

# Theoretical Modeling of the Halo Mass Function: a Path-Integral Approach to the Excursion Set Theory



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# Cosmic Structures and Dark Matter

## Dark Matter:

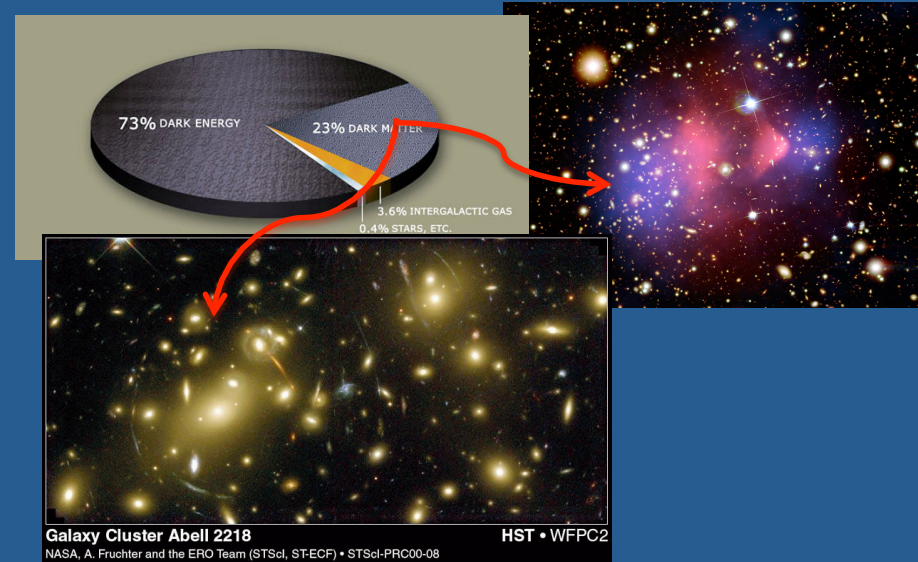
- Foster matter clustering
- Resides in virialized clumps

## The Role of Halos:

- Building blocks of cosmic structure formation
- Shape baryon distribution

## Halo Properties & Mass Function:

- Statistics of the initial density field
- Non-linear gravitational collapse
- Nature DM particles
- Underlying cosmology



## Numerical N-body Simulations Results:

- Limited to Poisson-Vlasov systems (CDM)
- Still require theoretical understanding

$$f(\sigma) = A \left[ \left( \frac{\sigma}{b} \right)^{-a} + 1 \right] e^{-c/\sigma^2}$$

(Tinker et al. 2008)

# Outline

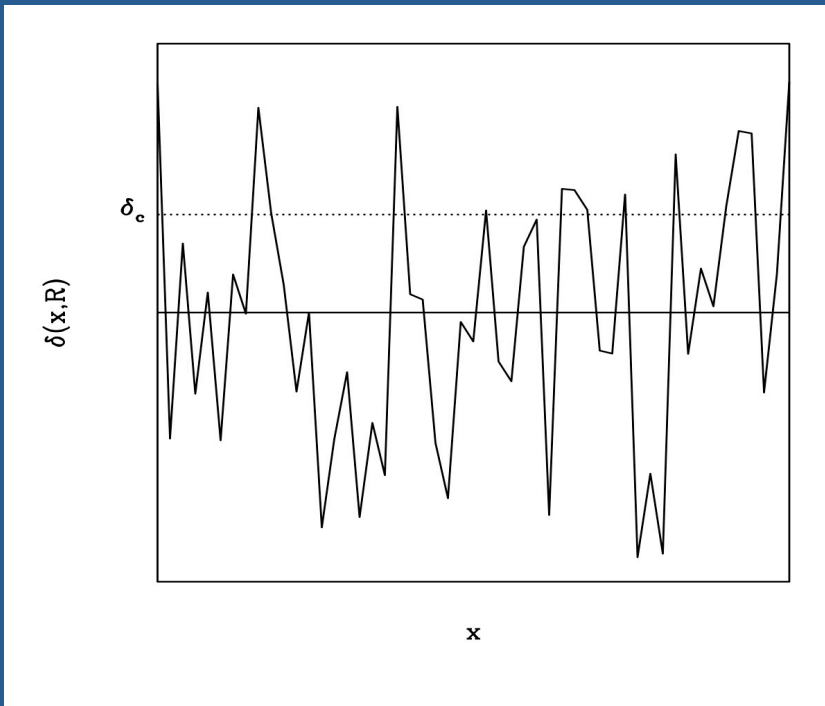
- Press-Schechter Formalism
- The Excursion Set Theory
- Halo Collapse Model
- Halo Mass Definition & Non-Markovianity
- Path-Integral Approach
- Comparison with N-body results

# Reference List

- Bond et al., ApJ 379, 440 (1991)
- Percival, MNRAS 327, 1313 (2001)
- Maggiore & Riotto, ApJ 711, 907 (2010)
- Maggiore & Riotto, ApJ 717, 515 (2010)
- Maggiore & Riotto, ApJ 717, 526 (2010)
- De Simone, Maggiore & Riotto, MNRAS 412, 2587 (2011)
- Ma et al., MNRAS 411, 2644 (2011)
- Corasaniti & Achitouv, PRL in press (arXiv:1012.3468)
- Corasaniti & Achitouv, PRD submitted (on arXiv next week)
- Achitouv & Corasaniti, (on arXiv next week)

# Press-Schechter Formalism

## Smoothed Linear Density Field:



- Filtering scale and mass

$$M = \bar{\rho} V(R)$$

- Statistics of the smoothed density field

$$\Pi(\delta, R)$$

- Linearly extrapolated collapse threshold

$$\delta_c$$

- #-halos with mass  $dM$

$$\frac{dn}{dM} = \frac{1}{V} \frac{dF}{dM}$$

- Fraction of mass in halos  $> M$

$$F_{PS}(R[M]) = \int_{\delta_c}^{\infty} d\delta \Pi(\delta, R[M])$$

# Press-Schechter Mass Function

## Gaussian Field:

- Variance of the smoothed field

$$\sigma^2(R) \equiv S(R) = \frac{1}{2\pi^2} \int k^2 P(k) |\tilde{W}(k, R)|^2 dk$$

- PDF

$$\Pi(\delta, S[R]) = \frac{1}{\sqrt{2\pi S[R]}} e^{-\delta^2/2S[R]}$$

- Spherical collapse threshold

$$\delta_c$$

- Fraction of mass in halos

$$F_{PS}(R) = \int_{\delta_c}^{\infty} d\delta \Pi(\delta, S(R)) = \frac{1}{2} \text{Erfc} \left[ \frac{\delta_c}{\sigma(R)\sqrt{2}} \right]$$

- Halo Mass Function

$$\frac{dn}{dM} = \frac{\bar{\rho}}{M^2} \frac{d \ln \sigma^{-1}}{d \ln M} f(\sigma)$$

with

$$f_{PS}(\sigma) = 2\sigma^2 \frac{dF_{PS}}{dS} = \frac{1}{\sqrt{2\pi}} \frac{\delta_c}{\sigma} e^{-\delta_c^2/(2\sigma^2)}$$

# Cloud-in-Cloud Problem

- Asymptotic Behavior:**
- In the limit  $R \rightarrow 0$  all mass must be in collapsed structures,  $F(0)=1$
  - In the PS calculation half of the mass is miscounted

$$F_{PS}(R) = \frac{1}{2} \text{Erfc} \left[ \frac{\delta_c}{\sigma(R)\sqrt{2}} \right] \xrightarrow{R \rightarrow 0} \frac{1}{2}$$

**Problem:** In the PS approach there is no mass ordering

- No distinction between different configurations (cloud-in-cloud)



- Halo of mass M embedded in a larger non-collapsed region contributes to mass function at mass M



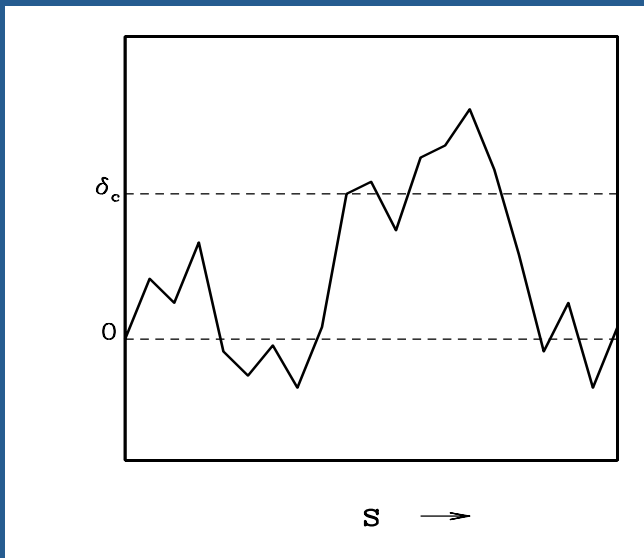
- Region M embedded in halo  $M' > M$ . M should not be counted in mass function since already included at  $M'$

# Excursion Set Theory

## Key point:

$$\delta(x, R) = \frac{1}{(2\pi)^3} \int d^3k \delta(k) \tilde{W}(k, R) e^{-ikx}$$

- random walks start at  $R = \infty$  ( $S=0$ ) with  $\delta = 0$  evolving toward smaller  $R$  (larger  $S$ )



- At any point  $x$ ,  $\delta$  performs a random walk as function of  $R$

- Langevin Equation:

$$\frac{\partial \delta}{\partial R} = \zeta(R) \quad \text{and} \quad \zeta(R) = \frac{1}{(2\pi)^3} \int d^3k \delta(k) \frac{\partial \tilde{W}}{\partial R} e^{-ikx}$$

-  $\zeta(R)$  depends on  $\Pi(\delta)$  and  $W(x, R)$

- Halos of mass  $M$  corresponds to trajectories crossing the threshold at  $S(M)$

- cloud-in-cloud solved by requiring first crossing



# Excursion Set Mass Function

## Stochastic Problem:

- Computation of the probability distribution of random walks with absorbing boundary,  $\Pi(\delta, \delta_c, S)$
- Multiplicity function obtained from the first-crossing rate  $\mathcal{F}(S) = dF/dS$

$$\int_0^S \mathcal{F}(S') dS' = 1 - \int_{-\infty}^{\delta_c} \Pi(\delta, \delta_c, S) d\delta$$



$$\frac{dF}{dS} = -\frac{\partial}{\partial S} \int_{-\infty}^{\delta_c} \Pi(\delta, \delta_c, S) d\delta$$

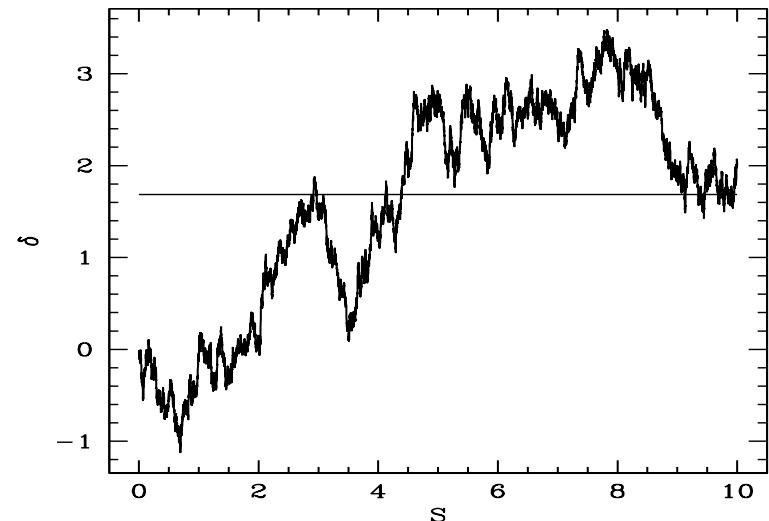
Sharp-k filter:  $\tilde{W}(k, R) = \theta(1/R - k)$

- Markovian random walks

$$\frac{\partial \delta}{\partial S} = \eta(S)$$

with

$$\begin{aligned} \langle \eta(S) \rangle &= 0 \\ \langle \eta(S) \eta(S') \rangle &= \delta_D(S - S') \end{aligned}$$



# Extended Press-Schechter

Fokker-Planck Equation:

$$\frac{\partial \Pi}{\partial S} = \frac{1}{2} \frac{\partial^2 \Pi}{\partial \delta^2}$$

with

$$\Pi(\delta, 0) = \delta_D(\delta)$$

$$\Pi(\delta_c, S) = 0$$

$$\Pi(\delta, S) = \frac{1}{\sqrt{2\pi S}} \left[ e^{-\delta^2/(2S)} - e^{-(2\delta_c - \delta)^2/(2S)} \right]$$

for  $\delta < \delta_c$

$$\frac{dF}{dS} = -\frac{\partial}{\partial S} \int_{-\infty}^{\delta_c} \Pi(\delta, \delta_c, S) d\delta$$



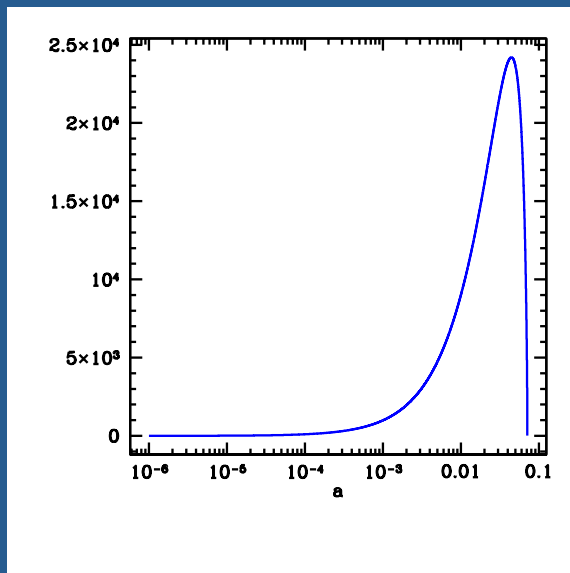
$$f_{EPS}(\sigma) = \sqrt{\frac{2}{\pi}} \frac{\delta_c}{\sigma} e^{-\delta_c^2/(2\sigma^2)}$$

# Halo Collapse Model

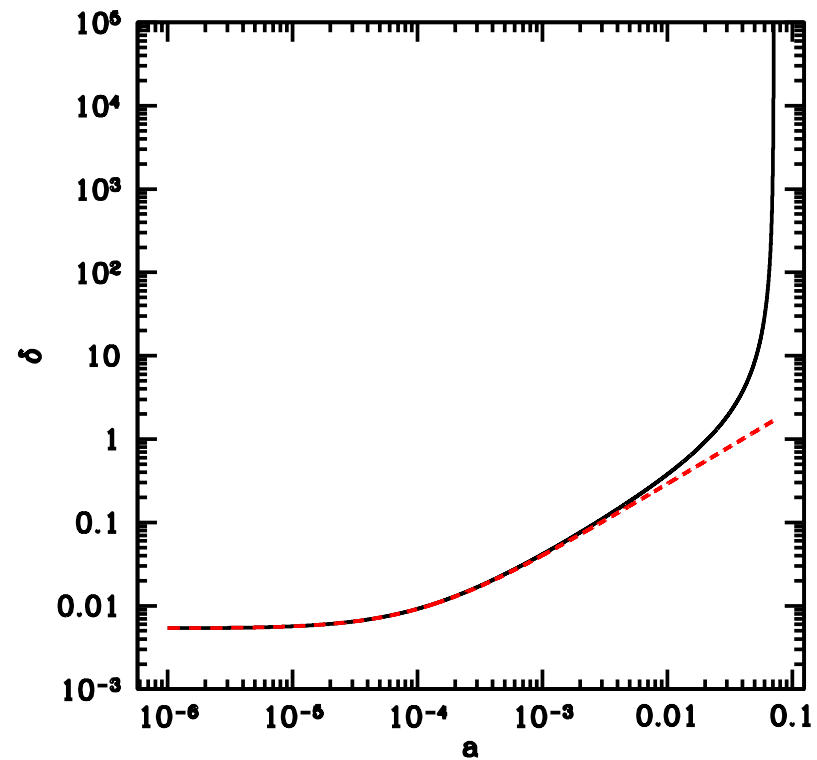
## Spherical Collapse (Gunn & Gott, 1973)

- Top-hat perturbation in FRW background,  $y=R/R_i$
- Dynamics independent of  $R_i$

$$\ddot{y} = -\frac{4}{3}\pi G\bar{\rho}_m^i(1+\delta_m^i)\frac{1}{y^2}$$



$$1 + \delta_m = (1 + \delta_m^i) y^3 \left( \frac{a}{a_i} \right)^3$$



# Non-Spherical Halo Collapse

## Ellipsoidal Collapse

- Initial Gaussian fluctuations are non-spherical (Doroshkevich, 1970)
- Ellipsoidal halos collapse and shear

$$\frac{d^2 a_i}{dt^2} = \frac{8}{3} \pi G \bar{\rho}_\Lambda a_i - 4 \pi G \bar{\rho}_m a_i \left[ \frac{1}{3} + \frac{\Delta(t)}{3} + \frac{b'_i(t)}{2} \Delta(t) + \lambda'_i(t) \right]$$

$$b'_i = -\frac{2}{3} + a_1 a_2 a_3 \int_0^\infty \frac{d\tau}{(a_i^2 + \tau) \prod_{m=1}^3 (a_m^2 + \tau)^{1/2}}$$

$$\lambda_1 = \frac{\delta}{3} (1 - 3e + p)$$

$$\lambda_2 = \frac{\delta}{3} (1 - 2p)$$

$$\lambda_3 = \frac{\delta}{3} (1 + 3e + p)$$

- Dynamics dependent of initial size of the collapsing region
- Critical Overdensity is mass dependent (e.g. Eisenstein & Loeb, 1995)

# “Fuzzy” Barrier

## Stochastic Barrier

- Ellipsoidal parameters are random variables with characteristic probability distribution

- e.g. for a Gaussian density field

$$g(e, p | \delta) = \frac{1125}{\sqrt{10\pi}} e(e^2 - p^2) \left(\frac{\delta}{\sigma}\right)^5 \exp\left[-\frac{5}{2} \frac{\delta^2}{\sigma^2} (3e^2 + p^2)\right]$$

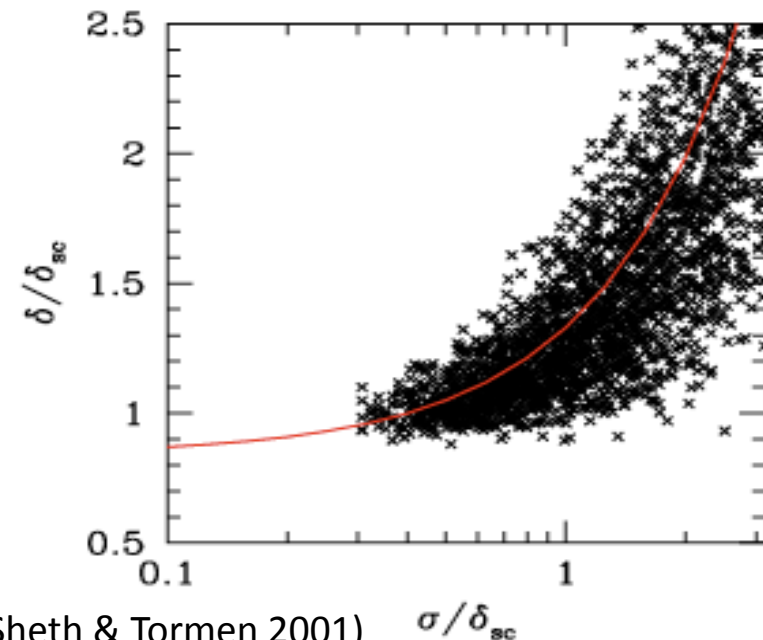
- Density threshold is a random variable

- e.g.  $\langle B(S) \rangle = \delta_c [1 + \beta (S/S_*)^\gamma]$

(Sheth, Mo & Tormen 2001)

- Stochastic barrier model: specify the moments of barrier's PDF

$\langle B(S) \rangle$  &  $\langle (B(S) - \langle B(S) \rangle)^2 \rangle, \dots$



(Sheth & Tormen 2001)

# Stochastic Barrier and Excursion Set

## Non-Spherical Collapse

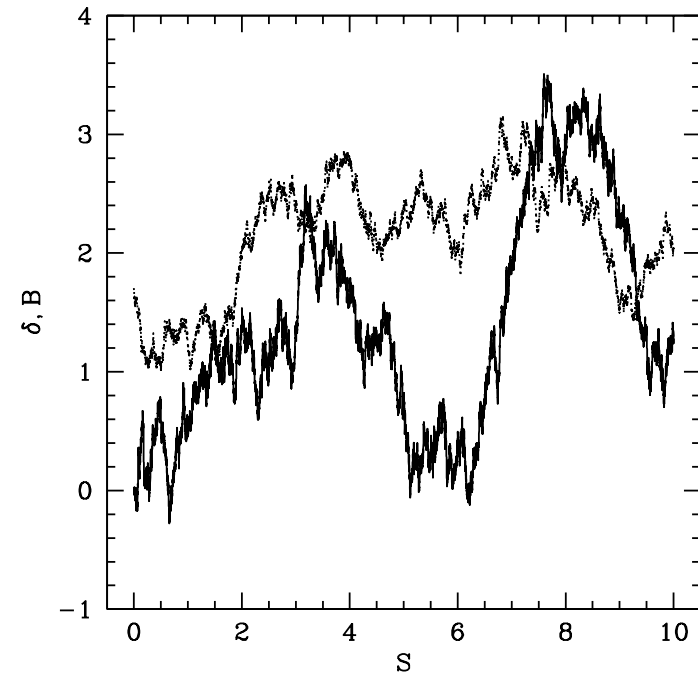
- Barrier performs a random walk characterized by non-spherical collapse predicted correlators
- e.g.  $\langle B(S) \rangle = \delta_c + \beta S$  &  $\langle (B(S) - \langle B(S) \rangle)^2 \rangle^{1/2} = \sqrt{D_B} \sigma$
- Introduce:  $Y = B - \delta$

## Sharp-k filter case:

$$\frac{\partial Y}{\partial S} = \beta + \eta(S)$$

$$\langle \eta(S) \rangle = 0$$

$$\langle \eta(S) \eta(S') \rangle = (1 + D_B) \delta_D(S - S')$$



$$\frac{\partial \Pi}{\partial S} = -\beta \frac{\partial \Pi}{\partial Y} + \frac{1 + D_B}{2} \frac{\partial^2 \Pi}{\partial Y^2}$$

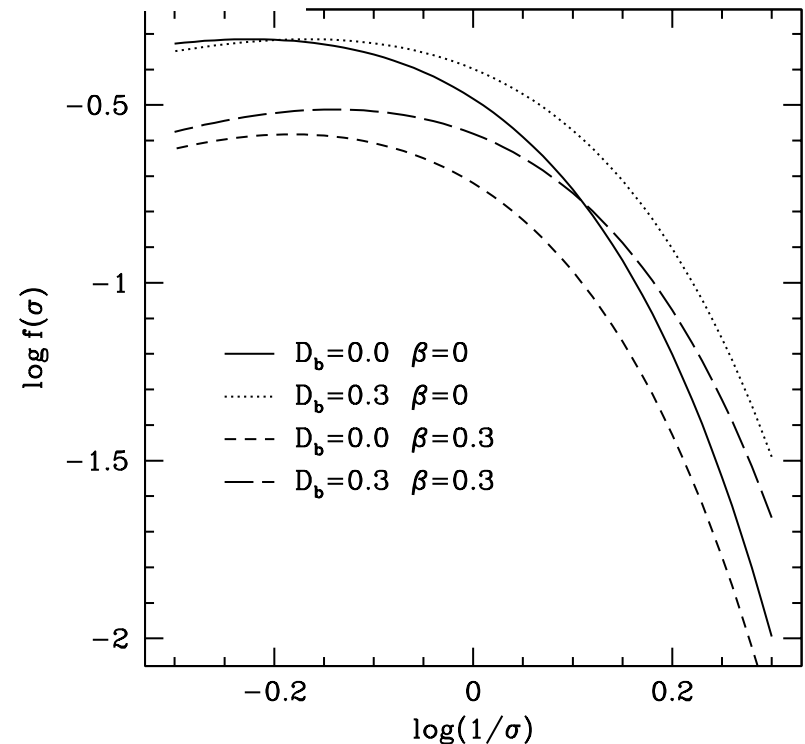
$$\Pi(Y, 0) = \delta_D(Y - \delta_c)$$

$$\Pi(0, S) = 0$$

$$\Pi(Y, S) = \frac{e^{\frac{\beta}{1+D_B}(Y-Y_0-\beta S/2)}}{\sqrt{2\pi S(1+D_B)}} \left[ e^{-\frac{(Y-Y_0)^2}{2S(1+D_B)}} - e^{-\frac{(Y+Y_0)^2}{2S(1+D_B)}} \right]$$

## Multiplicity Function:

$$f(\sigma) = \sqrt{\frac{2}{\pi}} \frac{\delta_c}{\sigma \sqrt{1+D_B}} e^{-\frac{(\delta_c + \beta \sigma^2)^2}{2\sigma^2(1+D_B)}}$$







# Filter Function and Halo Mass Definition

**Mass and smoothing scale:**  $M(R) = \rho V(R)$  with  $V(R) = \int d^3x W(x, R)$

- Unambiguously define only for sharp-x filter:  $W(x, R) = \theta(|\vec{x}| - R)$
- Generic filters define M up to a normalization constant
- Sharp-k leaves M undefined

## N-body Halo mass:

- Mass depends on Halo detection algorithm

- e.g SOD  $M_{\Delta} = \frac{4}{3} \pi R_{\Delta}^3 \bar{\rho} \Delta$  corresponding to  $W(x, R_{\Delta}) = \theta(|\vec{x}| - R_{\Delta})$

## Realistic Filtering: Sharp-x

- Leads to correlated random walks since

$$\tilde{W}(k, R) = 3 \frac{\sin(kR) - (kR) \cos(kR)}{(kR)^3}$$

# Top-Hat Filter and Correlated Random Walks

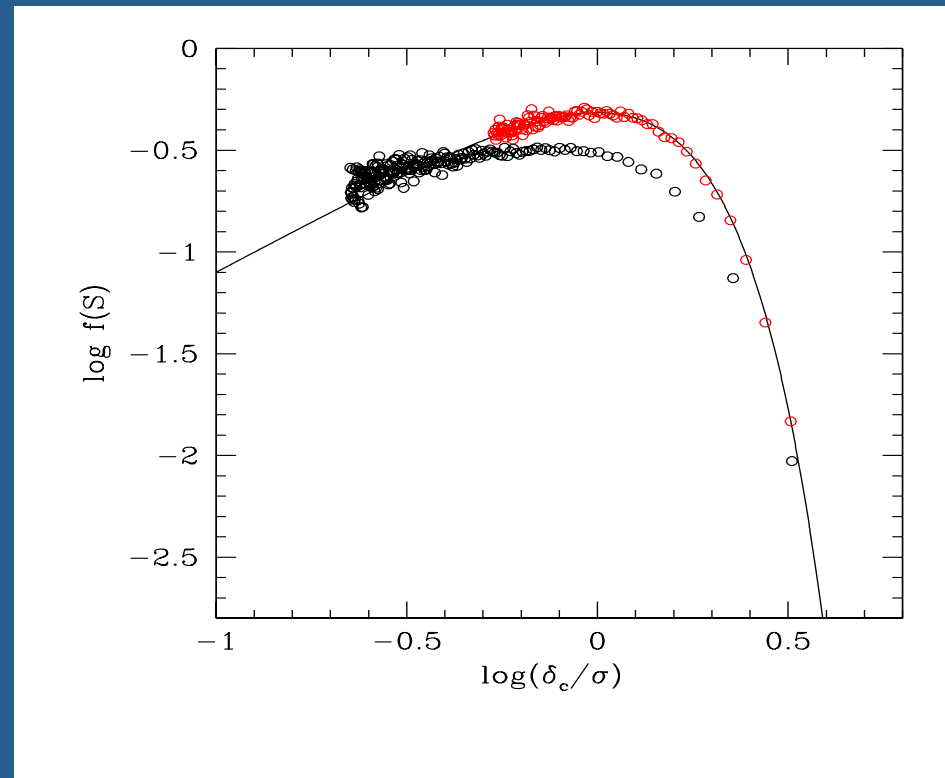
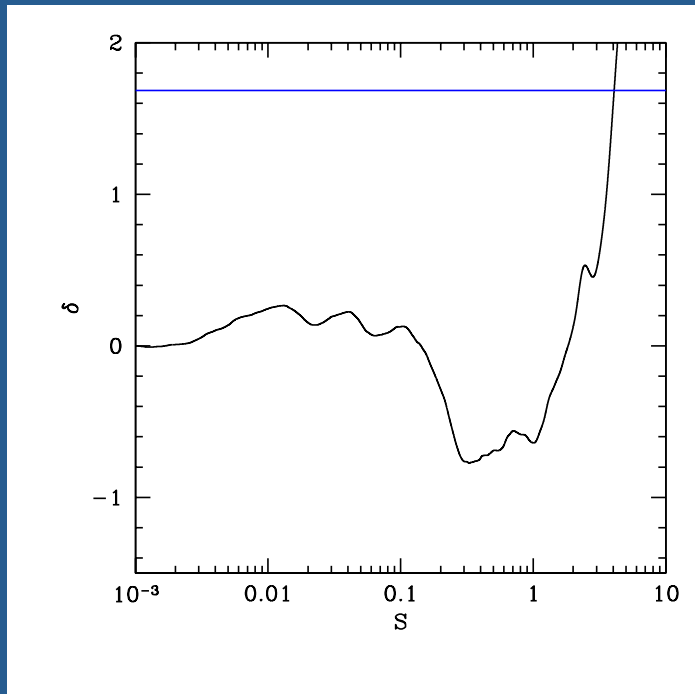
Langevin Equation:

$$\frac{\partial \delta}{\partial \ln k} = Q(\ln k) \tilde{W}(k, R)$$

$$\langle Q(\ln k) \rangle = 0$$

$$\langle Q(\ln k) Q(\ln k') \rangle = \Delta^2(k) \delta_D(\ln k - \ln k')$$

Spherical Collapse Case:



# Correlation Function

Generic Filter:

$$\langle \delta[R(S)]\delta[R(S')] \rangle = \frac{1}{2\pi^2} \int_0^\infty dk k^2 P(k) T^2(k) \tilde{W}[k, R(S)] \tilde{W}[k, R(S')]$$

Sharp-x Filter:

$$\langle \delta(S)\delta(S') \rangle = \int_0^S ds \int_0^{S'} ds' \langle \eta(s)\eta(s') \rangle = \min(S, S')$$

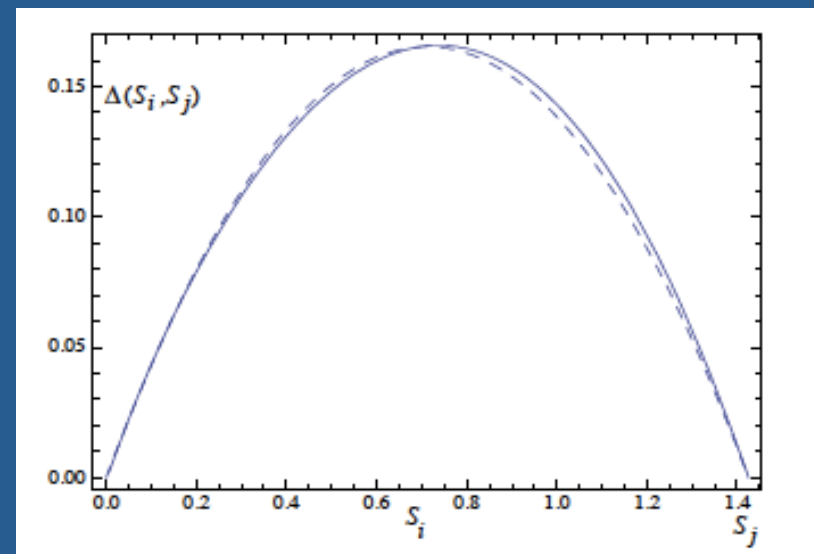
Introduce:

$$\Delta(S, S') = \langle \delta(S)\delta(S') \rangle - \min(S, S')$$

For LCDM power-spectrum:

$$\Delta(S, S') \cong \kappa \frac{S(S' - S)}{S'}$$

where  $\kappa \cong 0.47$



# Path-Integral Approach to Excursion Set

## Discrete Random Walks

- Trajectory over a discrete “time” interval  $\{Y_0, Y_1, \dots, Y_n\}$  with  $S_k = k \varepsilon$  and  $k=1, \dots, n$

## Ensemble Probability Density

$$p(Y_0, \dots, Y_n, S_n) = \langle \delta_D[Y(S_1) - Y_1] \cdot \dots \cdot \delta_D[Y(S_n) - Y_n] \rangle = \int D\lambda e^{i \sum \lambda_i Y_i} \left\langle e^{-i \sum \lambda_i Y(S_i)} \right\rangle$$

## Partition function

$$e^Z = \left\langle e^{-i \sum \lambda_i Y(S_i)} \right\rangle \quad \text{with} \quad Z = \sum_{p=1}^{\infty} \frac{(-i)^p}{p!} \sum_{i_1=1}^n \dots \sum_{i_p=1}^n \lambda_{i_1} \dots \lambda_{i_p} \langle Y(S_{i_1}) \dots Y(S_{i_p}) \rangle_c$$

## Connected Correlators

$$\langle Y(S_i) \rangle_c \equiv \bar{B}(S_i) = \delta_c + \beta S_i$$
$$\langle Y(S_i) Y(S_j) \rangle_c = (1 + D_B) \min(S_i, S_j) + \Delta(S_i, S_j)$$

# $\kappa$ -expansion around Markovian solution

## Probability Distribution

$$\Pi_{\varepsilon}(Y_0, Y_n, S_n) = \int_0^{\infty} dY_1 \dots \int_0^{\infty} dY_{n-1} p(Y_0, \dots, Y_n, S_n)$$

## Expansion to $O(\kappa)$

$$\Pi_{\varepsilon}(Y_0, Y_n, S_n) = \int_0^{\infty} dY_1 \dots \int_0^{\infty} dY_{n-1} \int D\lambda \left( 1 - \frac{1}{2} \sum_{i,j} \lambda_i \lambda_j \Delta_{ij} \right) e^{i \sum_k \lambda_k [Y_k - \bar{B}_k]} e^{-i \sum_{n,m} \lambda_n \lambda_m A_{nm}}$$

## Isolate non-Markovian correction

$$\Pi_{\varepsilon}(Y_0, Y_n, S_n) = \Pi_{\varepsilon}^M(Y_0, Y_n, S_n) + \Pi_{\varepsilon}^{\kappa^{(1)}}(Y_0, Y_n, S_n)$$

# Memory and Memory-of-Memory Terms

## First Order Correction

$$\Pi_{\varepsilon}^{k(1)}(Y_0, Y_n, S_n) = \Pi_{\varepsilon}^m(Y_0, Y_n, S_n) + \Pi_{\varepsilon}^{m-m}(Y_0, Y_n, S_n)$$

$$\Pi_{\varepsilon}^m(Y_0, Y_n, S_n) = - \sum_{i=1}^{n-1} \Delta_{in} \partial_n \left[ \Pi_{\varepsilon}^{M,f}(Y_0, 0, S_i) \Pi_{\varepsilon}^{M,f}(0, Y_n, S_n - S_i) \right]$$

$$\Pi_{\varepsilon}^{m-m}(Y_0, Y_n, S_n) = \sum_{i < j} \Delta_{ij} \left[ \Pi_{\varepsilon}^{M,f}(Y_0, 0, S_i) \Pi_{\varepsilon}^{M,f}(0, 0, S_j - S_i) \Pi_{\varepsilon}^{M,f}(0, Y_n, S_n - S_j) \right]$$

- Markovian solution around the barrier  $\Pi_{\varepsilon}^{M,f}$

## Continuous Limit

$$\sum_{i=1}^{n-1} \rightarrow \lim \frac{1}{\varepsilon} \int_0^S dS_i$$

&

$$\sum_{i < j} \rightarrow \lim \frac{1}{\varepsilon^2} \int_0^S dS_i \int_{S_i}^S dS_j$$

# Corrections to Mass Function

## Memory

$$f_1^m(\sigma) = -2\sigma^2 \frac{\kappa Y_0}{(1+D_B)^2} \frac{\partial}{\partial S} \int_0^\infty dY_n \partial_n \left\{ Y_n e^{\frac{\beta}{1+D_B}(Y-Y_0-\beta S/2)} \operatorname{Erfc} \left[ \frac{Y_0 + Y_n}{\sqrt{2S(1+D_B)}} \right] \right\} = 0$$

## Memory-of-Memory

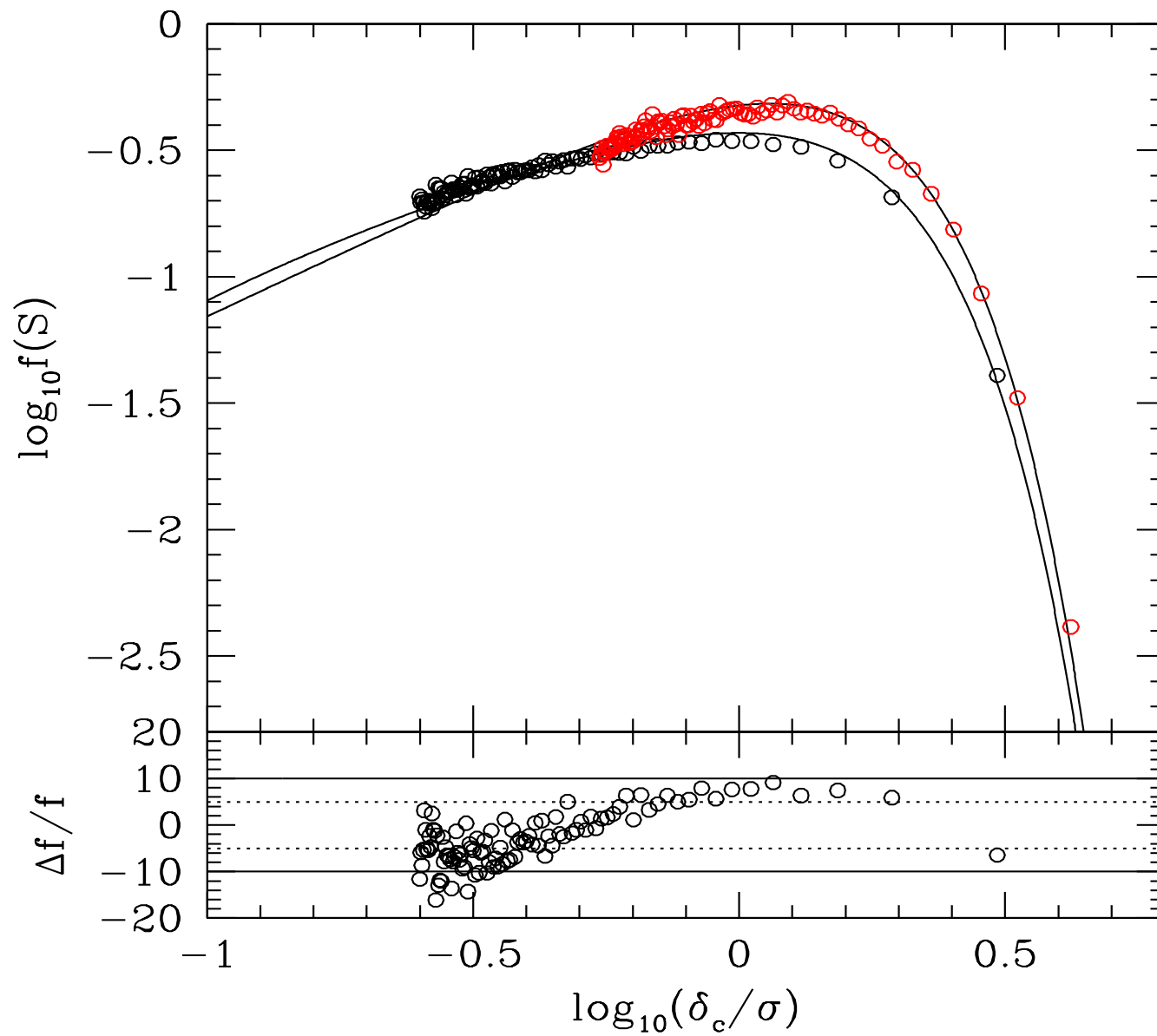
$$f_{1,\beta=0}^{m-m}(\sigma) = -\tilde{\kappa} \frac{\delta_c}{\sigma} \sqrt{\frac{2a}{\pi}} \left[ e^{-\frac{a\delta_c^2}{2\sigma^2}} - \frac{1}{2} \Gamma \left( 0, \frac{a\delta_c^2}{2\sigma^2} \right) \right]$$

$$a = \frac{1}{1+D_B}$$

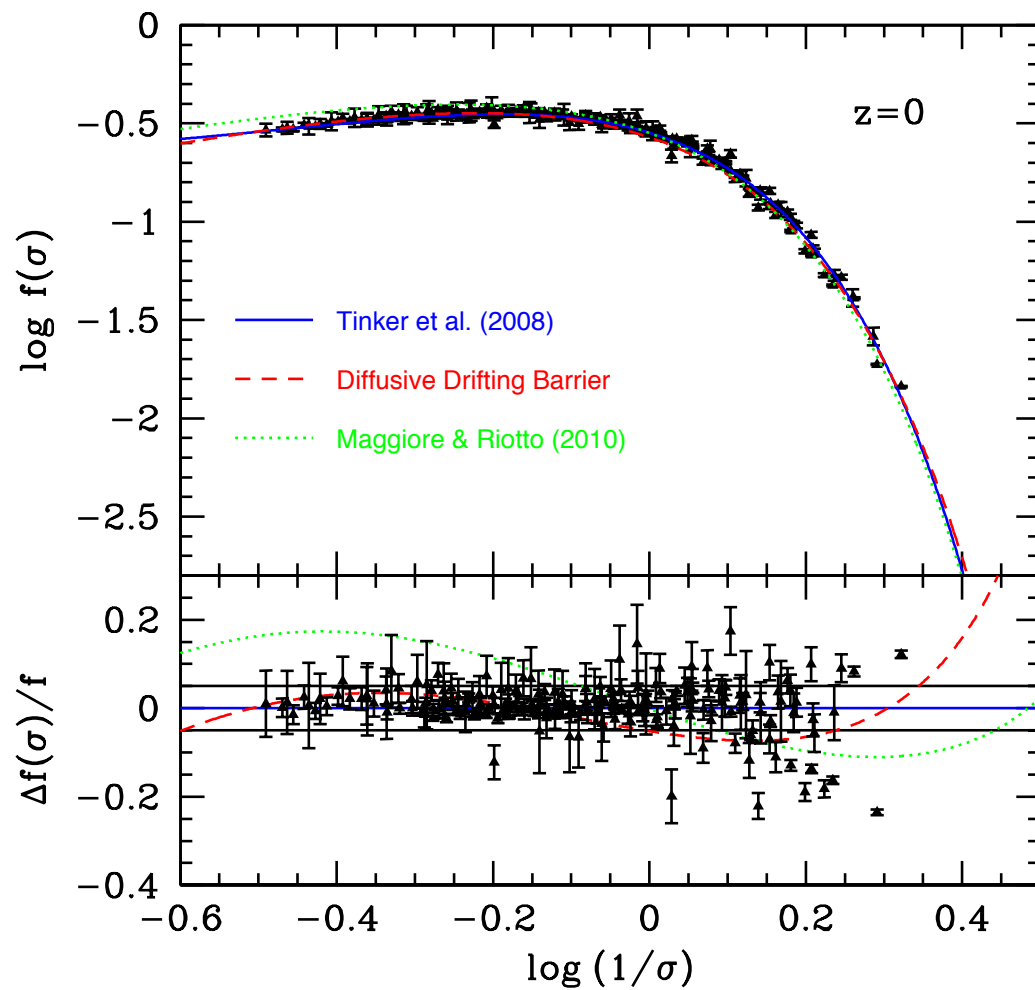
$$f_{1,\beta^{(1)}}^{m-m}(\sigma) = -\beta a \delta_c \left[ f_{1,\beta=0}^{m-m}(\sigma) + \tilde{\kappa} \operatorname{Erfc} \left( \frac{\delta_c}{\sigma} \sqrt{\frac{a}{2}} \right) \right]$$

$$\tilde{\kappa} = \kappa a$$

$$f_{1,\beta^{(2)}}^{m-m}(\sigma) = \beta^2 a \delta_c \tilde{\kappa} \left\{ a \delta_c \operatorname{Erfc} \left( \frac{\delta_c}{\sigma} \sqrt{\frac{a}{2}} \right) + \right. \\ \left. + \sigma \sqrt{\frac{a}{2\pi}} \left[ e^{-\frac{a\delta_c^2}{2\sigma^2}} \left( \frac{1}{2} - \frac{a\delta_c^2}{\sigma^2} \right) + \frac{3}{4} \frac{a\delta_c^2}{\sigma^2} \Gamma \left( 0, \frac{a\delta_c^2}{2\sigma^2} \right) \right] \right\}$$







Calibration to Tinker et al.

- $\beta=0.057$
- $D_B=0.294$

If you are interested in non-CDM mass function: what to do?

**1) Solve linear perturbation theory and estimate  $S(R)$**

- Derive form of 2-pt correlator induced by the filter

**2) Write down non-linear collapse equation (spherical or ellipsoidal models) and estimate PDF( $\delta_{\text{coll}}$ )**

- Infer moments of the barrier statistics

**3) Solve Markovian System  $\{\delta, B\}$**

**4) Compute non-Markovian corrections using Path-Integral**

# Conclusions

- DM Halo mass function is crucial in modern cosmology, physical understanding is needed
- Excursion Set is a self-consistent framework for a theoretical modeling of MF
- Path-Integral Methods allows analytical computation of generic filter corrections and consistent comparison with N-body results
- Focus on implementing physical models of halo collapse and properties of invisible components