



# The challenges of relic neutrino direct detection and status of the **PTOLEMY** experiment

Princeton Tritium Observatory for Light, Early-universe,  
Massive-neutrino Yield (PTOLEMY)

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Princeton University

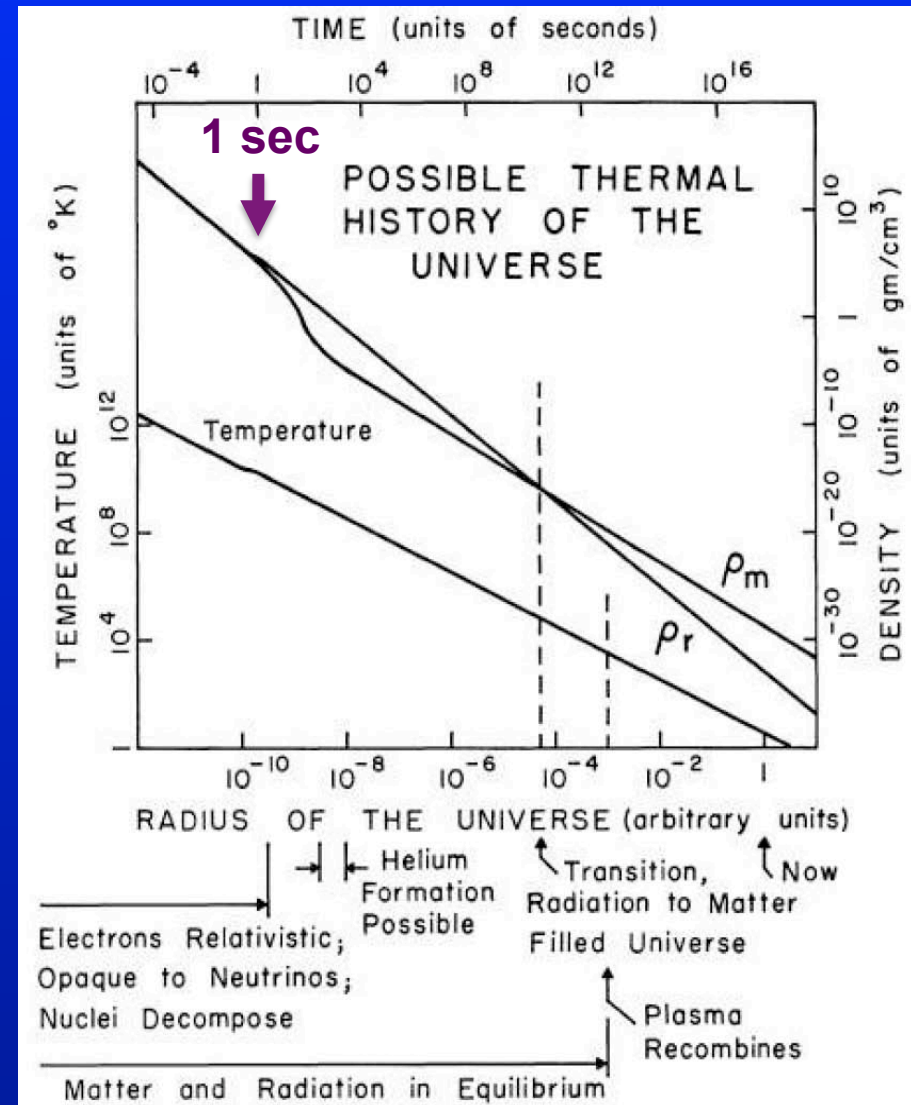
Chalonge Meudon Workshop 2014

5 June 2014

# Looking Back in Time



- The Universe was not always as cold and dark as it is today – there are a host of landmark measurements that track the thermal history of the universe
- Few measurements, however, reach back as far in time as  $\sim 1$  second after the Big Bang
  - At  $\sim 1$  second the hot, expanding universe is believed to have become transparent to neutrinos
  - In the present universe, relic neutrinos are predicted to be at a temperature of 1.9K ( $1.7 \times 10^{-4}$  eV) and to have an average number density of  $\sim 56/\text{cm}^3$  per lepton flavor



Dicke, Peebles, Roll, Wilkinson (1965)



# Big Bang Prediction I



When the mean free path exceeds the horizon size, the neutrinos decouple from matter

$$\lambda_v \sim \frac{1}{\sigma_v n_e} \sim \frac{1}{(G_F^2 T^2) T^3}$$

$$\lambda_h \sim \frac{1}{\sqrt{G\rho}} \sim \frac{M_{Pl}}{T^2}$$

$$\frac{\lambda_h}{\lambda_v} \sim \left( \frac{T}{T_{vd}} \right)^3 \sim M_{Pl} G_F^2 T^3$$

$$T_{vd} \sim M_{Pl}^{-1/3} G_F^{-2/3} \sim 1 \text{ MeV}$$

Neutrinos decouple before  $e^+e^-$  annihilation,  $e^+e^-$  heats up photons  
2 photon + 7/8 ( 2 electron + 2 positron )  $\rightarrow$  2 photon

Relic neutrino temperature in lock step with photons and both drop at the same rate with the Hubble expansion

$$T_\nu(t) = T_\nu(t_{vd}) \frac{a(t_{vd})}{a(t)} = \left( \frac{4}{11} \right)^{1/3} T_{CMB}$$

$$T_\nu \sim 1.95 \text{ K}$$

# Big Bang Prediction II



Relic neutrino number density follows photon number density

$$n_\nu = \left(\frac{3}{4}\right)\left(\frac{4}{11}\right)n_\gamma = 112/\text{cm}^3$$

per neutrino species  
(neutrino+antineutrino)

Present-day relic neutrinos are distributed in velocity according to their original relativistic Fermi-Dirac distribution at one second after the Big Bang

$$g(p_\nu) = \frac{1}{1 + e^{p_\nu/T}}$$

in the relativistic limit

$$E \approx p_\nu + \frac{m^2}{2p_\nu}$$

$$\langle v_{rms} \rangle \propto T/m_\nu \text{ instead of } \propto \sqrt{T/m_\nu}$$

Relic velocity depends on mass  $\langle v_{rms} \rangle = 160 \text{ km/s } (1 \text{ eV}/m_\nu)$



# Timing of Neutrino Decoupling



- Timing is essential to Big Bang predictions
  - Neutrino decoupling at  $\sim 1$  MeV ( $\sim 1$  sec)
    - Weak interactions constantly regenerate neutrons
  - Neutron-Proton mass difference  $m_n - m_p \sim 1.3$  MeV
  - Deuterium ( $\rightarrow$  Helium) at  $\sim 0.07$  MeV ( $\sim 132$  sec) compared to neutron lifetime of  $\sim 886.7$  sec
    - $n/p \sim 0.15 * 0.74 \sim 0.11$  at the start of nucleosynthesis

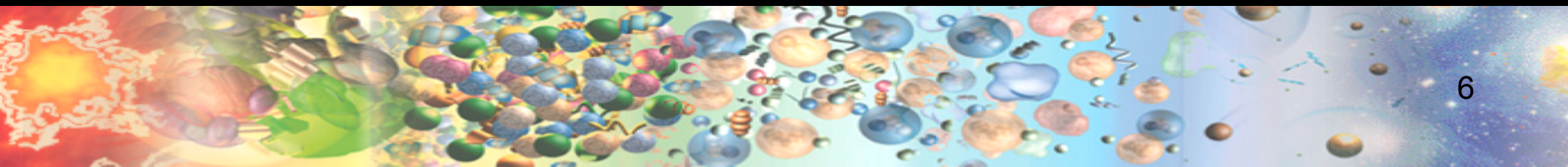
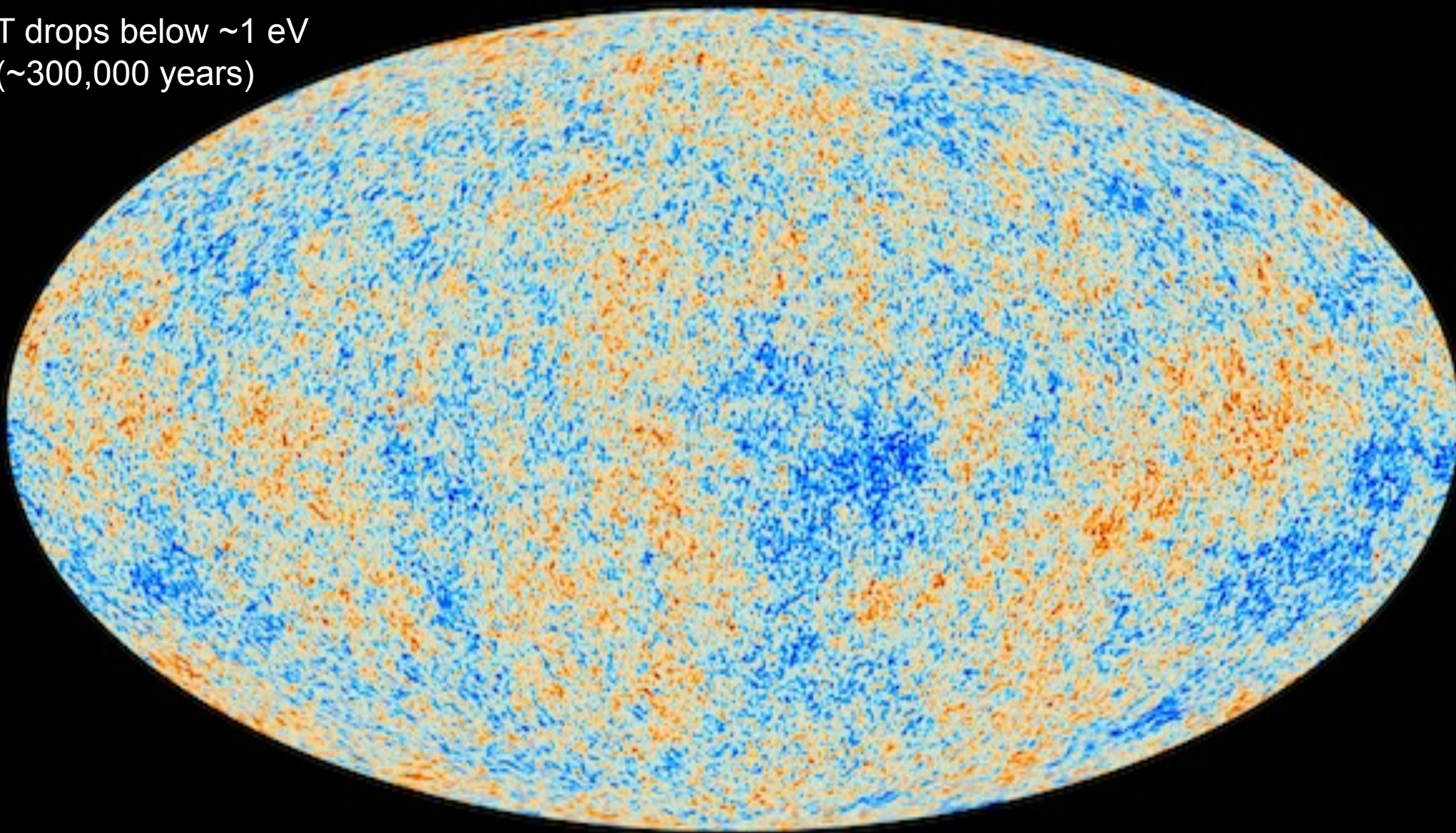
Not much wiggle room for the standard BBN prediction



# Cosmic Background Radiation

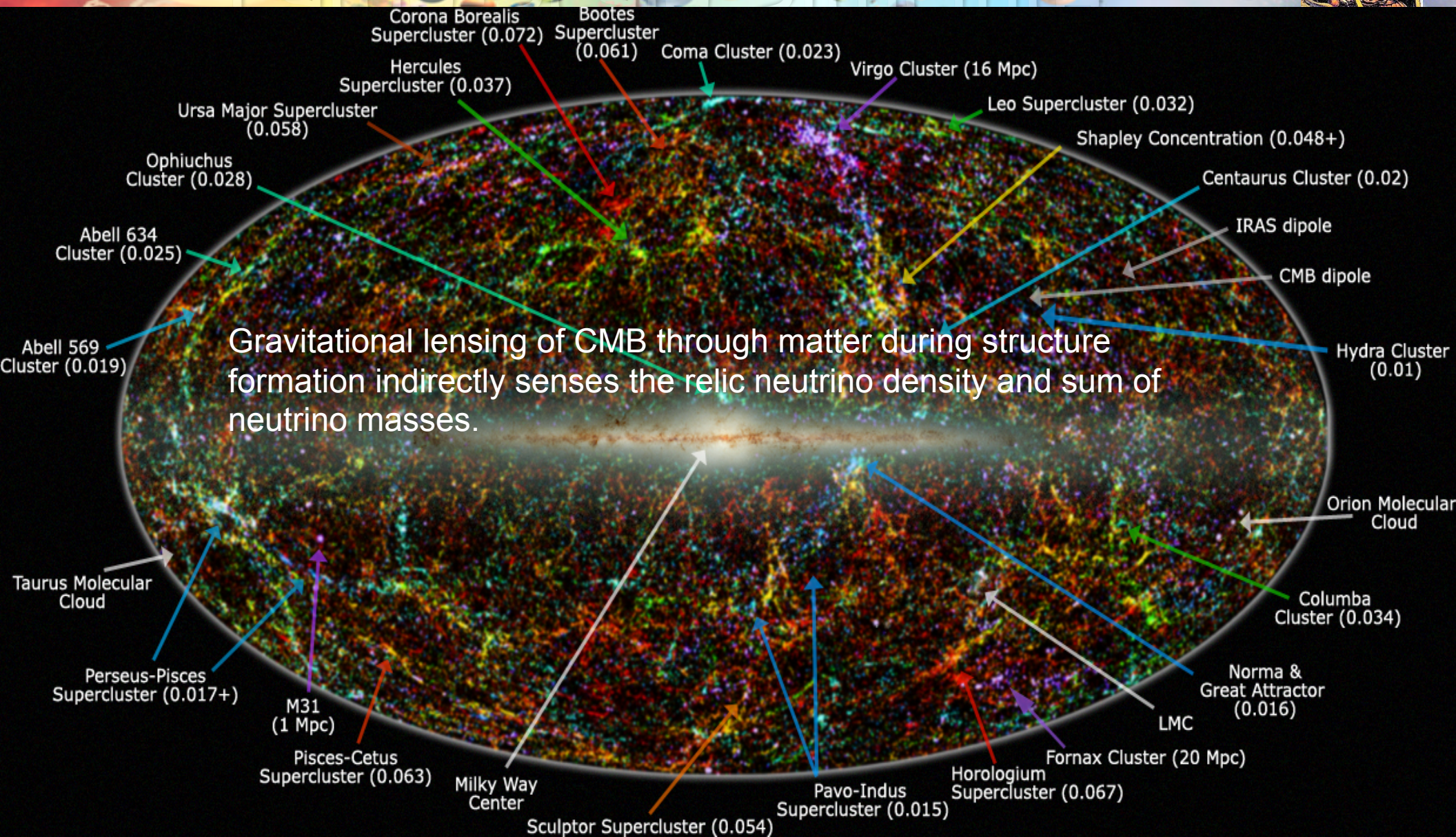


T drops below  $\sim 1$  eV  
( $\sim 300,000$  years)



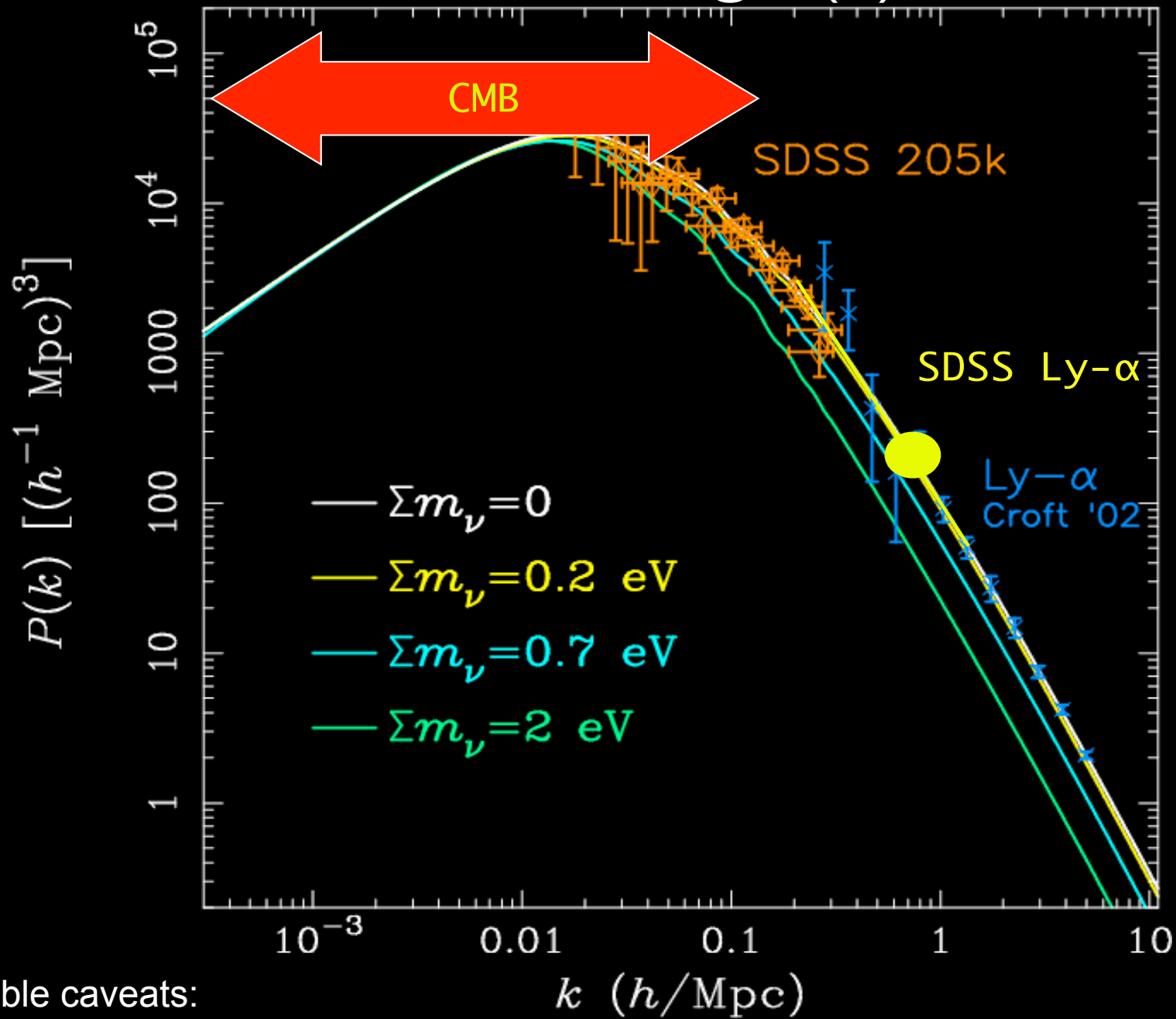


# Large-Scale Structure





# Measuring $P(k)$

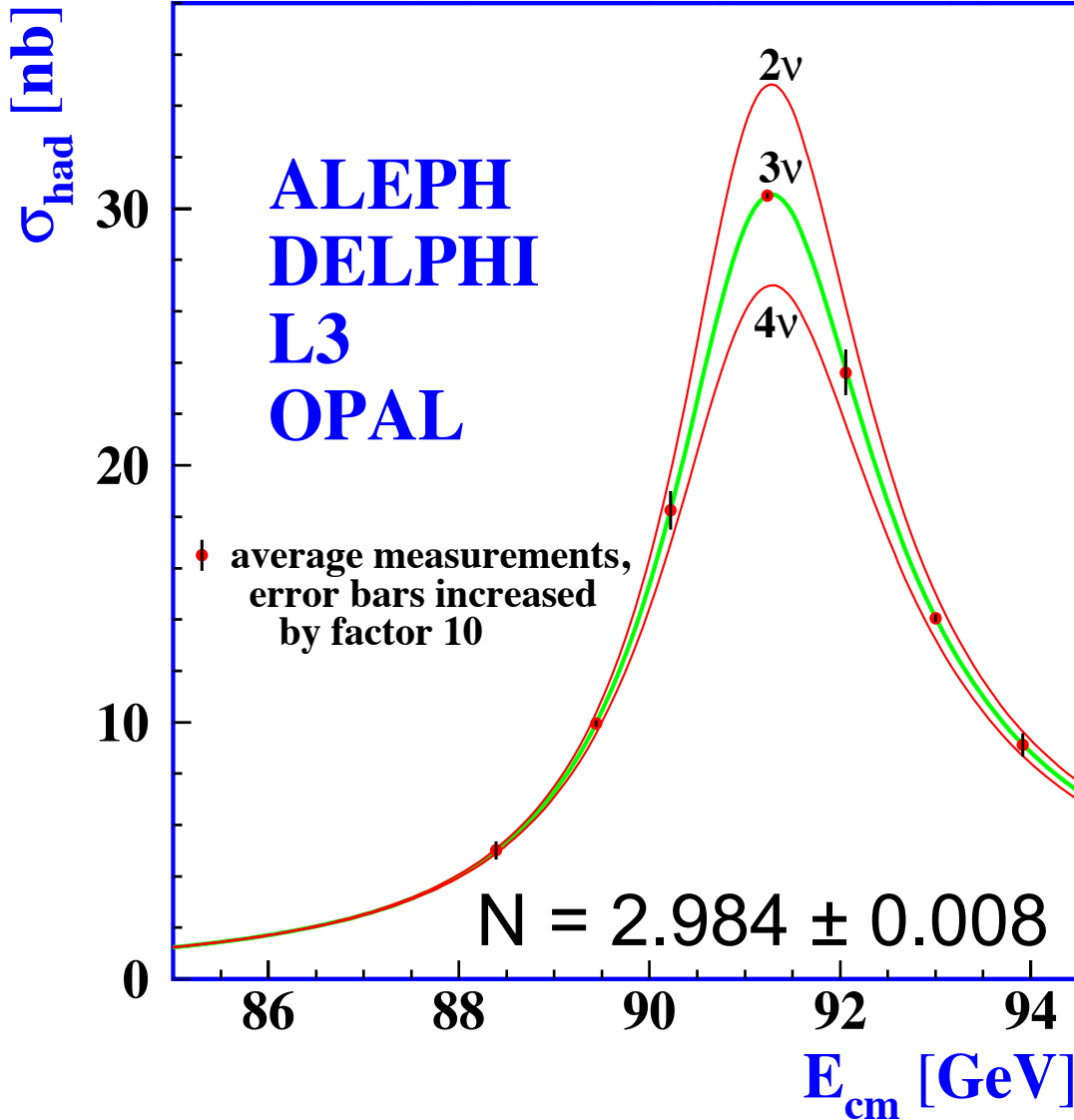


Some notable caveats:

Bounds depend on cosmology assumptions, such as the dark energy contribution to the equation of state. One can also have a delay in the matter-radiation transition from dark radiation (the number of relativistic degrees of freedom above  $N=3.04$ ).



# Neutrino Counting



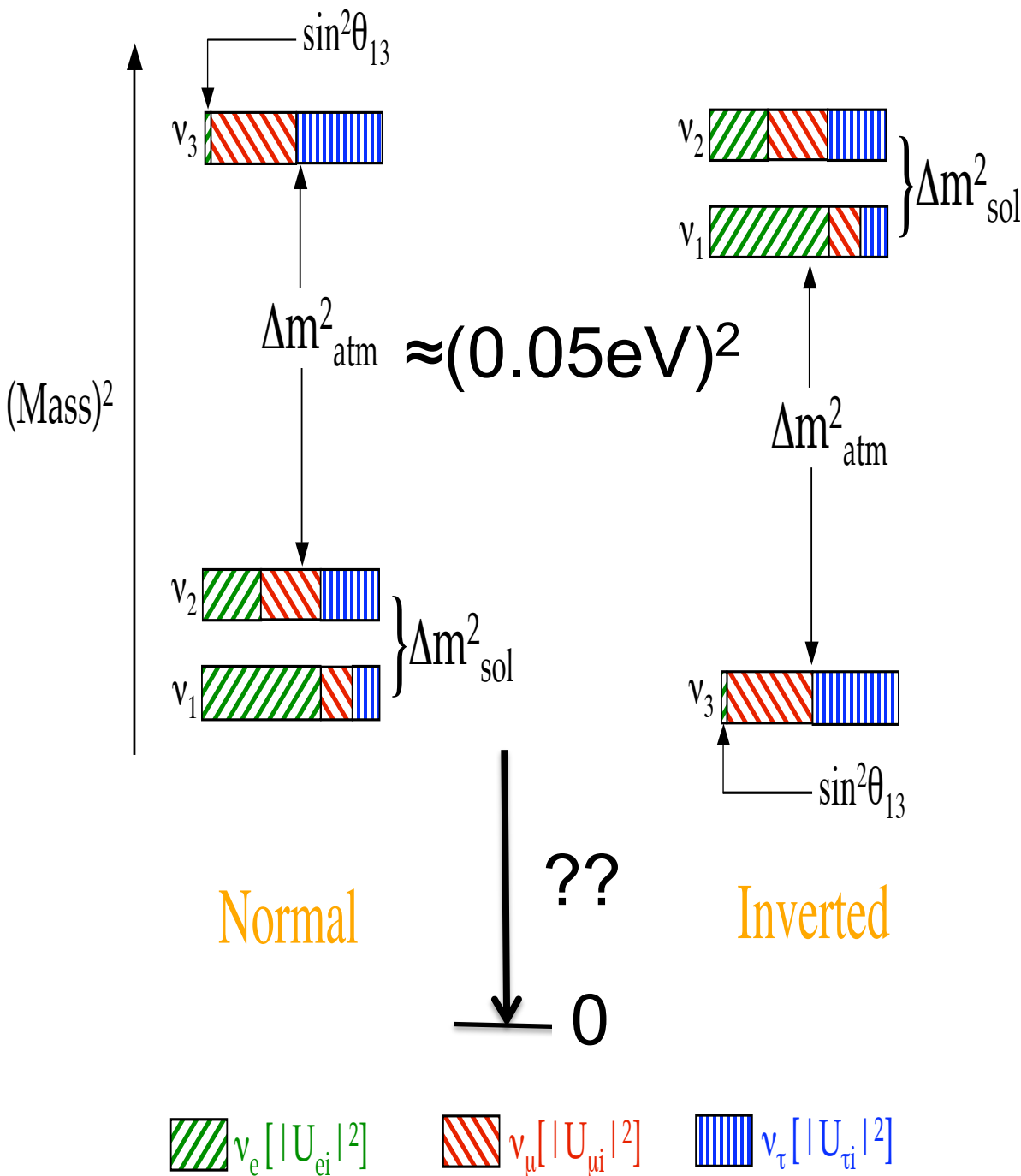
Produce  $\sim 1\text{M}$  Z bosons at an  $e^+e^-$  collider

Scan the line shape in center-of-mass energy

Count the number of hadronic Z decays

Compute the total width from visible decays and add an invisible width scaled by the SM couplings to neutrinos

# Neutrino Masses from Oscillations



An incredible phenomenon appeared when neutrinos were measured from different sources: solar, atmospheric, reactor, accelerator.

A neutrino created with a definite lepton flavor (in this case, electron or muon) would arrive with a lower probability to be detected with the same flavor and a non-zero probability to have mixed into another flavor.

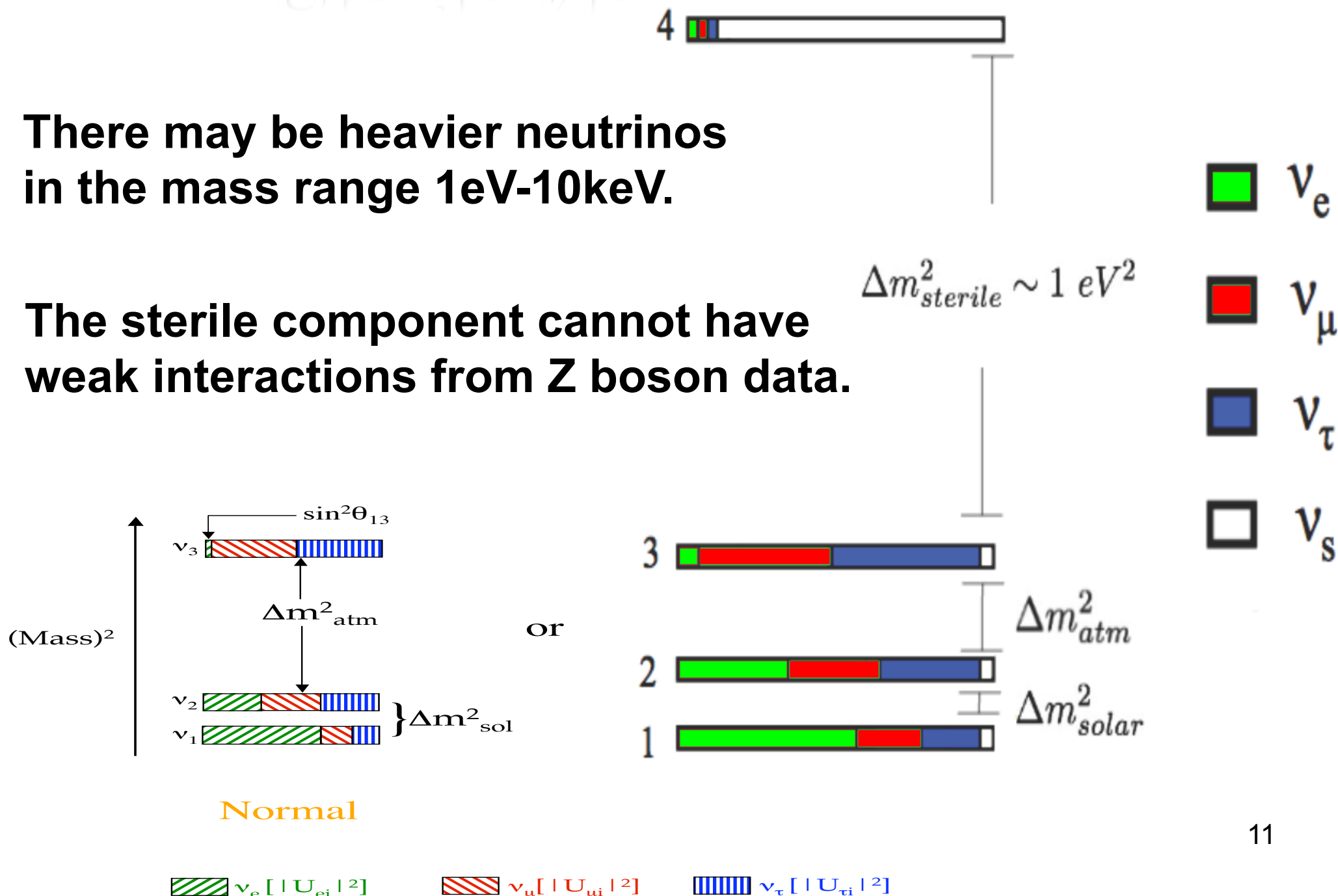


# Sterile Neutrinos?



There may be heavier neutrinos in the mass range 1eV-10keV.

The sterile component cannot have weak interactions from Z boson data.





# Relic Neutrino Detection

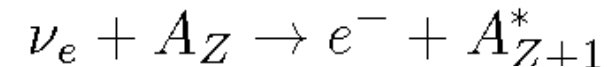
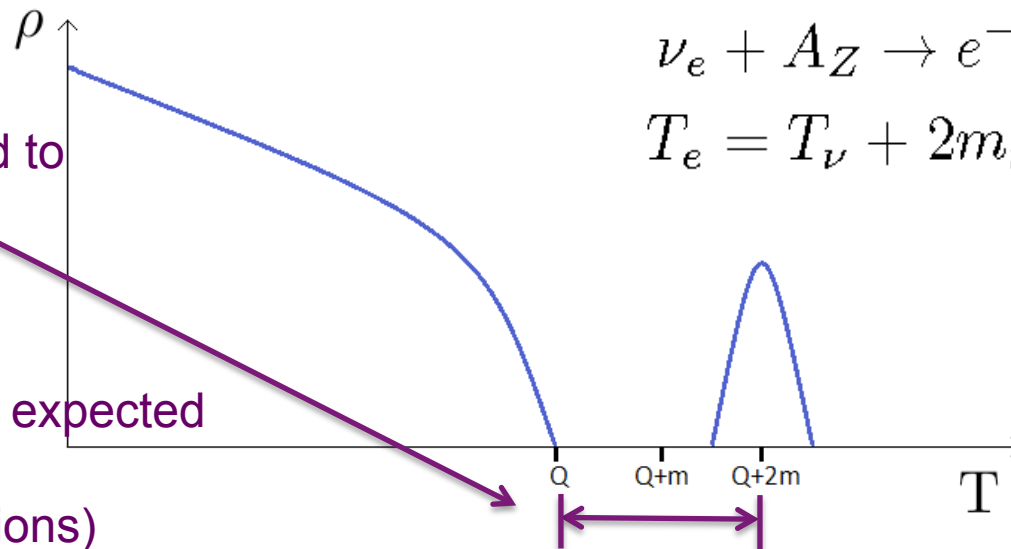


- Basic concepts for relic neutrino detection were laid out in a paper by Steven Weinberg in 1962 [*Phys. Rev.* 128:3, 1457]
  - Look for relic neutrino capture on tritium by measuring electrons at

## What do we know?

Gap ( $2m$ ) constrained to  $< \sim 0.6\text{eV}$  from Cosmology

(some electron flavor expected with  $2m > 0.1\text{eV}$  from neutrino oscillations)



$$T_e = T_\nu + 2m_\nu + Q$$

Figure 1: Emitted electron density of states vs kinetic energy for neutrino capture on beta decaying nuclei. The spike at  $Q + 2m$  is the CNB signal

# Hydrogen (Isotope) Bonding



Tritium experiments typically use diatomic tritium  $T^2$  where the bond strength is approximately 4eV.

But what happens when one T atom decays?

Answer: The maximum  $He^3$  recoil energy is  $\sim 3eV$ .  $He^3$  stays bound to the remaining T to form a T- $He^3$  molecule – and can fall into a number of closely spaced rotational and vibrational excited states



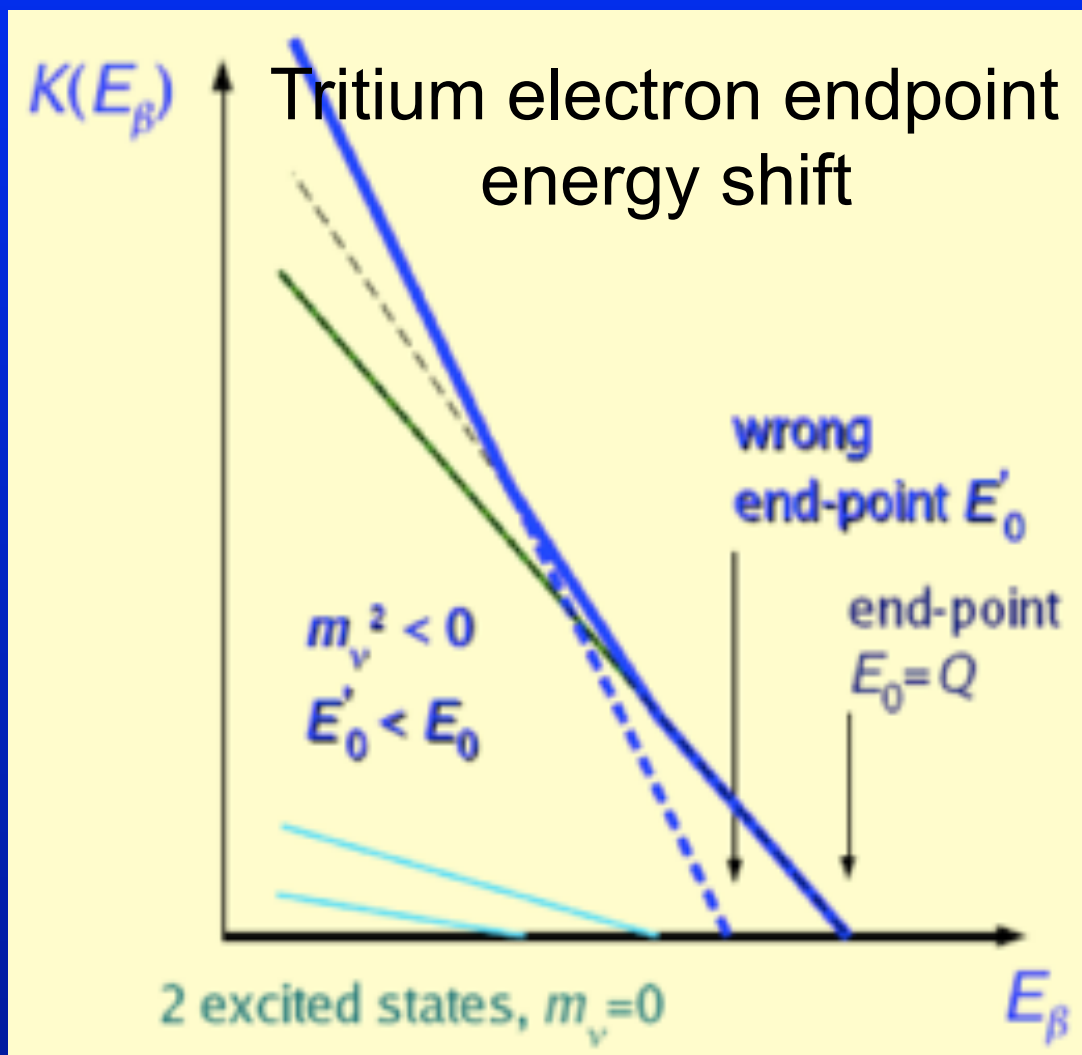
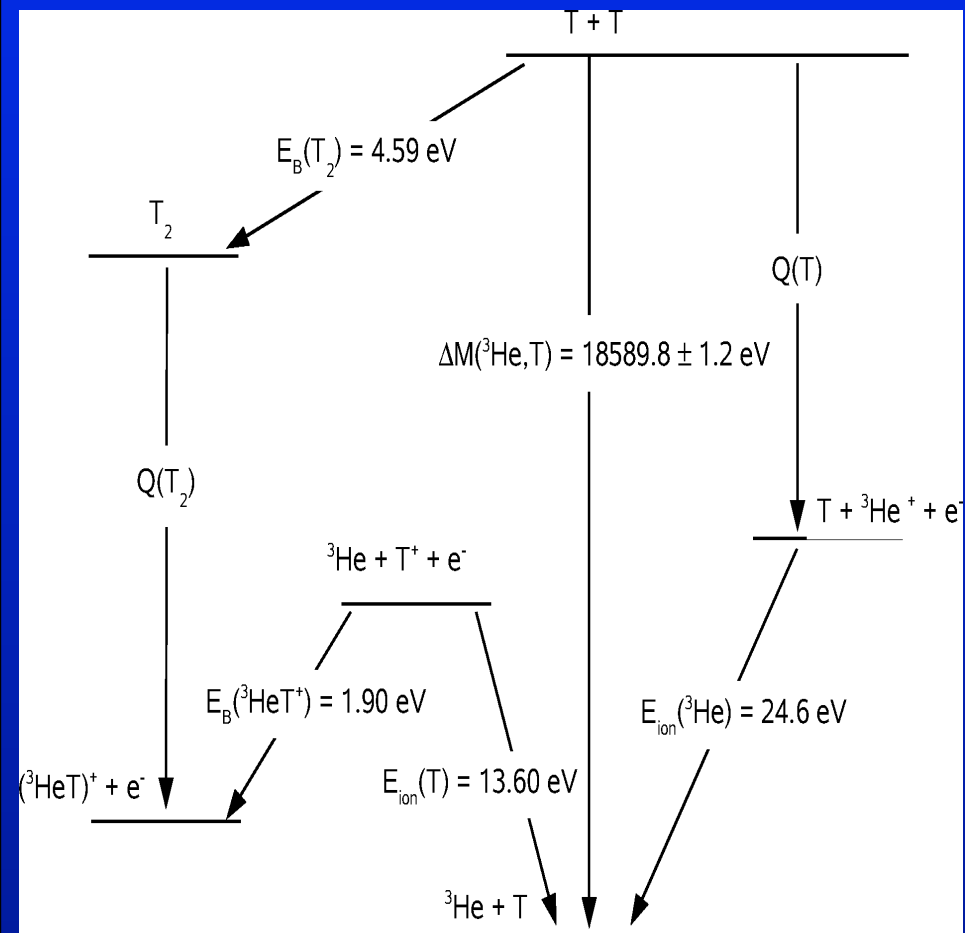
Quantum Mechanics tells us that the outgoing electron energy depends on the change in the binding energy of  $T^2$  to  $(THe^3)^*$



# He<sup>3</sup> Binding Energy Shift



## T-T → T-He<sup>3</sup> Level Diagram

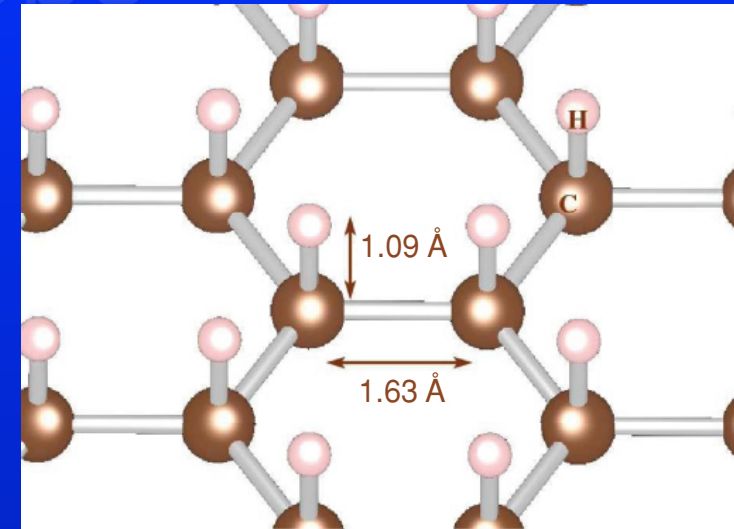


# Tritium on Graphene



- In the hunt for alternative energies, there has been a great focus on the development of Hydrogen storage systems

- Hydrogen binds to the surface of graphene in a solid form (6%wt) at room temperature, but with a weak enough binding that the hydrogen can be readily released



## Single-sided-hydrogenated Graphene

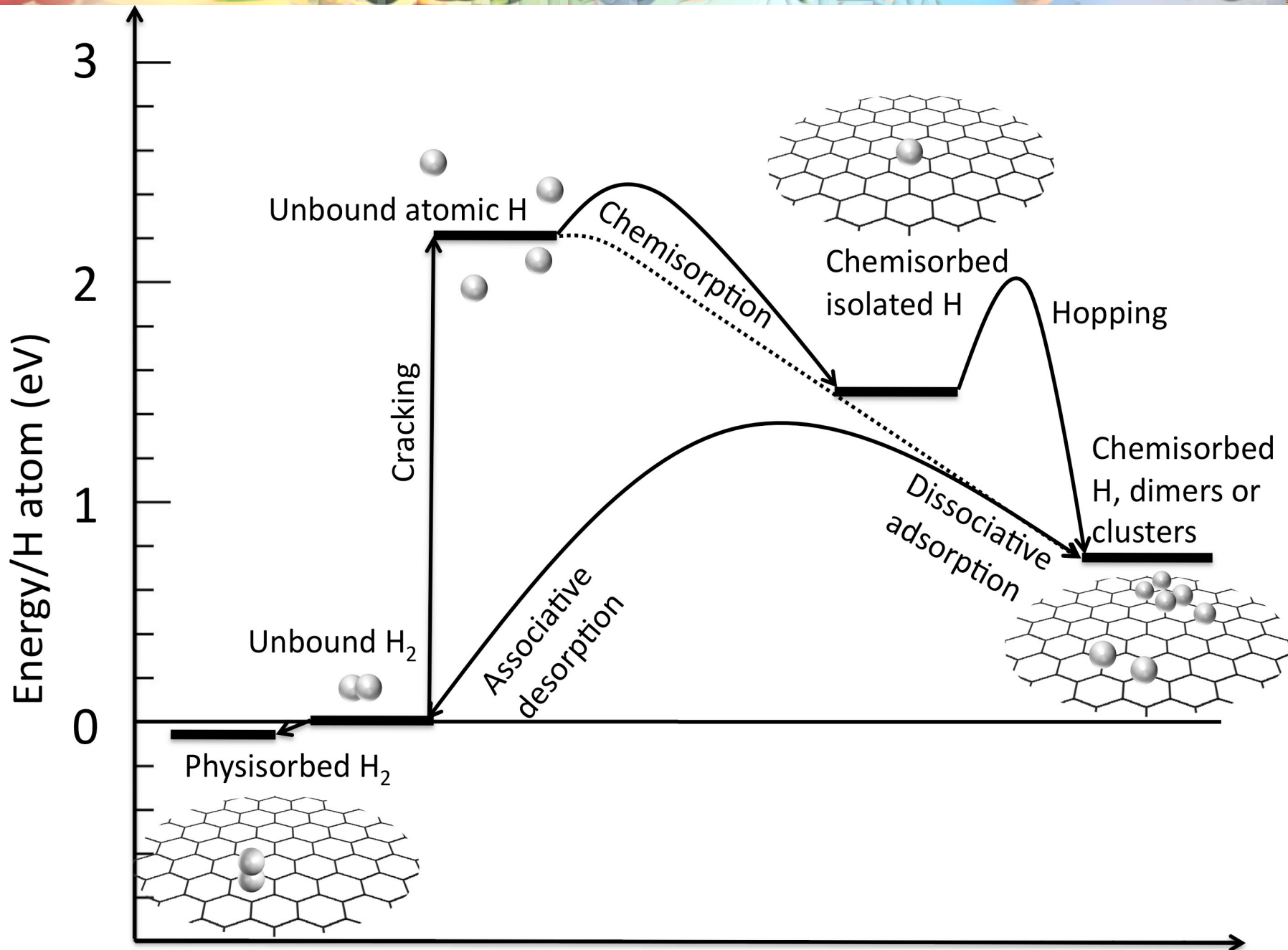
- Planar (uniform bond length)
- Semiconductor (~Si gap)
- Polarized tritium(?)

$\sim 3 \times 10^{13}$  T/mm<sup>2</sup> ( $\sim 80$ kHz of decays/mm<sup>2</sup>)

Different forms of hydrogenated graphene have a hydrogen binding energy less than 3eV – with no binding for He<sup>3</sup>.



# Chemist's View



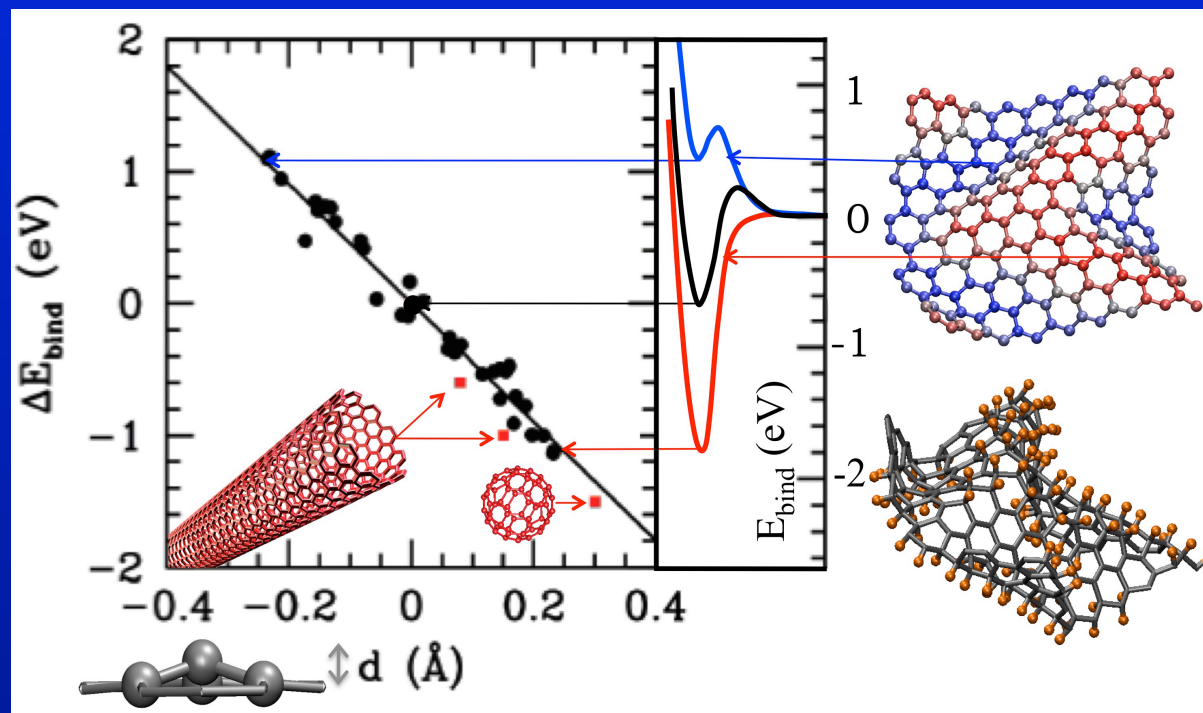
# When in Doubt, Measure!



- 1<sup>st</sup> proposal of the PTOLEMY project is to use the setup to measure hydrogen bond strength differences using tritium beta decay energy measurements and a high resolution microcalorimeter

**New C-H bond measurement technique:**

May help accelerate Hydrogen fuel cell/storage research





# Relic Neutrino Capture Rates



- Target mass: 100 grams of tritium ( $2 \times 10^{25}$  nuclei)
- Capture cross section  $\times (v/c) \sim 10^{-44}$  cm<sup>2</sup> (flat up to 10 keV)
- (Very Rough) Estimate of Relic Neutrino Capture Rate:  
( $56 \nu_e/\text{cm}^3$ ) ( $2 \times 10^{25}$  nuclei) ( $10^{-44}$  cm<sup>2</sup>) ( $3 \times 10^{10}$  cm/s) ( $3 \times 10^7$  s)  $\sim 10$  events/yr

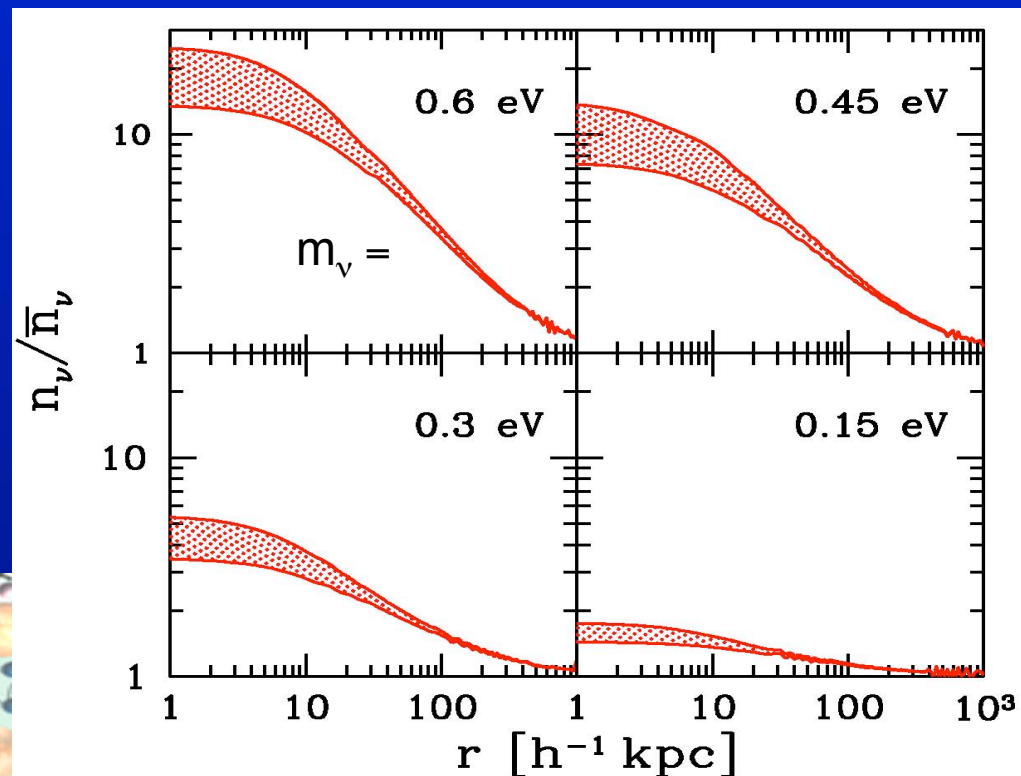
Lazauskas, Vogel, Volpe: J.Phys.G G35 (2008) 025001.

Cocco, Mangano, Messina: JCAP 0706 (2007) 015

$$\sigma(v/c) = (7.84 \pm 0.03) \times 10^{-45} \text{ cm}^2$$

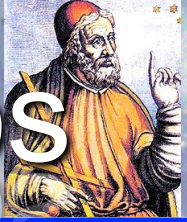
Gravitational clumping could potentially increase the local number of relic neutrinos.

For low masses  $\sim 0.15$  eV, the local enhancement is  $\sim \times 1.5$



Ringwald and Wong (2004)

# Dirac versus Majorana Neutrinos



Long, Lunardini, Sabancilar: arXiv:1405.7654

“Detecting non-relativistic cosmic neutrinos by capture on tritium: phenomenology and physics potential” → **Factor of 2 difference in capture rate**

Relic neutrino capture rate on tritium depends on whether neutrinos are Dirac or Majorana:

Neutrinos decouple at relativistic energies

Helicity is conserved as the universe expands and the relic neutrinos become non-relativistic

Dirac: initially left-handed chiral=helical neutrinos and right-handed chiral=helical anti-neutrinos are active → cooldown → ~half of left-handed helical Dirac neutrinos are right-handed chiral (non-active) → Factor of 2 drop in present-day capture rate on tritium for Dirac

Majorana: initially left-handed chiral=helical neutrinos=antineutrinos and right-handed chiral=helical neutrinos=antineutrinos are active → cooldown → No change in present-day capture rate (heavy neutrino components are decoupled)

**First Majorana/Dirac test outside of neutrinoless double-beta decay?**

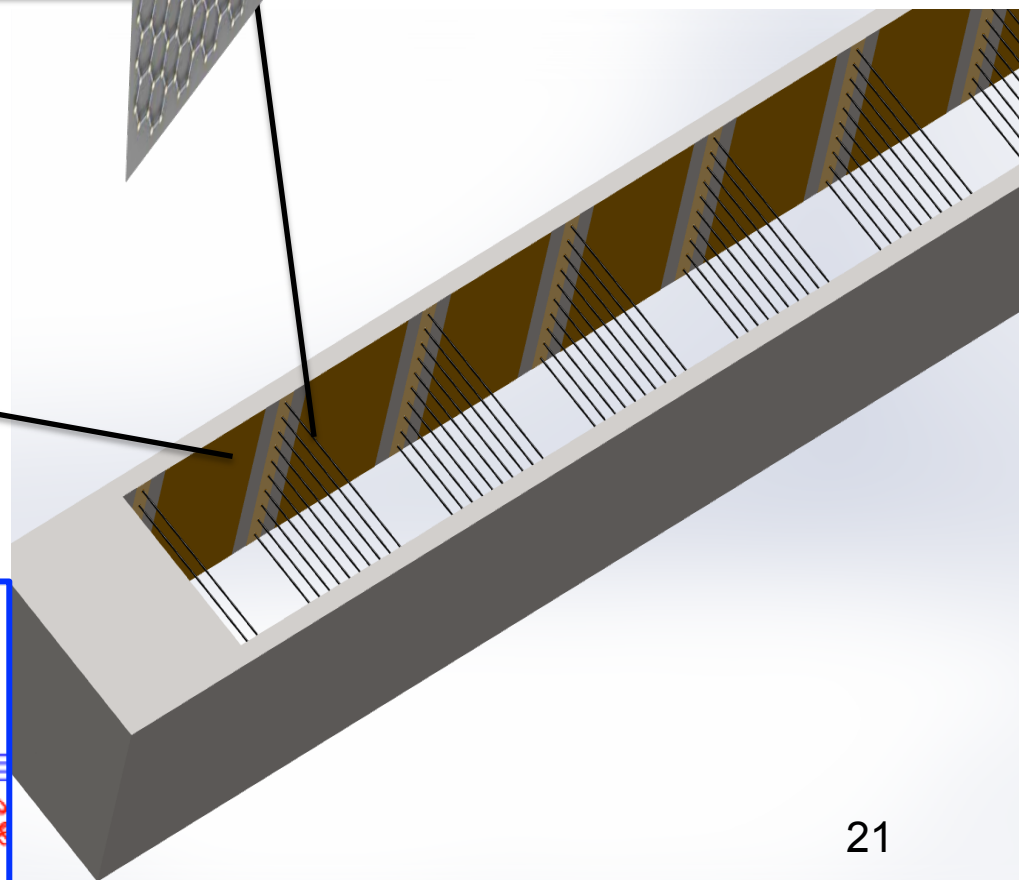
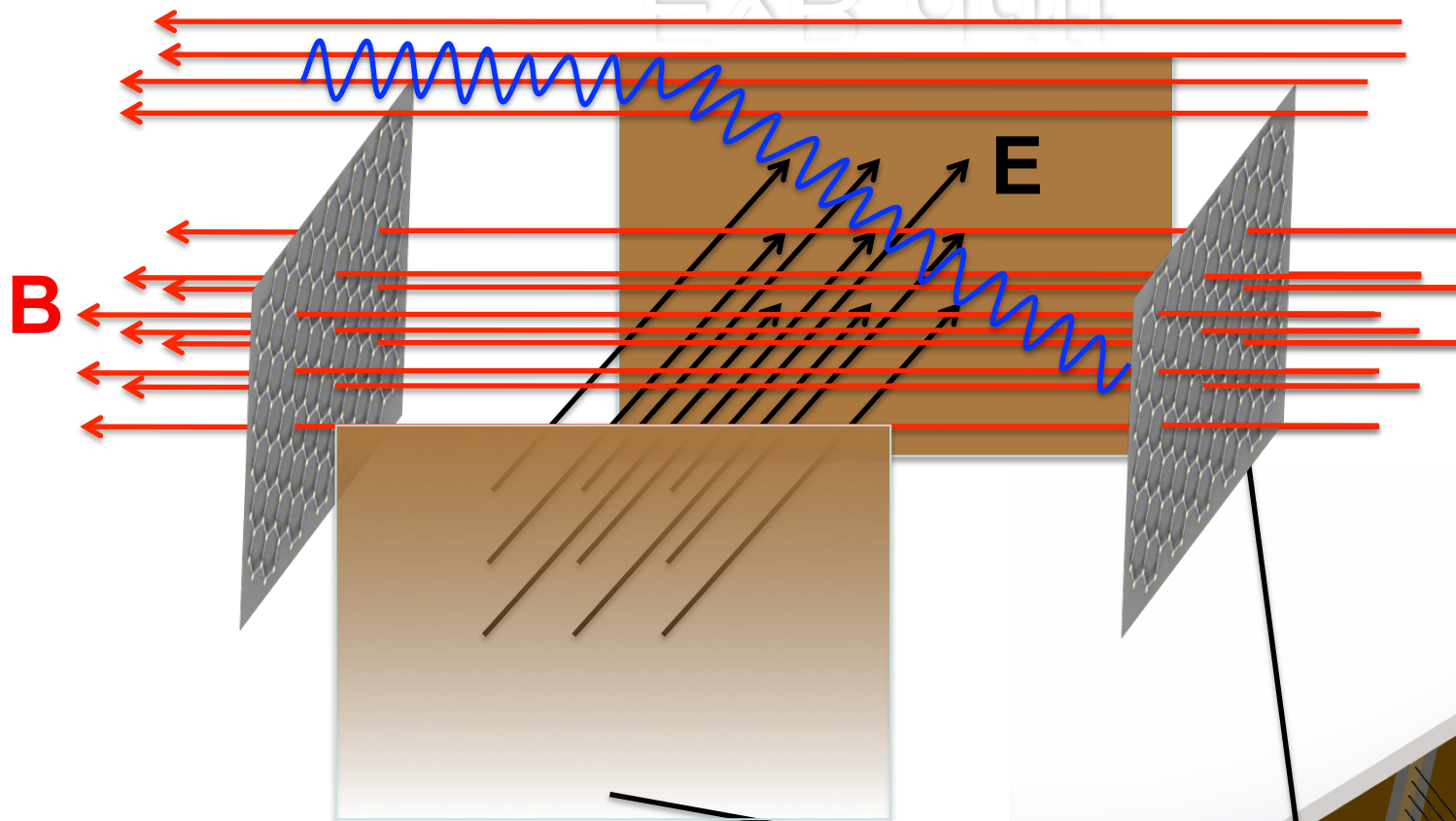


# THE Challenge

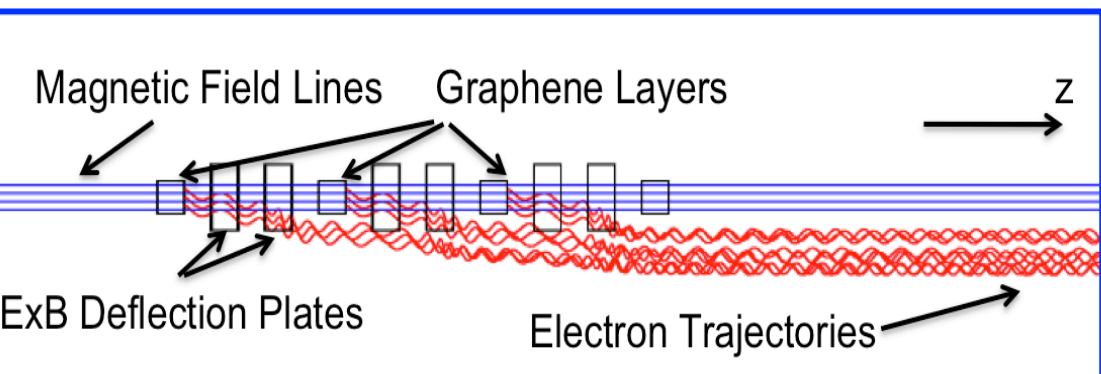


- The largest and nearly insurmountable problem of relic neutrino detection is to provide a large enough surface area to hold at least 100 grams of weakly bound atomic tritium
  - The trajectory of the outgoing electrons from tritium decay must have a clear vacuum path to the calorimeter (up to one or two atomic layers of carbon)
  - Need approximately  $10^6$  m<sup>2</sup> of expose surface area, that's ~200 football fields
  - Cannot be achieve with a flat planar surface – needs nanotechnology to solve

# ExB drift

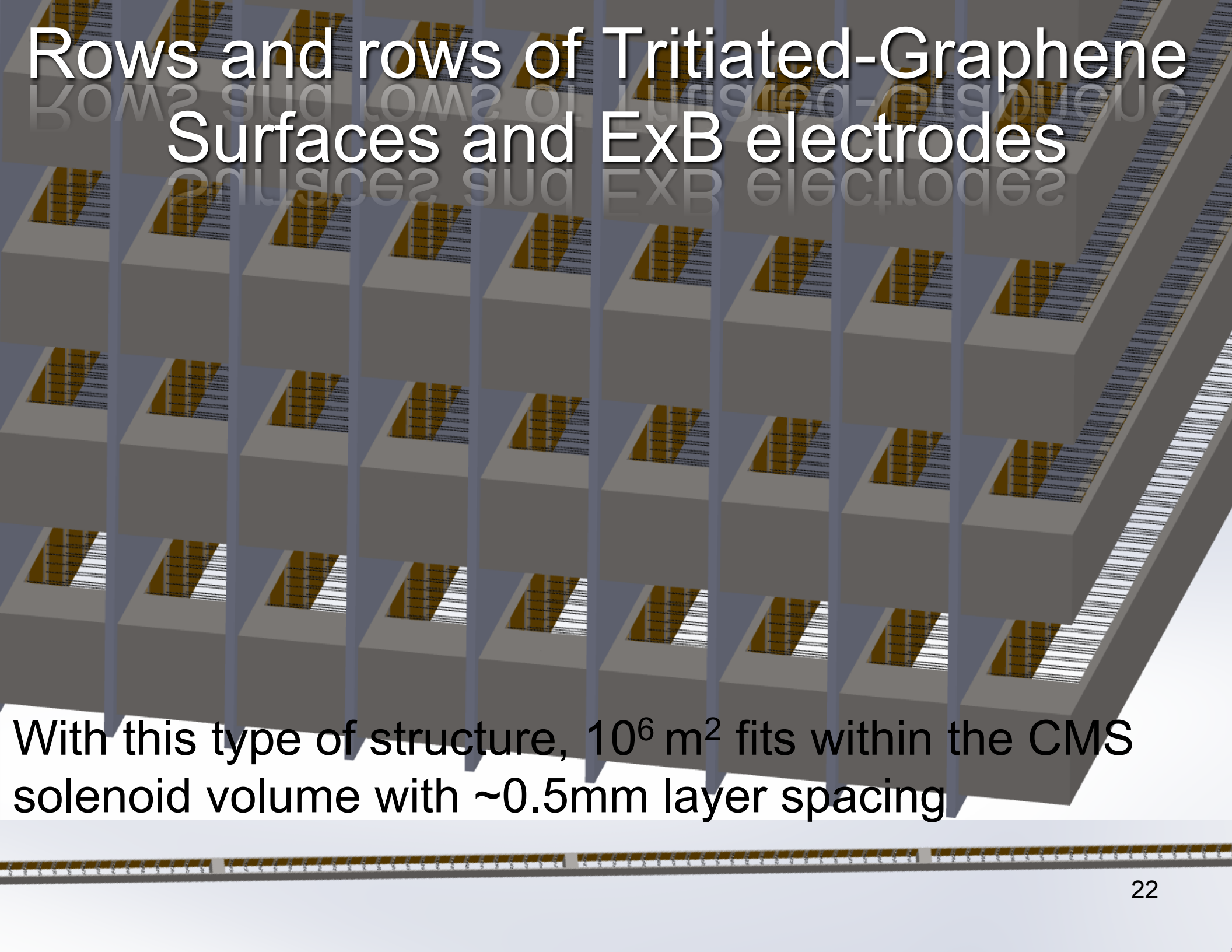


COMSOL calculation:





# Rows and rows of Tritiated-Graphene Surfaces and ExB electrodes

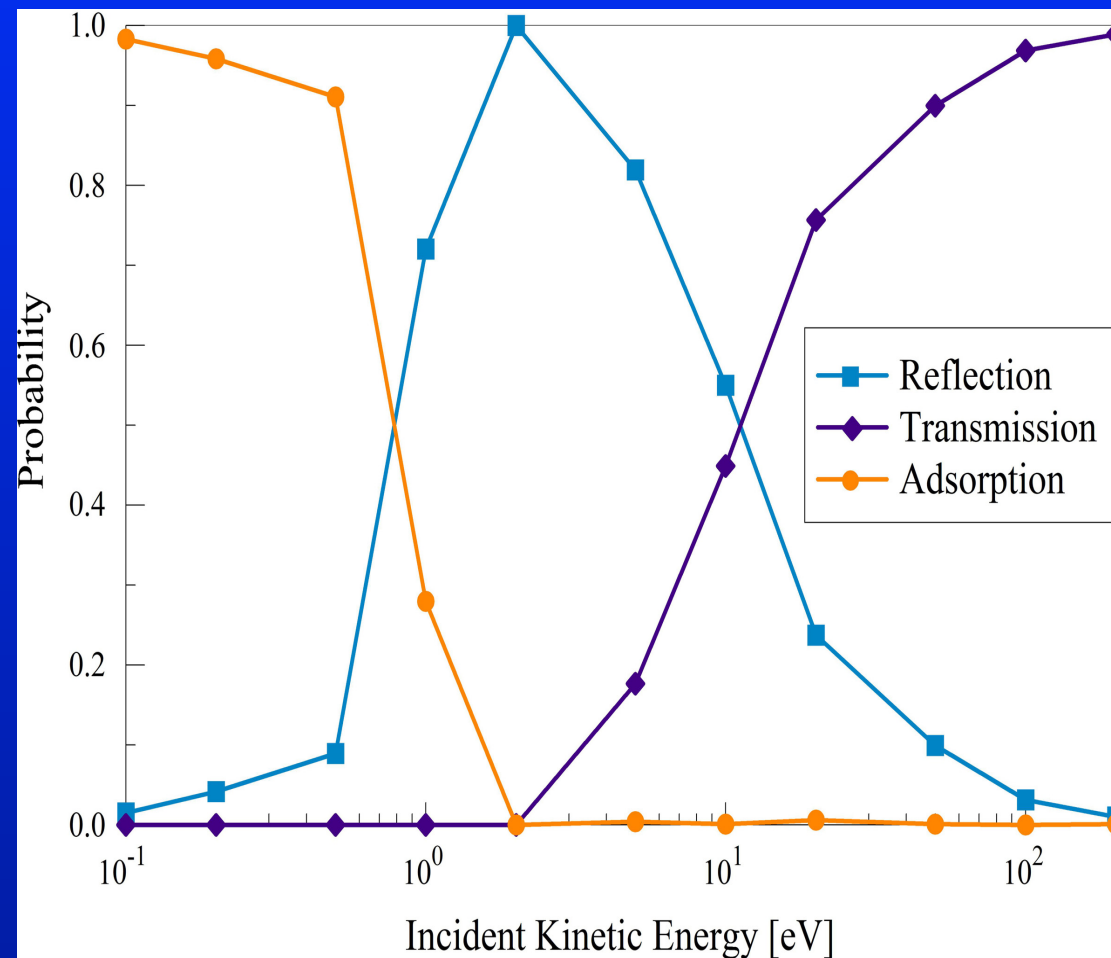
A 3D perspective view of a grid of tritiated-graphene surfaces and ExB electrodes. The structure consists of multiple rows of rectangular blocks, each containing a tritiated-graphene surface and an ExB electrode. The blocks are arranged in a regular grid pattern, with a small gap between them. The surfaces are shown in a perspective view, highlighting the depth of the structure.

With this type of structure,  $10^6 \text{ m}^2$  fits within the CMS solenoid volume with  $\sim 0.5 \text{ mm}$  layer spacing

# Tritium Loading

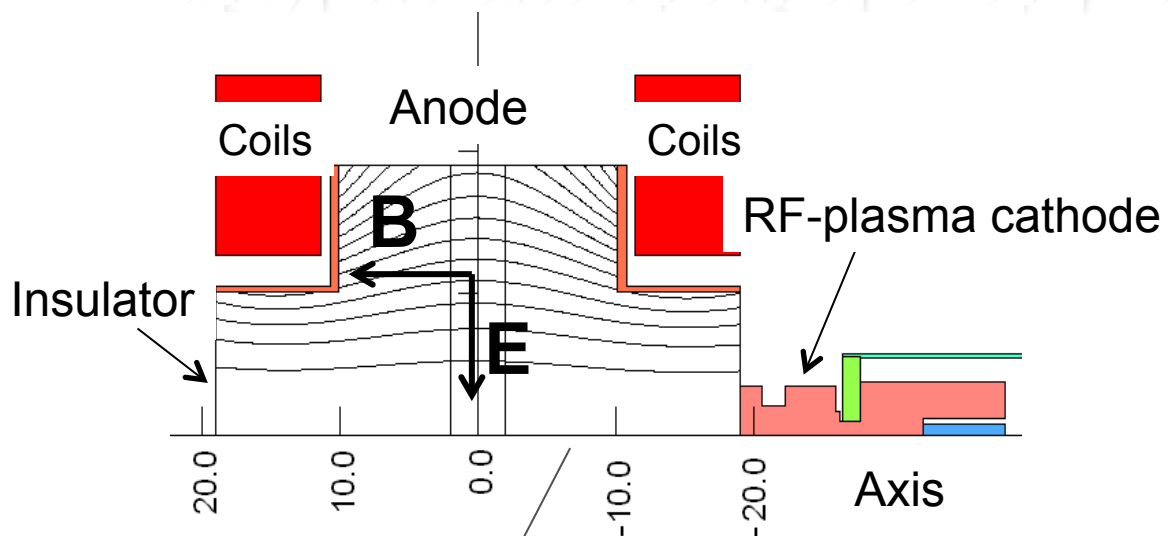


- The most common form of hydrogen loading is done at high pressure (~100atm) which is prohibitive for large surface areas.
- Ultra-low proton “beams” with  $T < 1\text{eV}$  bombarding the graphene surface have near unity probability to be absorbed onto the surface
- Above 2eV, the adsorption probability drops off rapidly

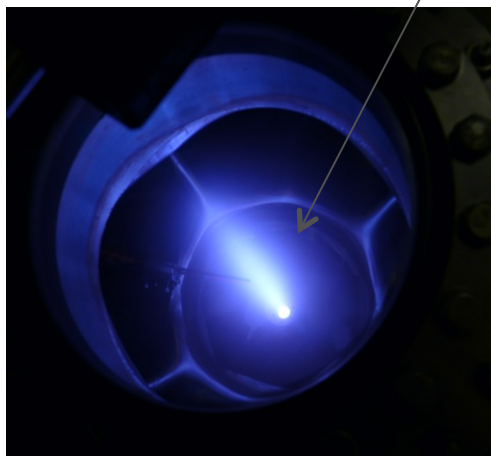




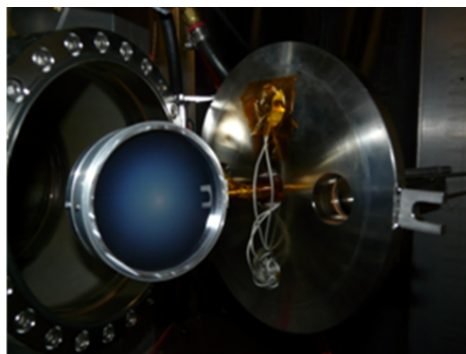
# Cold Magnetized Plasma for Processing of Nanomaterials at Low Pressure



## Source operation



## Si wafer immersed in the plasma source



DC  $E \times B$  fields applied in a 20 cm  $\times$  50 cm st. steel chamber with ceramic side walls.

**Plasma cathode:** 2 MHz, 50-200 W Ferromagnetic ICP

## Diagnostics:

Langmuir probes, emissive probes, optical emission and laser diagnostics of plasma

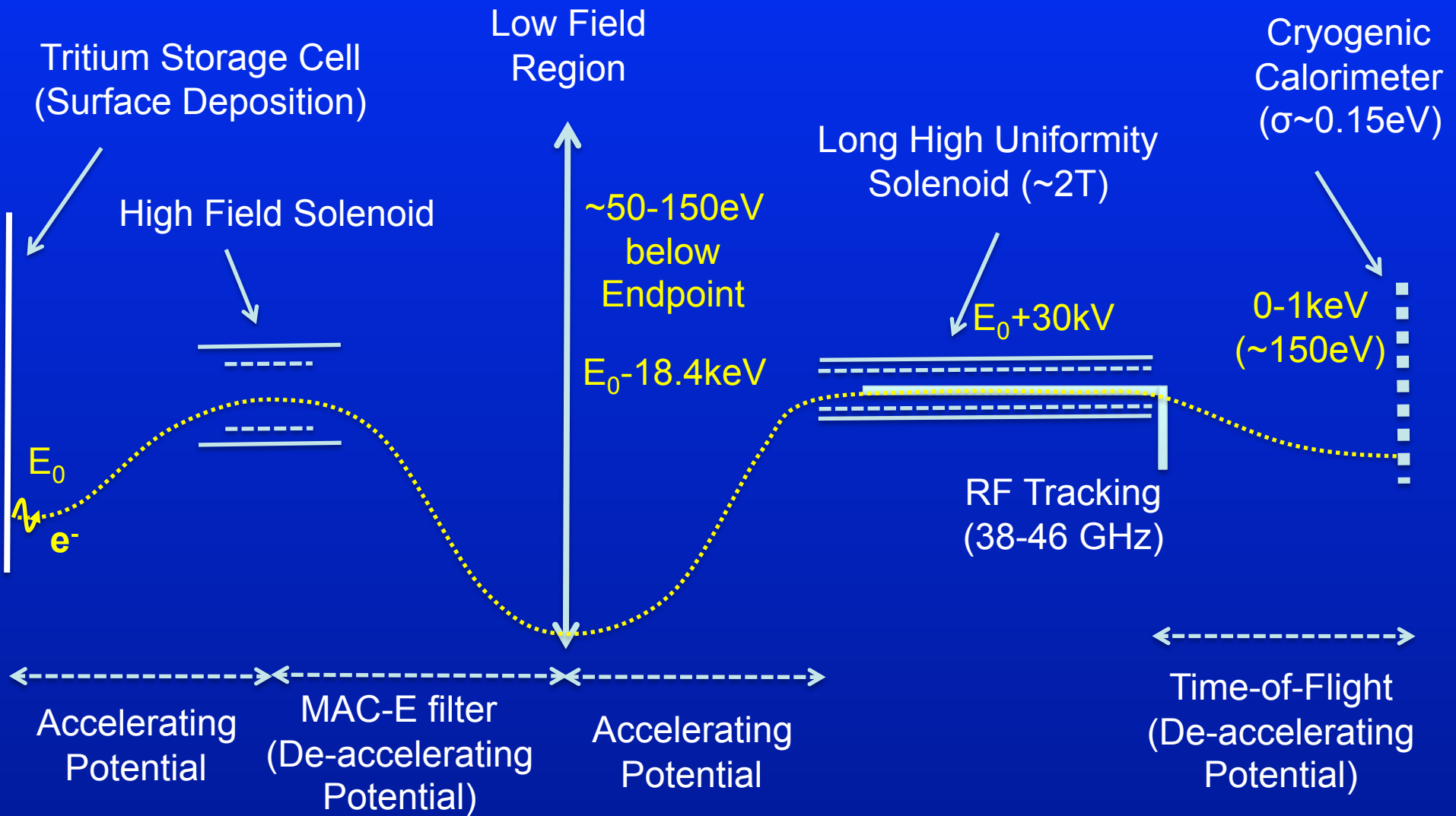
*Raitses et al., DOE PSC Meeting, 2013*

We will use cold magnetized plasma,  $T_e < 1$  eV, for the hydrogenation of graphene

# PTOLEMY Experimental Layout



Princeton Tritium Observatory for Light, Early-universe, Massive-neutrino Yield

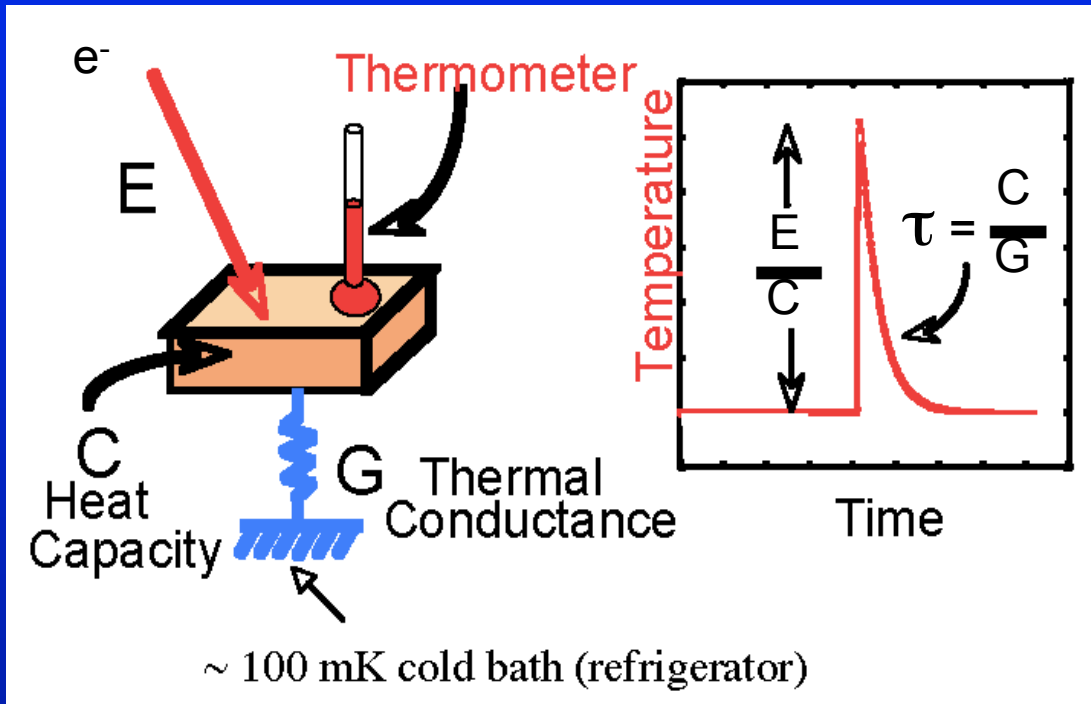




# Transition-Edge Sensors for Calorimetry

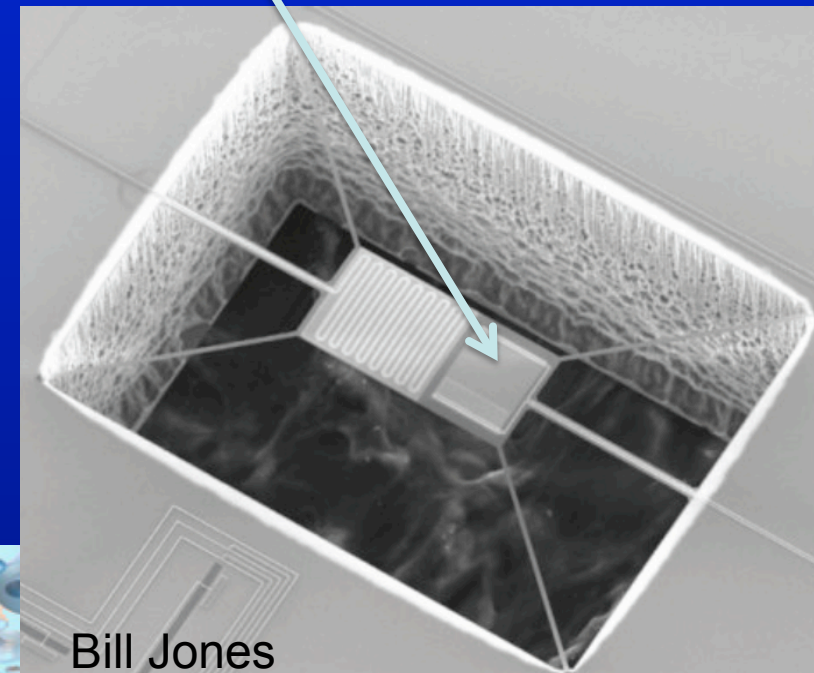


- ANL Group (Clarence Chang) estimates  $\sim 0.55\text{eV}$  at  $1\text{keV}$  and  $\sim 0.15\text{eV}$  at  $0.1\text{keV}$  operating at  $70\text{-}100\text{mK}$



100eV electron can be stopped with very small C

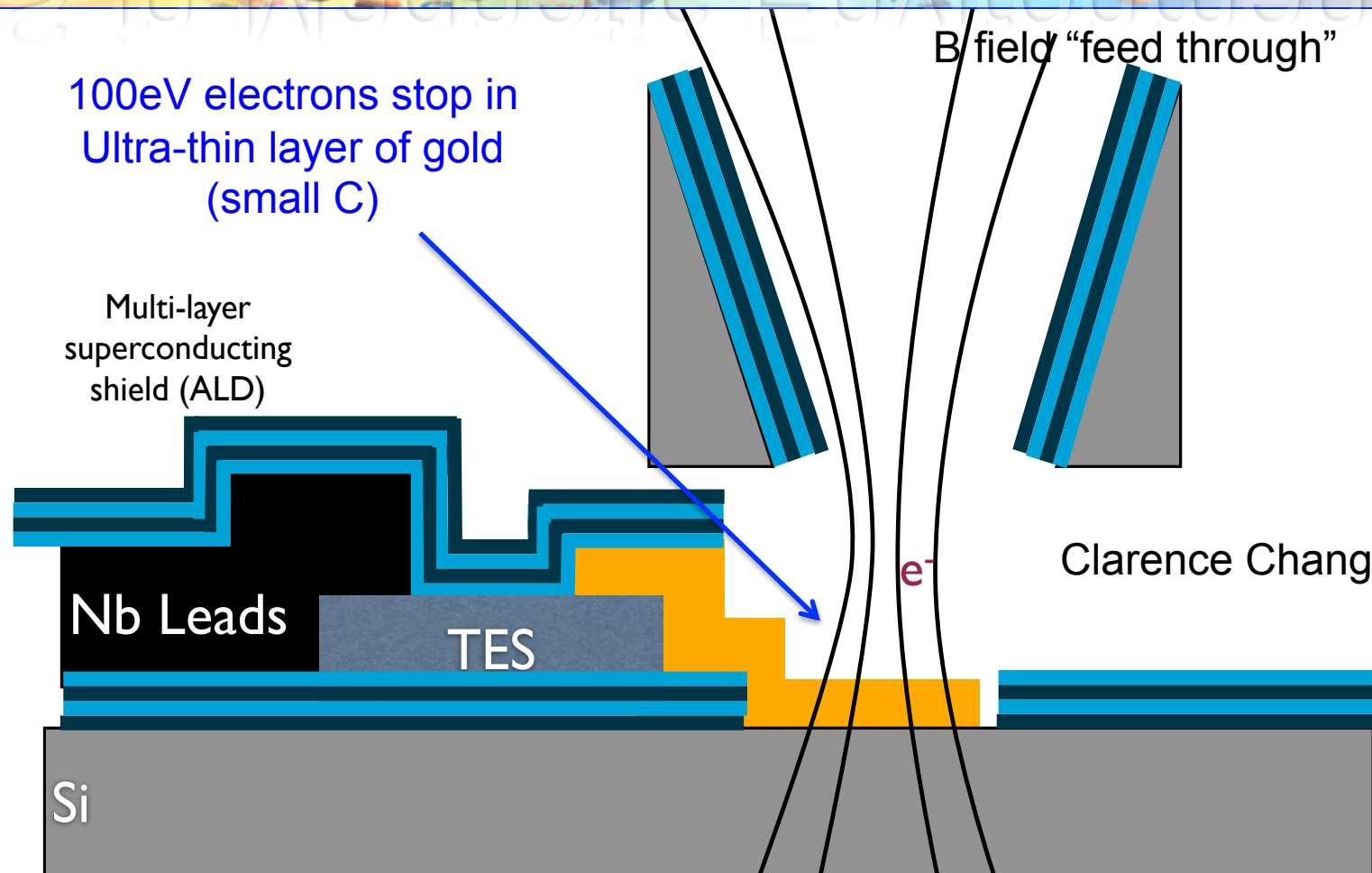
(example) SPIDER Island TES



Bandwidths of  $\sim 1\text{ MHz}$  to record  $\sim 10\text{ kHz}$  of electrons hitting the individual sensors

Bill Jones

# TES in Magnetic Environment



G set by e-ph coupling  $\propto T^5$   
 $\Delta E = 0.15 \text{ eV @ } 100 \text{ eV}$   
Operating at 70-100mK



# Calorimeter Energy Resolution



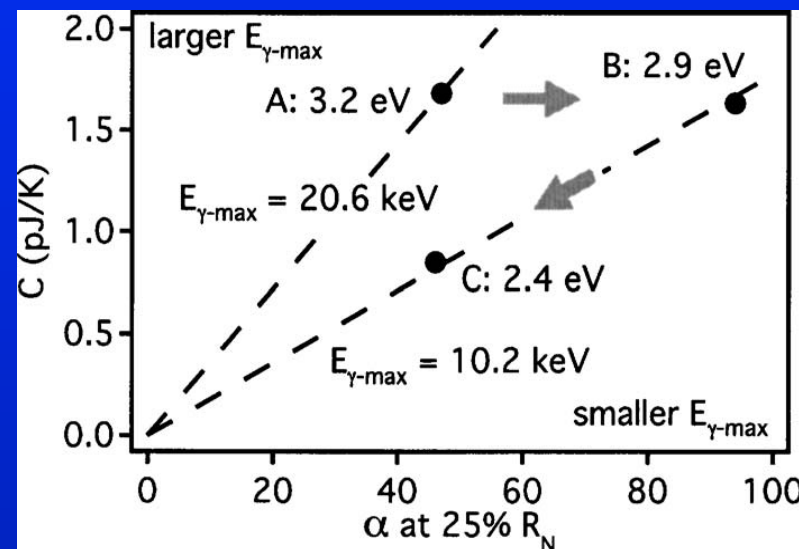
$$\Delta E_{\text{FWHM}} = 2.355 \sqrt{(4k_b T_c^2 C / \alpha) \sqrt{(1 + M^2) n / 2}}$$

Applied Physics Letters 87, 194103 (2005);  
doi: 10.1063/1.2061865

$(C/\alpha)$  scaled down by a factor of 100

Keep  $\alpha$  large, but not too large to keep M small

Electron energy at calorimeter:	Thickness of Gold Absorber:
600 eV	9.64 nm
400 eV	6.63 nm
200 eV	3.82 nm
100 eV	2.39 nm
10 eV	0.68 nm



$$\alpha \propto \frac{1}{\Delta T_{\text{width}}}$$

- Thickness of Gold absorber can be 5 nm (~40 atomic layers), corresponding to  $C_p$  of approximately 0.04 pJ/K per  $\text{mm}^2$
- Transition-edge steepness ( $1/\alpha$ ) controlled by normal regions and magnetic field.

Au is not ideal as it doesn't stick well for thin layers,  
Alternative materials would be studied (15nm of Bi)

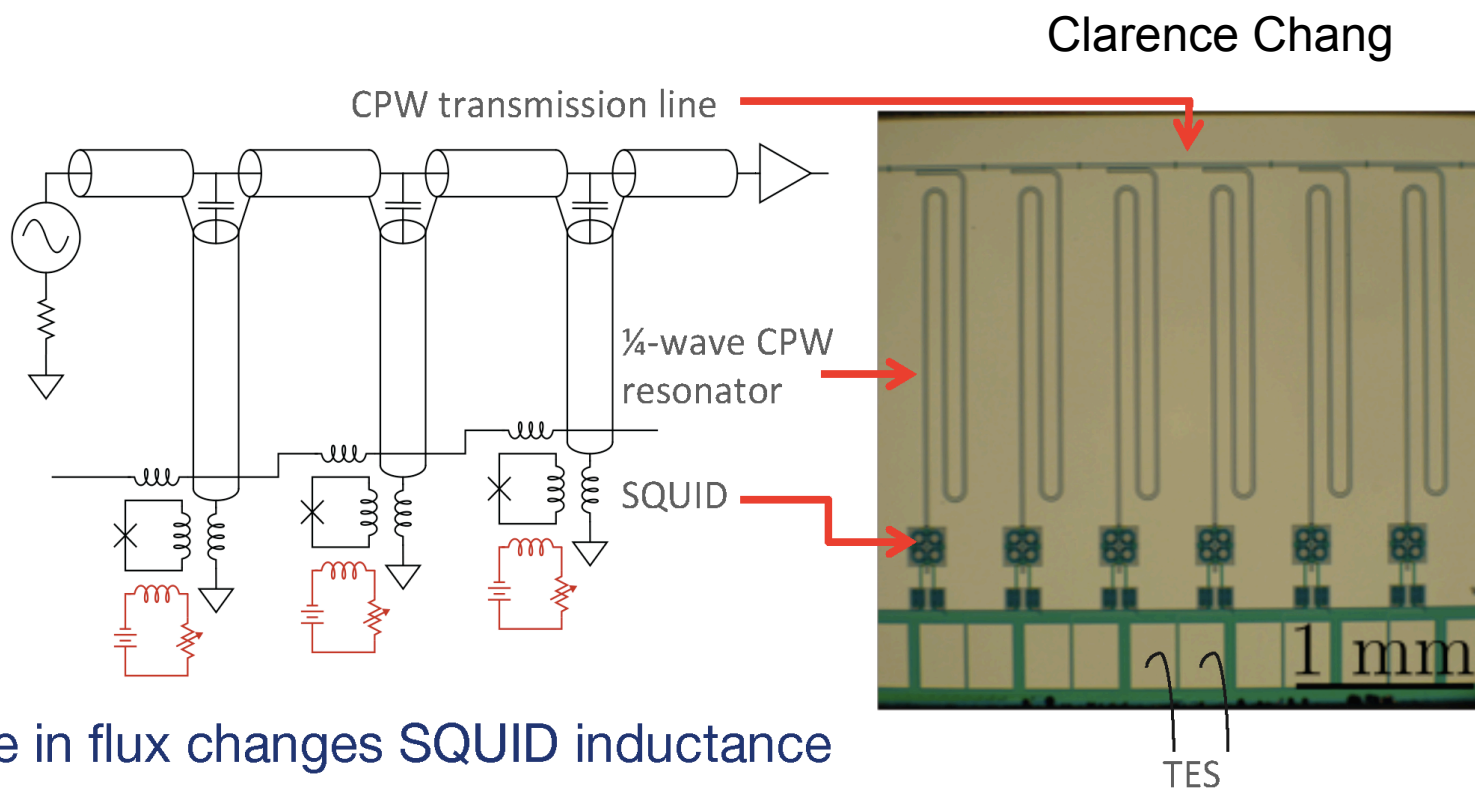
Clarence Chang

Backscatter from calorimeter can be efficiently collected by placing calorimeter surface at a +50V minimum (advantage of having only atomic electron backscatter)

# Highly Multiplexed SQUID Readout



## Microwave-readout Massive SQUID Multiplexer



- Change in flux changes SQUID inductance
- at 1-10 GHz, can support ~1 MHz of bandwidth with ~1000 channels per line
- Originally developed for CMB measurements, recently demonstrated successful operation with X-ray u-cals

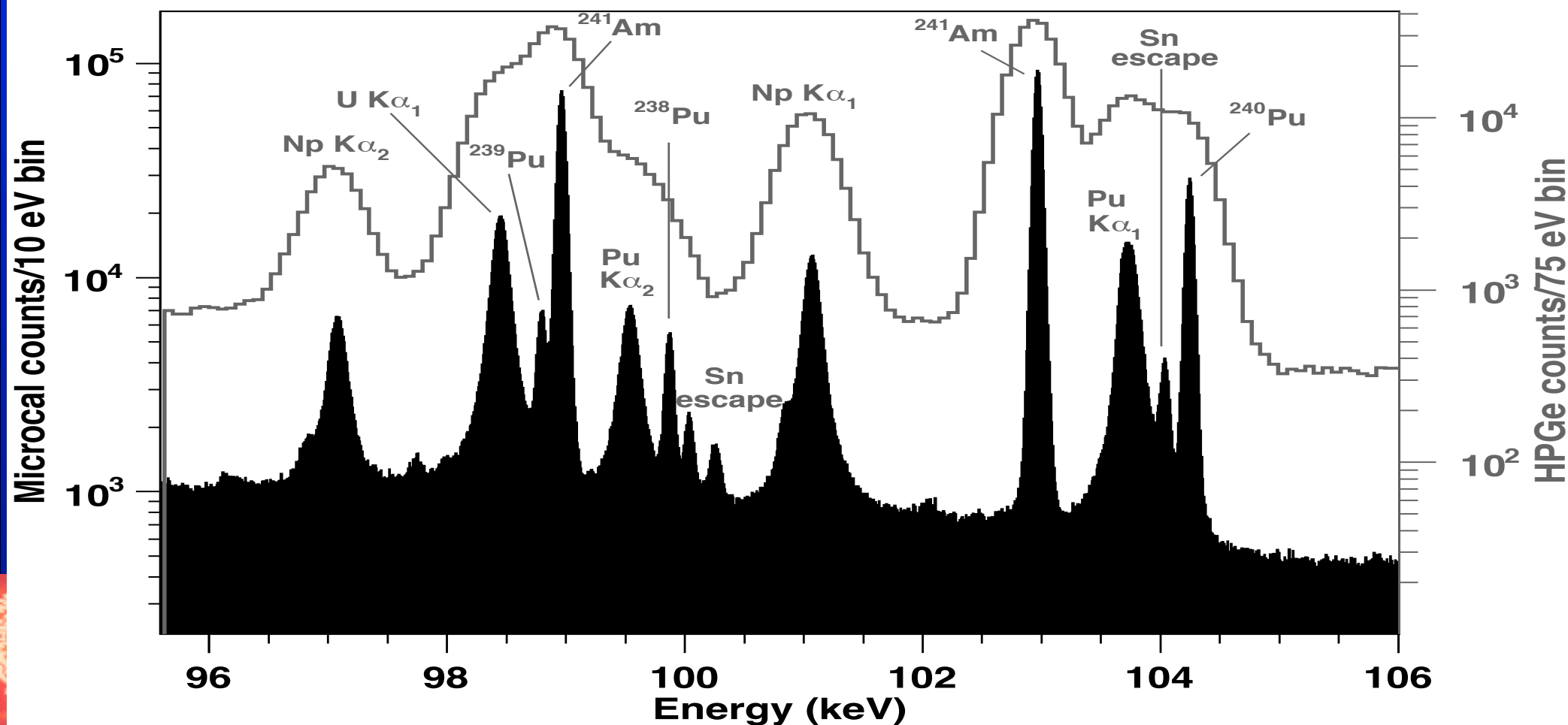
Kent Irwin



# Microcalorimeter Resolution



- Cryogenic microcalorimeters promise to greatly improve many research areas with vastly higher precision and data collection



# Calorimeter Rate



- Rate suppression is achieved with a MAC-E filter  
 $2 \times 10^{-10}$  suppression at 10eV, that scales as the endpoint distance cubed  $(E-E_0)^3$ 
  - At 0.1eV and zero neutrino mass, there are roughly  $10^5$  more background expected than signal – a finite neutrino mass is required to introduce a 2m gap between signal and background
  - At 0.001eV, the rate of signal and background are approximately equal
- 10eV endpoint window for 100g of tritium corresponds to  $(2 \times 10^{25})(2 \times 10^{-10}) / (4 \times 10^8 \text{s}) = 10 \text{ MHz}$ 
  - $10^3$ - $10^4$  calorimeter channels (for 10kHz bandwidth)



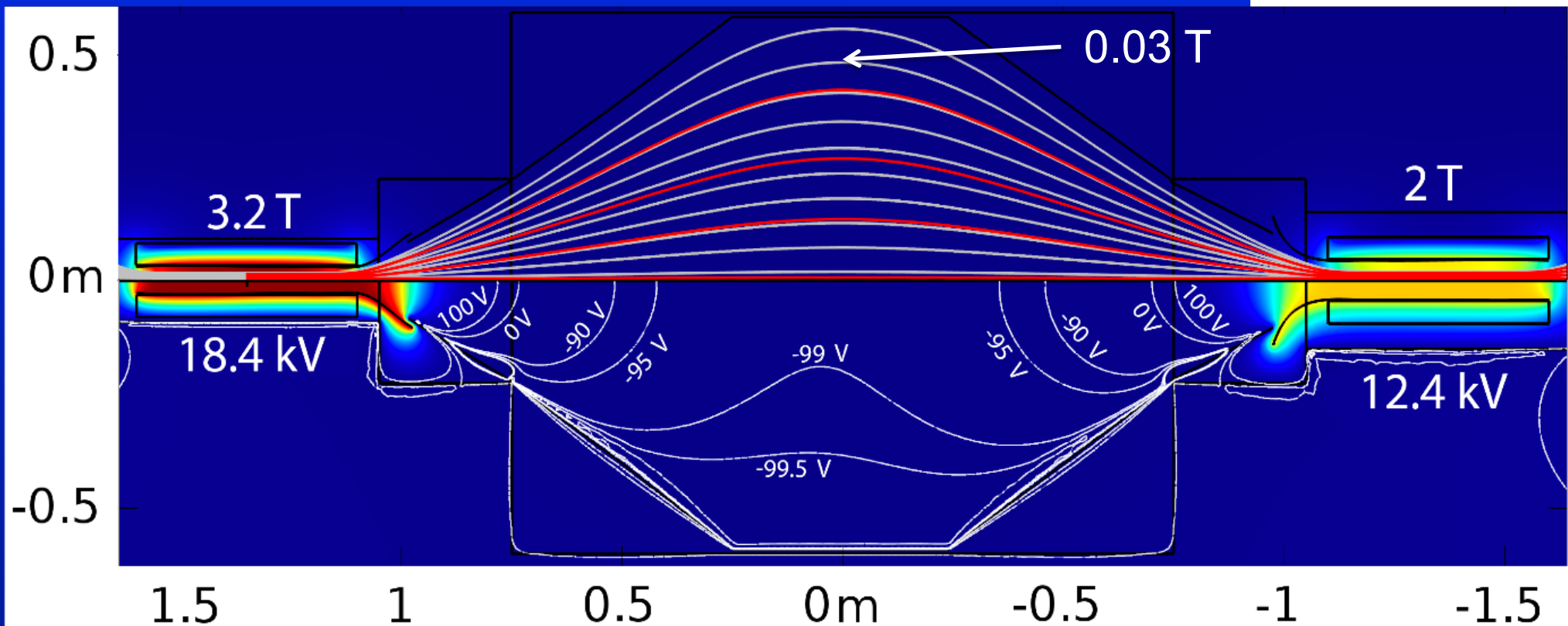
# PTOLEMY MAC-E filter



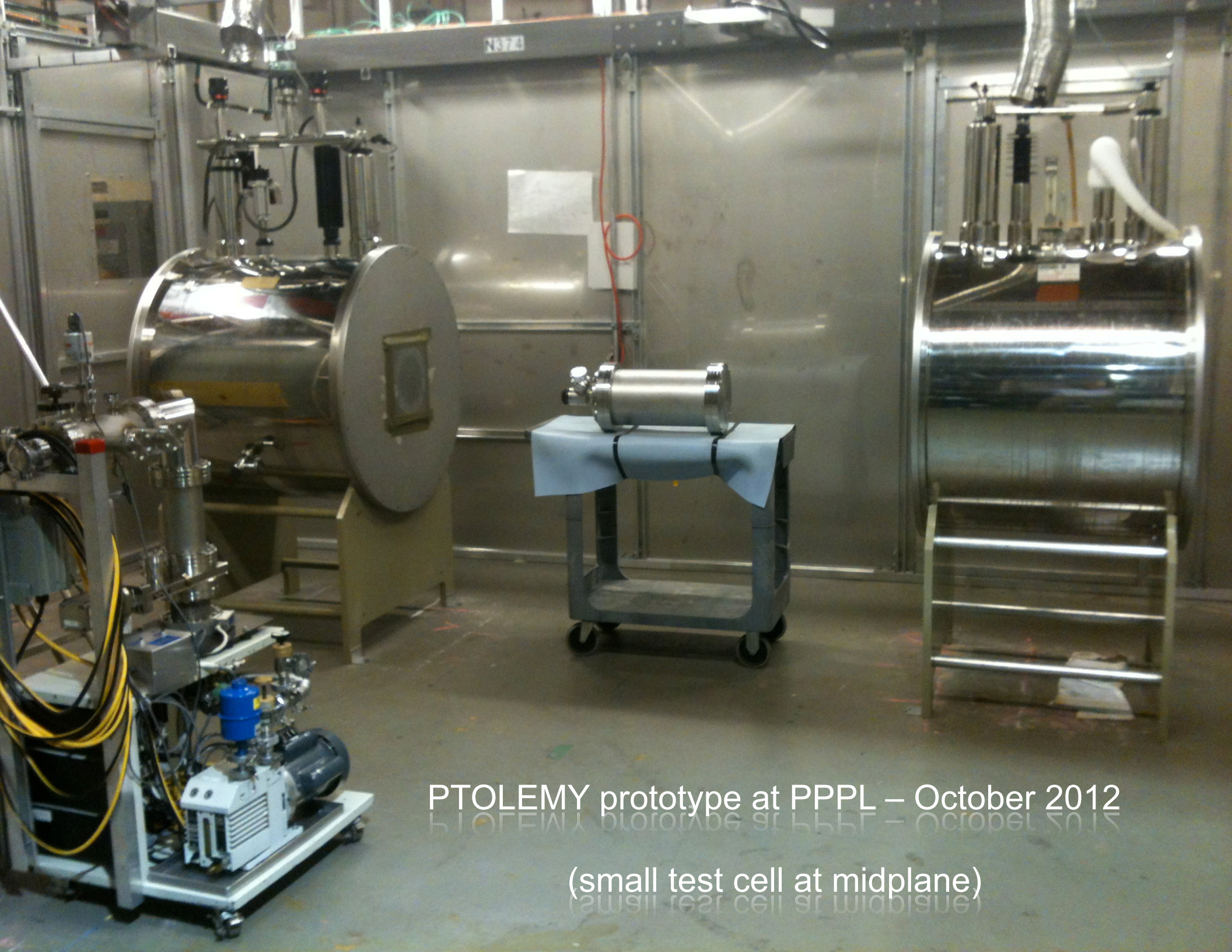
- MAC-E filter cutoff of  $10^{-2}$  to  $10^{-3}$  precision on electron energy

- Energy window below endpoint needed for  $2\pi$  acceptance
- Voltage of filter cut-off accurate to  $\sim 1\text{eV}$

Adiabatic Invariant:  $\mu = \frac{E^\perp}{B}$







PTOLEMY prototype at PPPL – October 2012

(small test cell at midplane)





**WARNING**  
**STRONG MAGNETIC**  
**FIELD AREA**  
**DO NOT BRING MAGNETIC OR**  
**FERROUS OBJECTS INTO THE LAB**

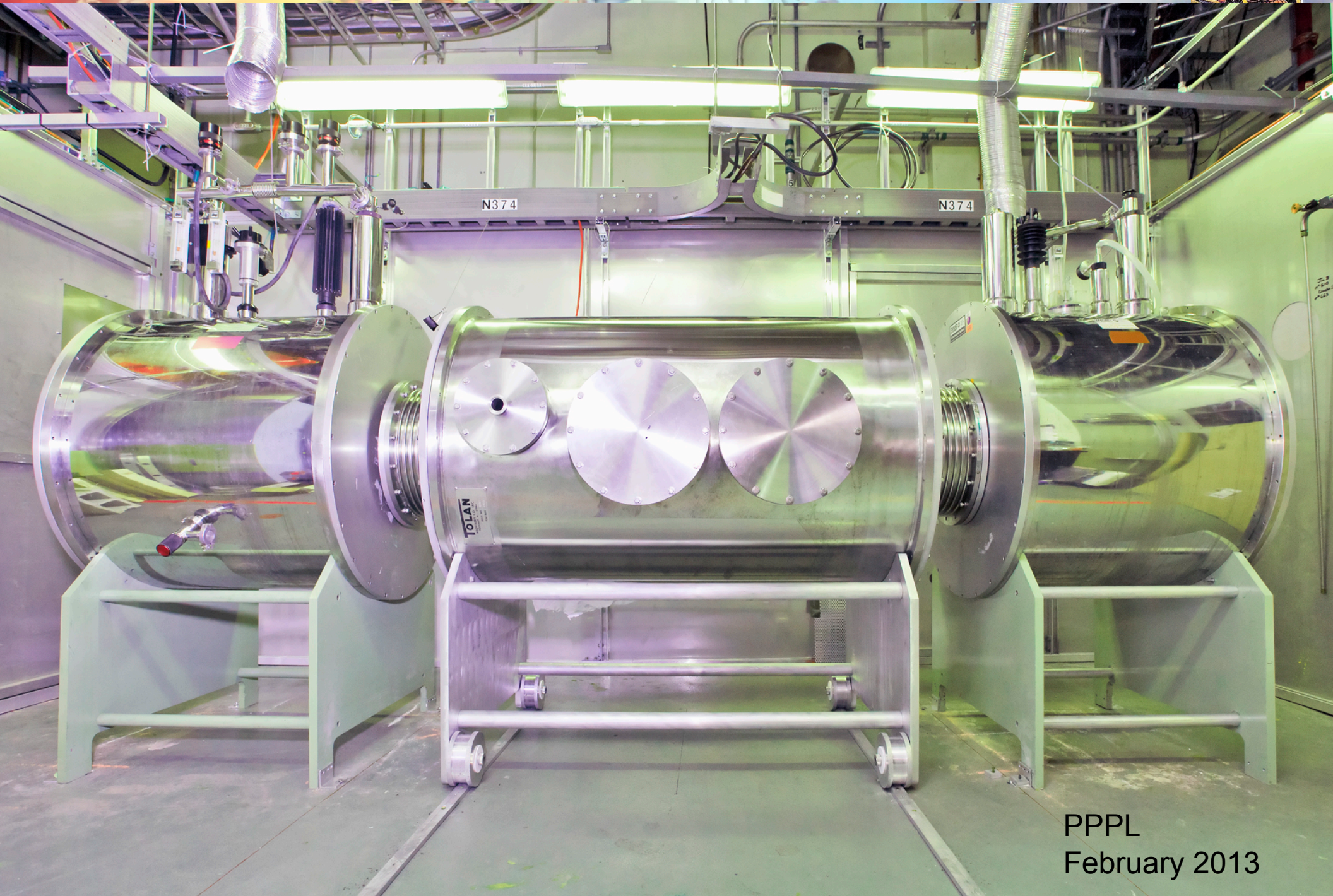
Place all ID badges, wallets, cell phones, keys, watches, and other magnetically susceptible objects in the containers provided on the table outside the PTOLEMY lab door.

All objects not placed in the bins and brought into the magnetic fields will fly to the magnets and be stuck within the bore.

PTOLEMY prototype at PPPL – January 2013  
(large vac-tank ready for install)



# PTOLEMY





# MAC-E Filter Size

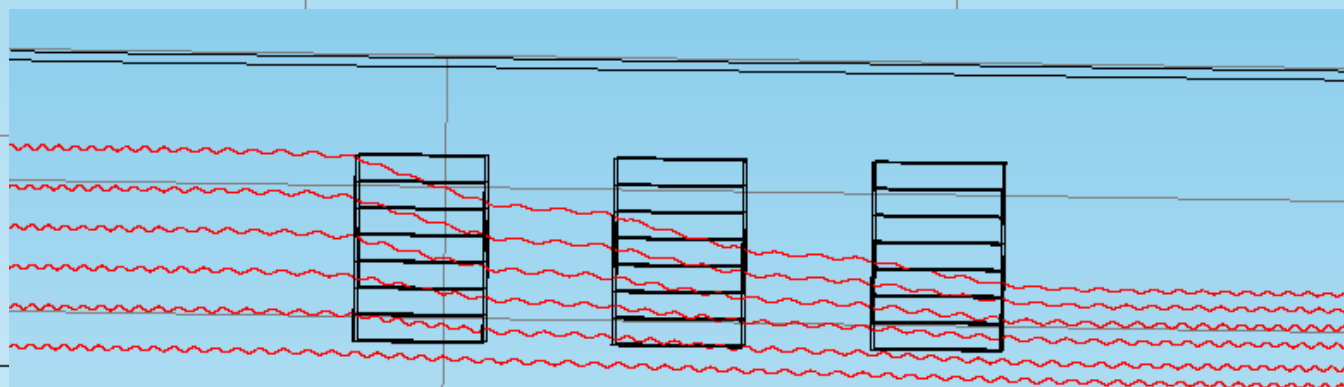


Limiting electron trajectory  
(hits outer radius of filter)

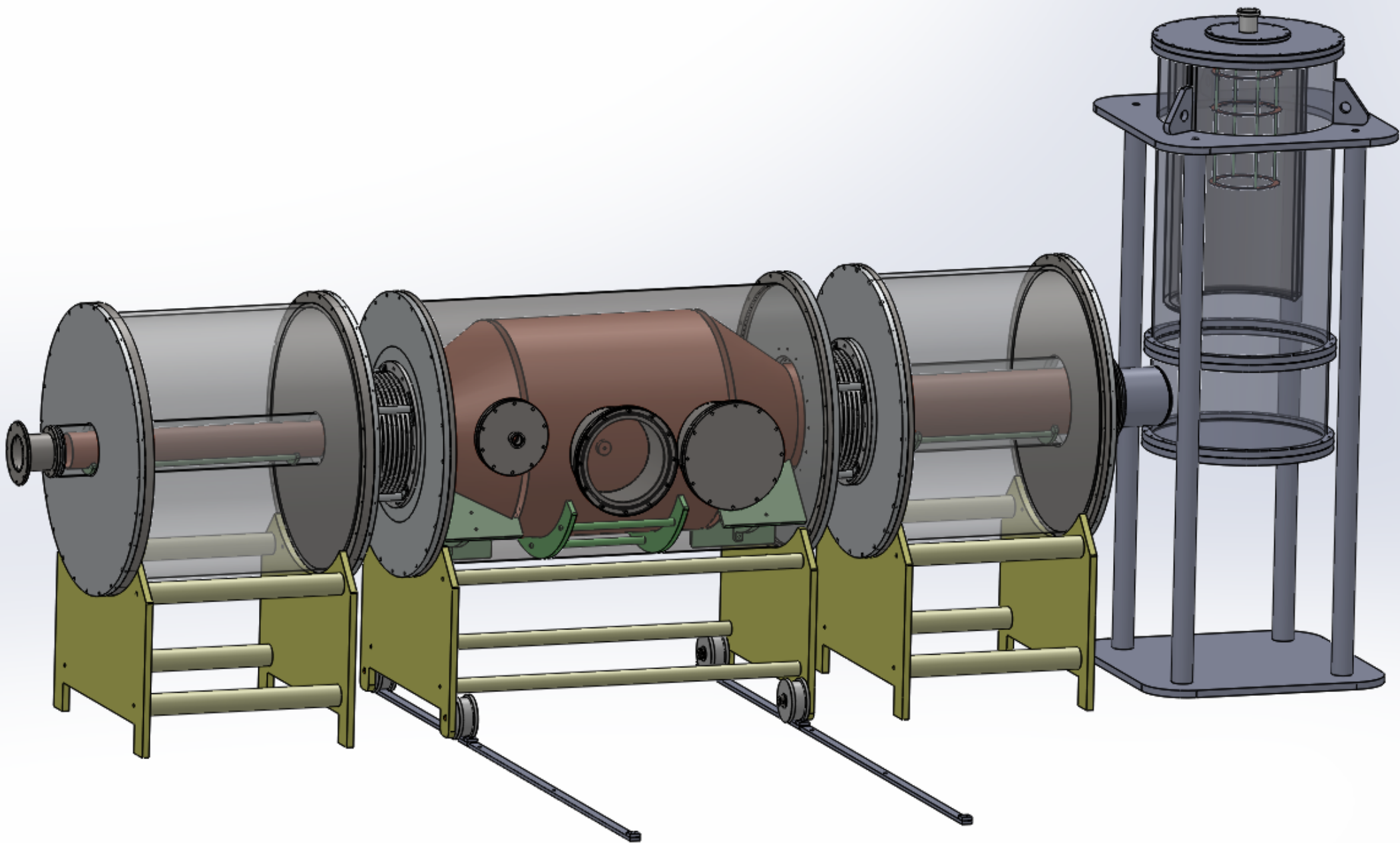
ExB focus before calorimeter  
(more compact calorimeter)

Calorimeter

ExB focus before  
entering MAC-E filter  
(same filter power with  
smaller central radius)



# Calorimeter Interface to MAC-E Filter

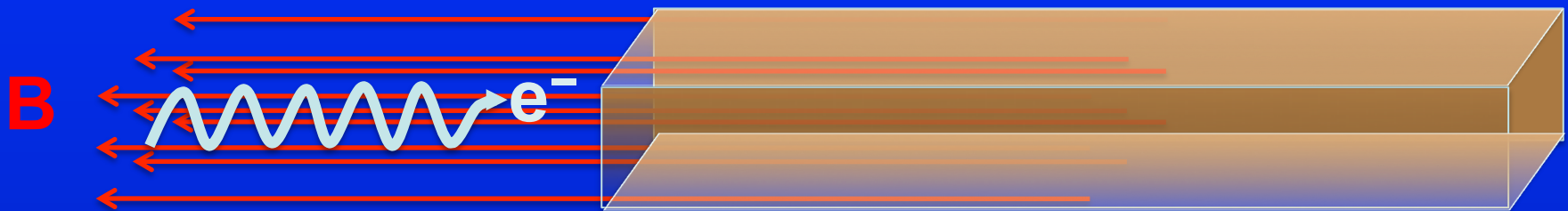




# Semi-relativistic Electron Identification



## Concept Developed by Project 8



- RF tracking and time-of-flight

- Thread electron trajectories (magnetic field guide lines) through a waveguide with  $\sim$ wide bandwidth (few  $\times 10^{-5}$ ) to identify cyclotron RF signal in transit times of order  $0.2\mu\text{sec}$ 
  - Currently using WMAP (Norm Jarosik) HEMT amplifiers with 1K/GHz noise and operating in the Q-Band range 38-46 GHz ( $\sim 1.9\text{T}$ )
  - Accelerate electrons to  $E_0 + 30\text{keV}$  in antenna region to increase electron cyclotron radiation – record in long uniform field (few  $\times 10^{-5}$ )

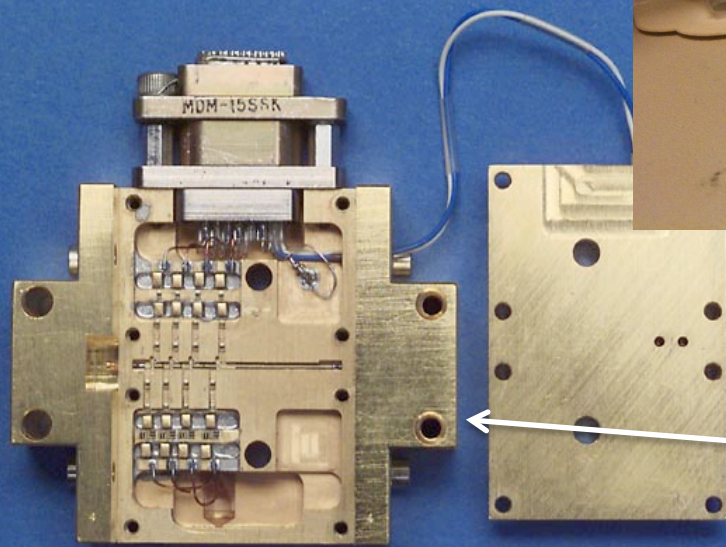
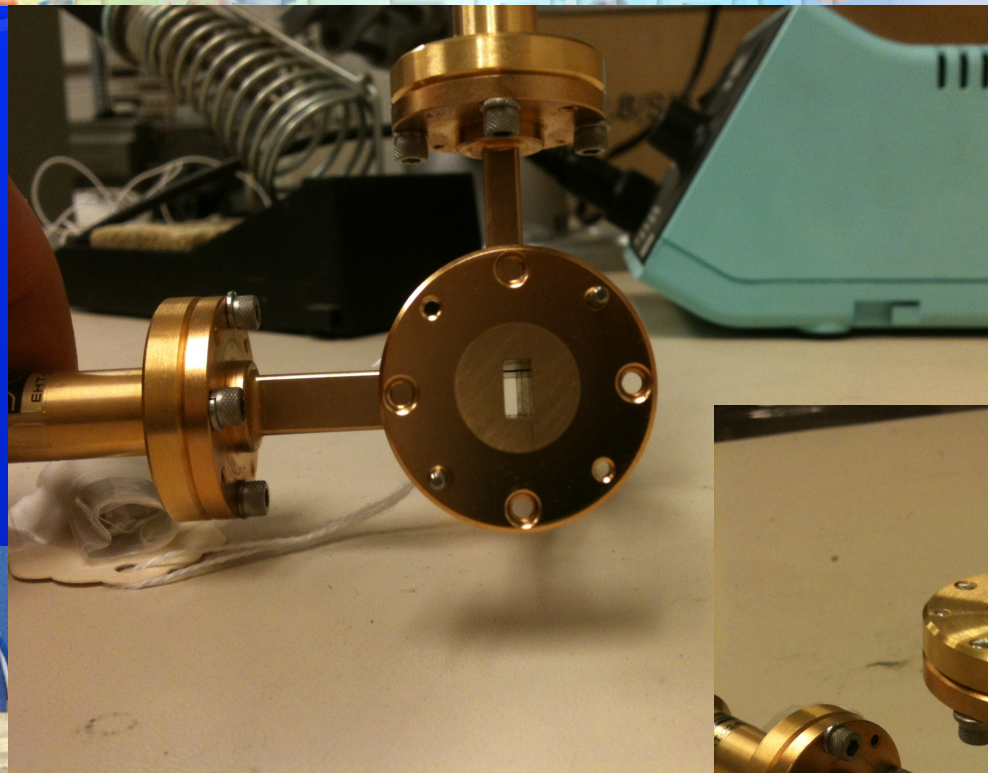


# RF Tracker Element



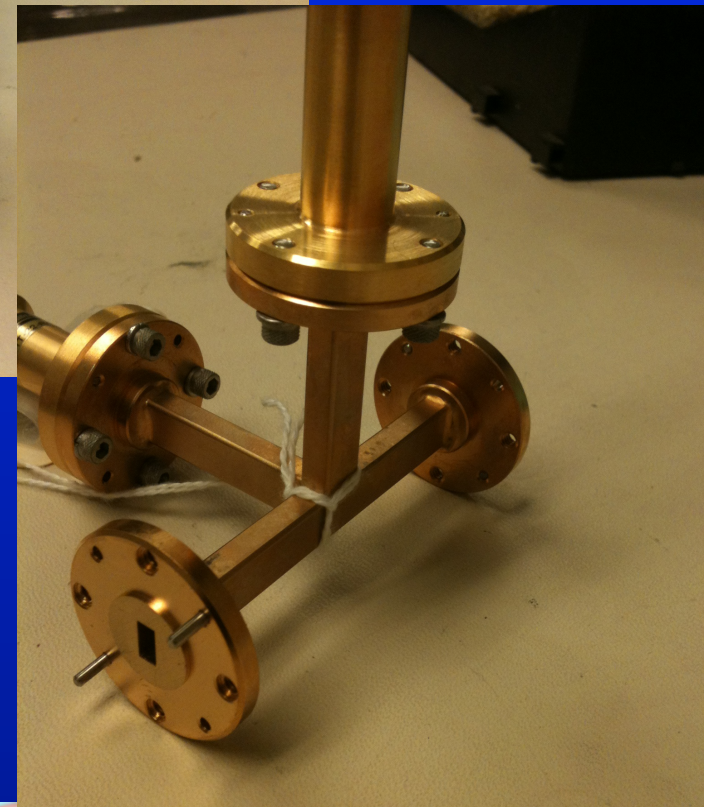
Readout Orthogonal to  
Electron Trajectory

Q-Band (38-46 GHz)  
Magic Tee Waveguide  
Junction



Q-11  
25 mm

Q-Band (38-46 GHz)  
WMAP Amplifier

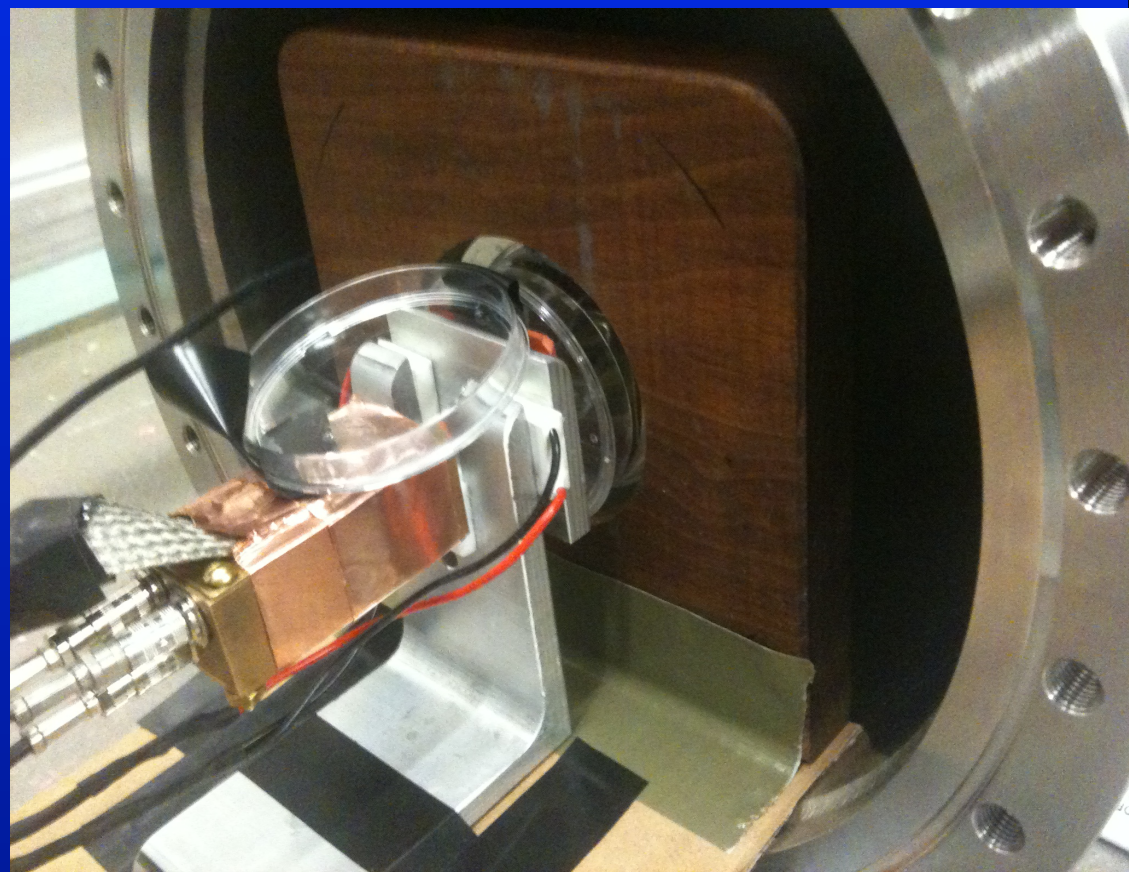
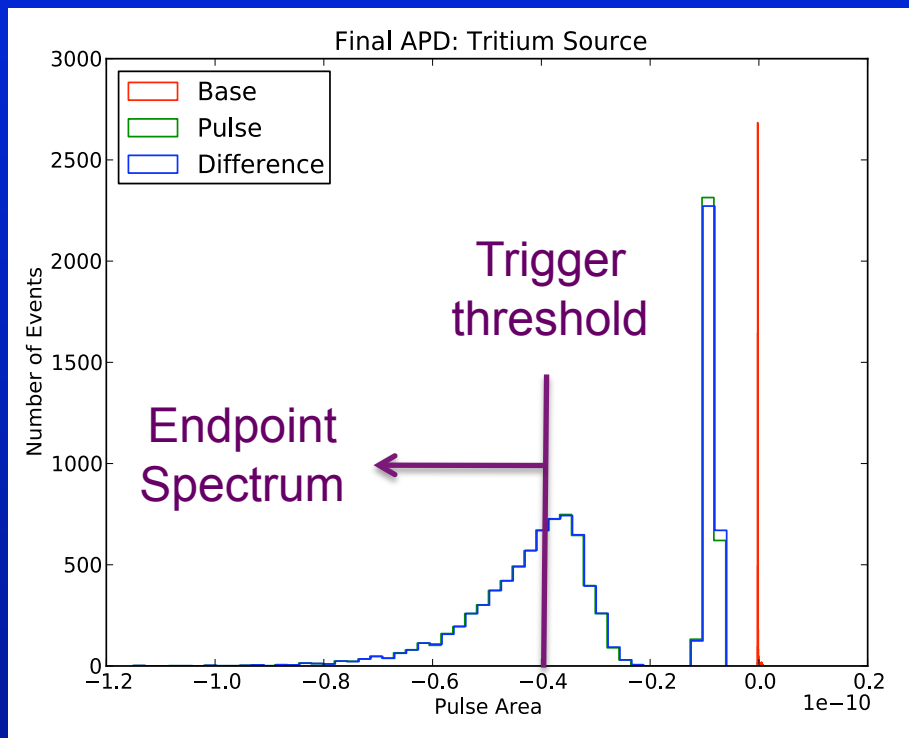




# Tritium Tag Detector



- For studying antenna data, a windowless APD is used to tag the tritium decay from a tritium disk source
  - Trigger on APD and record antenna (50 GHz mixed down to ~10 MHz bandwidth)





# What about Carbon-14?



- Take a biological sample from centuries ago exposed to atmospheric carbon

—————→ Now

With a half-life of ~5700 years  
levels of  $C^{14}$  of  $10^{-12}$  in  $2 \times 10^{25}$  nuclei  
will produce 100 Hz of decays

In a window of 0.5 eV ( $Q=156\text{keV}$ ), biological levels of  $C^{14}$  are four orders of magnitude too much radiation for a relic neutrino experiment with a graphene substrate. Fortunately, underground carbon sources have  $10^{-18}$  levels of  $C^{14}$  (achieved in Borexino).



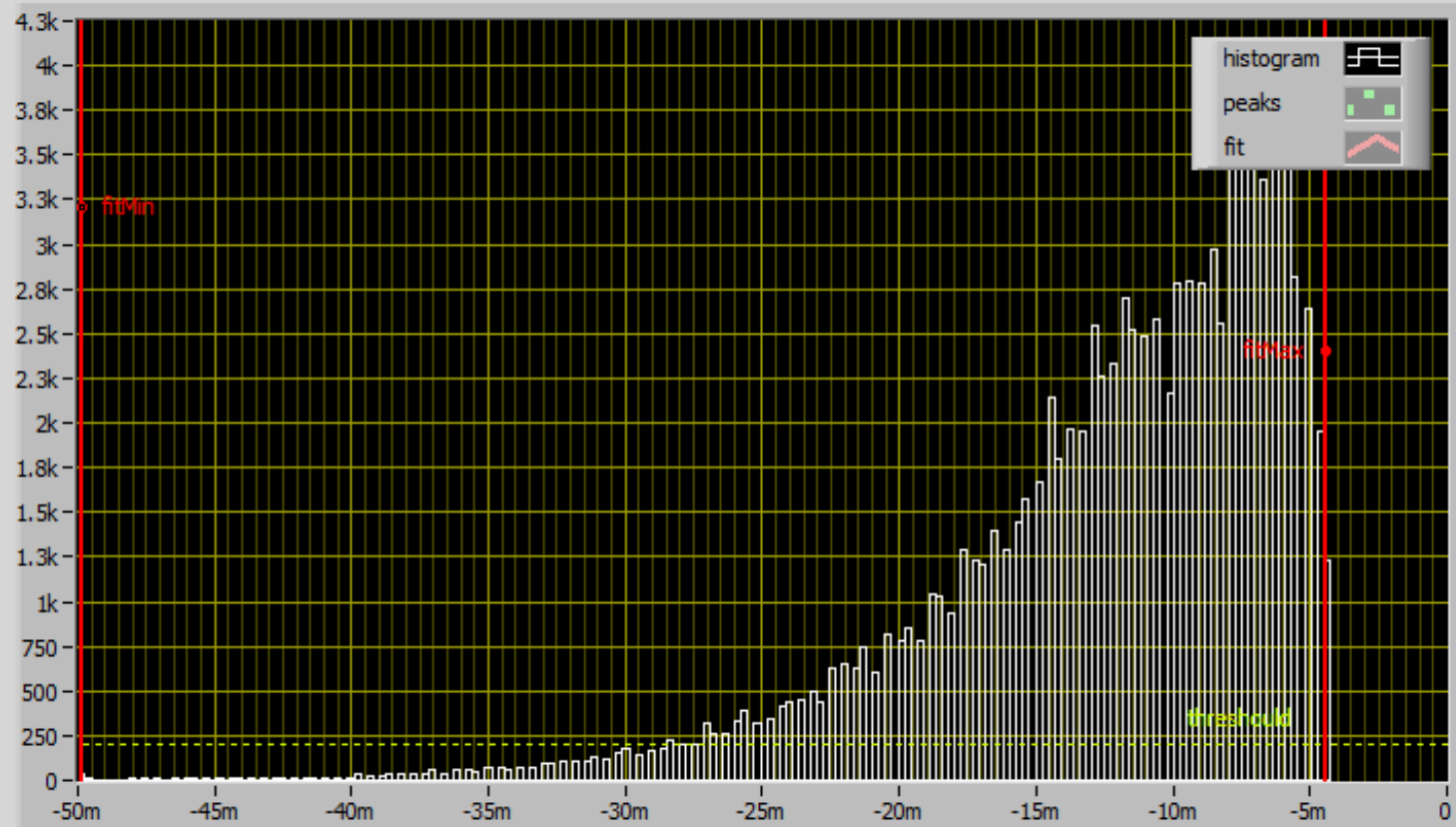
# Carbon-14 Spectrum



Read & Filtering    Waveforms & Cursors    **Histogram & Fitting**

**Print** **CLOSE**

Data source File 1



**Fill histo**    **Add** ?    **Sa**    **Lo**

amplitude (min)  Negate

200 bins    228u

100000 entries

timeDiffAmpl 0

Peaks threshold    pwidth

5 %    3

1st peak relaxed search

# 45    **to Excel**

Fit Window Full Scale

gaussian (LabView)

Least Square    **Run Fit**

Use stat weights

0

0	ampl
0E+0	center
0E+0	stdDev

Fit Points

Stat Points

Mean

Std. Dev



**simu**

expFitCoef

0    ampl

0    1/tau

0    offset

gaussInitials

0

ampl

center

stdDev

**Cursors:**

	X	Y
fitMin	-49.9m	3.2k
fitMax	-4.44m	2.4k
threshold	-4.44m	200

value

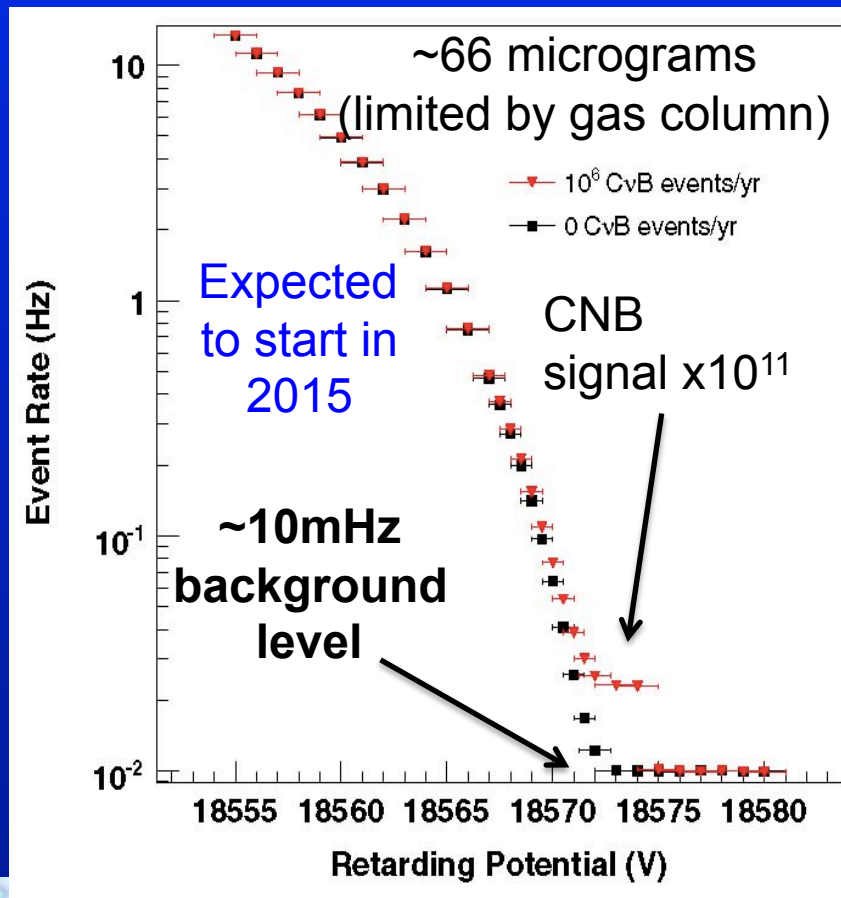
entries

numOfPeaks 2

# Karlsruhe TRItium Neutrino (KATRIN)



- Uses large uniform geometry to achieve  $\sim 0.2\text{eV}$  cut-off sensitivity – “Cut and Count” experiment
  - PTOLEMY Goal:  $10\text{mHz} \rightarrow \text{sub-}\mu\text{Hz}$  Background





# Annual Modulation of Cosmic Relic Neutrinos



B. Safdi, M. Lisanti, et al.  
<http://arxiv.org/pdf/1404.0680.pdf>

If CMB rest frame = relic neutrino rest frame,  
 the direction and velocity of the Sun is known relative  
 to the rest frame ( $\langle v \rangle \approx 370$  km/s)  
 represented by the unbound neutrino wind.

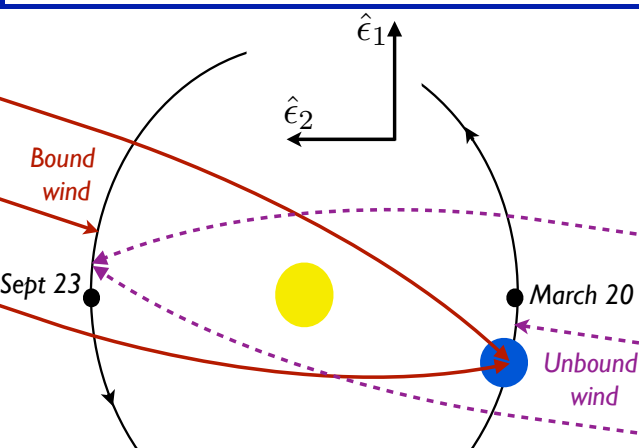
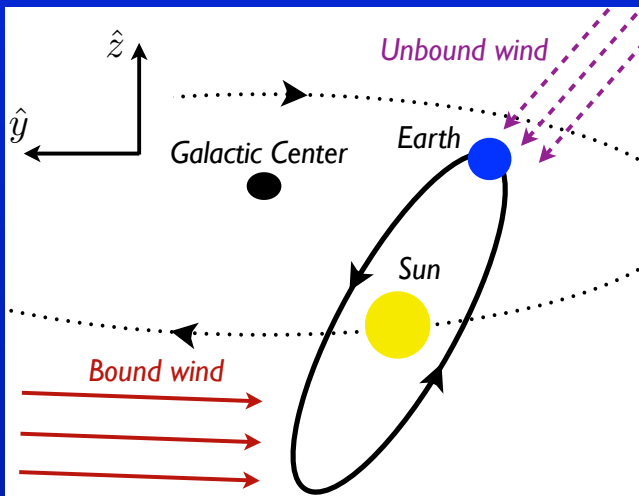
Sensitivity to relic neutrino  
 velocity and direction  
 through  
 annual modulation amplitude  
 (0.1-1%) and phase

Additional sensitivity to neutrino velocity  
 and direction for a polarized tritium

$$d\sigma_{\text{NCB}} \propto \left[ 1 + a \frac{\mathbf{p}_e \cdot \mathbf{p}_\nu}{E_e E_\nu} + A \frac{\hat{\mathbf{s}} \cdot \mathbf{p}_e}{E_e} + B \frac{\hat{\mathbf{s}} \cdot \mathbf{p}_\nu}{E_\nu} \right]$$

Tritium polarization calibrated using beta decay electrons.

Velocity sensitivity provides possibility to measure:  
 Relic Neutrino Rest Frame, and potentially,  
 Relic Neutrino Temperature (from velocity and mass)



# Neutrino Calibration Source



(quoted from work of) **V. Kornoukhov**

$^{51}\text{Cr}$  decays to  $^{51}\text{V}$  (Vanadium) through electron capture



Neutrino lines are 746(81%), 751(9%), 426(9%), 431(1%) keV.

Direct neutrino detection at short distances considered for neutrino oscillation experiments into sterile neutrinos.

Capture rate calculation (B. Safdi)

$\sigma(^{51}\text{Cr}) / \sigma(\text{relic neutrinos}) \sim 19$  (weighted)

$N_{\text{events}} \sim 40 \text{ events} / 100\text{g-years} (1 \text{ m} / R)^2 (M_{\text{source}} / \text{MCi})$

Relativistic neutrinos  $\rightarrow$  clear angular dependence from tritium polarization

Neutrino scattering off of atomic electrons also a signal for the neutrino flux

Also considering Titanium thin films embedded with tritium to produce a high mass (kg), low resolution neutrino capture target for solar neutrinos and for a relic sterile neutrino search.



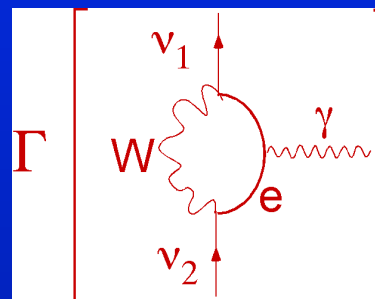
# Relic Sterile Neutrinos



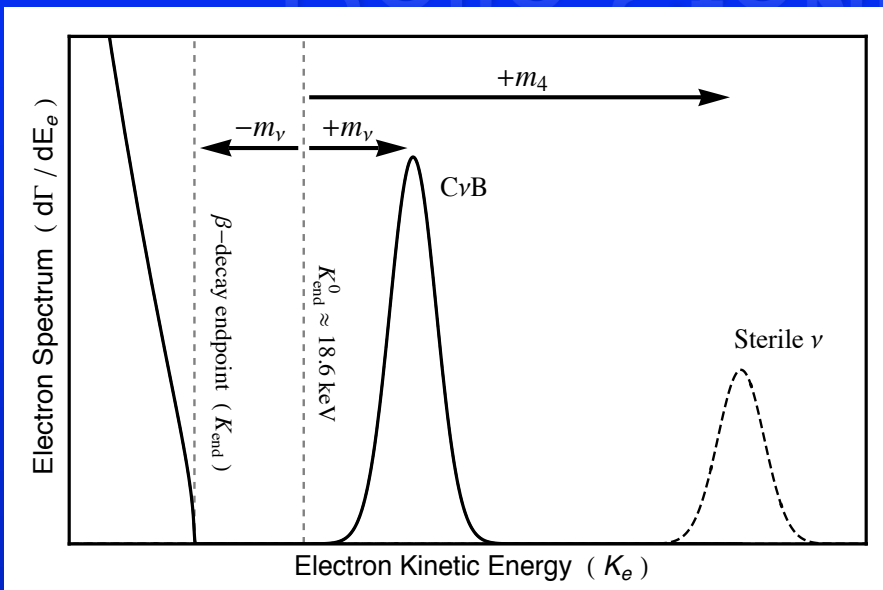
Long, Lunardini, Sabancilar: arXiv:1405.7654

Either a Majorana or a Dirac neutrino can decay radiatively:

$$\nu_2 \rightarrow \nu_1 + \gamma$$



$$\Gamma \propto G_F^2 [m(\nu_2)]^5$$



For  $m(\nu_2) \sim 0.5$  eV,

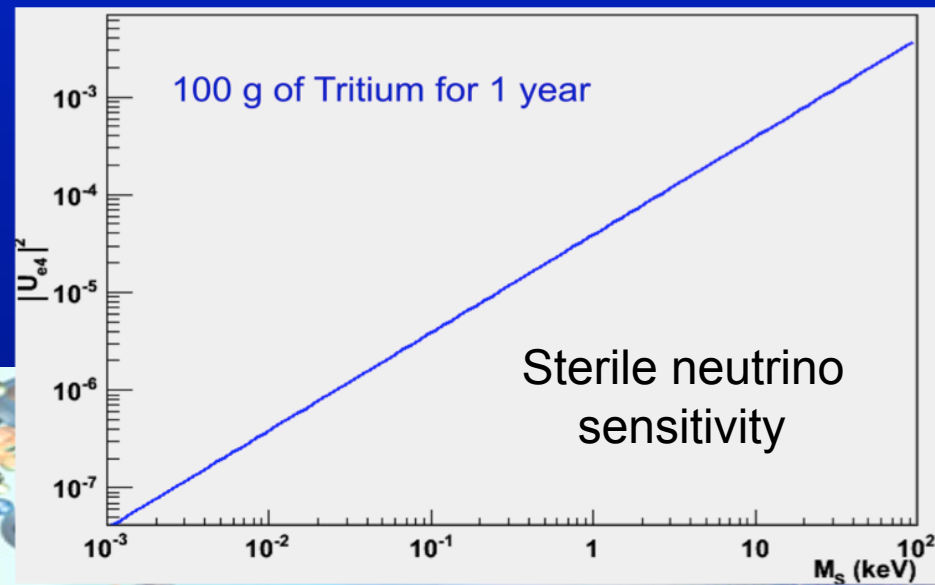
$$|U_{e4}|^2 \sim 0.03, f_c \sim 20, \Delta N_{\text{eff}} = 1$$

$$N_{\text{capture}} \sim 5 \text{ events/100g-year}$$

For  $m(\nu_2) \sim 1$  keV,

$$|U_{e4}|^2 \sim 10^{-4}, f_c \sim 5000 (\rho_{\text{DM}} \sim 0.3 \text{ GeV/cm}^3),$$

$$N_{\text{capture}} \sim 4 \text{ events/100g-year (10kg?)}$$



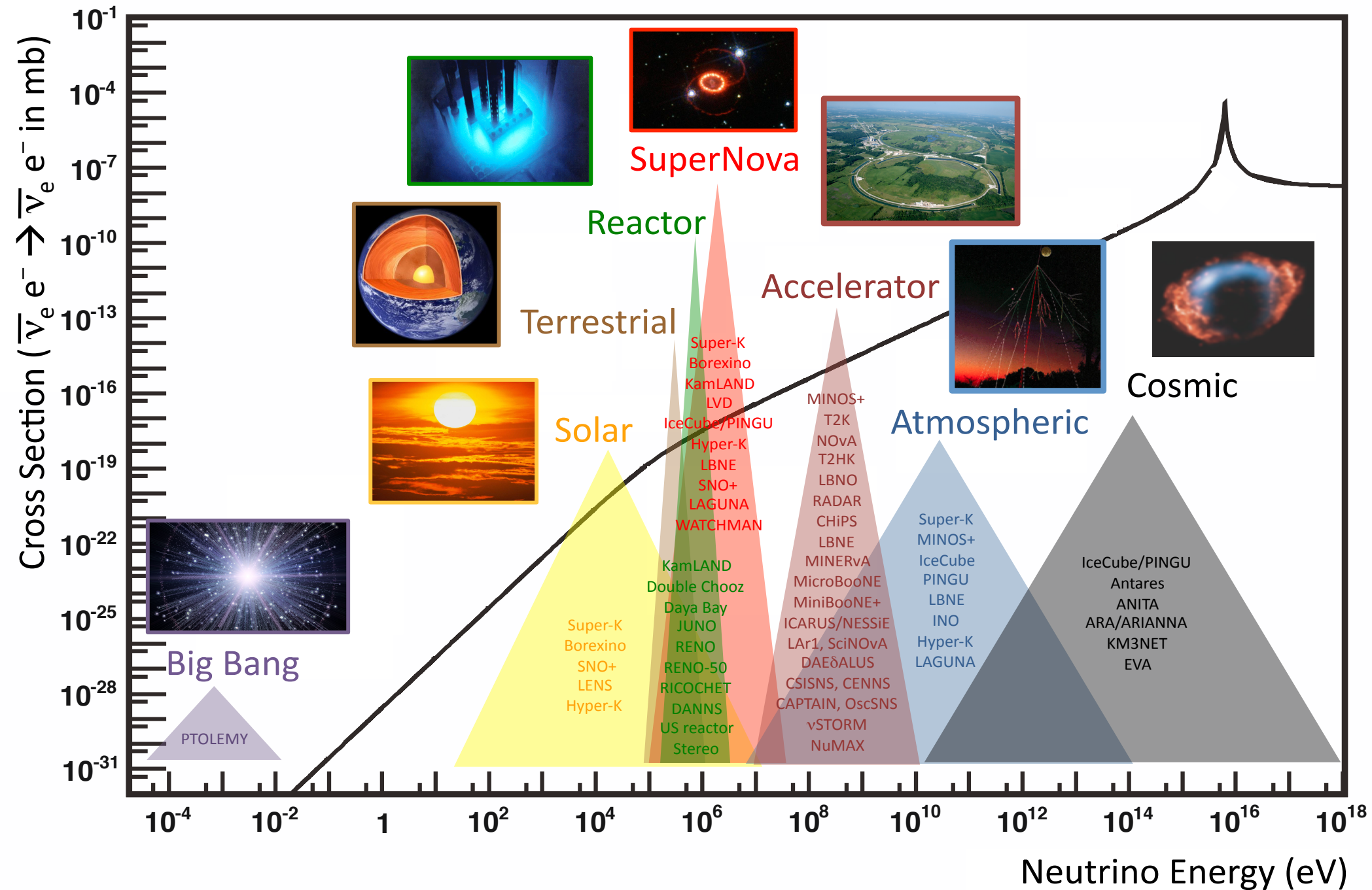
# Outlook 2-3 years



- **PTOLEMY** will operate with:
  - Completed MAC-E filter (1% cut-off)
  - Collect tritium spectra in 50-150eV of endpoint
  - 100 micrograms (1 Cu) of tritium with 1m<sup>2</sup> area (~300 layer cell)
  - 0.15eV energy resolution at 100eV
  - Demonstrate RF-tagged electron identification
  - Measure tritium cell systematics to sub-eV
- **Physics**
  - 1<sup>st</sup> direct constraint on relic neutrino density ( $10^6$  above nominal)
  - Competitive resolution performance on neutrino mass (systematics will be measured)
  - Early universe relic sterile neutrino limits (up to ~10keV) for a range ( $|U_{e4}|^2 \sim 10^{-4}-10^{-6}$ ) of sterile neutrino electron flavor content



# Overview of Neutrino Experiments



# What can Relic Neutrino Density tell us?



- Are there experimental outcomes that are inconsistent with Big Bang cosmology? **Yes!**
  - Too many cold neutrinos with no visible mass separation from the end-point (no galactic clumping factor) would contradict the initial conditions of Big Bang nucleosynthesis (present day H, D, He, Li abundances)
- Are there outcomes that are inconsistent with the Standard Model of particle physics? **Yes!**
  - No neutrino detection (exclusion of the relic neutrino density below prediction) could mean that neutrinos have a finite lifetime
- Are there possibilities for discovering new physics? **Yes!**
  - Alternative dark matter candidates such as keV sterile neutrinos may have a non-zero electron flavor content and would appear as a mass peak above the end-point

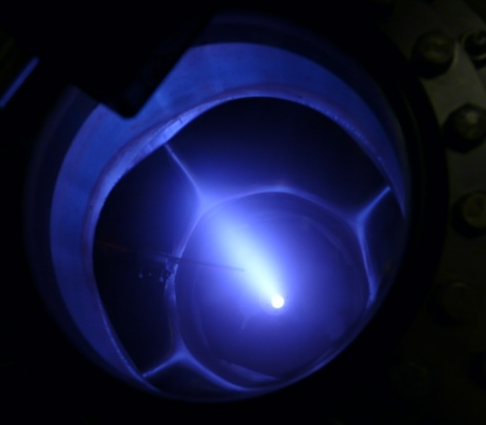
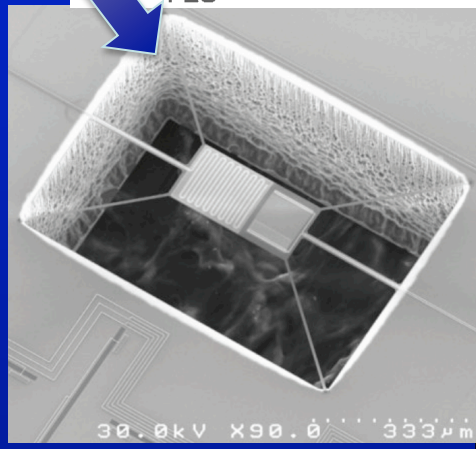
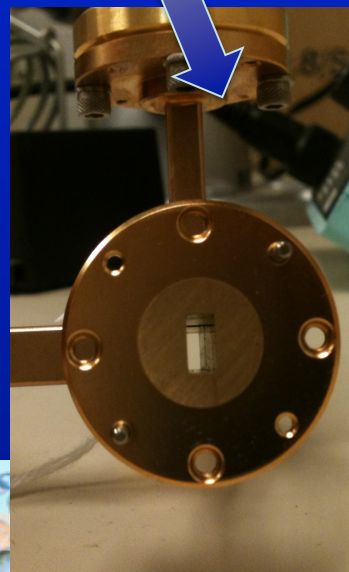
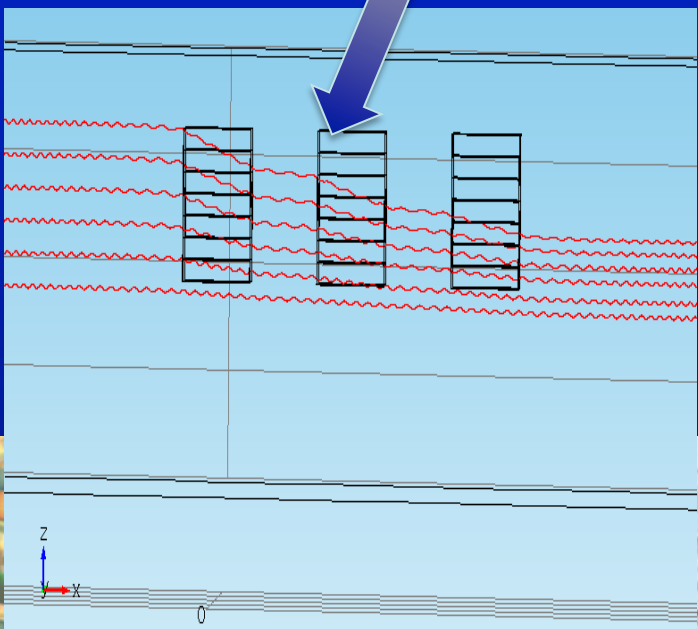
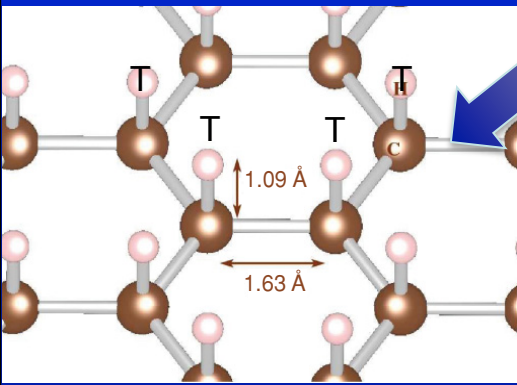
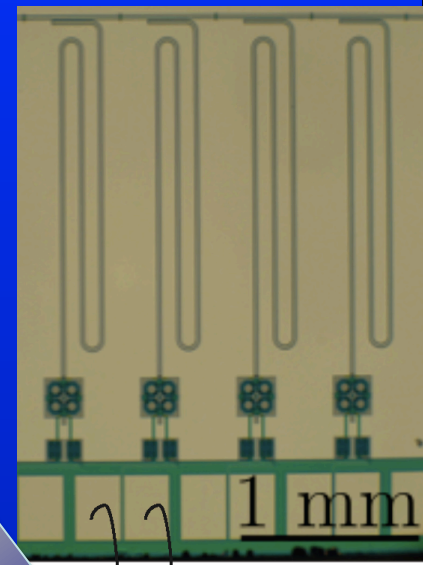
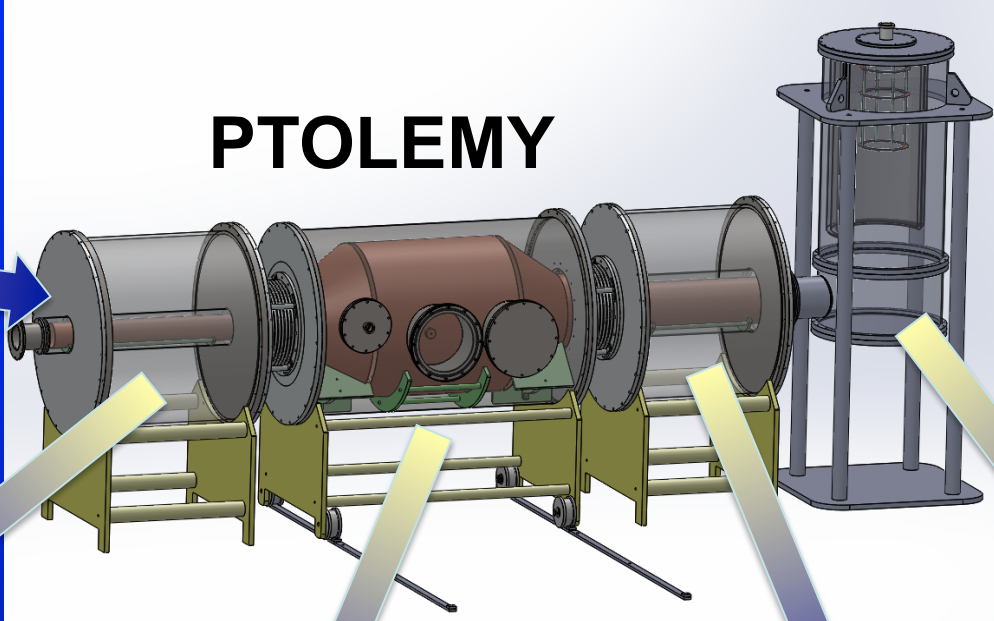
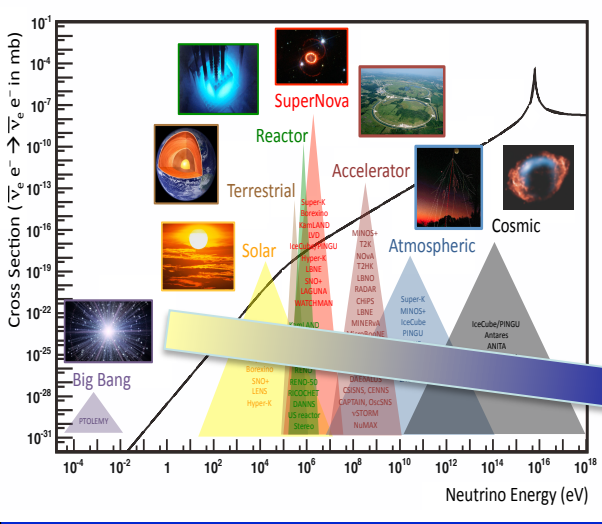


# What can Relic Neutrino Density tell us?

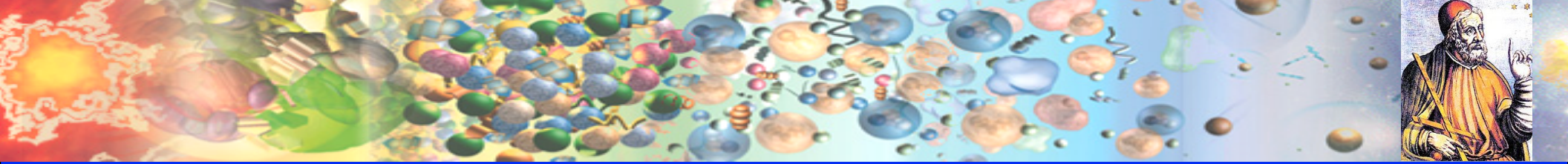


- Is there a possibility to make long-term contributions to the understanding of the Universe?
  - Absolutely! We believe that we live in a sea of 14 billion year old neutrinos all around us (the oldest relics in the Universe) – is it true?
  - When one opens a new frontier of exploration, there is no telling what will be found and learned

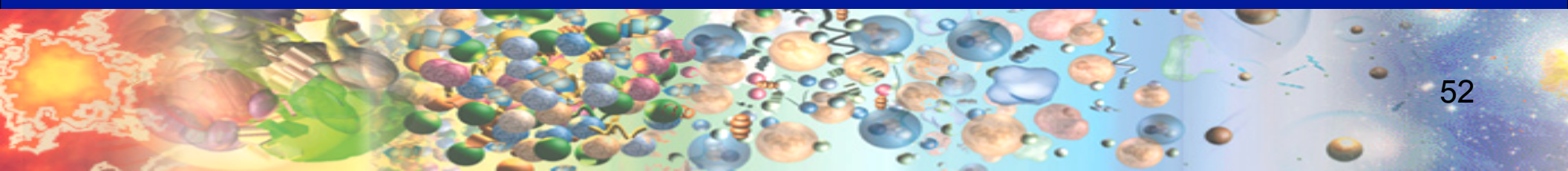
# Summary







# Backup Slides



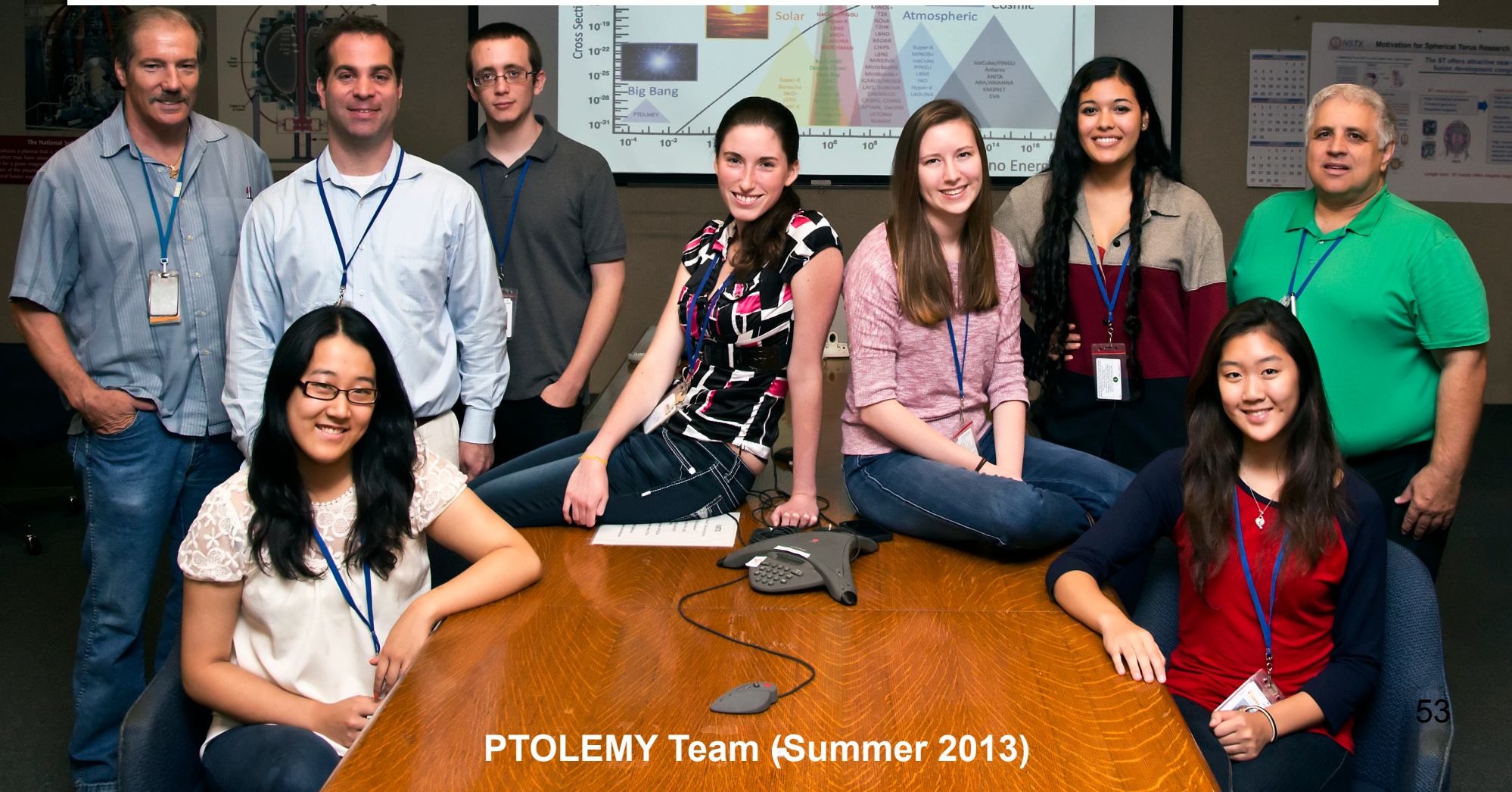


Development of a Relic Neutrino Detection Experiment at PTOLEMY:  
Princeton Tritium Observatory for Light, Early-Universe, Massive-Neutrino Yield

S. Betts<sup>1</sup>, W. R. Blanchard<sup>1</sup>, R. H. Carnevale<sup>1</sup>, C. Chang<sup>2</sup>, C. Chen<sup>3</sup>, S. Chidzik<sup>3</sup>, L. Ciebiera<sup>1</sup>, P. Cloessner<sup>4</sup>, A. Cocco<sup>5</sup>, A. Cohen<sup>1</sup>, J. Dong<sup>1</sup>, R. Klemmer<sup>3</sup>, M. Komor<sup>3</sup>, C. Gentile<sup>1</sup>, B. Harrop<sup>3</sup>, A. Hopkins<sup>1</sup>, N. Jarosik<sup>3</sup>, G. Mangano<sup>5</sup>, M. Messina<sup>6</sup>, B. Osherson<sup>3</sup>, Y. Raitses<sup>1</sup>, W. Sands<sup>3</sup>, M. Schaefer<sup>1</sup>, J. Taylor<sup>1</sup>, C. G. Tully<sup>3</sup>, R. Woolley<sup>1</sup>, and A. Zwicker<sup>1</sup>

<sup>1</sup>Princeton Plasma Physics Laboratory

arXiv: 1307.4738



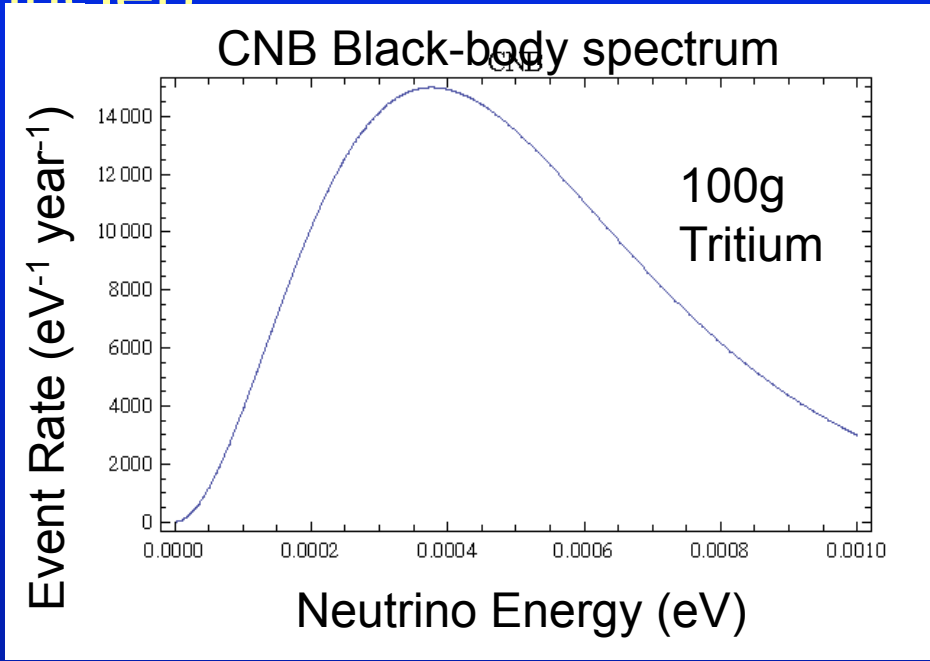
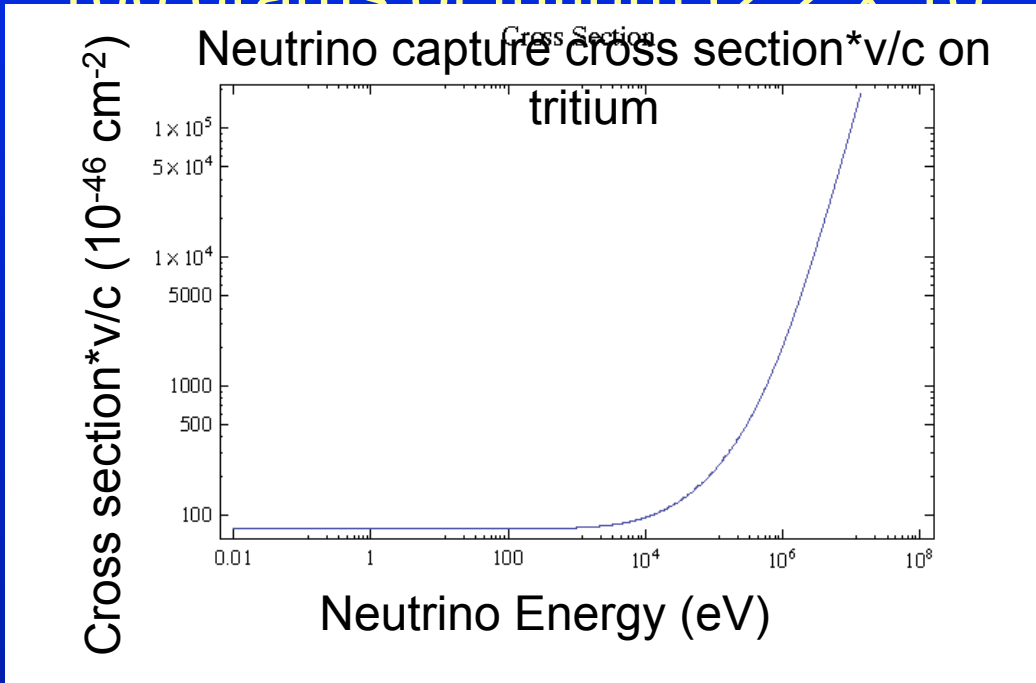
PTOLEMY Team (Summer 2013)



# Neutrino Interaction Rates



- 1 SNU = 1 neutrino interaction per second for  $10^{36}$  target nuclei
- 100 grams of tritium ( $2.2 \times 10^{25}$  nuclei)



$$\int \sigma(p_\nu) v_\nu f_\nu(p_\nu) \frac{d^3 p_\nu}{(2\pi)^3}$$

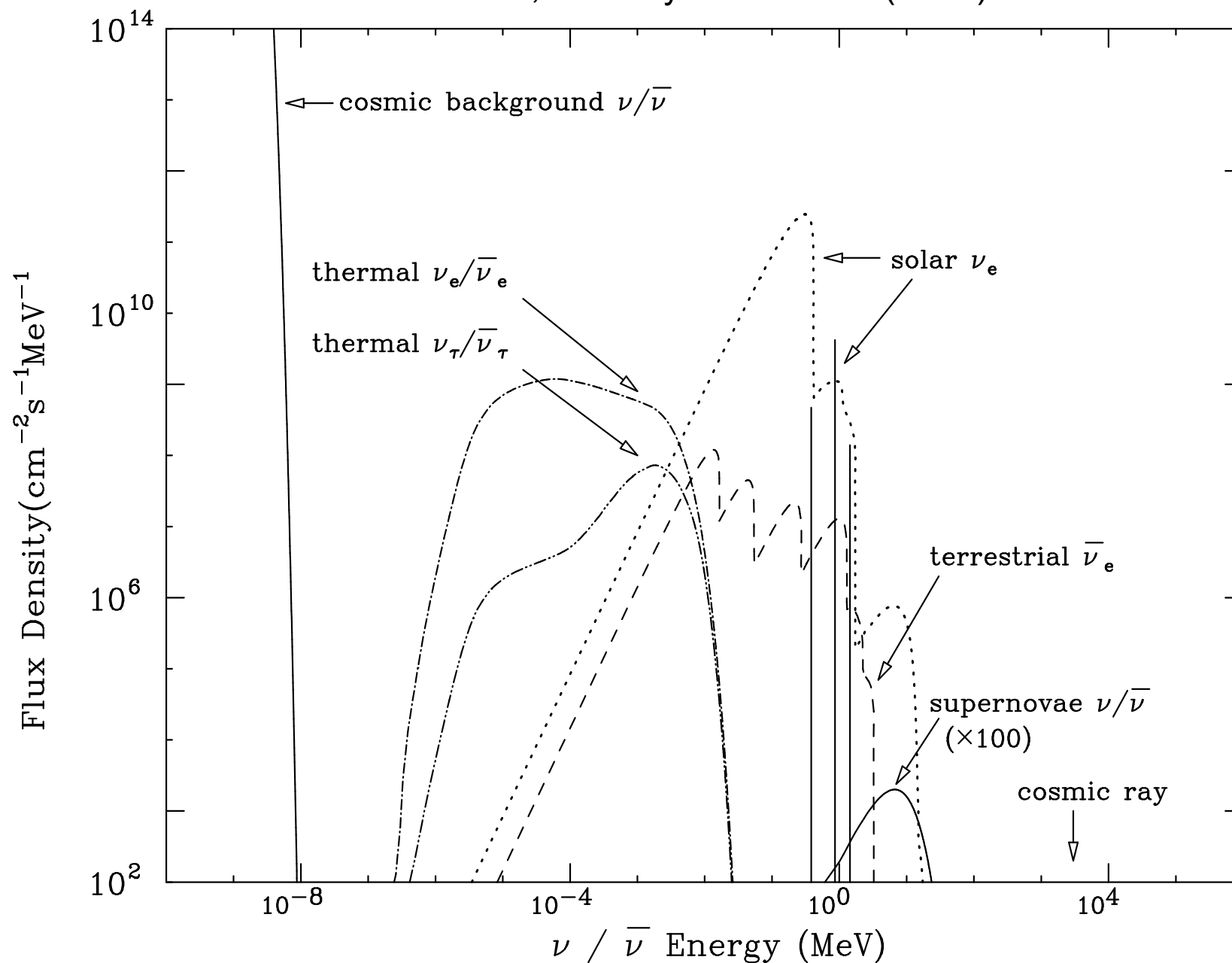
$9.51 \pm 0.03$  events/year ( $13600 \pm 50$  SNU)

Laurentiu Rodina

# Neutrino Flux



Haxton, Lin: Phys.Lett. B486 (2000) 263-271



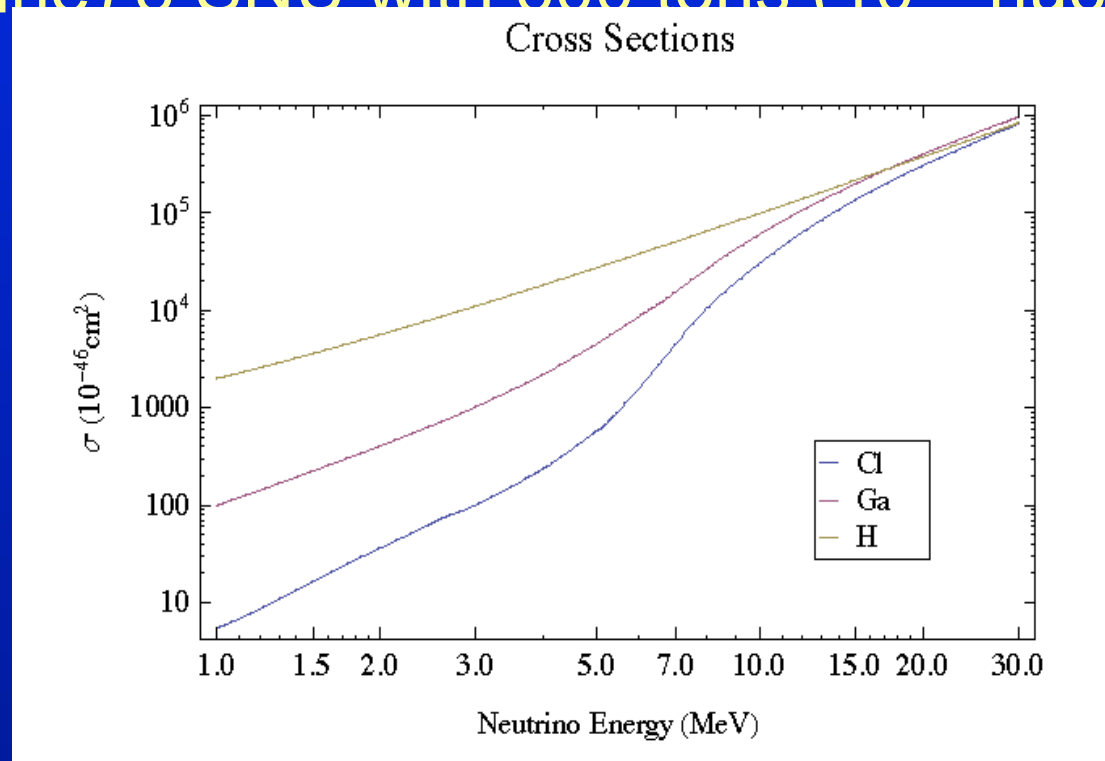


# Solar Neutrino Capture Experiments



- PTOLEMY ~3618 SNU with 100g ( $10^{25}$  nuclei) 2.5 evts/year
- Gallex 70 SNU with 30 tons ( $10^{29}$  nuclei) 1200 evts/year
- Homestake (Chlorine) 8 SNU with 600 tons ( $10^{31}$  nuclei) 2500 evts/year

Hard to compete with  
Tritium for sub-MeV  
neutrino energies



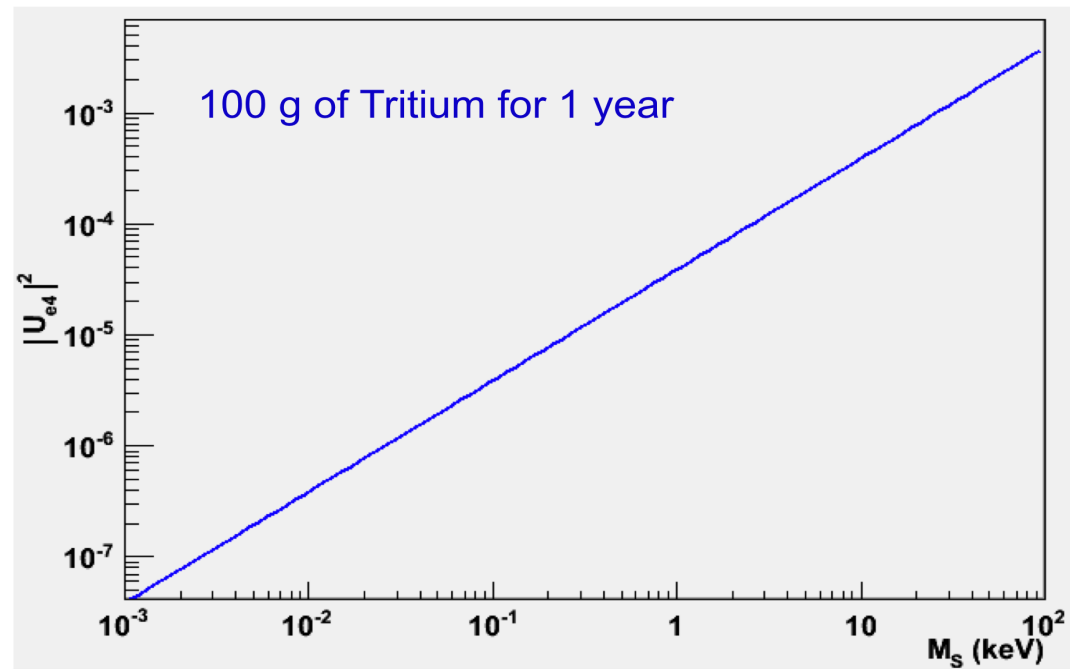
# Sterile Neutrino Search



Using  $\nu$  capture...

If Dark Matter is made by sterile neutrino  $\rightarrow \rho_S \sim \frac{0.4 \times 10^6}{M_S [\text{keV}]} \text{ cm}^{-3}$

Looking beyond the beta decay endpoint energy (background free region)





# PTOLEMY Conceptual Design

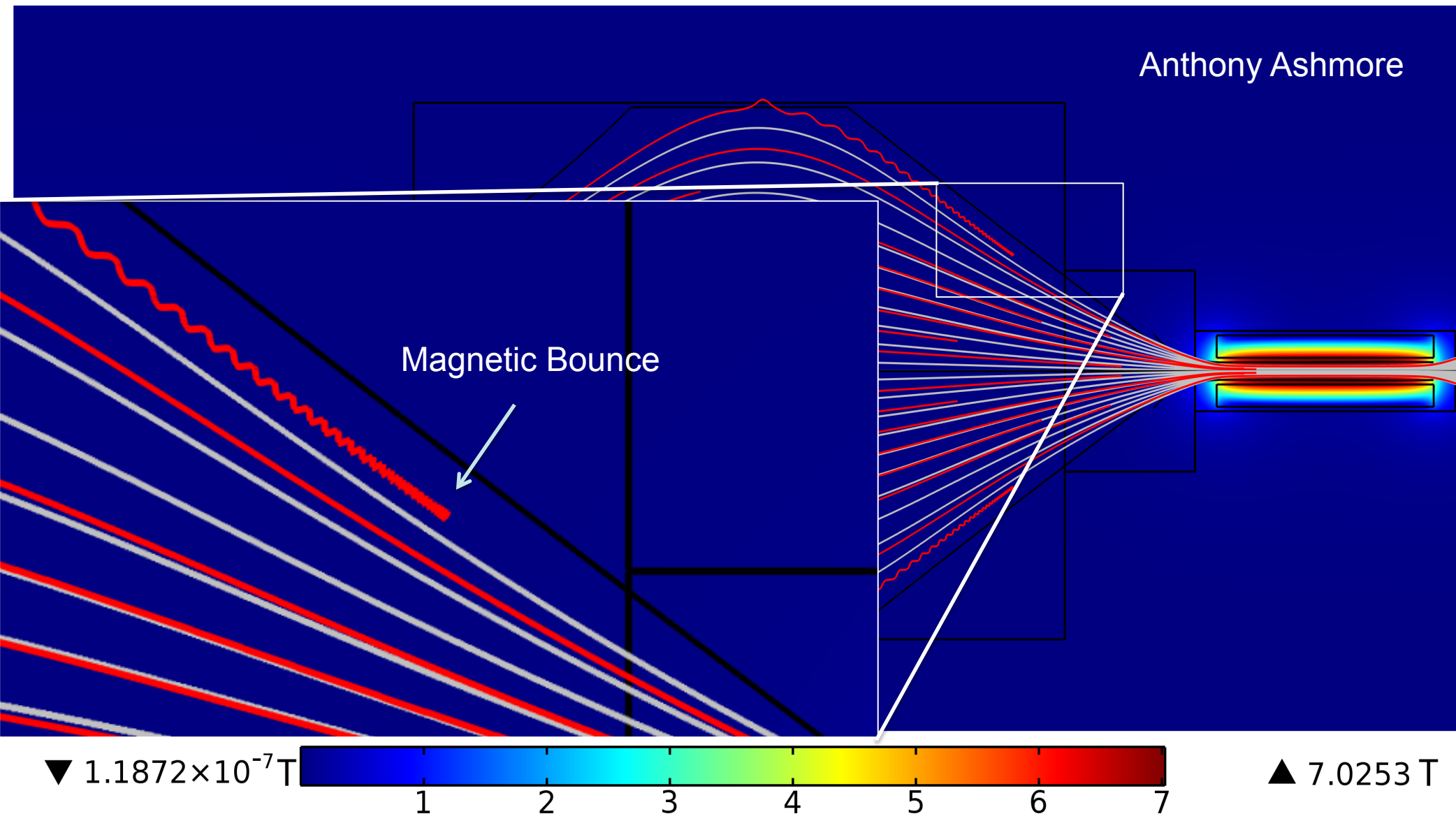


- High precision on endpoint
  - Cryogenic calorimetry energy resolution
  - **Goal: 0.1eV resolution**
- Signal/Background suppression
  - RF tracking and time-of-flight system
  - **Goal: sub-microHertz background rates above endpoint**
- High mass, high resolution tritium target
  - Surface deposition (tenuously held) on conductor in vacuum
  - **Goal: for CNB: maintains 0.1eV signal features with high efficiency**
  - **For sterile nu search: maintains 10eV signal features w/ high eff.**
- Scalable mass/area of tritium source and detector
  - **Goal: relic neutrino detection at 100g**
  - **Sterile neutrino (w/ % electron flavor) at ~1g**

# Trajectory Calculations



Anthony Ashmore

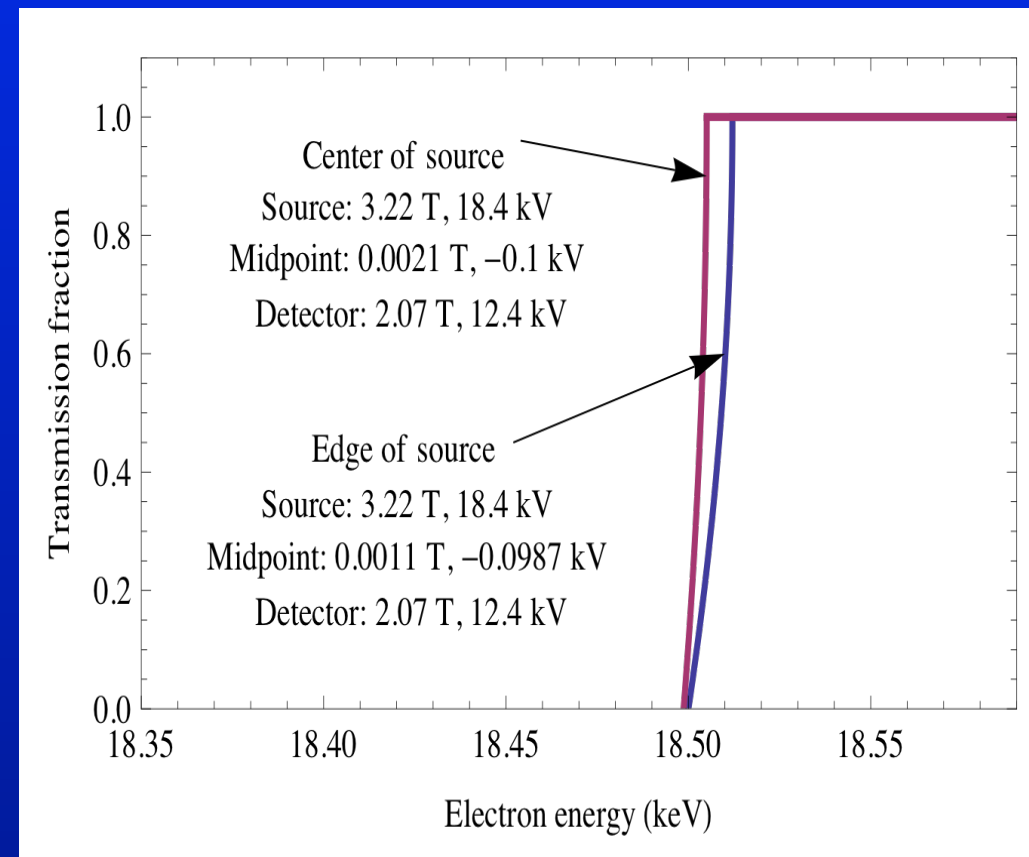
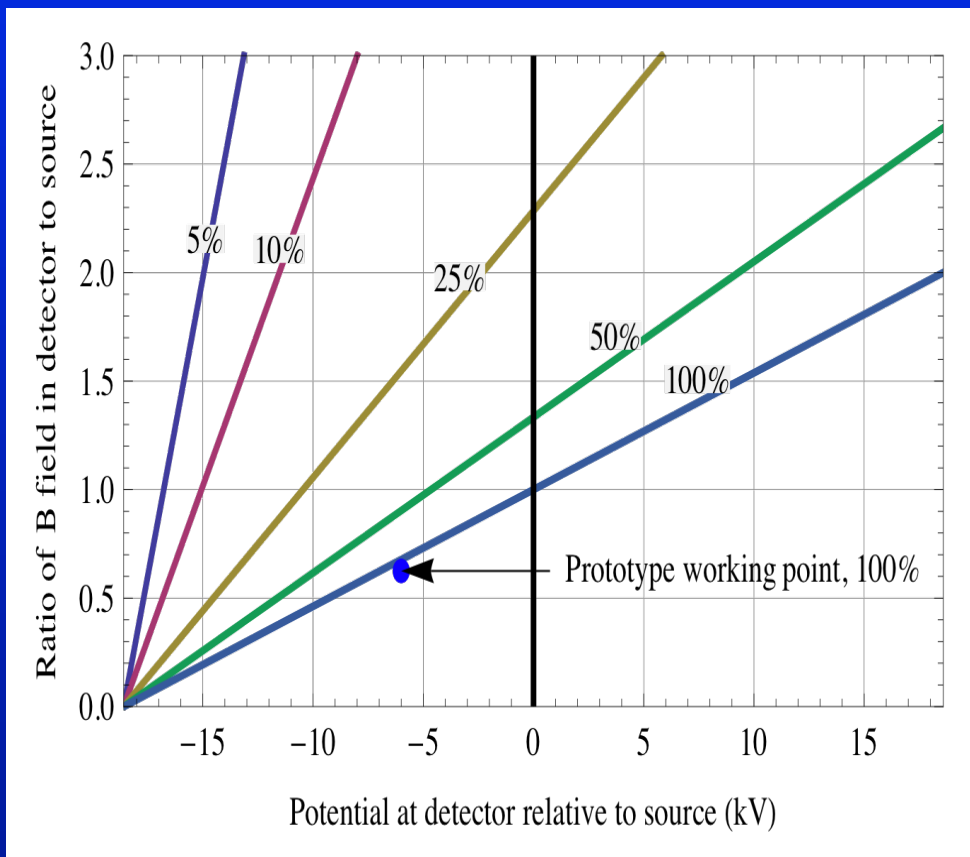




# Cut-off Uniformity and Decay Acceptance



- In order to avoid magnetic bounce, electrons must be accelerated back up in going from mid-plane to detector
- Different trajectories have different cut-off precisions



Anthony Ashmore

# Calibration and Backgrounds



- High precision (0.1eV) electron gun
  - Off-axis directionality needed for RF antenna calibration
  - Investigating possibility of a single or multiple high precision guns situated outside of the magnetic field of the tritium target plate with a “switch yard” of input spigots to provide in situ calibration peaks for every calorimeter channel and electron trajectory
- Vacuum studied with residual gas analyzer (RGA)
- Several possibilities for background estimation
  - sideband data-driven background estimation below MAC-E filter cutoff
  - out-of-time tracking-calorimeter coincidence
  - (vacuum-)scattered electron trajectory analysis
  - varying source strength tiles (null sources)
- NMR calibration for magnetic field uniformity in RF tracker



# Calibration System



Precision e-gun  
“switchyard”

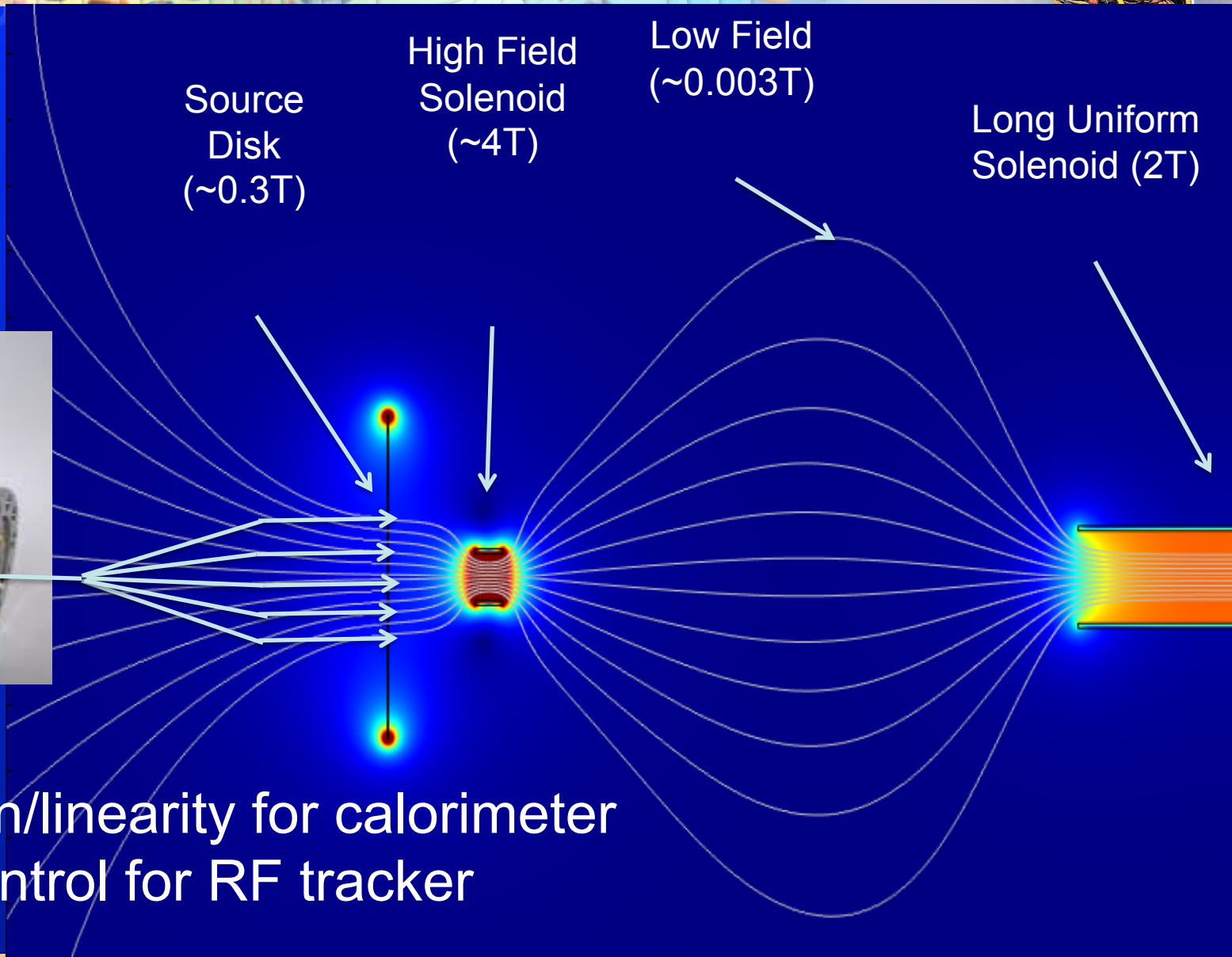


Source  
Disk  
(~0.3T)

High Field  
Solenoid  
(~4T)

Low Field  
(~0.003T)

Long Uniform  
Solenoid (2T)



Energy resolution/linearity for calorimeter  
Angular control for RF tracker

# TFTR Carbon tiles



- PPPL already produced tritium-graphene samples in the 90's in TFTR
  - These have been analyzed with Raman scattering
- SRNL are providing tritium loaded graphene samples
  - Loaded at 100 atm

“Hot spots” of TFTR carbon tile indicate tritium

