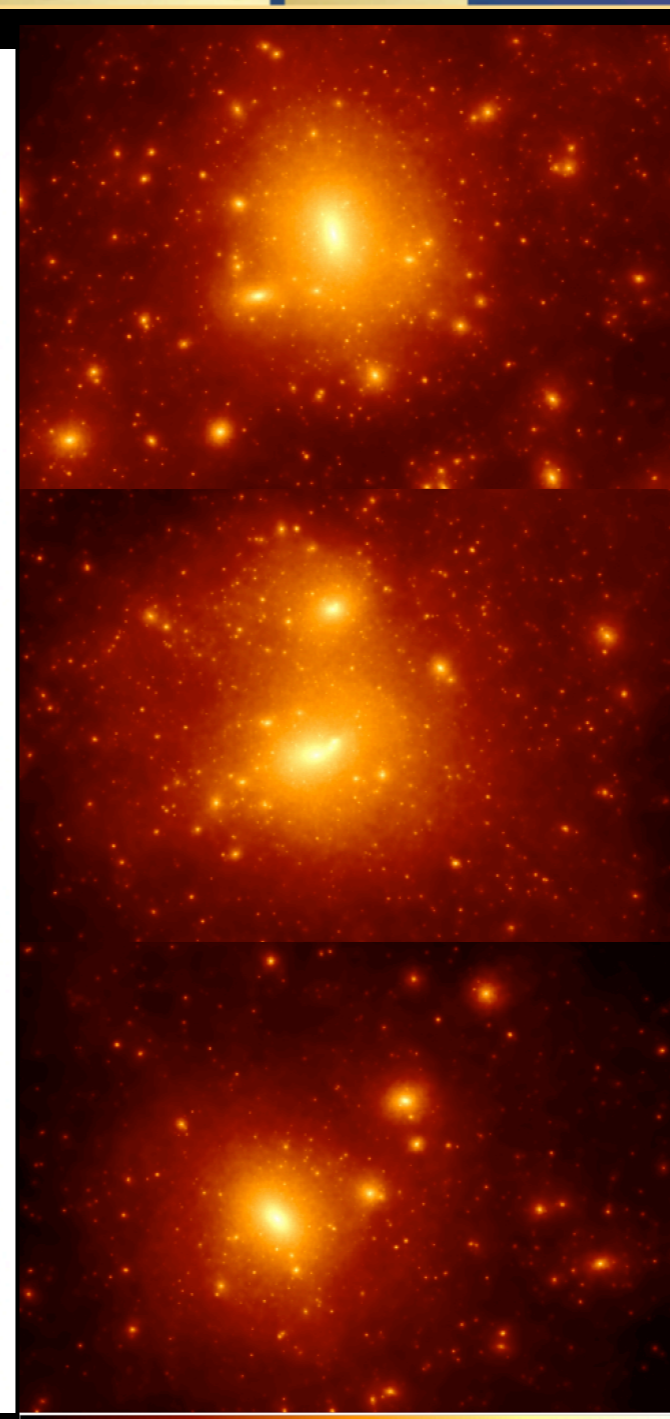
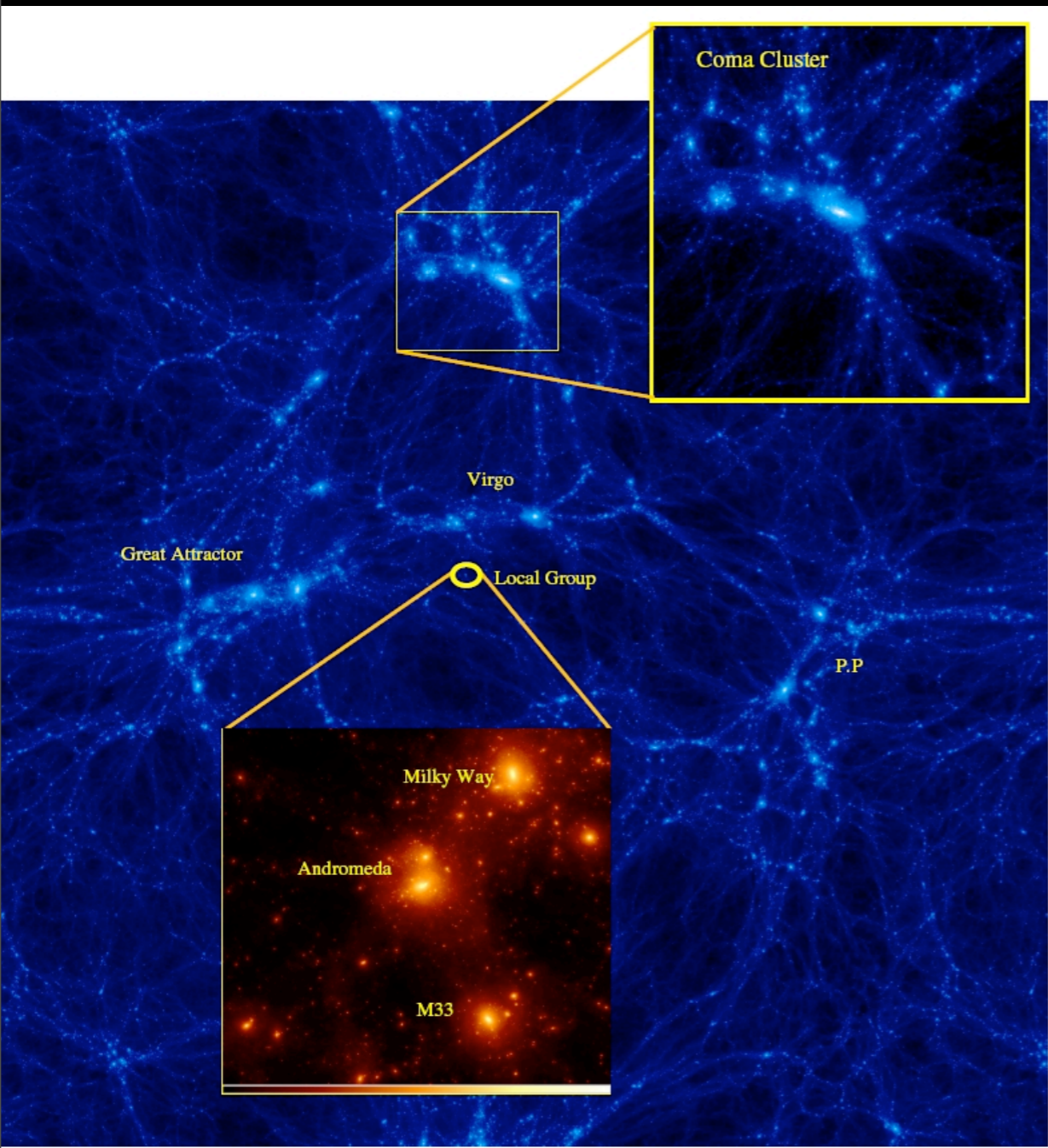
Talks

Articles

Image Gallery

Movies

<http://www.clues-project.org/>



Stefan Gottloeber
 Anatoly Klypin
 Gustavo Yepes
 YH

The CLUES' WMAP3 LG

Cosmology: WMAP3 ($\sigma_8=0.75$,
 $\Omega_m=0.24$, $h=0.73$)

- LG: MW, M31, M33
- $D(\text{MW}-\text{M31})=0.9 \text{ Mpc}/h$
- $D(\text{LG}-\text{Virgo})=9.7 \text{ Mpc}/h$

	$M_{\text{vir}}, M_{\text{stellar}}$ [$1.e11\text{Msun}/h$]	R_{vir} [kpc/h]	V_{max} [km/s]
M31	5.7	174	128
(BDM)	0.14		182
MW	4.6	162	131
(BDM)	0.12		155
M33	2.2	127	112
(BDM)	0.06		118

Simulations

- Code: GADGET-2
- Halo finder: AHF
- Box: 64 Mpc/h
- zoom: $R=2 \text{ Mpc}/h$
effective 4096^3
- PDM: $DM=2.5e5 \text{ Msun}/h$
- BDM: $DM=2.1e5 \text{ Msun}/h$
gas= $4.4e4 \text{ Msun}/h$
stars= $2.2e4 \text{ Msun}/h$

The CLUES' WMAP3 LG

Cosmology: WMAP3 ($\sigma_8=0.75$,
 $\Omega_m=0.24$, $h=0.73$)

- LG: MW, M31, M33
- $D(\text{MW}-\text{M31})=0.9 \text{ Mpc}/h$
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MW	4.6	162	131
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M33	2.2	127	112
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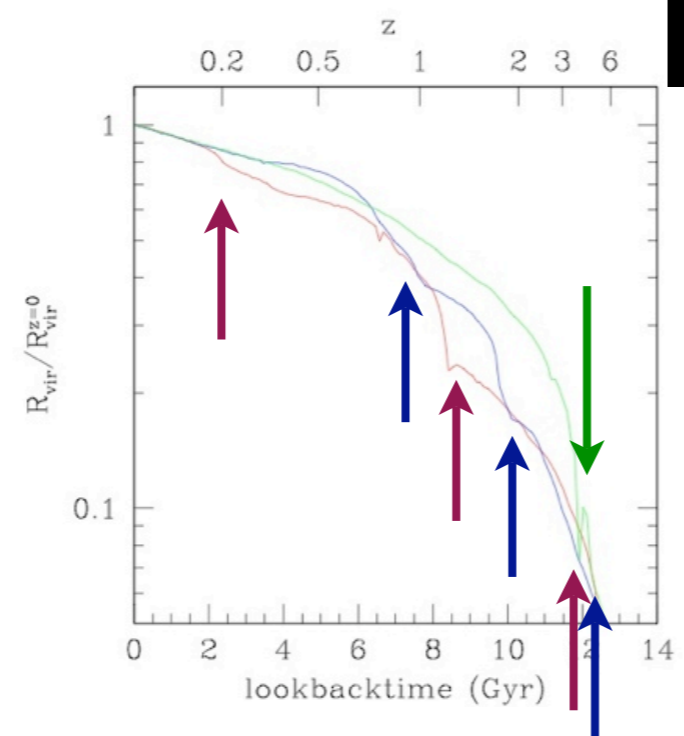
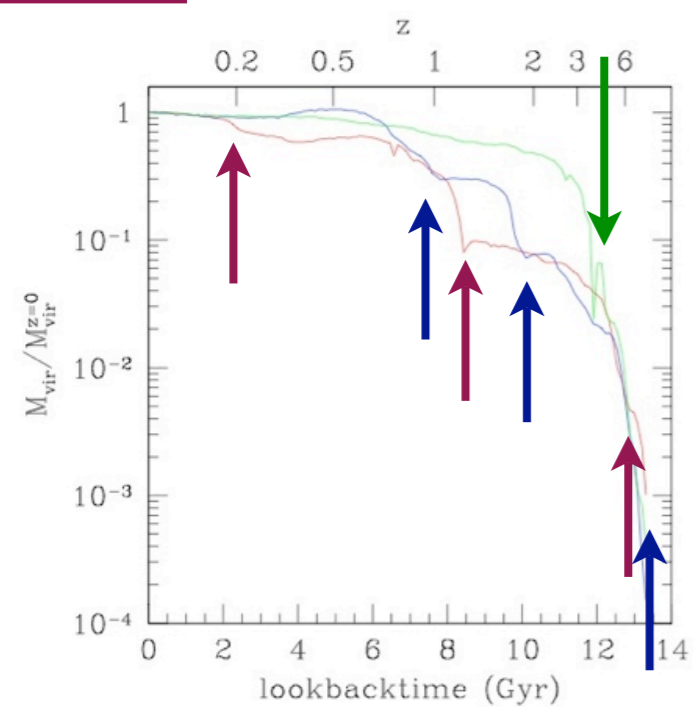
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- BDM: $DM=2.1e5 \text{ Msun}/h$
gas= $4.4e4 \text{ Msun}/h$
stars= $2.2e4 \text{ Msun}/h$

Work in progress ...

VIR

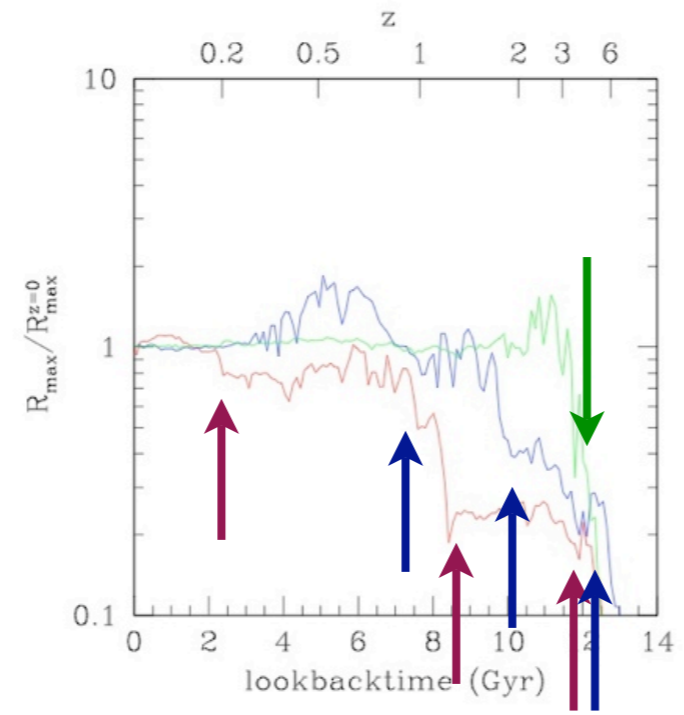
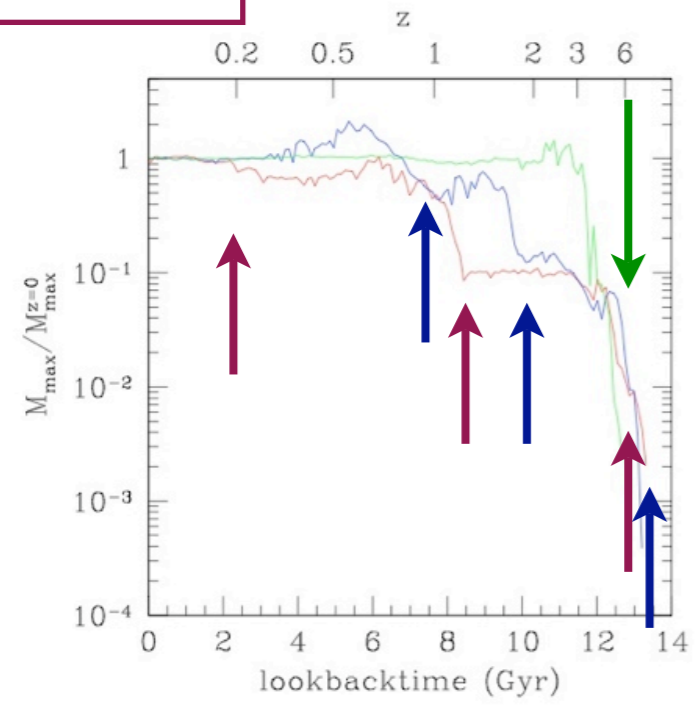
Mass Accretion History



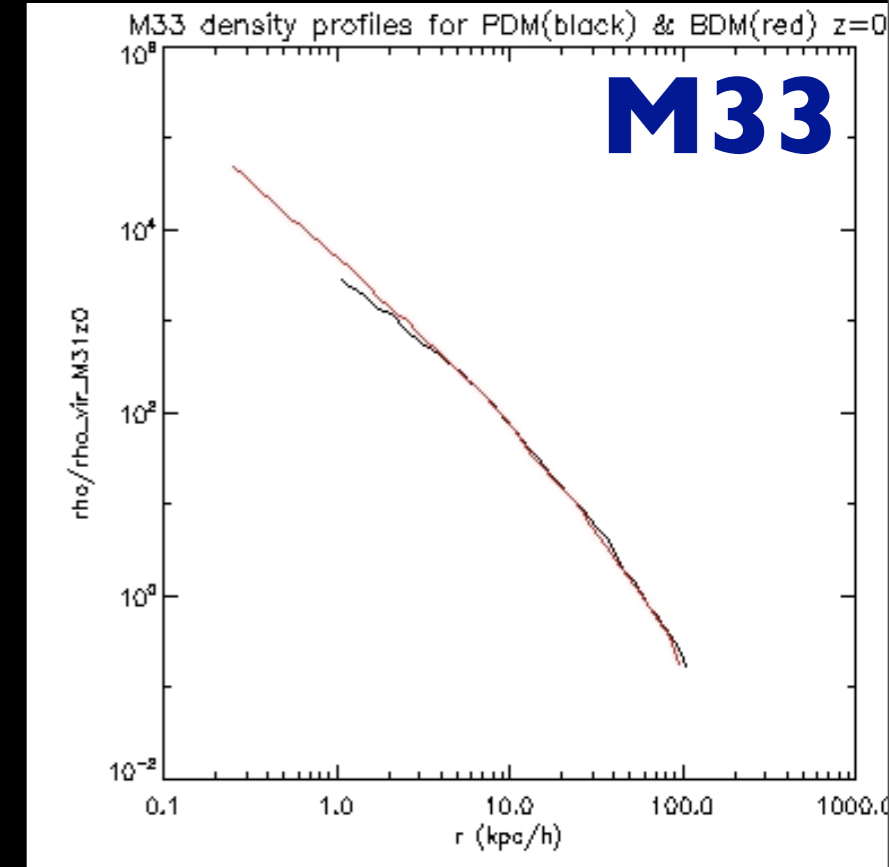
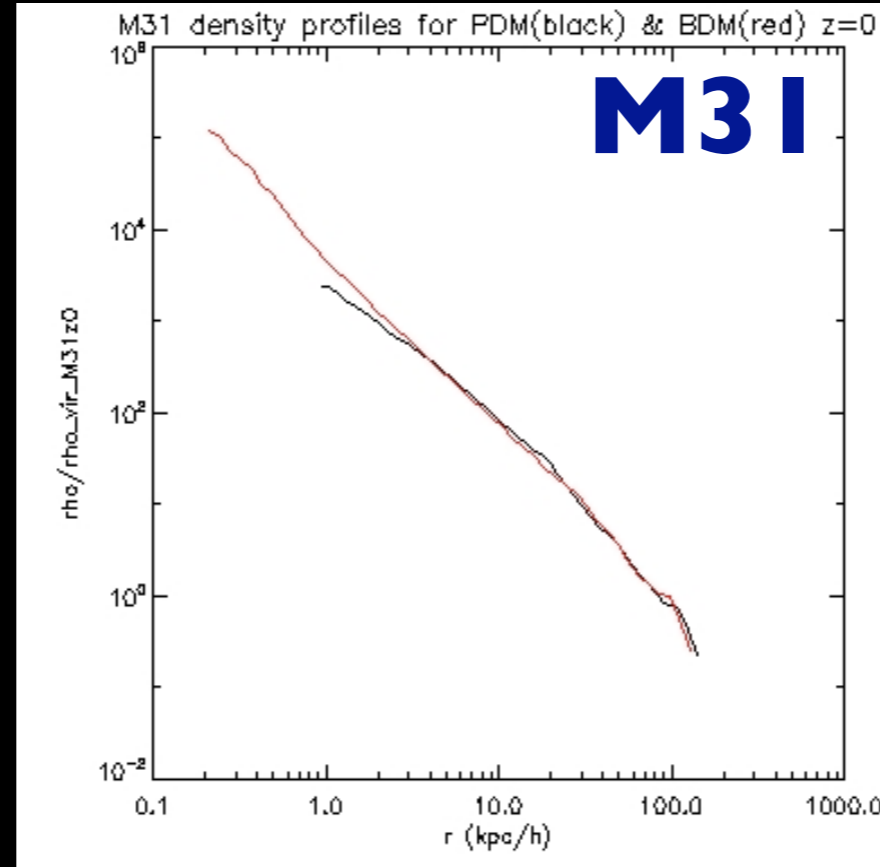
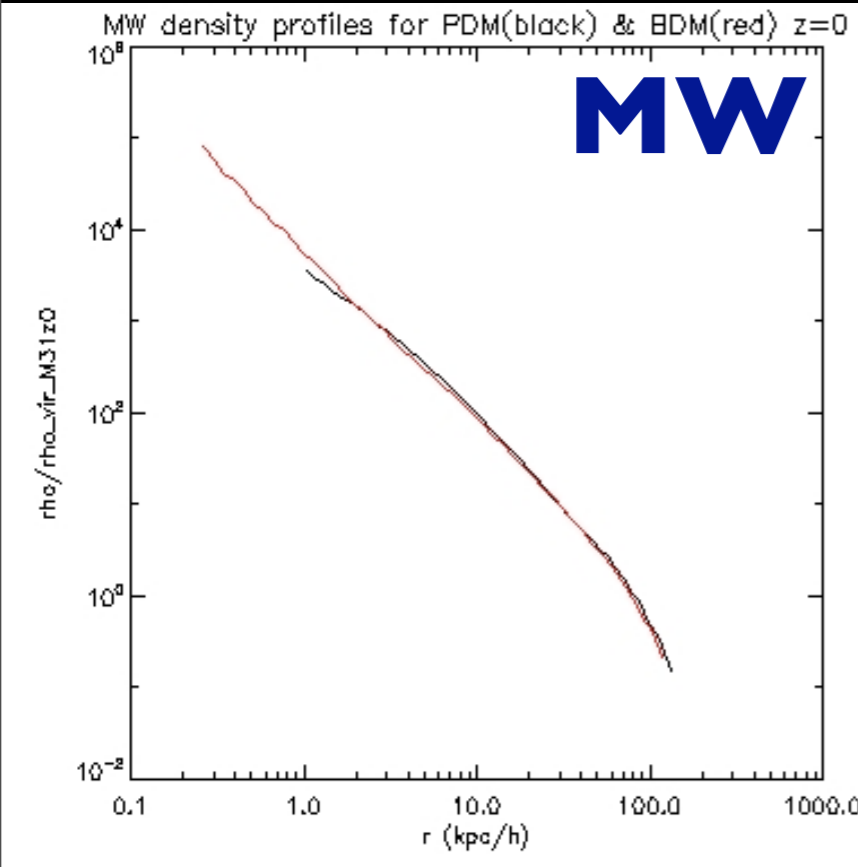
M31 ———
 MW ———
 M33 ———

Arrows indicate epochs of major mergers

MAX



DM density profiles (z=0):



Apparent disagreement with Romano-Diaz et al - no flattening of $\rho(r)$

- Different codes (numerics, resolution, ...)
- Different physics (feedback, ...)
- Different halos:
 - ◆ mass
 - ◆ merging history

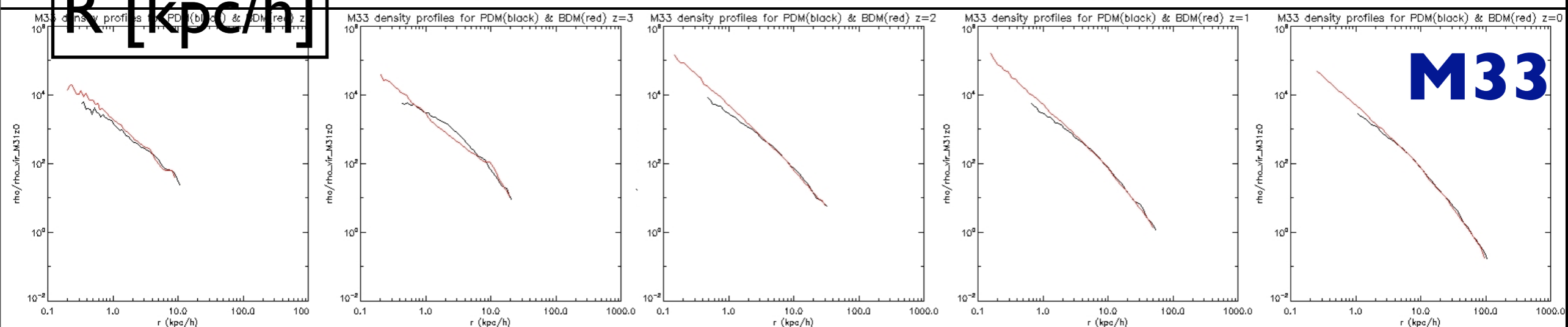
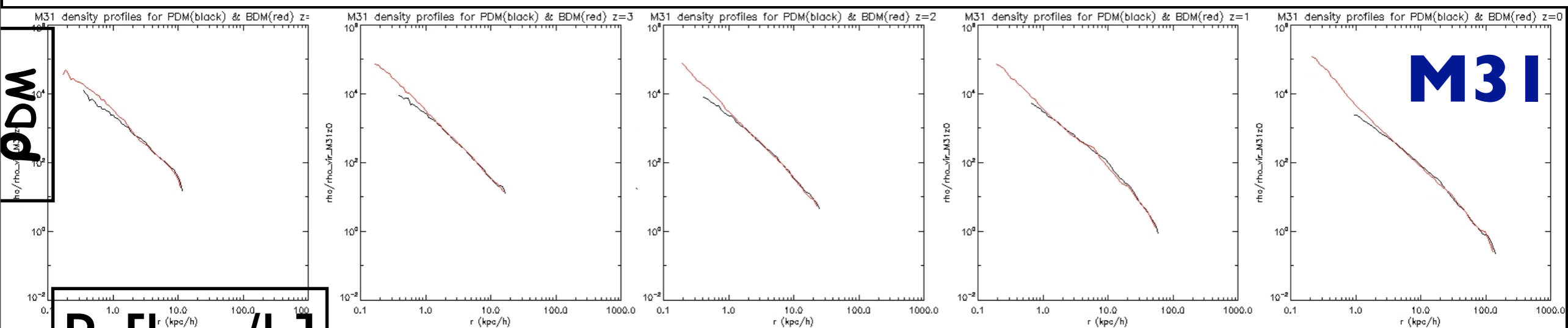
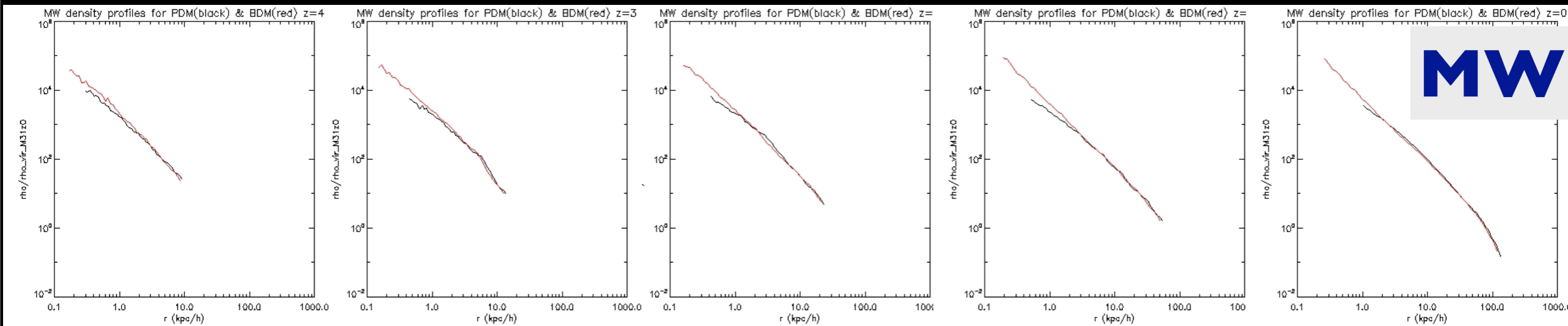
Z= 4

3

2

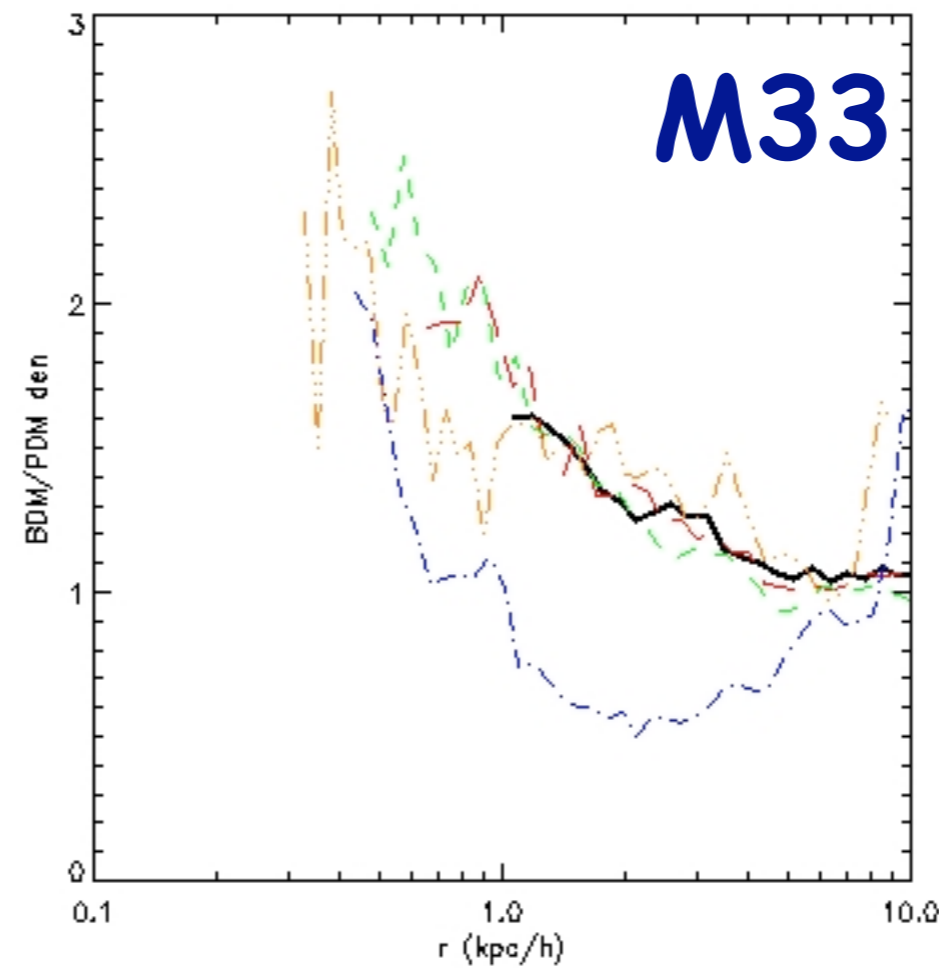
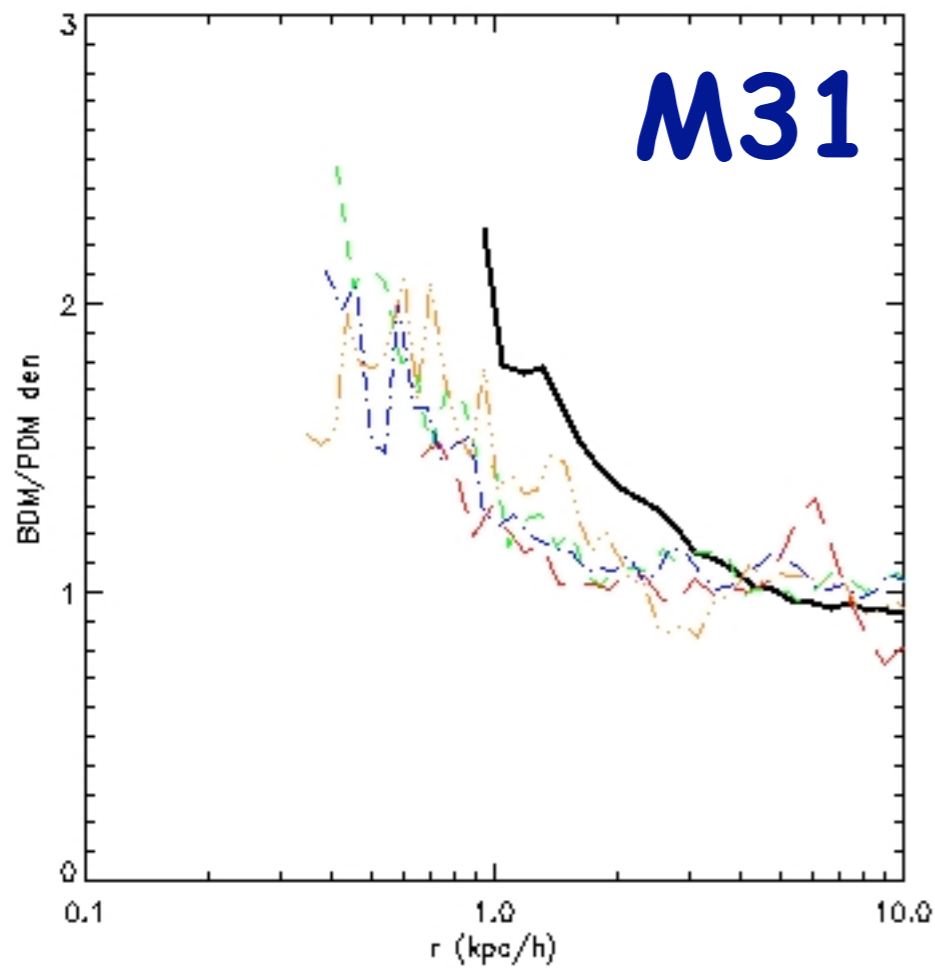
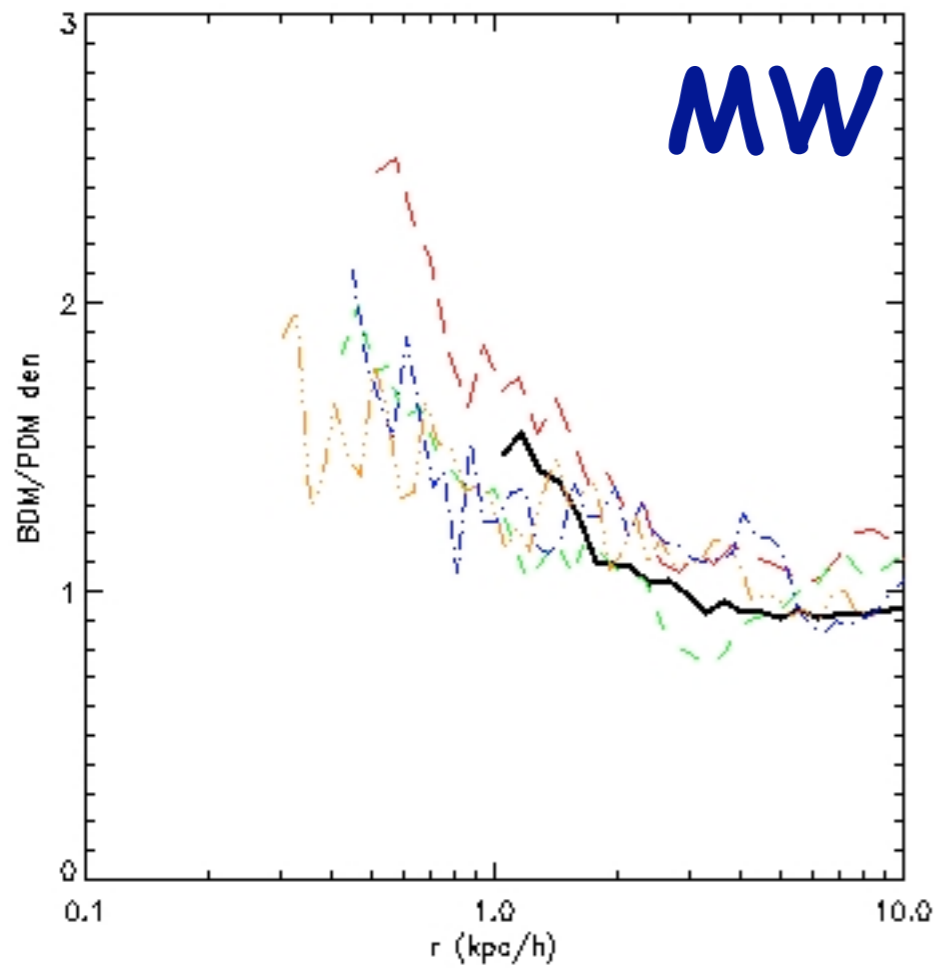
1

0



Density Ratio

- z=0 Black
- z=1 Red
- z=2 Green
- z=3 Blue
- z=4 Yellow



Meudon (June 2010)

Z= 4

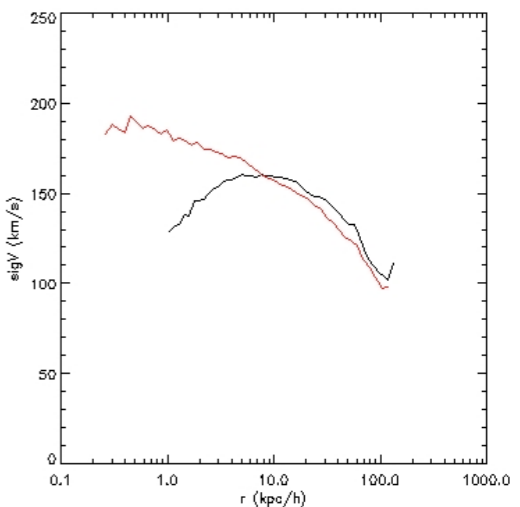
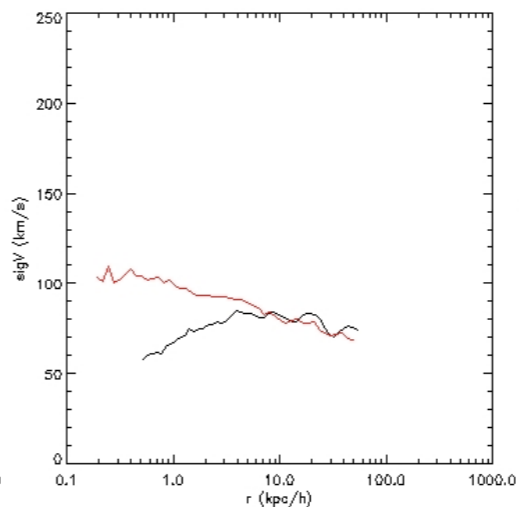
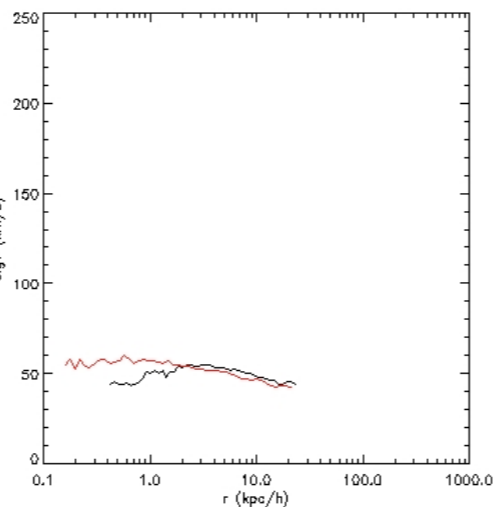
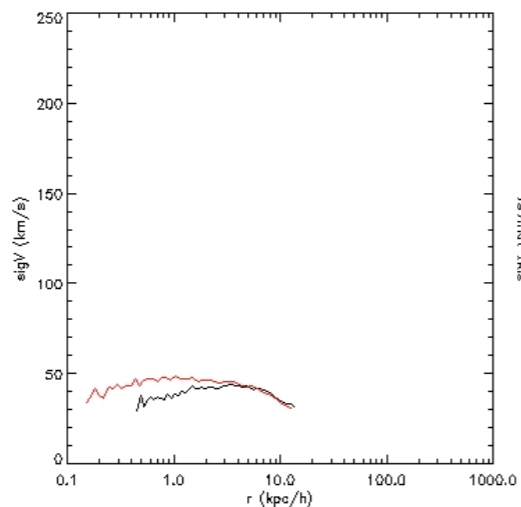
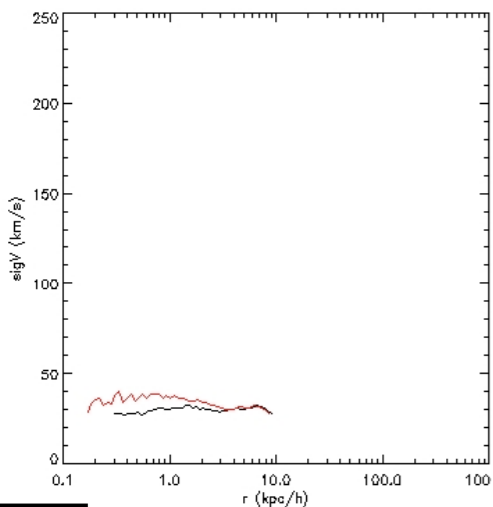
3

2

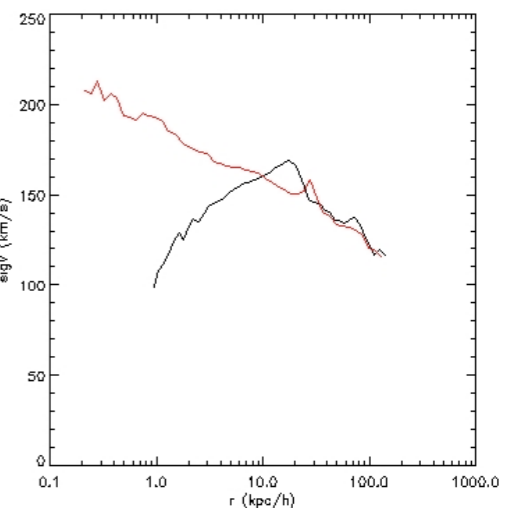
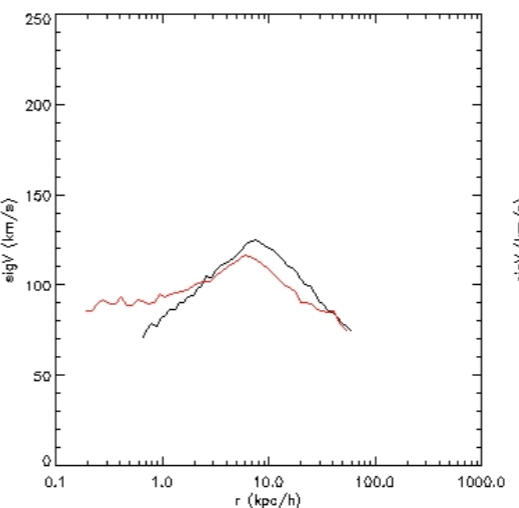
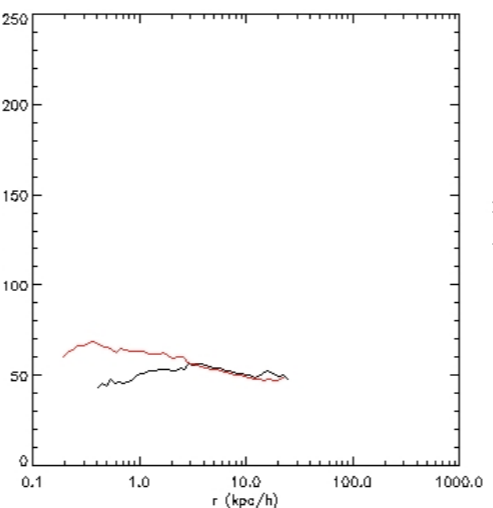
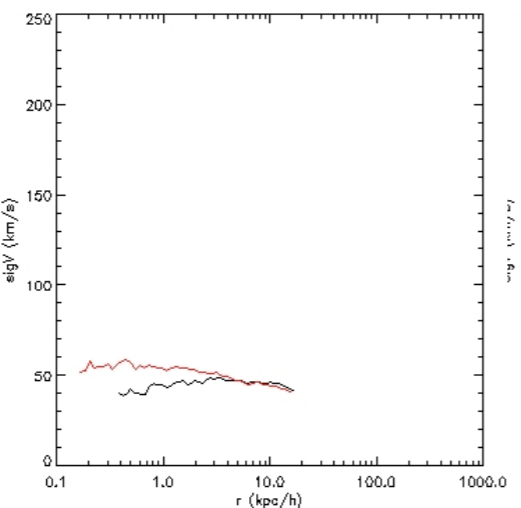
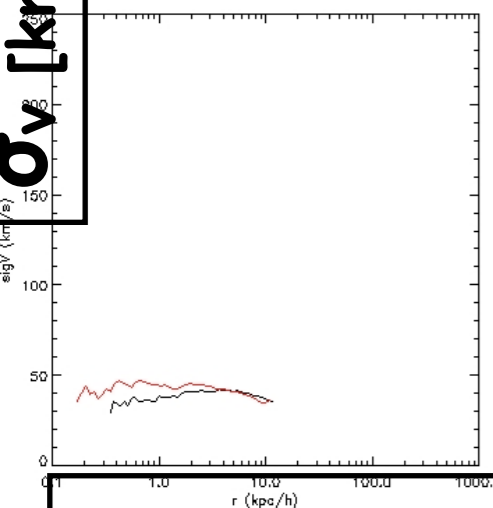
1

0

MW

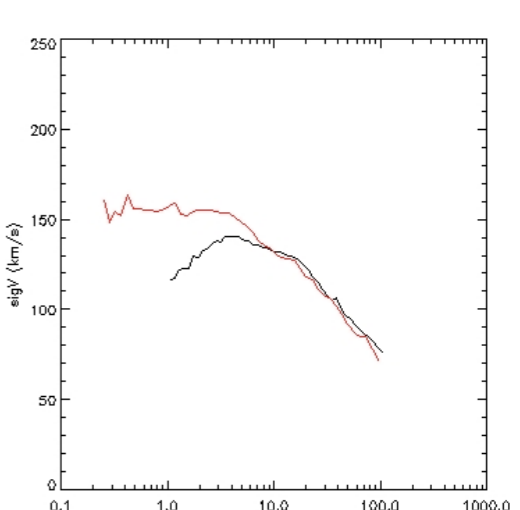
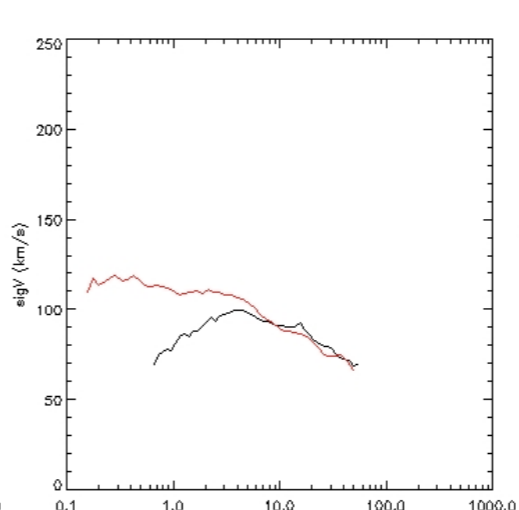
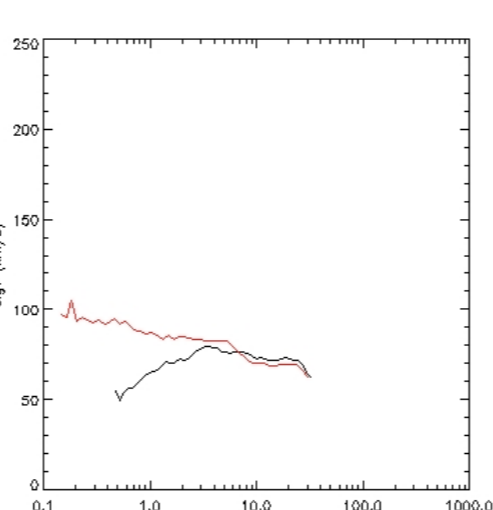
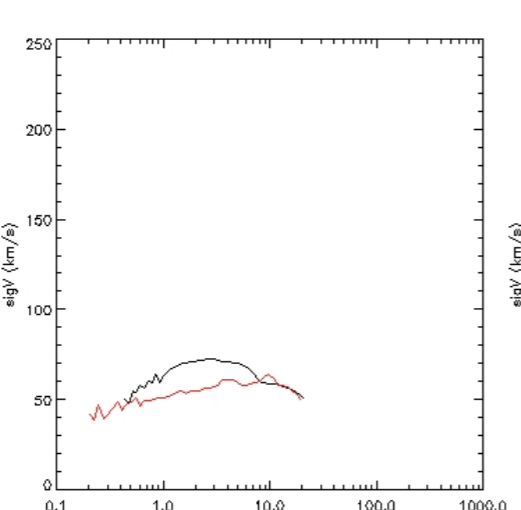
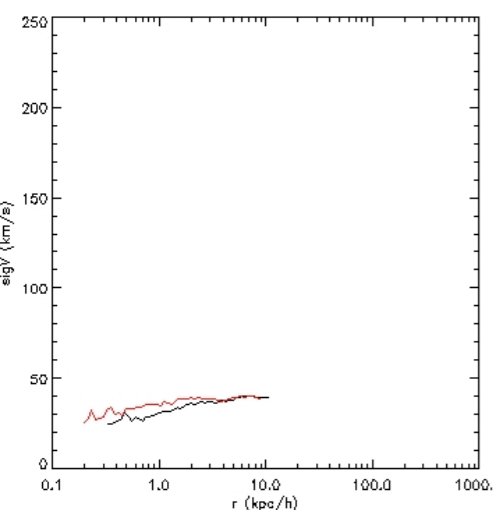


σ_V [km/s]



R [kpc/h]

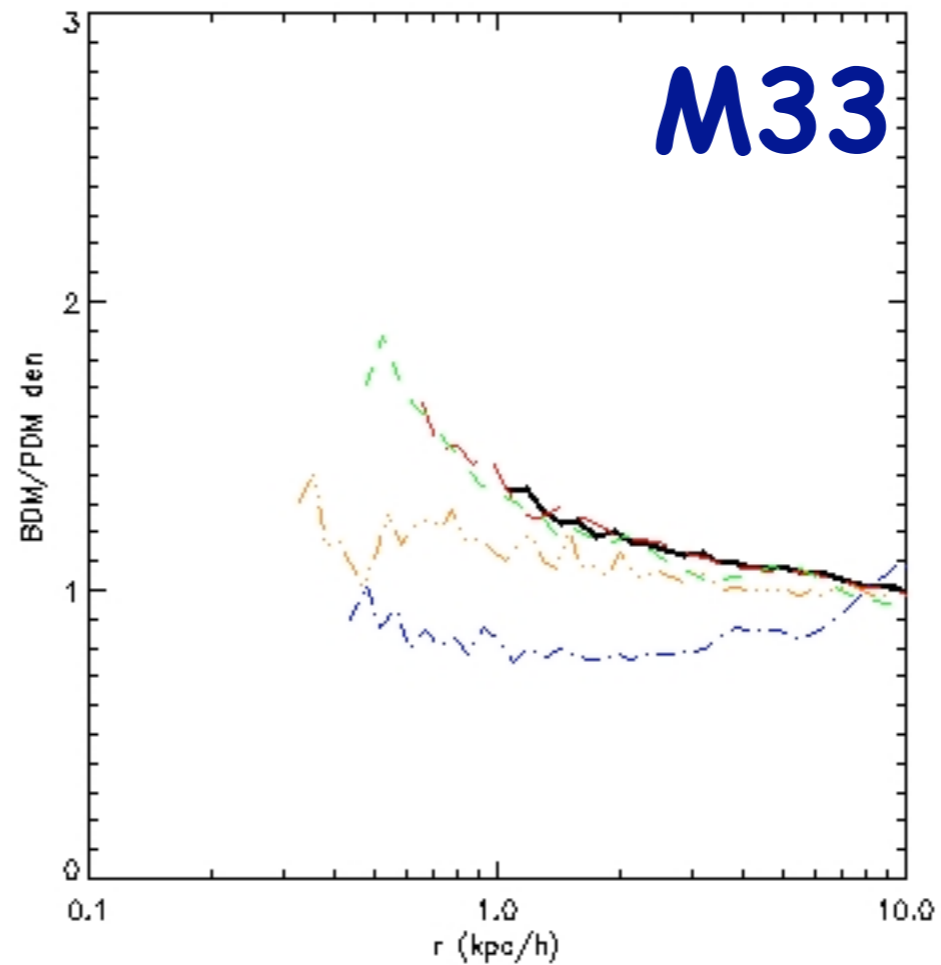
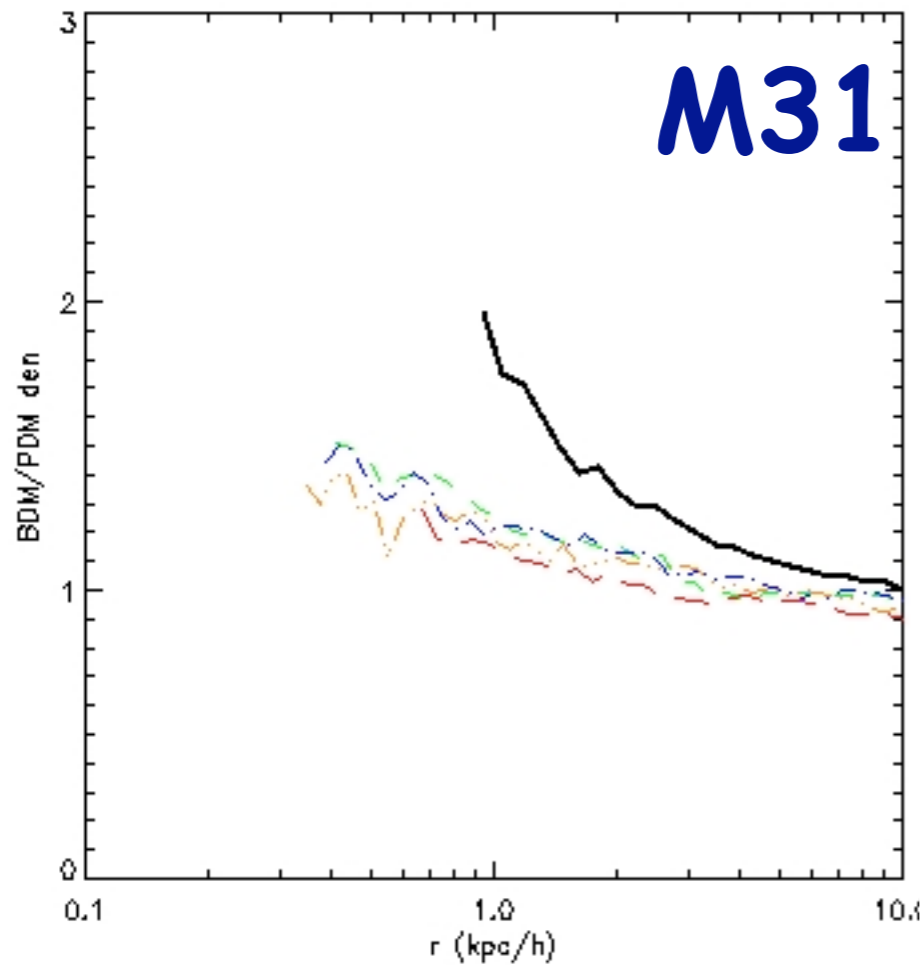
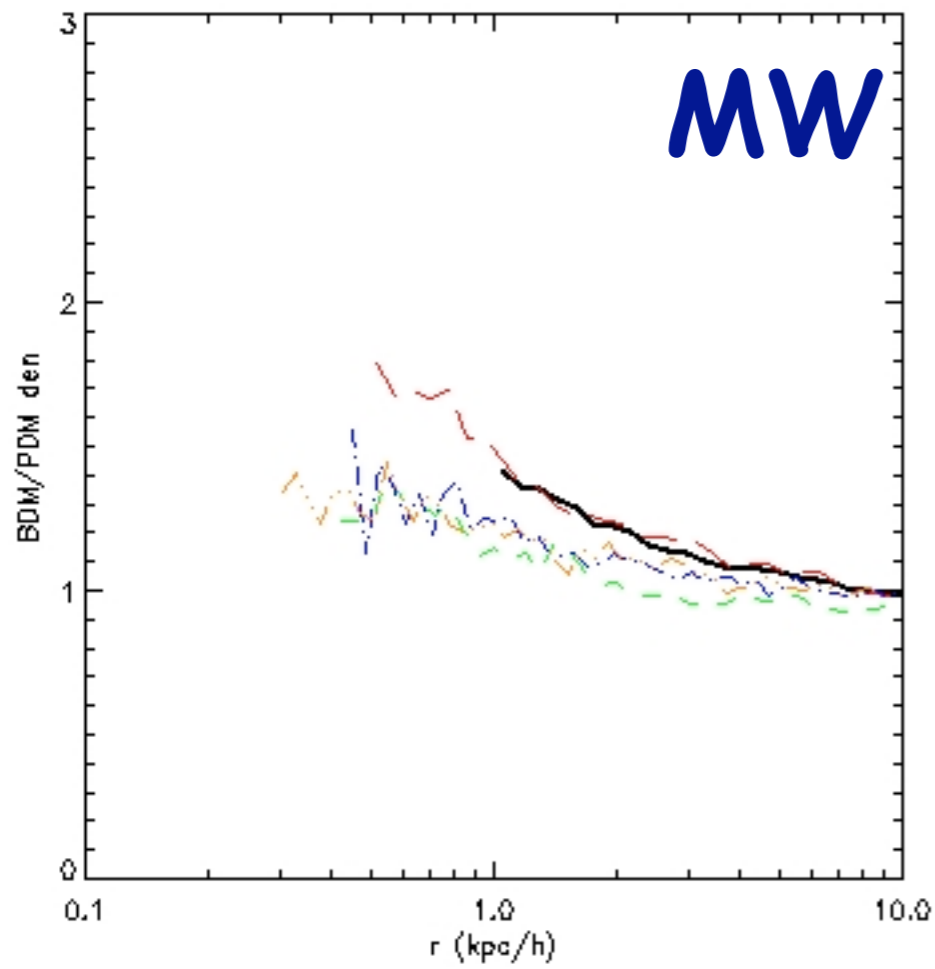
M31



M33

Velocity dispersion Ratio

- z=0 Black
- z=1 Red
- z=2 Green
- z=3 Blue
- z=4 Yellow

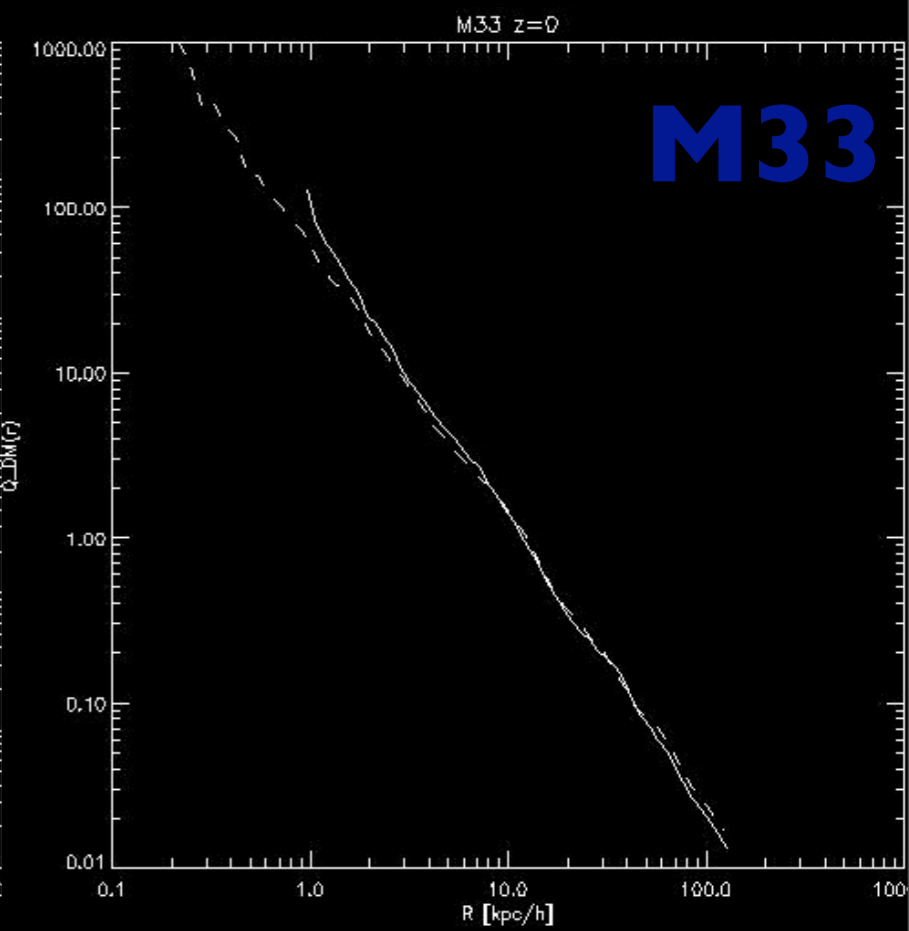
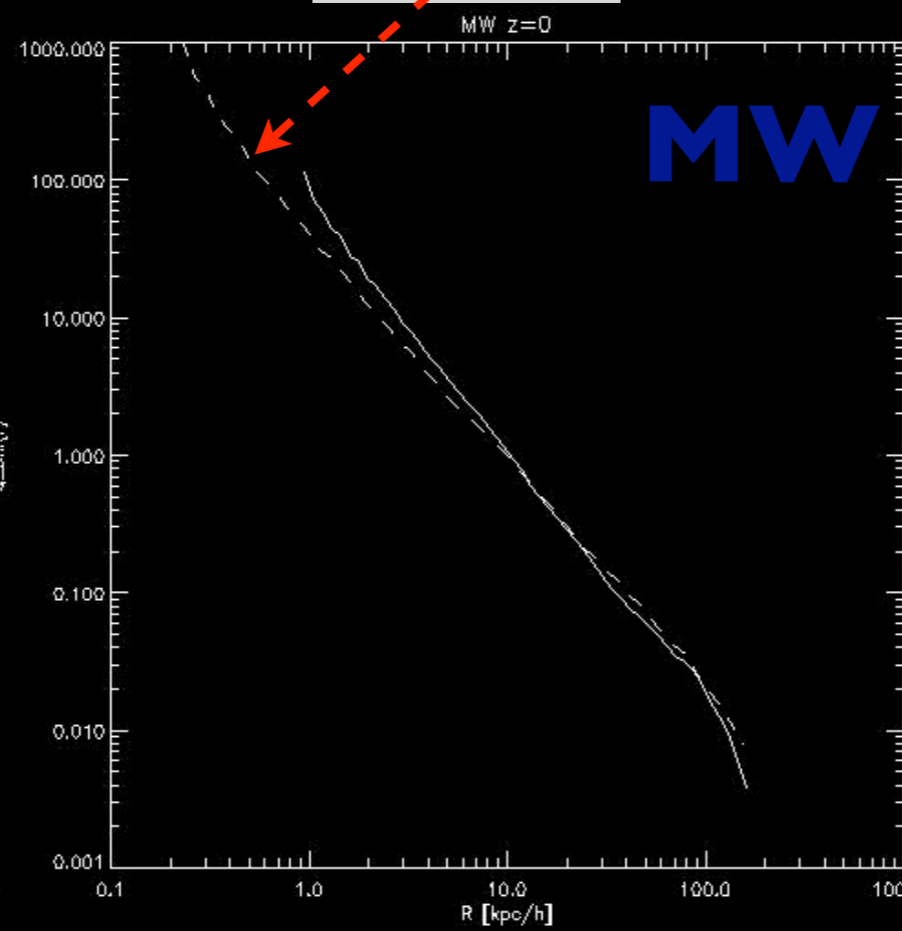
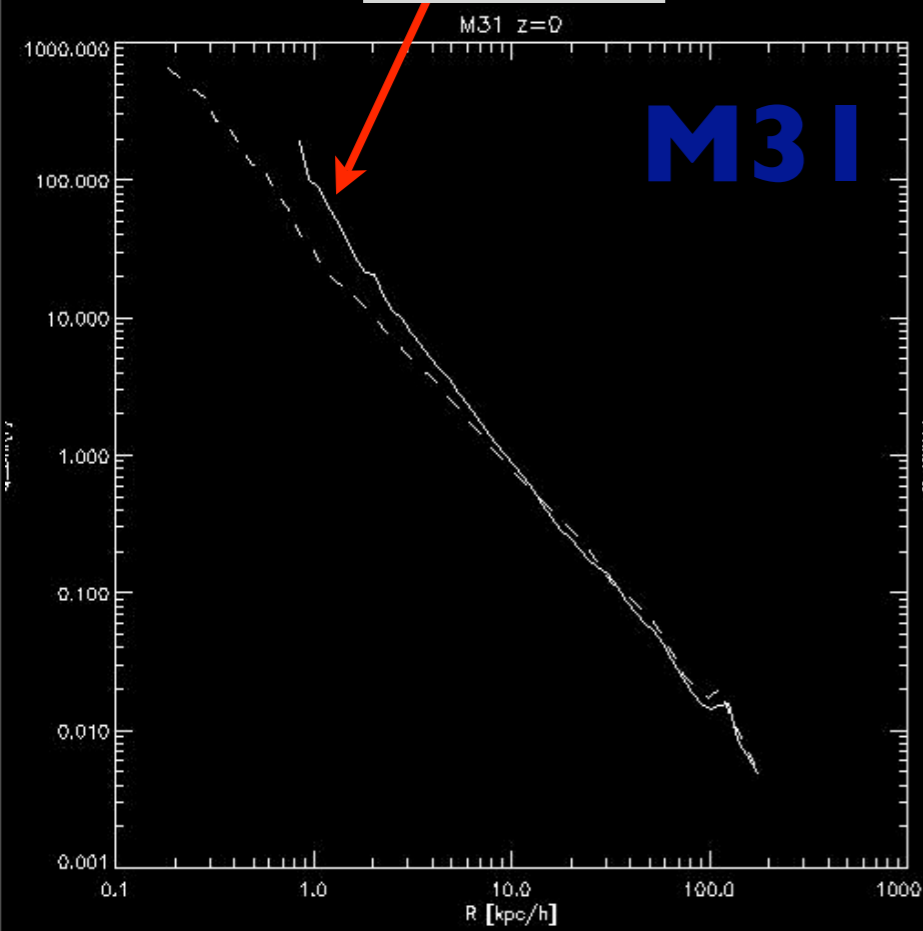


Meudon (June 2010)

$$Q(r) = \frac{\rho}{\sigma_v^3}$$

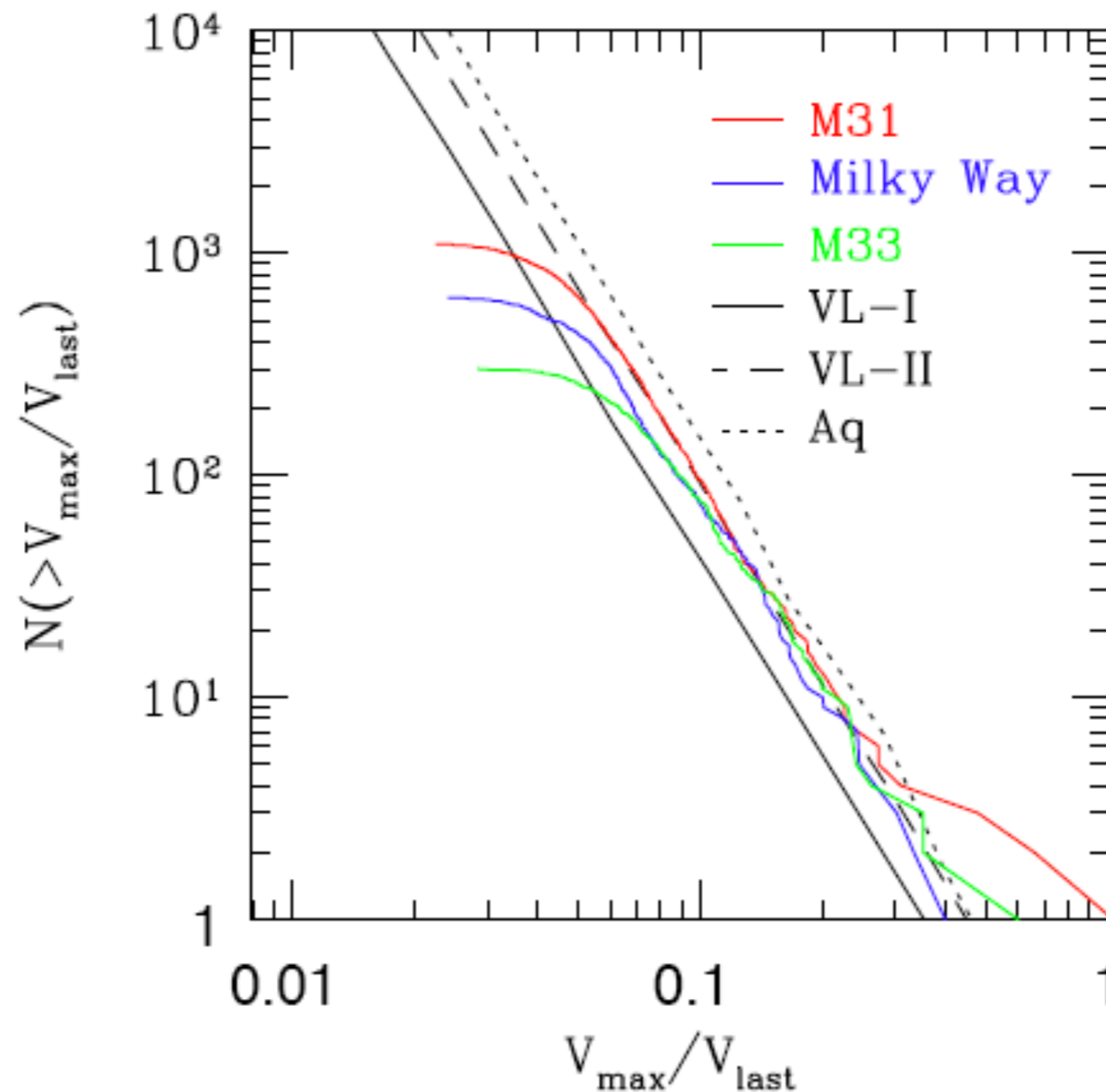
PDM

BDM



Self-Similar Nature of (PDM) Subhalos Abundance

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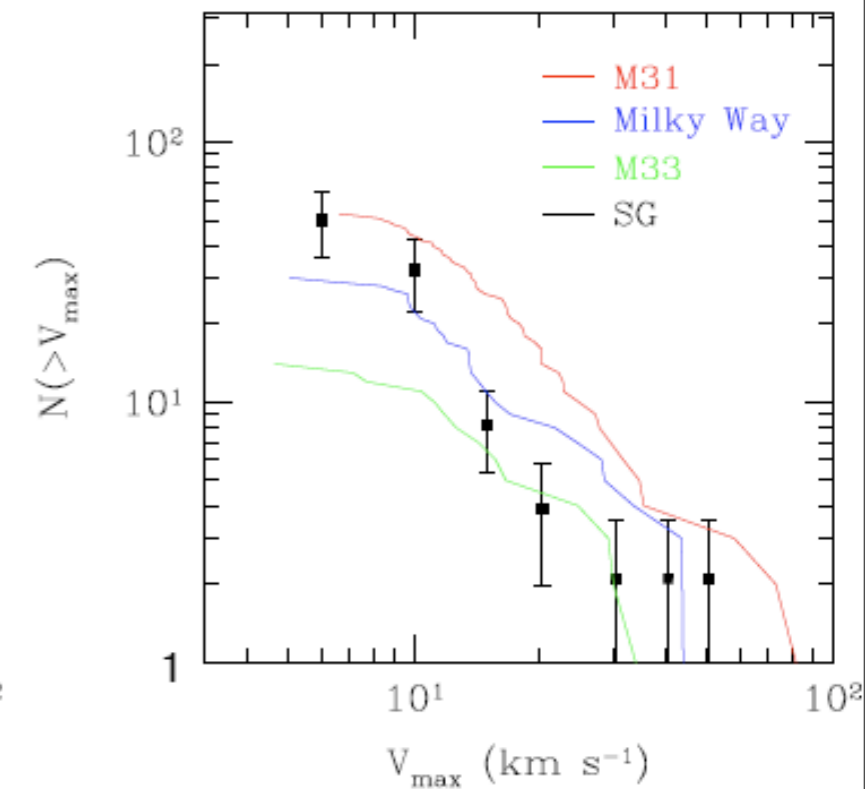
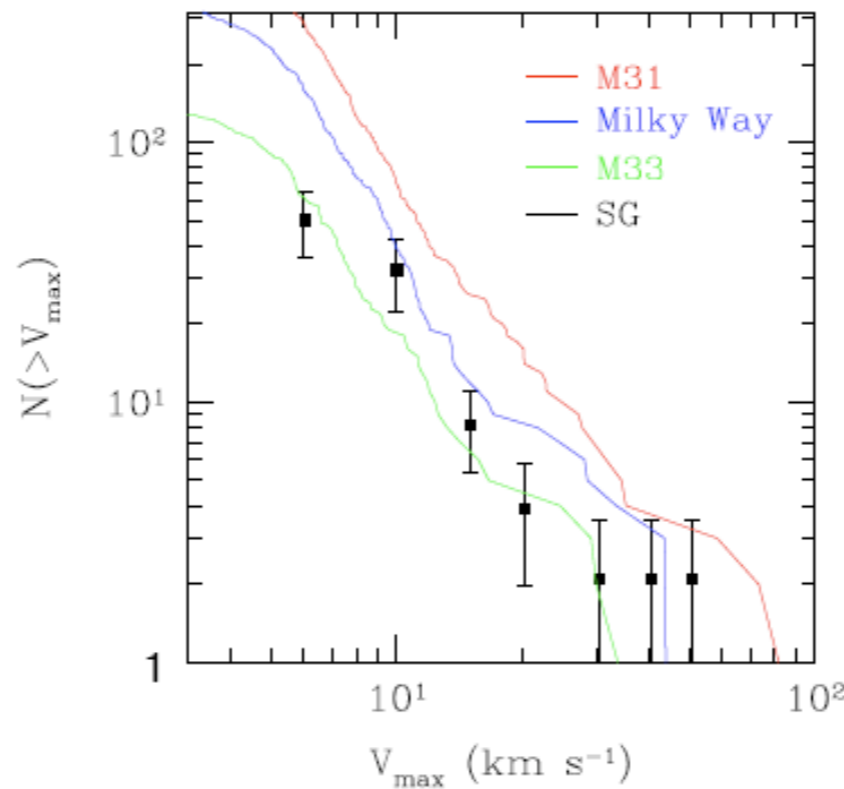
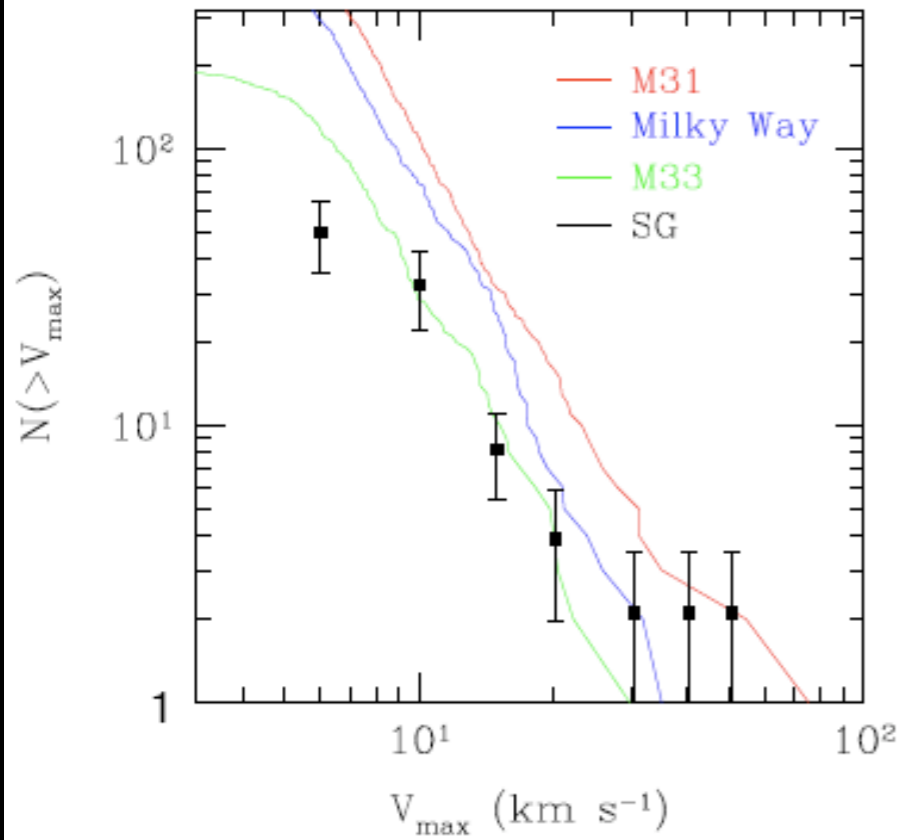


Substructures: PDM vs BDM

PDM - subhalos

BDM - subhalos

BDM - satellites

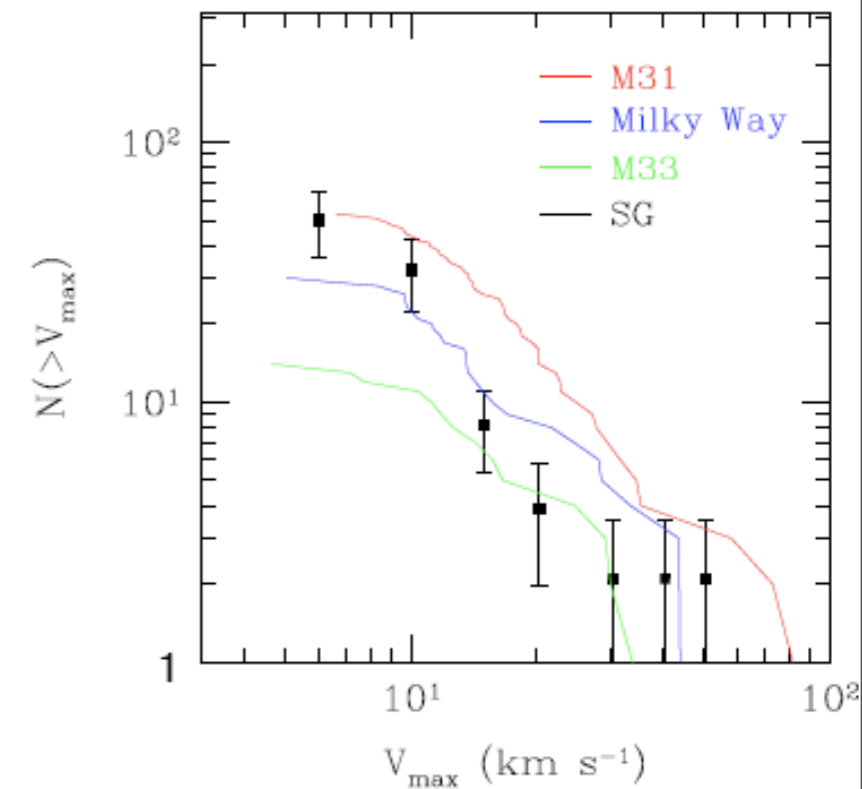
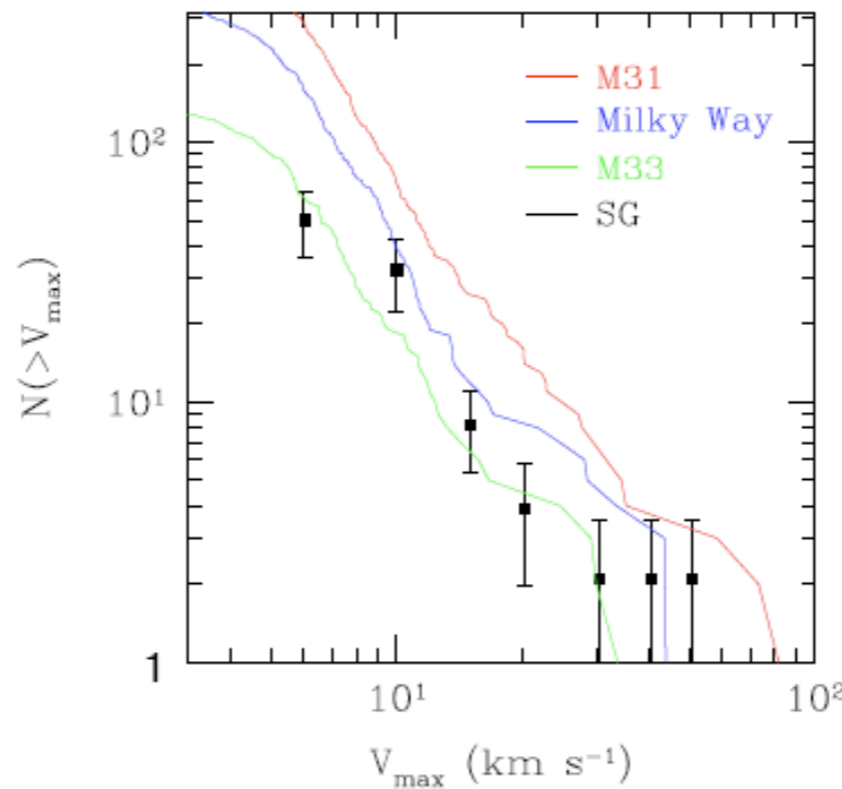
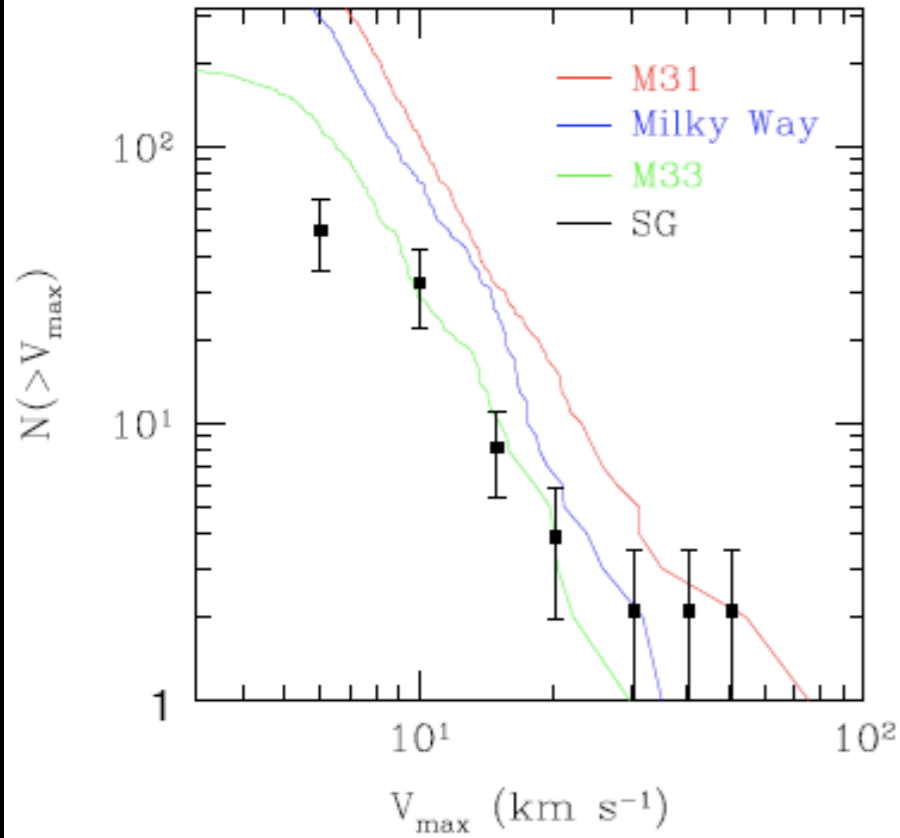


Substructures: PDM vs BDM

PDM - subhalos

BDM - subhalos

BDM - satellites



PDM & BDM V_{\max} functions are similar

$$\left(\frac{dV_M}{dt}\right)_{\text{DF}} = -\frac{4\pi G^2 \ln \Lambda \rho M}{V_M^3} \times \left(\text{erf}(X) - \frac{2X}{\sqrt{\pi}} \exp[-X^2] \right) V_M$$

$$X \equiv V_M / \sqrt{2}\sigma$$

$$\Lambda \equiv bV_M^2 / G(M + m)$$

$$\theta \equiv M_{\text{tot}} / M$$

In virial equilibrium:

$$\tau_{\text{DF}} / \tau_{\text{dyn}} \approx 0.75\theta / \ln \theta$$

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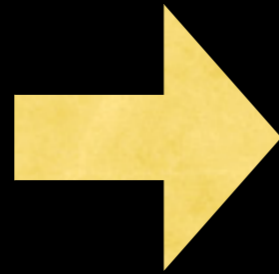
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$\tau_{\text{DF}} / \tau_{\text{dyn}}$ is expected to be (roughly) scale independent!

Subhalos mass function is scale-independent in $M_{\text{subhalo}} / M_{\text{host}}$

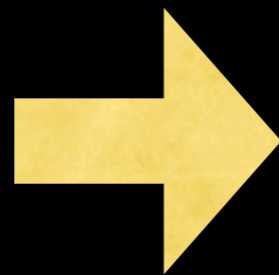
3rd Interim Report: Adiabatic Contraction, Dynamical Friction and (BDM) central density profile

Romano-Diaz et al
(2008, 2009, 2010)
Johansson, Naab &
Ostriker (2009)



In massive ($> \text{few } 10^{12} M_{\text{sun}}$) halos,
dynamical friction by substructures
efficiently reduces the central DM
density and flattens the inner
density cusp.

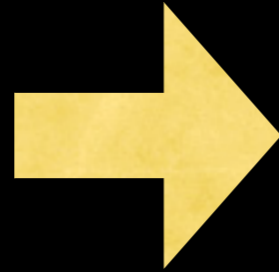
CLUES WMAP3 LG
simulations
Pedrosa, Tissera &
Scannapieco (2009,
2010) & Tissera et
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In less massive ($\leq 10^{12} M_{\text{sun}}$) halos,
dynamical friction by substructures
is less effective - reduces the
central DM density but does not
flatten the inner density cusp.

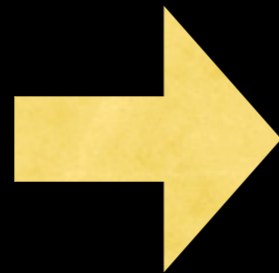
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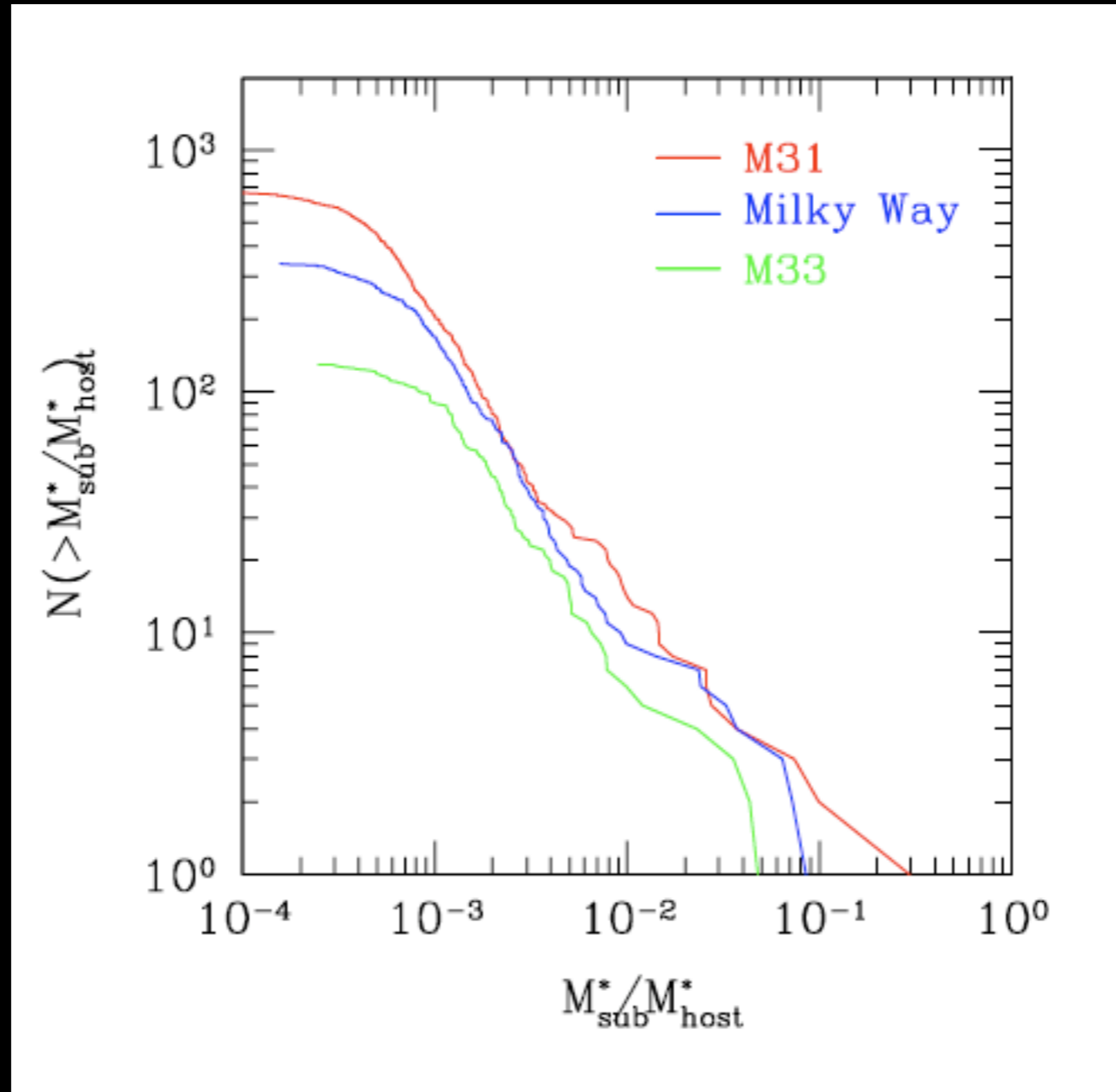
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WHY?

Substructures: baryonic (stellar) mass function



The self-similarity does not hold for the baryonic component

But ...

- $M_{\text{substructure}}$ scales with M_{host} - the more massive is the host halo the more massive are the substructures.
- A conjecture: More massive substructures are more baryon rich.
- It follows that in more massive hosts the dynamical friction brings in more stellar rich substructures to the center, whose orbital energy is pumped into the DM distribution.

SUMMARY

- Adiabatic contraction works - in all simulations the central BDM DM density profile exceeds the PDM profile
- Dynamical friction on substructures:
 - ◆ Dynamical friction is the mechanism that transfers energy to the DM
 - ◆ Key point - substructures need to be baryon rich, so as to replace DM by baryons
 - ◆ The magnitude of the effect depends on the details of baryonic physics: star formation, feedback, ...
 - ◆ The effect depends on the halo (mass, merging history?)
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THANK YOU