

## Outline

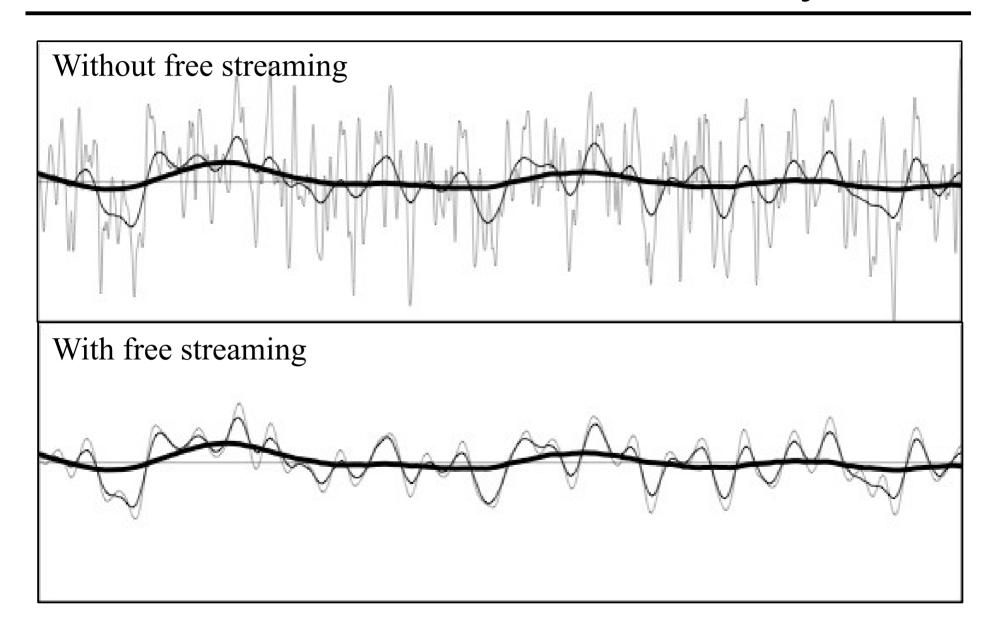
# Structure Formation and Free Streaming

Warm Dark Matter - Simulations

Modeling the Halo Mass Function

**Predictions for Cold Dark Matter** 

## Structure formation: linear density field



## Structure formation: linear density field

Equation of perturbation:

$$\frac{d^2\delta}{dt^2} + 2\frac{\dot{a}}{a}\frac{d\delta}{dt} = \left[4\pi G\rho_{\rm b}(t) - \frac{\sigma^2(t)k^2}{a^2}\right]\delta$$

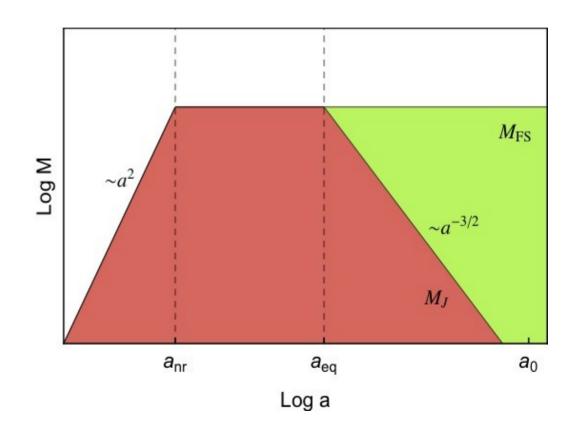
$$\delta = \frac{\rho - \rho_{\rm b}}{\rho_{\rm b}}$$

Jeans criterion:

$$\lambda_{
m J}(t) = \sqrt{rac{\pi\sigma^2(t)}{G
ho_{
m b}(t)}}$$

Free streaming criterion:

$$\lambda_{\mathrm{FS}}(t) = \int_0^{t_{\mathrm{eq}}} \frac{\sigma(t)}{a(t)} dt$$



## Structure formation: dark matter

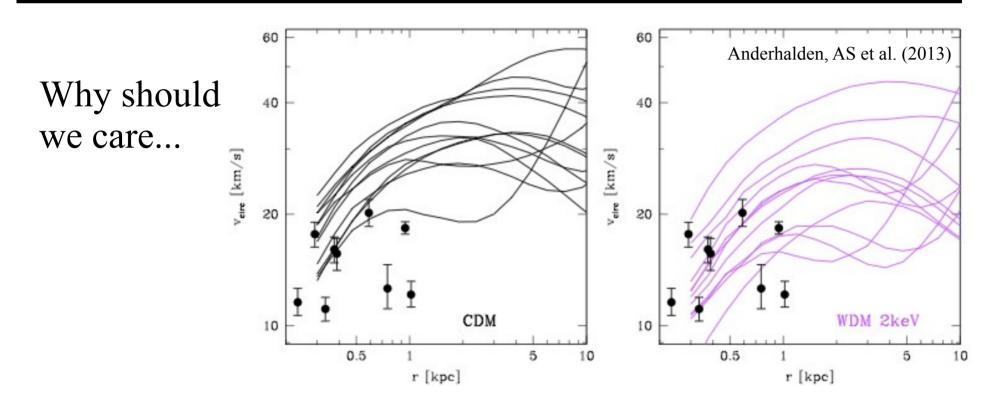
Dark Matter Candidates and smallest haloes:

• CDM (WIMP): 
$$\sim 10^{-6} M_{\odot}/h$$
 (earth-mass microhaloes)

• CDM (Axion): 
$$\sim 10^{-13} M_{\odot}/h$$

• WDM (sterile neutrino):  $\sim 10^8 M_{\odot}/h$  (dwarf galaxies)

### Structure formation: dark matter



- ... because dark matter is a mystery!
- ... because of inconsistencies in small scales structure formation.
- → WDM effects could be visible in the sky!

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# Structure Formation and Free Streaming

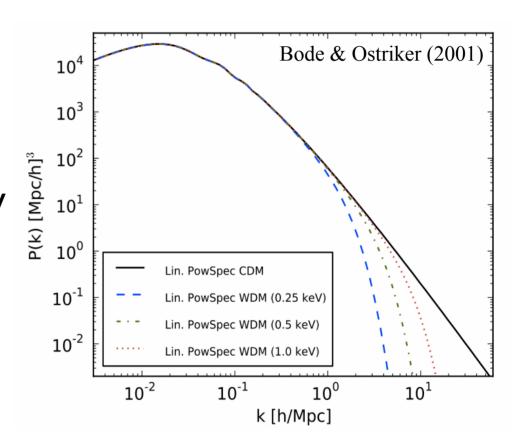
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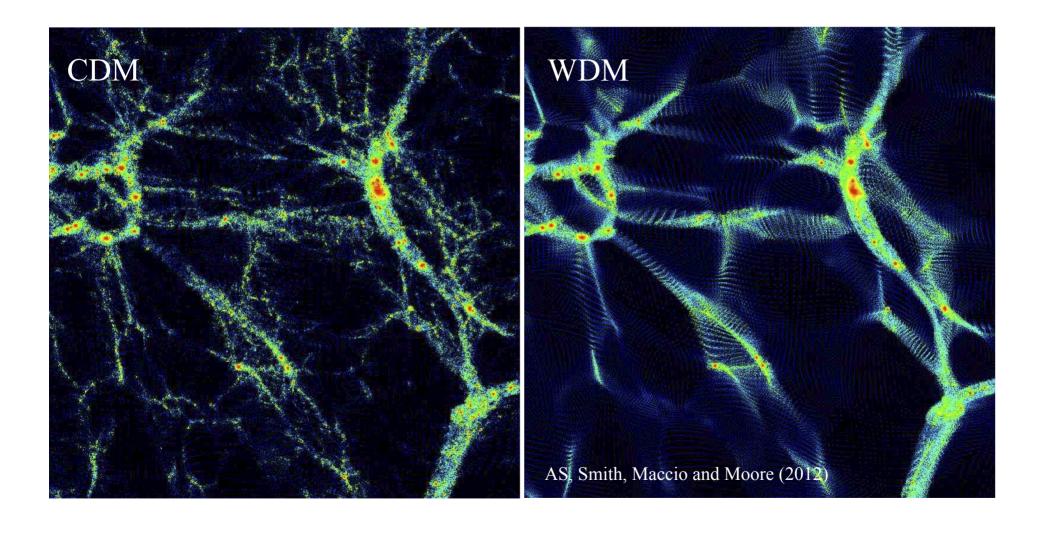
**Predictions for Cold Dark Matter** 

## WDM Simulations: Setup

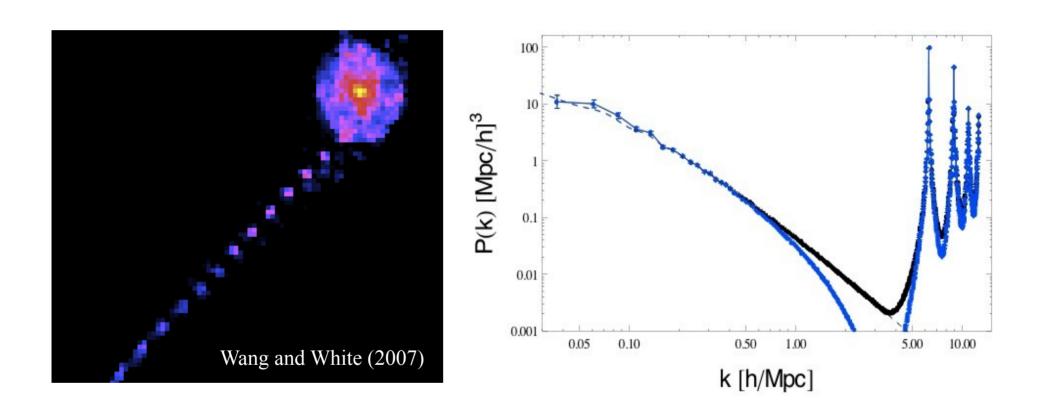
- ICs with WDM power spectrum
- No thermal velocities
- $m = \{0.25, 0.5, 1, \infty\} \text{ keV}$
- $L = \{16, 64, 256\} \text{ Mpc/h}$
- $N = \{256^3, 512^3, 1024^3\}$



## WDM Simulations: Pictures

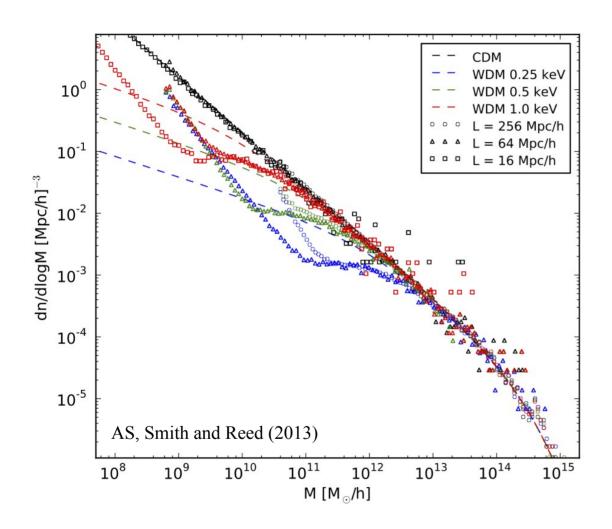


### WDM Simulations: Artifacts



Artificial clumping induced by initial grid. Simulations do not converge!

## WDM Simulations: massfct



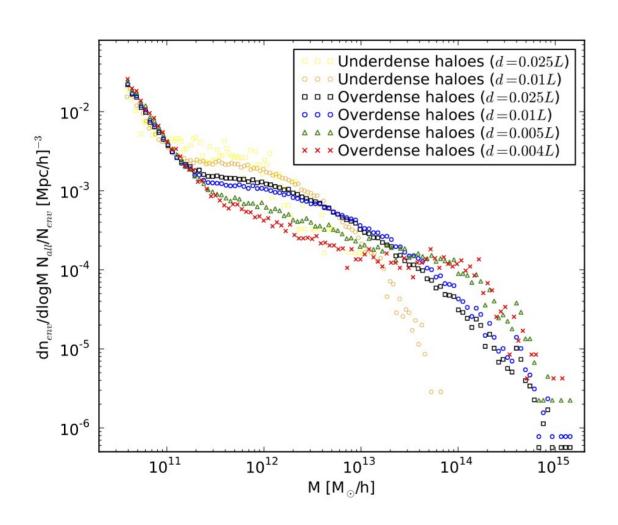
Suppressed WDM mass function.

Artificial upturn (power law).

Resolution  $\sim N^{1/3}$ .

Can artifacts be removed?

### WDM Simulations: conditional massfct



Mass function and environment:

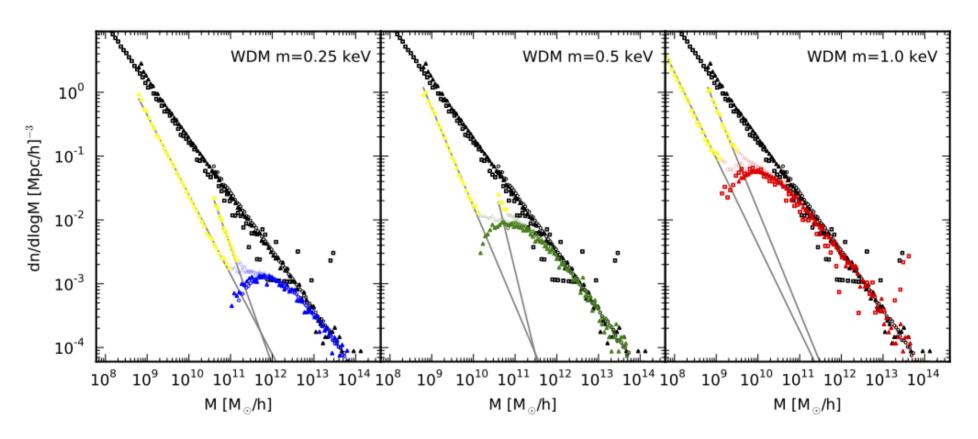
Underdense: no neighbour within distance *d*.

Overdense: at least one neighbour within *d*.

Upturn is shifted to smaller masses in underdense environemnts.

→ Artificial upturn is independent of environment!

## WDM Simulations: corrected massfct



Subtracting artificial power law.

Mass function turns over!

Roughly convergent.

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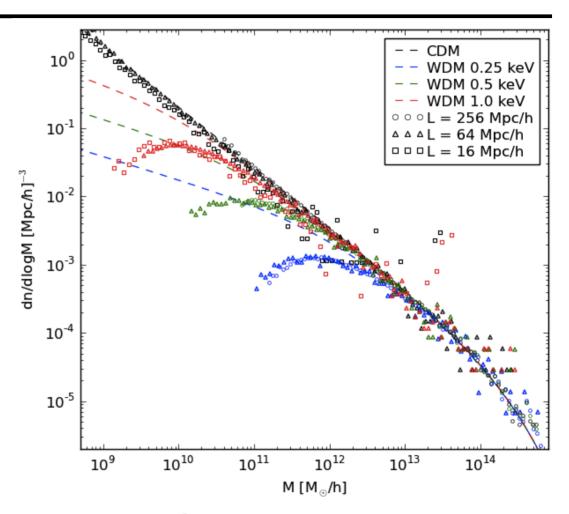
## Mass Function: Sheth-Tormen model

$$\frac{dn}{d\log M} = -\frac{1}{2}\frac{\bar{\rho}}{M}f(\nu)\frac{d\log\sigma^2}{d\log M}$$

$$f(\nu) = A\sqrt{\frac{2q\nu}{\pi}} \left[ 1 + (q\nu)^{-p} \right] e^{-q\nu/2},$$

$$\sigma^2(R) = \int \frac{d\mathbf{k}^3}{(2\pi)^3} P_{\text{Lin}}(k) W^2(kR)$$

$$\nu = \frac{\delta_c^2}{\sigma^2(M)}$$



Choice of window function:

Tophat: 
$$W_{\text{TH}} = 3 \left[ \sin y - y \cos y \right] / y$$

Sharp-k: 
$$W_{SK} = \Theta(1-y)$$
  $y = kR$ 

$$\lim_{M \to 0} \frac{dn}{d \log M} \propto M^{-1/3} = \infty$$

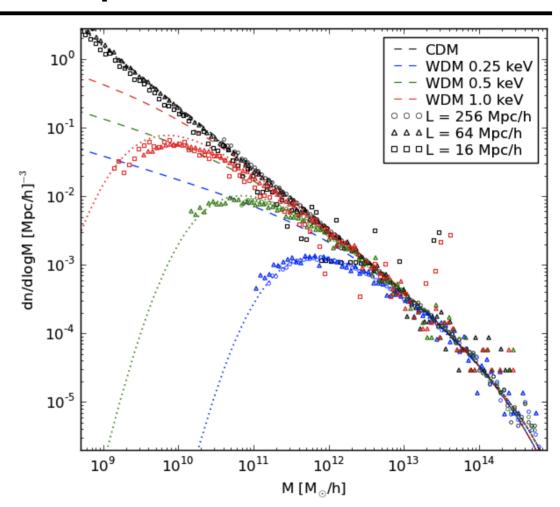
$$\lim_{M\to 0}\frac{dn}{d\log M}\propto M^{6-n/3}=0$$

$$\frac{dn}{d\log M} = -\frac{1}{2}\frac{\bar{\rho}}{M}f(\nu)\frac{d\log\sigma^2}{d\log M}$$

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#### Parameters:

Tophat: A = 0.322, p = 0.3, q = 0.707.

 $M = \frac{4\pi}{3}\bar{\rho}R^3$ 

Sharp-k: A = 0.322, p = 0.3, q = 1.0.

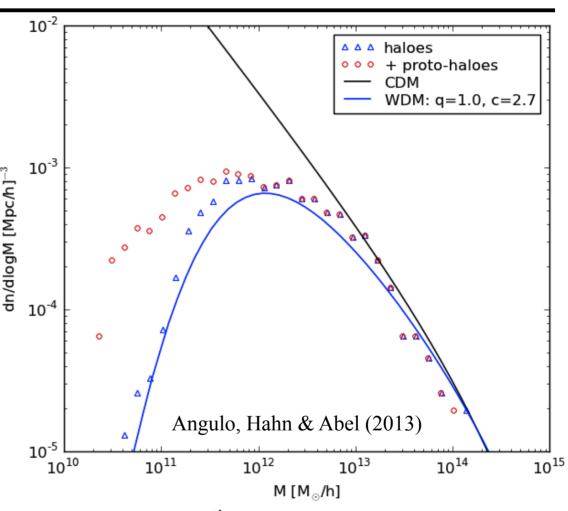
 $M = \frac{4\pi}{3}\bar{
ho}[cR]^3, \ \ c = 2.7$ 

$$\frac{dn}{d\log M} = -\frac{1}{2}\frac{\bar{\rho}}{M}f(\nu)\frac{d\log\sigma^2}{d\log M}$$

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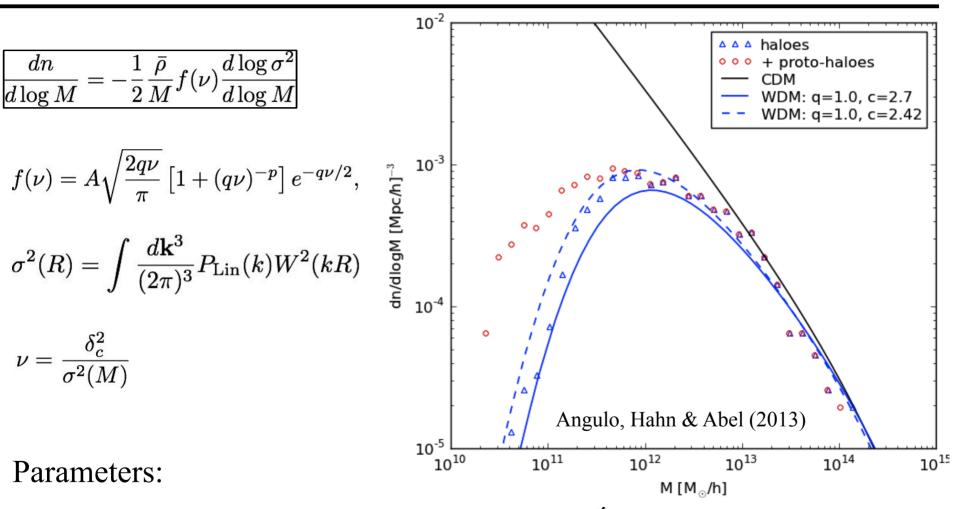
 $M = \frac{4\pi}{2}\bar{\rho}[cR]^3, \ c = 2.7$ 

$$\frac{dn}{d\log M} = -\frac{1}{2}\frac{\bar{\rho}}{M}f(\nu)\frac{d\log\sigma^2}{d\log M}$$

$$f(\nu) = A\sqrt{\frac{2q\nu}{\pi}} \left[1 + (q\nu)^{-p}\right] e^{-q\nu/2},$$

$$\sigma^2(R) = \int \frac{d\mathbf{k}^3}{(2\pi)^3} P_{\text{Lin}}(k) W^2(kR)$$

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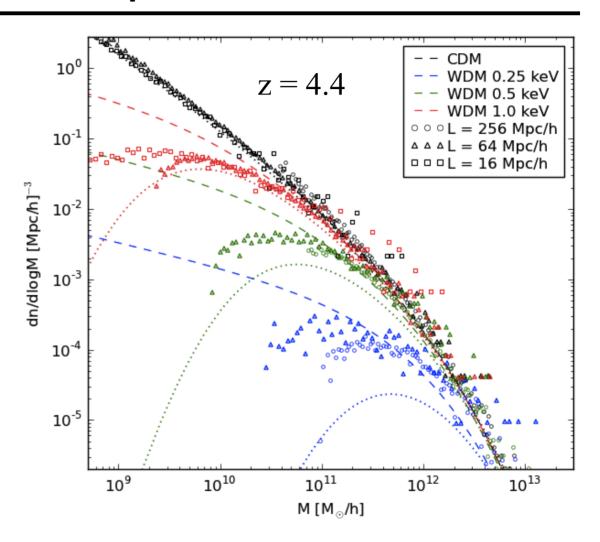
 $M=\frac{4\pi}{2}\bar{\rho}R^3$ 

Sharp-k: A = 0.322, p = 0.3, q = 1.0.

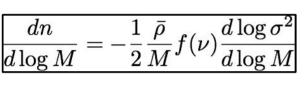
 $M = \frac{4\pi}{2}\bar{\rho}[cR]^3, \ c = 2.7$ 

Problem solved?

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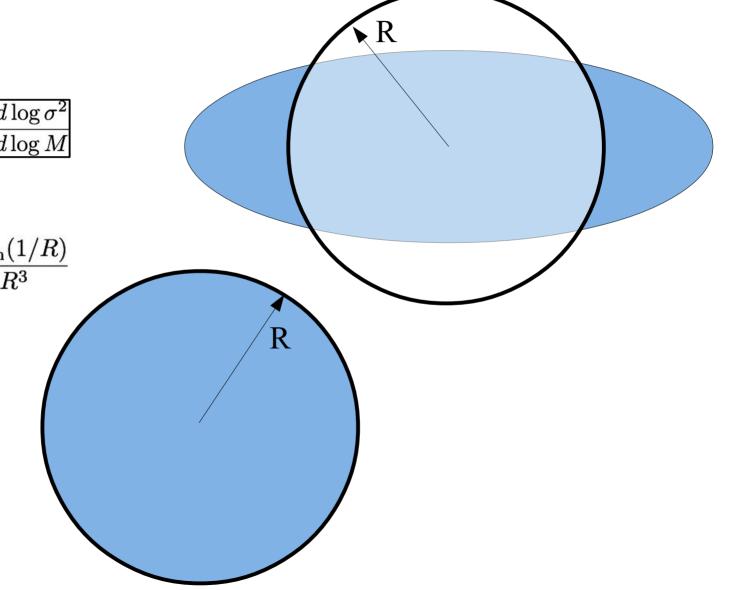


...not at high redshift!



$$\frac{d\log\sigma^2}{d\log M} = \frac{1}{6\pi^2\sigma^2} \frac{P_{\rm Lin}(1/R)}{R^3}$$

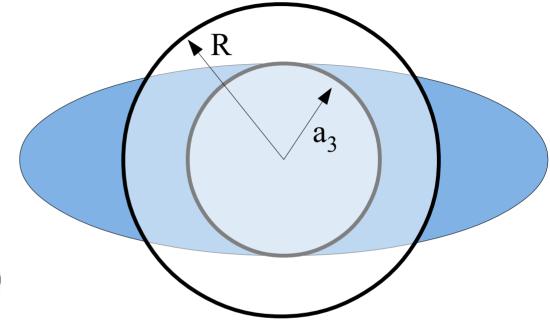
$$M=rac{4\pi}{3}ar
ho[cR]^3$$



$$rac{dn}{d\log M} = -rac{1}{2}rac{ar
ho}{M}f(
u)rac{d\log\sigma^2}{d\log M}$$

$$\frac{d\log\sigma^2}{d\log M} = \frac{1}{6\pi^2\sigma^2} \frac{P_{\mathrm{Lin}}(1/a_3)}{a_3^3} \frac{da_3}{dR} \xi(R)$$

$$R^3 = a_1 a_2 a_3 = \left(\frac{a_1}{a_3} \frac{a_2}{a_3}\right) a_3 = (\xi a_3)^3, \qquad a_3(R) = \frac{R}{\xi(R)}$$

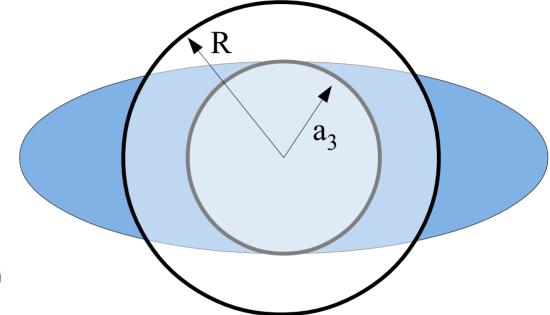


$$a_3(R) = \frac{R}{\xi(R)}$$

$$M=rac{4\pi}{3}ar
ho\left[cR(a_3)
ight]^3$$

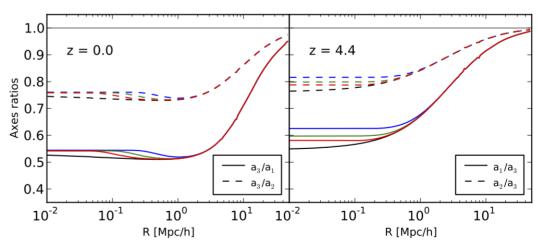
$$\frac{dn}{d\log M} = -\frac{1}{2} \frac{\bar{\rho}}{M} f(\nu) \frac{d\log \sigma^2}{d\log M}$$

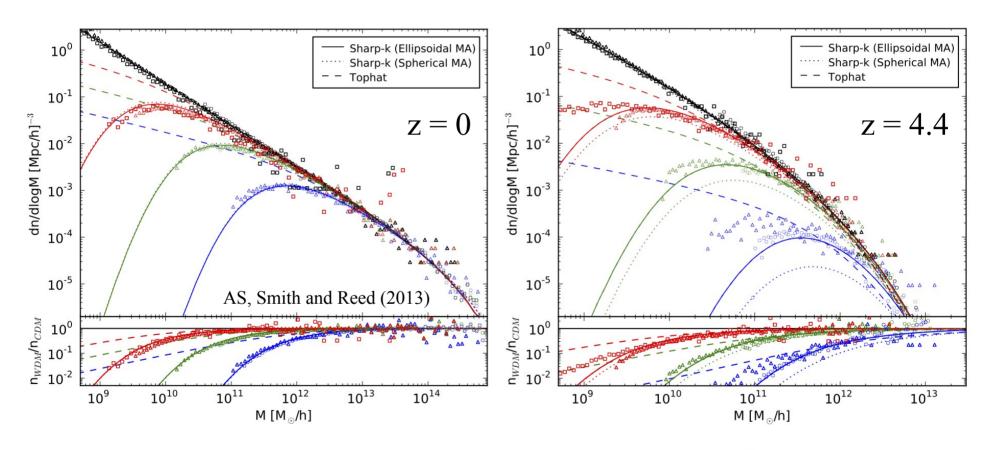
$$\frac{d \log \sigma^2}{d \log M} = \frac{1}{6\pi^2 \sigma^2} \frac{P_{\text{Lin}}(1/a_3)}{a_3^3} \frac{d a_3}{d R} \xi(R)$$



$$R^{3} = a_{1}a_{2}a_{3} = \left(\frac{a_{1}}{a_{3}}\frac{a_{2}}{a_{3}}\right)a_{3} = (\xi a_{3})^{3}$$

Shape of initial patches depend on scale and redshift (Bardeen et al 1986)





• Spherical sharp-k: A = 0.322, p = 0.3, q = 1.0.

• Ellipsoidal sharp-k: A = 0.322, p = 0.3, q = 0.75

$$M = \frac{4\pi}{3}\bar{
ho}[cR]^3, \ \ c = 2.7$$

$$M = rac{4\pi}{3}ar{
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ight]^3, \;\; c = 2.0$$

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# Cold Dark Matter: comparison

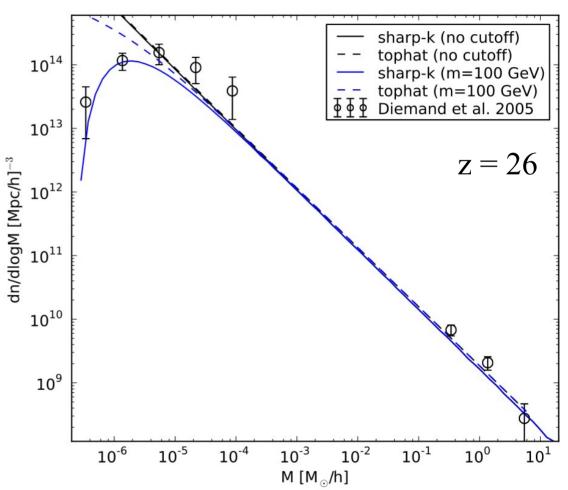
# Does model work for neutralino-CDM?

- Free streaming scale: 10<sup>15</sup> times smaller!
- Different cutoff

# Cold Dark Matter: comparison

Does model work for neutralino-CDM?

- Free streaming scale: 10<sup>15</sup> times smaller!
- Different cutoff



Diemand et al. (2006): m = 100 GeV, T<sub>dk</sub> = 28 MeV

