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# Non-Linear Small Scale Structure Formation in Non-Standard Dark Matter Cosmologies

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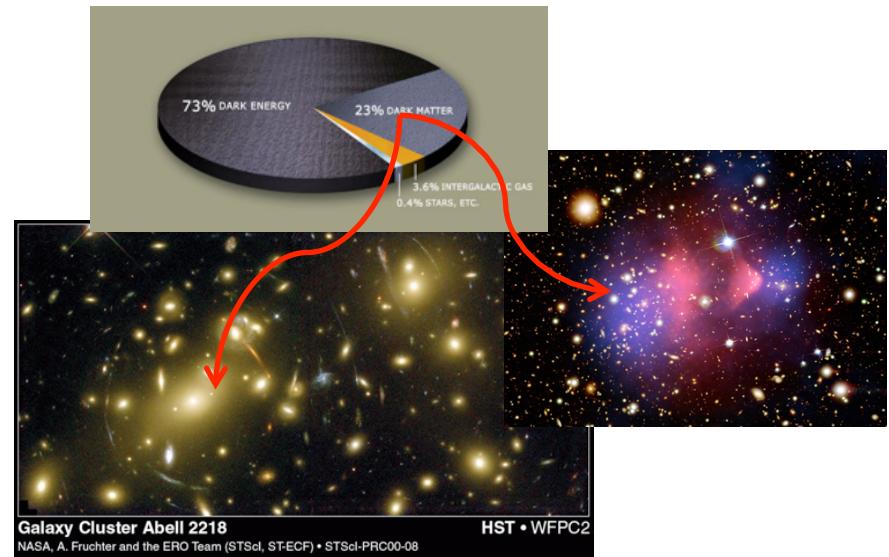
**Subinoy Das**

**Doddy Marsh**

# Standard Cosmological Scenario

## Dark Matter:

- Gravitational Collapse Initial Fluctuations
- Foster matter clustering
- Resides in virialized clumps



## Halos:

- Building blocks of cosmic structure formation
- Shape baryon distribution & formation of visible structures

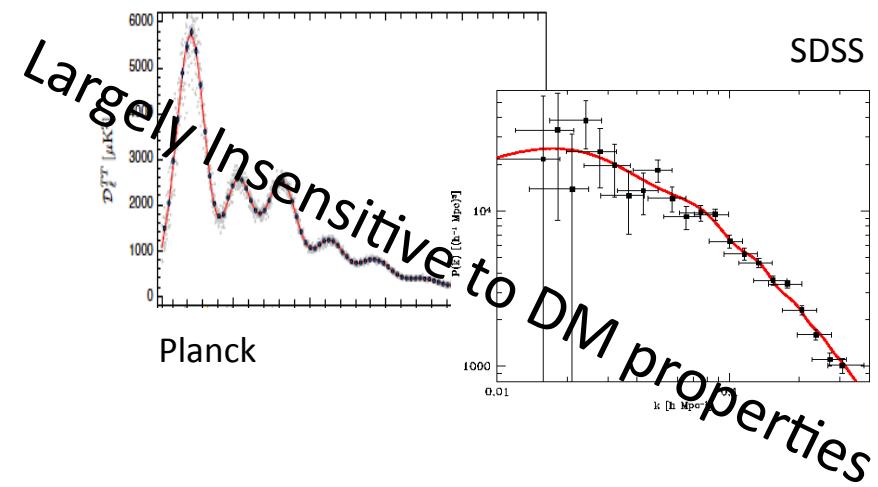
# CDM Paradigm

## Theoretical Bias:

- WIMP hypothesis/Neutralino

## Large Scales

- Successful Description CMB spectra
- Clustering of Matter from galaxy surveys



## Minimal Scenario - LCDM

- 6 parameter models:  $H_0$ ,  $\Omega_m(1-\Omega_\Lambda)$ ,  $\Omega_b$ ,  $\sigma_8$ ,  $n_s$ ,  $\tau$
- Censored Comments

# Small Scales & Beyond CDM

## DM Direct Searches:

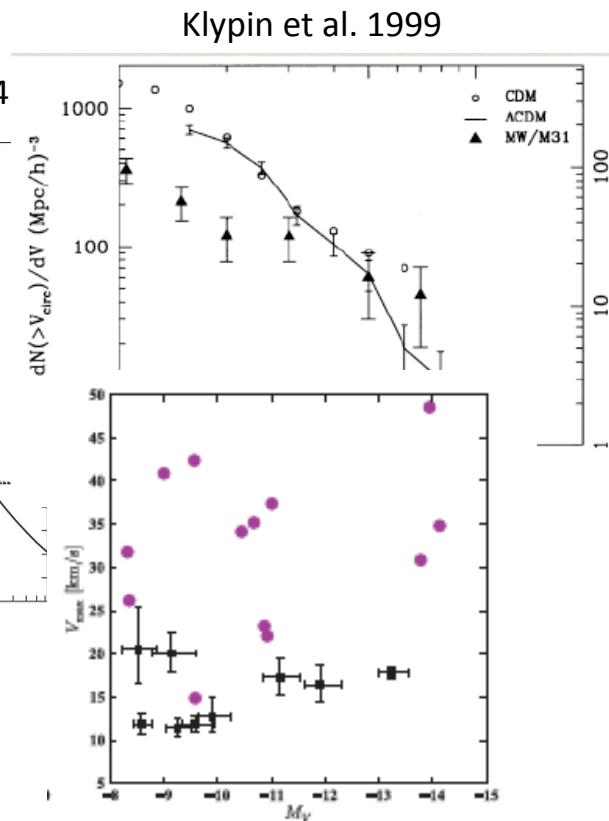
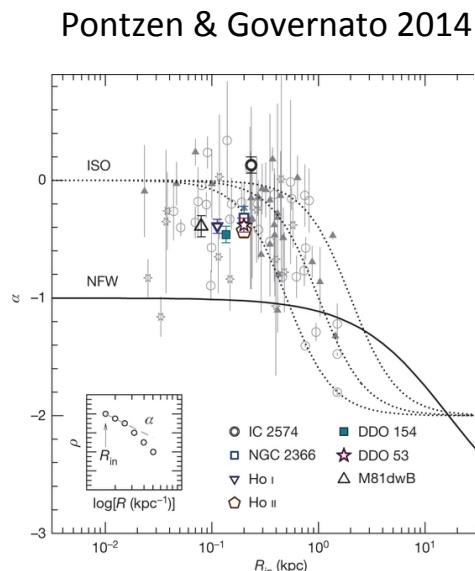
- Negative/Contrasting Results
- No signal at LHC

## CDM anomalies :

- Core vs Cusp Profiles
- Missing Satellites
- Too-big-too fail problem

## Complex Physics:

- Baryonic Feedback
- Observational Selection Effects
- Uncertainties of Milky-Way Mass



# Alternative Scenarios

## Warm Dark Matter

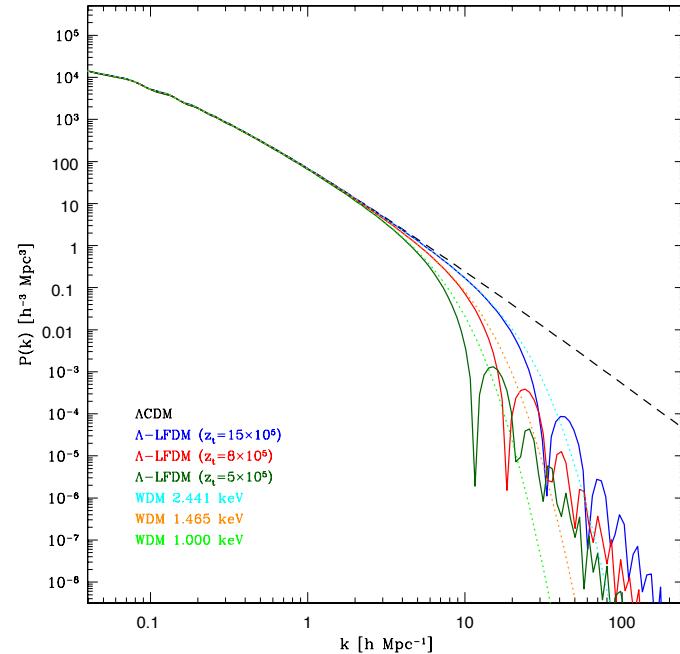
- Thermal Relic ,  $m_{\text{WDM}} \approx \text{keV}$
- Free-Stream  $\leq 100 \text{ kpc}$
- Small-Scale Power Spectrum Cut-off

## Self-Interacting DM

- DM scattering cross-section
- Interaction with radiation

## Late Forming DM & Ultra-Light Axions

- LFDM: Before matter/radiation equality, decay of scalar field coupled to relativistic particles ( $w \approx 1/3 \rightarrow w=0$ )
- ULA: Axion field transition from vacuum to matter ( $w=-1 \rightarrow w=0$ )



# Observational Constraints

## Warm Dark Matter\*

- $m_{\text{WDM}} < 0.1 \text{ keV}$  to core the profiles e.g. Maccio et al. 2012
  - $1.5 < m_{\text{WDM}} [\text{keV}] < 2$  give too-big-to-fail Lowell et al. 2012, 2014
  - $m_{\text{WDM}} > 3.3 \text{ keV}$  Lyman- $\alpha$  power spectrum at  $z > 2$  Viel et al. 2012
- CUM GRANO SALIS*

## Self-Interacting DM

- Lower density sub-halos and core profiles Vogelsberger et al. 2012; Zavala et al. 2013
- Low mass halo abundances unaltered

## Late Forming DM

- Evade Lyman-alpha constraints ( $z_t > 5 \cdot 10^5$ ) Agarwal, Corasaniti, Das & Rasera 2015
- Suppressed halo abundances
- Flatter profiles than LCDM (not cored)

# N-body Simulations

## Monte Carlo Method

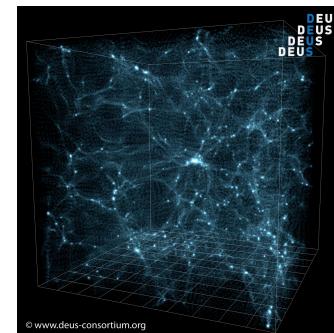
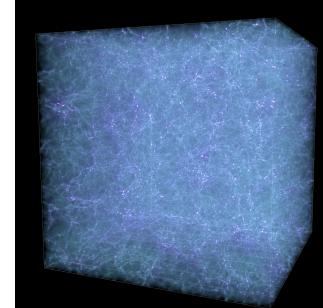
- Macro-particle sampling DM field
- Initial Conditions:
  - WDM thermal velocities negligible (?)
  - LFDM effectively collision-less
- Numerical integration trajectories

$$\frac{d\vec{p}_i}{da} = -\frac{\nabla\Phi}{\dot{a}}, \quad \frac{d\vec{x}_i}{da} = \frac{\vec{p}}{\dot{a}\ddot{a}}, \quad i = 1, N$$
$$\nabla^2\Phi = 4\pi G a^2 \bar{\rho}_m \delta_m$$

- Analyze final macro-particle distribution

## Simulation Characteristics

- Gravity Solver
- Cosmological Volume
- Mass Resolution
- Spatial Resolution



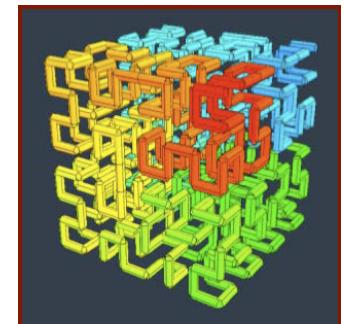
## Initial Conditions

- ZA displacements from input P(k)



## N-body Solver

- AMR, TreePM, etc.



$$L_{box}^3$$

$$m_p = \rho_c \Omega_m L_{box}^3 / N_p$$
$$\Delta x$$

# Artificial Fragmentation (Mass Segregation)

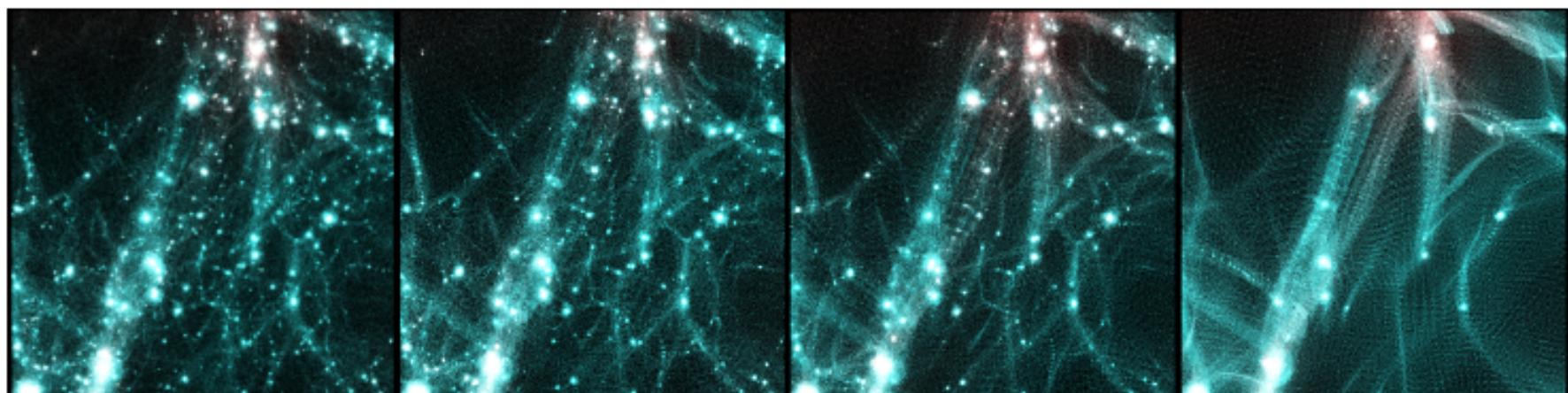
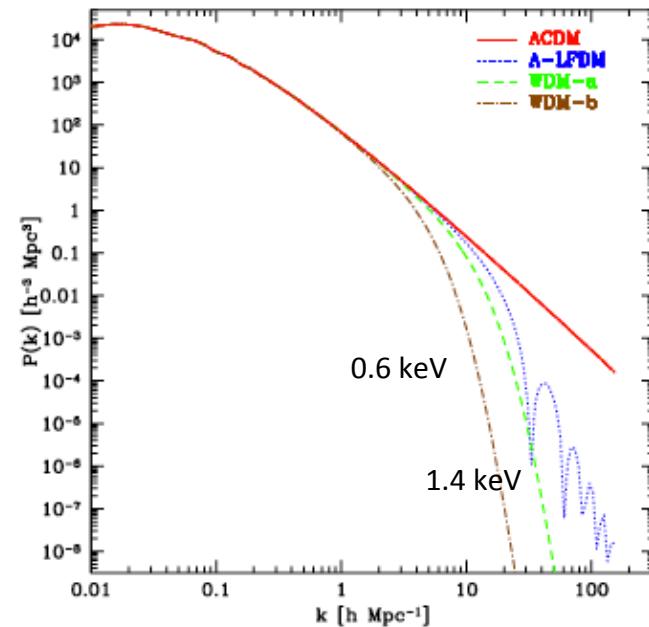
## Discretization Effect

- Sampling Poisson Noise ( $k > k_{\text{cut-off}}$ )
- Spurious Numerical Halos

Gotz & Sommer-Larsen 2002, 2003; Wang & White 2007

## Example

RAMSES    -  $N_p = 512^3$                             -  $m_p \sim 10^7 M_{\text{sun}} h^{-1}$   
AMR        -  $L_{\text{box}} = 27.5 \text{ Mpc } h^{-1}$     -  $dx_{\text{coarse}} \sim 54 \text{ kpc } h^{-1}$



# Spurious Halo Contamination

## Halo Mass Function

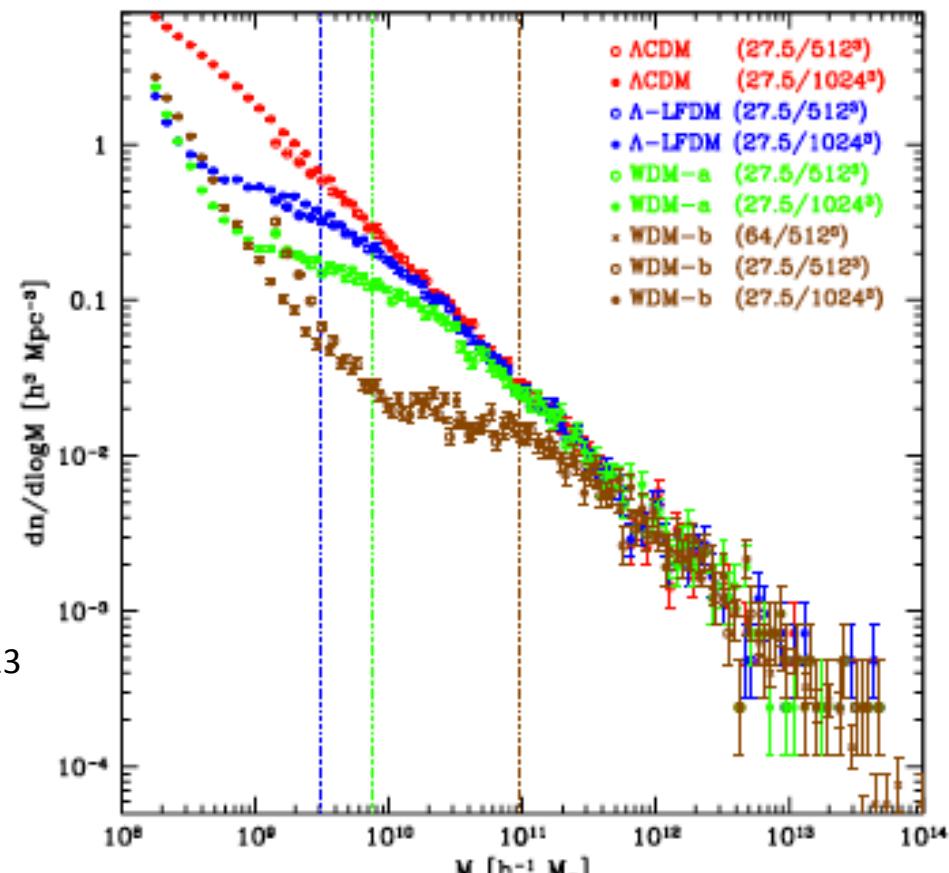
- $N_{\text{h-particles}} > 100$
- Upturn at  $M < M_*$
- Simulation Dependent Slope

## Proposed Cures

- Mass Cut:  $M_{\min} = 10.1 \rho d k_p^{-2}$   
Wang & White 2007
- Select Unflatten Proto-Halos  
in Initial Lagrangian Patch &  
Apply Mass Cut      Lowell et al. 2012
- Visual Inspection    Angulo, Hahn, Abel 2013
- Tessellation 6-d phase-space  
folding (reduce but doesn't solve)

Hahn, Abel, Kaehler 2013

- $N_p = 1024^3$
- $L_{\text{box}} = 27.5 \text{ Mpc } h^{-1}$
- $m_p \sim 10^6 M_{\text{sun}} h^{-1}$
- $d\chi_{\text{coarse}} \sim 26 \text{ kpc } h^{-1}$



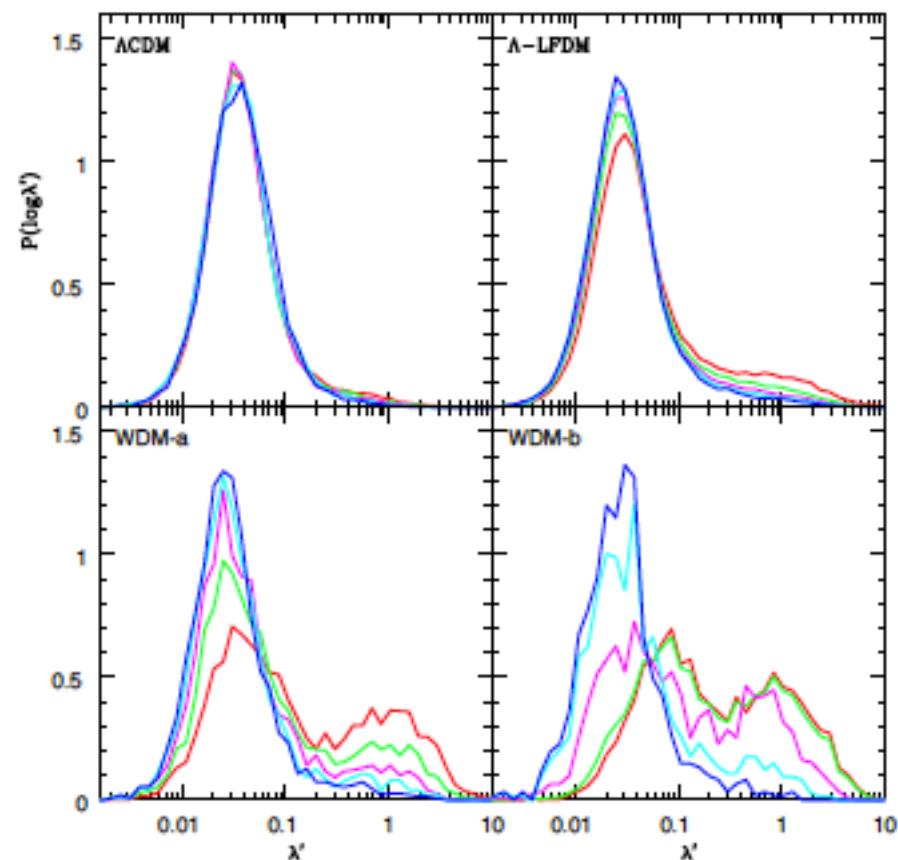
Agarwal & Corasaniti 2015

# Structural Properties of Halos

Agarwal & Corasaniti 2015

## Halo Spin

- Spin parameter  $\lambda' = \frac{J}{\sqrt{2}MVR}$
- $V = \sqrt{GM/R}$
- 8 bins:  $4 < M[10^9 M_{\text{sun}} h^{-1}] < 8$
- CDM: log-normal & mass independent
- non-CDM: deviations from lognormality/bimodality and mass dependent
- spurious halos have large spins

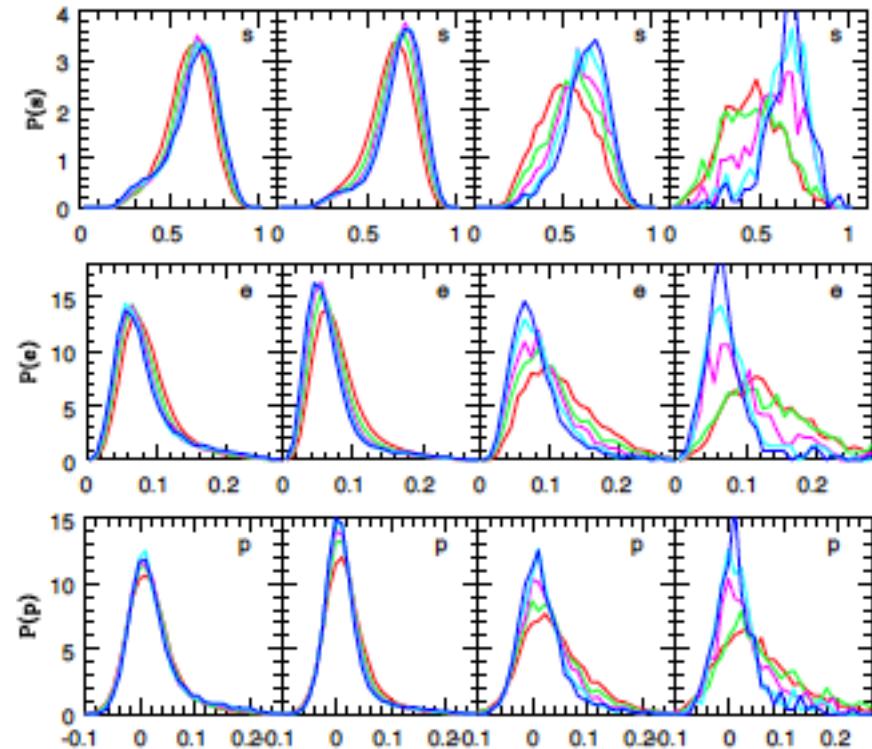


# Structural Properties of Halos

## Halo Shape

- Symmetric Mass Distribution Tensor
- sphericity, ellipticity & prolatness
- CDM: mass independent & elliptical halos
- non-CDM: mass dependent & highly non-spherical (elliptical & prolate, i.e. alignment with filaments)

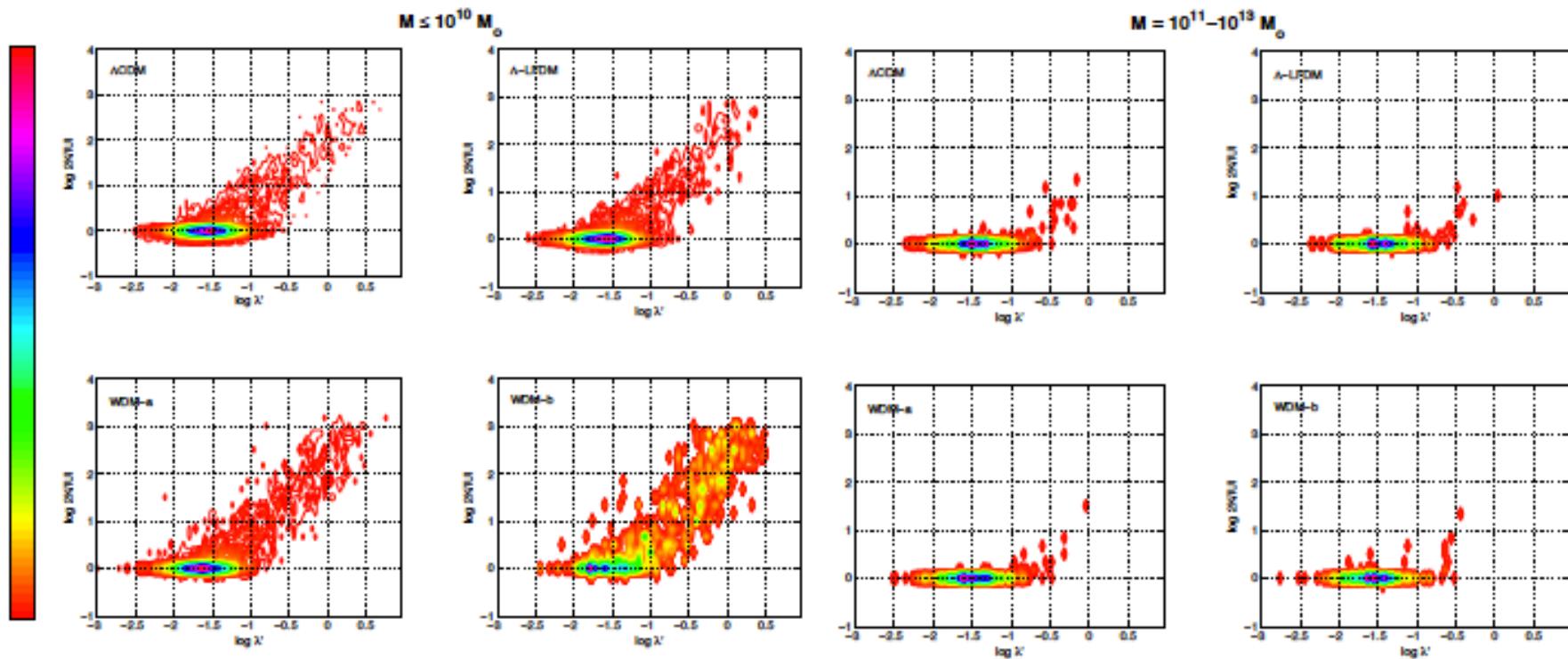
$$M_{\alpha\beta} = \frac{m_p}{M} \sum_{i=1}^{N_h} (r_{\alpha,i} - r_{\alpha,c})(r_{\beta,i} - r_{\beta,c})$$



# Halo Dynamical State

## Virial Condition

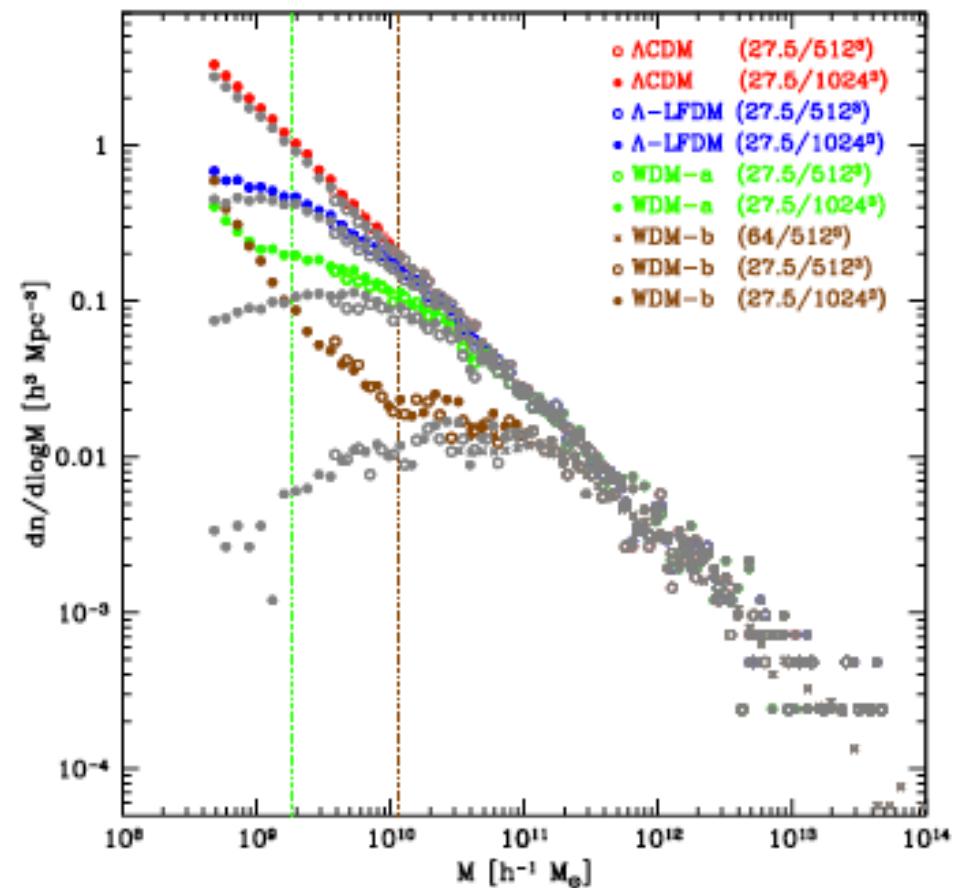
- proxy:  $\eta = 2 K / |E|$
- correlation  $\lambda' - \eta$  for  $\eta > 1$



# Virial State Selection

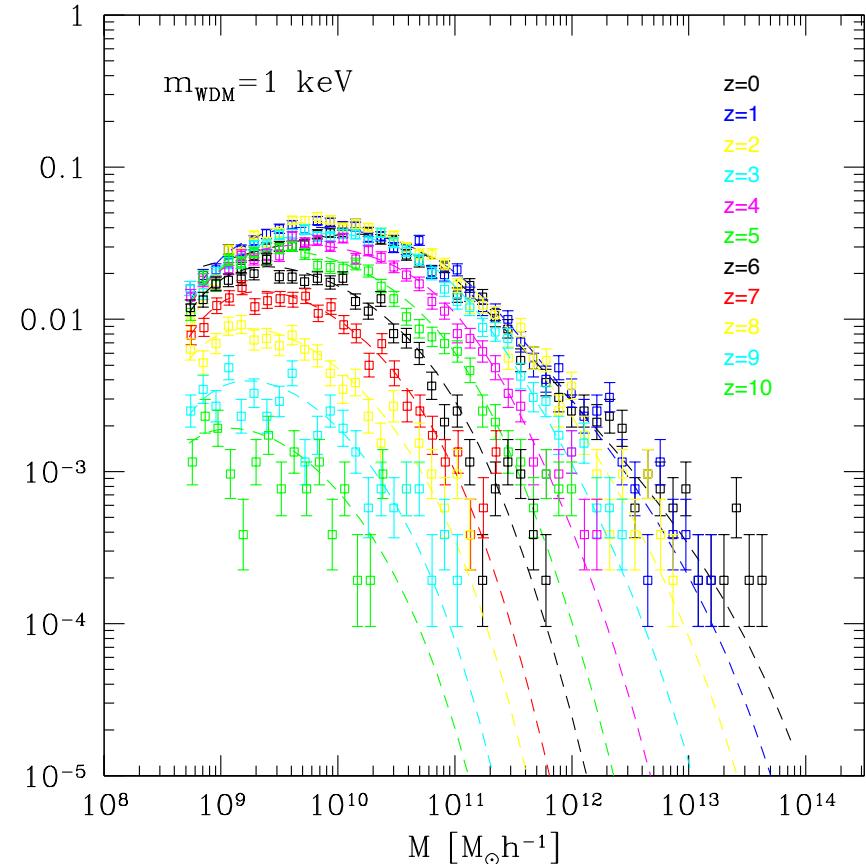
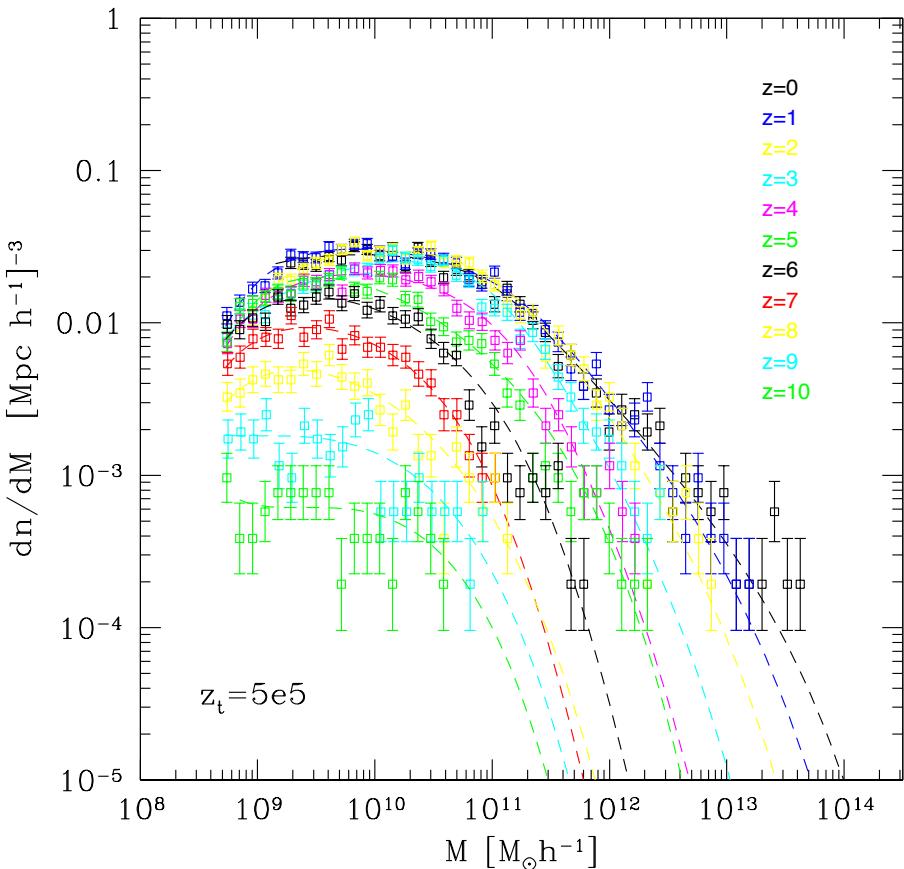
## Removing Spurious Halos

- $0 < \eta=2 K/|E| < 1.5$
- recover halo triaxial distribution
- recover spin log-normality
- recover suppressed mass function at low mass (mass resolution convergence)
- spurious halos still present with simple mass-cut at  $M_{\min}$
- mass range larger than mass cut



# Evolution of HMF in NDM models

Corasaniti et al. in preparation



- Low mass end saturates at  $z < 3$

- Fitting function:

$$\left( \frac{dn}{d \ln M} \right)_{NDM} = 10^{A_0 + A_1 M} \left( \frac{dn}{d \ln M} \right)_{LCDM} \left[ 1 - e^{-\frac{M}{m_*}} \right]^{-\alpha}$$

# Abundance of Field Dwarf-Galaxies z=0

## Velocity Function

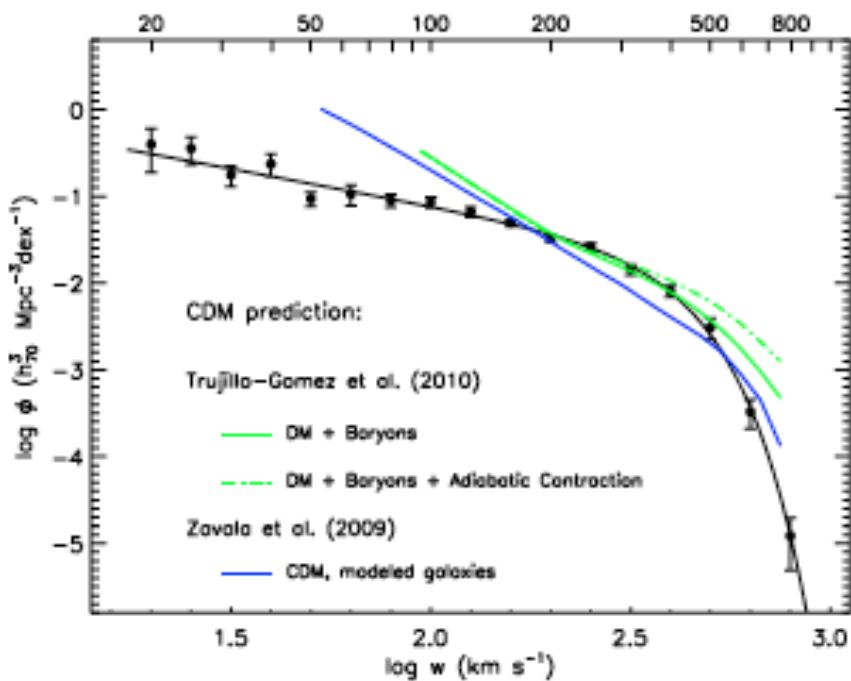
- Abundance of galaxies with given circular velocity
- More sensitive to halo dynamics than gas physics
- Mostly individual halos rather than satellite

## Caveats

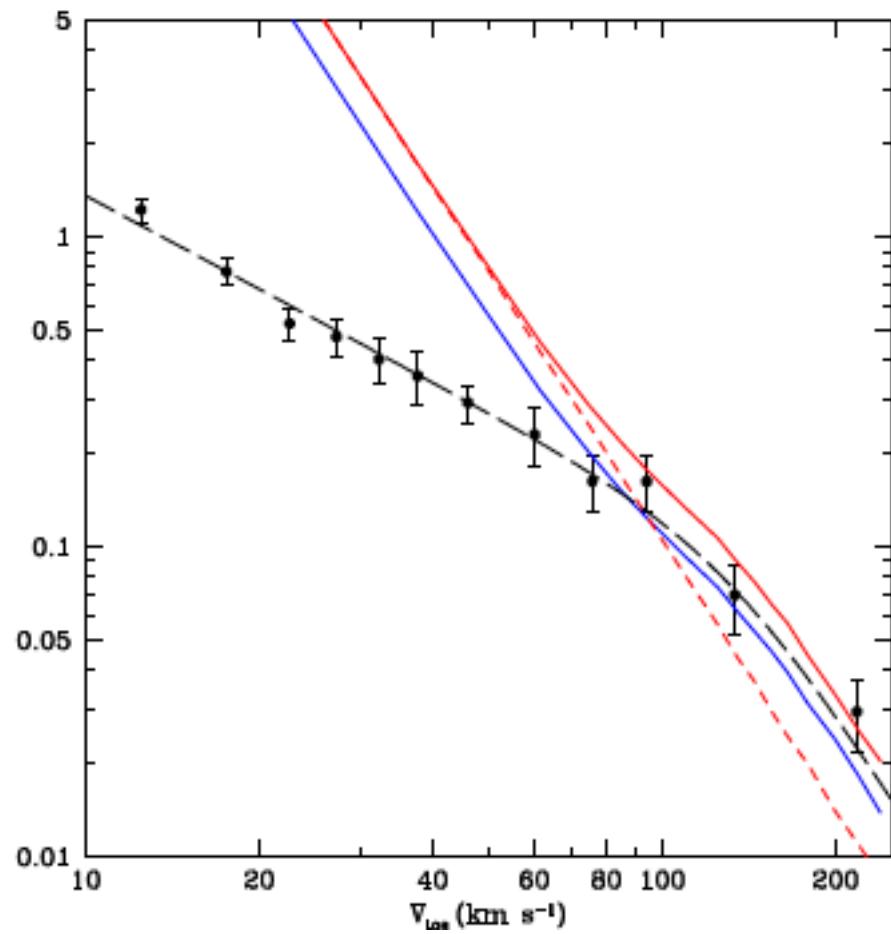
Theory: circular velocity / Observations: line-of-sight velocity

- Inclination effect
- Baryons at high-velocity end
- Selection

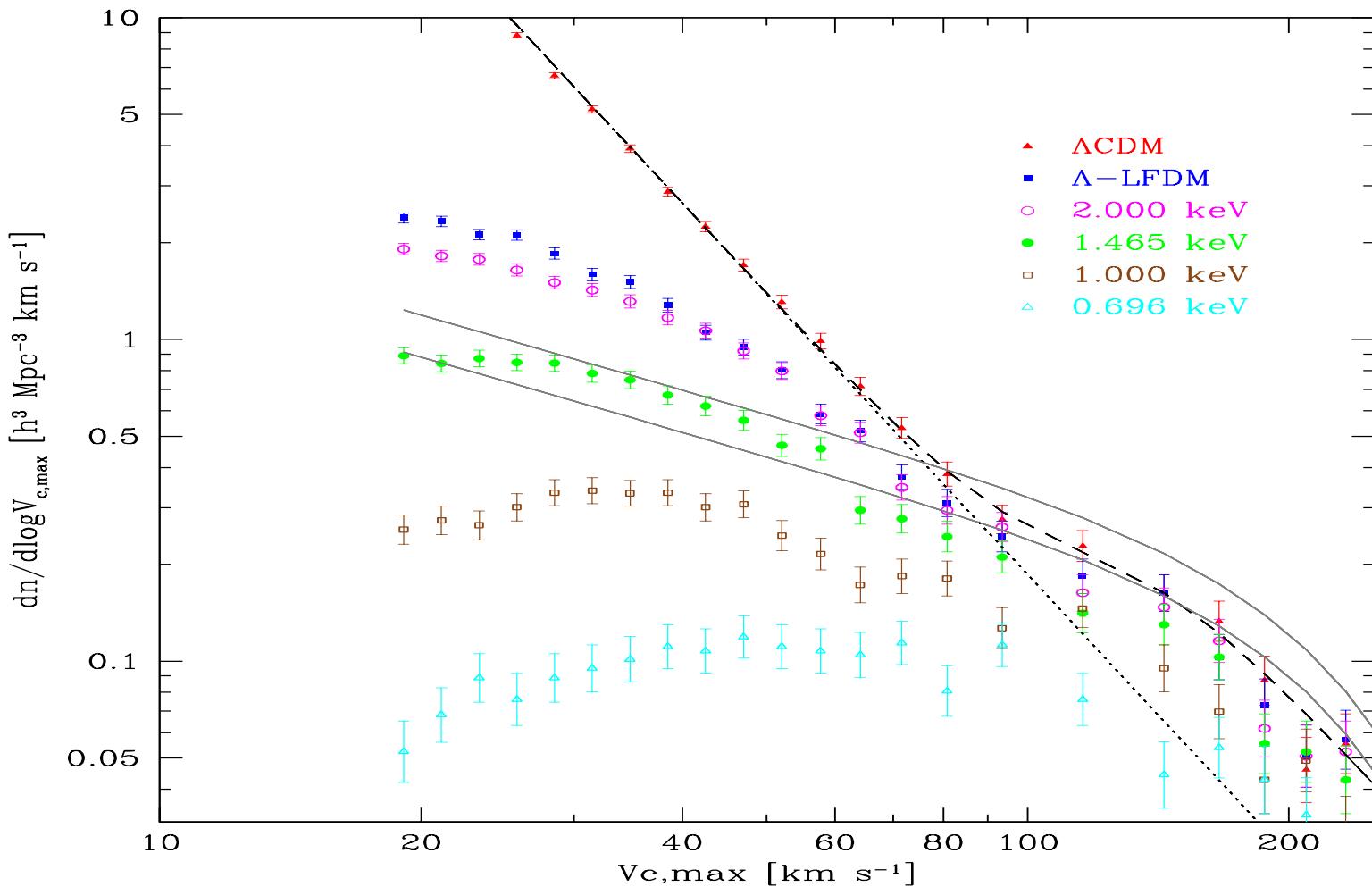
# CDM Overabundance



Klypin et al. 2015



# Velocity Function



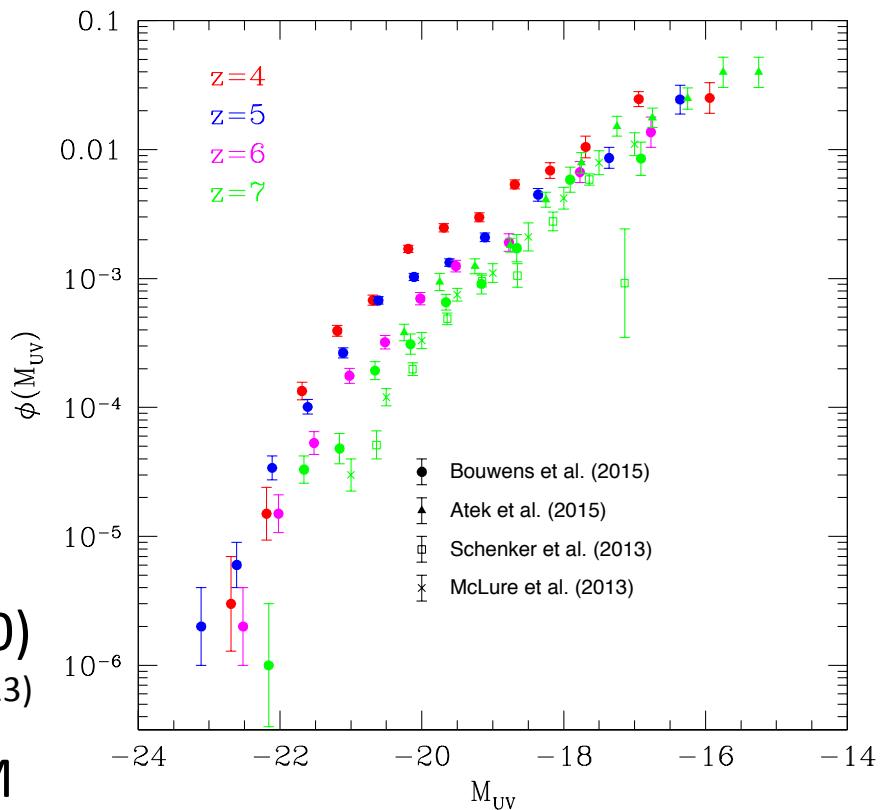
# High-Redshift Universe

## Galaxy Luminosity Function

- Probe low mass end HMF evolution
- Depends on galaxy formation
- Several estimates at  $z > 4$

## Constraints on NDM

- $m_{\text{WDM}} > 1 \text{ keV}$  (CLASH gal. den.  $z = 10$ )  
Pacucci, Mesinger & Haiman (2013)
- $m_{\text{WDM}} > 0.8 \text{ keV}$  from HAM of LCDM  
Schultz et al. (2014)
- $m_{\psi} > 10^{-22} \text{ eV}$  from HAM assuming parametrize  $L_{\text{UV}}(M)$  relation

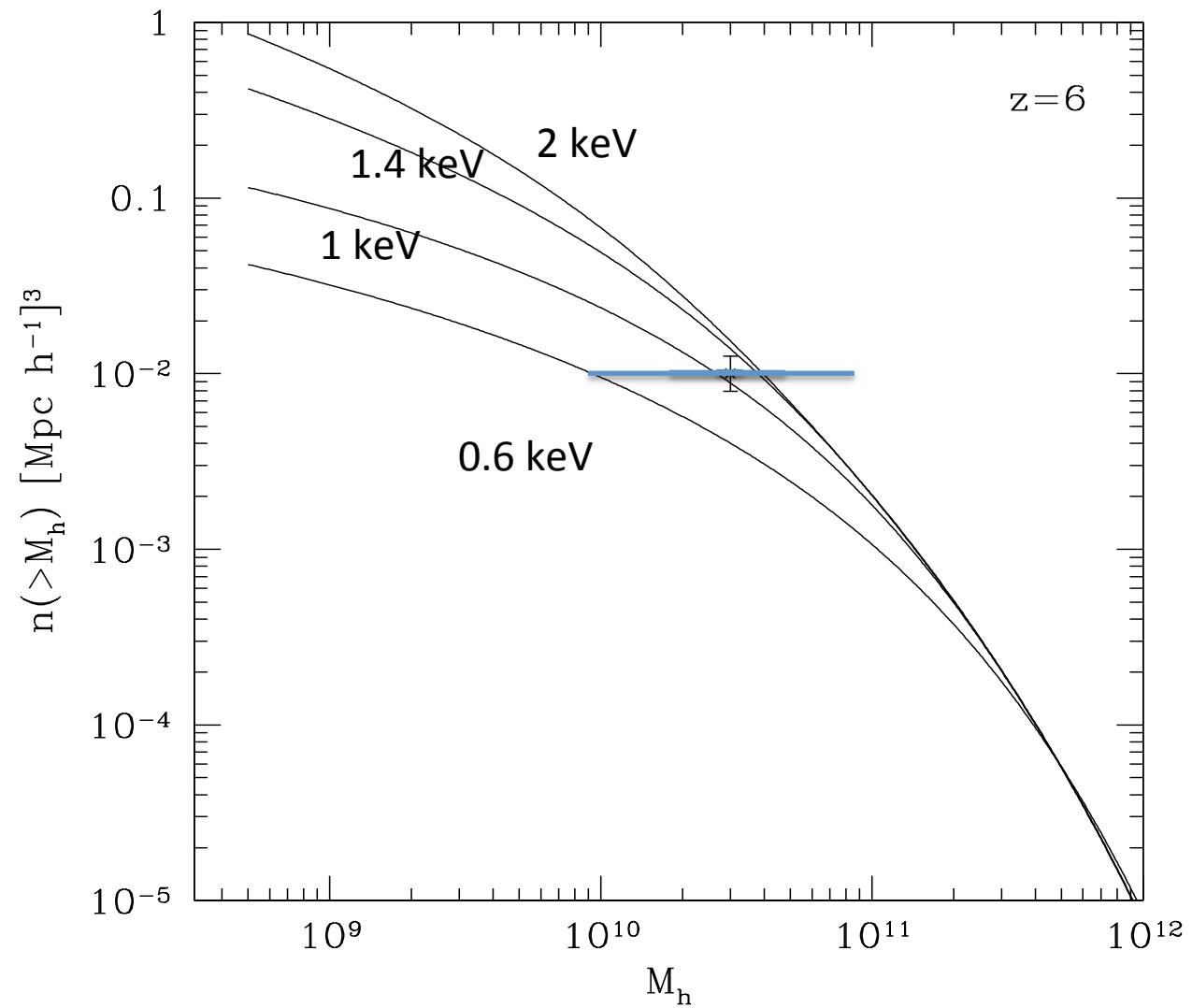


Schive et al. (2016)

# Cumulative Number Density

## How reliable?

- DM Halo mass is highly uncertain



# Abundance Matching

## Empirical Approach:

$$\int_M d\tilde{M} \frac{dn}{d\tilde{M}} = \int_{M_{UV}} d\tilde{M}_{UV} \phi(\tilde{M}_{UV})$$

- Rely on approximate analytical form of HMF
- Extrapolate LCDM relation to NDM models
- Redshift evolution may be driven by extinction

## What can we learn?

- **About galaxy formation in NDM models?**
- **What high-z can tell about NDM under minimal assumptions?**

# HAM & SFR

## UV-Luminosity to SFR

- Account for dust extinction
- Convert  $M_{UV}$  corrected to SFR (Kennicut relation)
- Derive SFR-density functions (see Mashian, Oesch, Loeb 2015 for LCDM)

## Extinction Correction

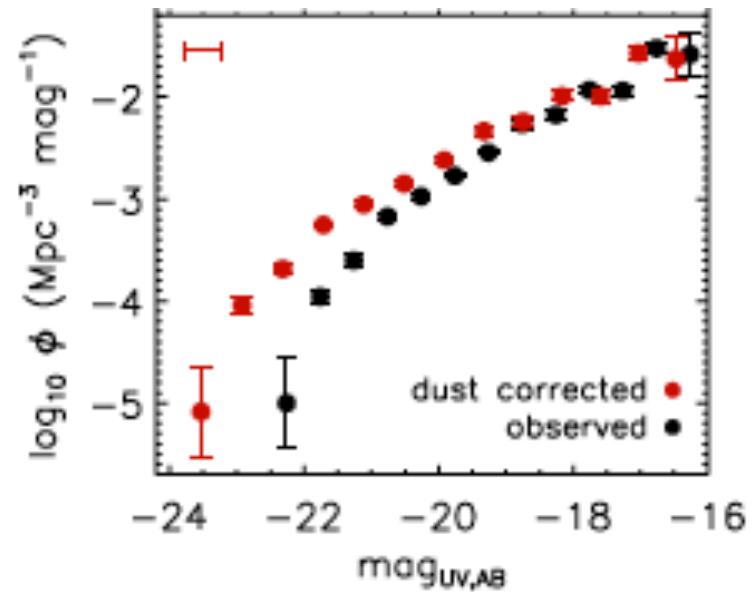
$$\langle A_{UV} \rangle = 4.43 + 0.79 \ln 10 \sigma_{\beta}^2 + 1.99 \langle \beta \rangle$$

Meurer et al. (1999)

$$\langle \beta(M_{UV}, z) \rangle = \begin{cases} [\beta_{M_0}(z) - C] e^{\frac{\beta'(z)(M_{UV} - M_0)}{\beta_{M_0}(z) - C}} & M_{UV} \geq M_0 \\ \beta'(z)[M_{UV} - M_0] + \beta_{M_0}(z) & \end{cases}$$

Tacchella et al. (2013), Mason, Trenti & Treu (2015)

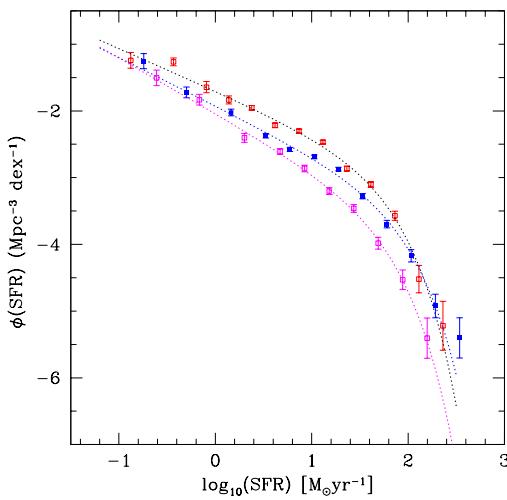
- Changes UV-mag bin size
- Shift toward higher luminosities



Smit et al. (2012)

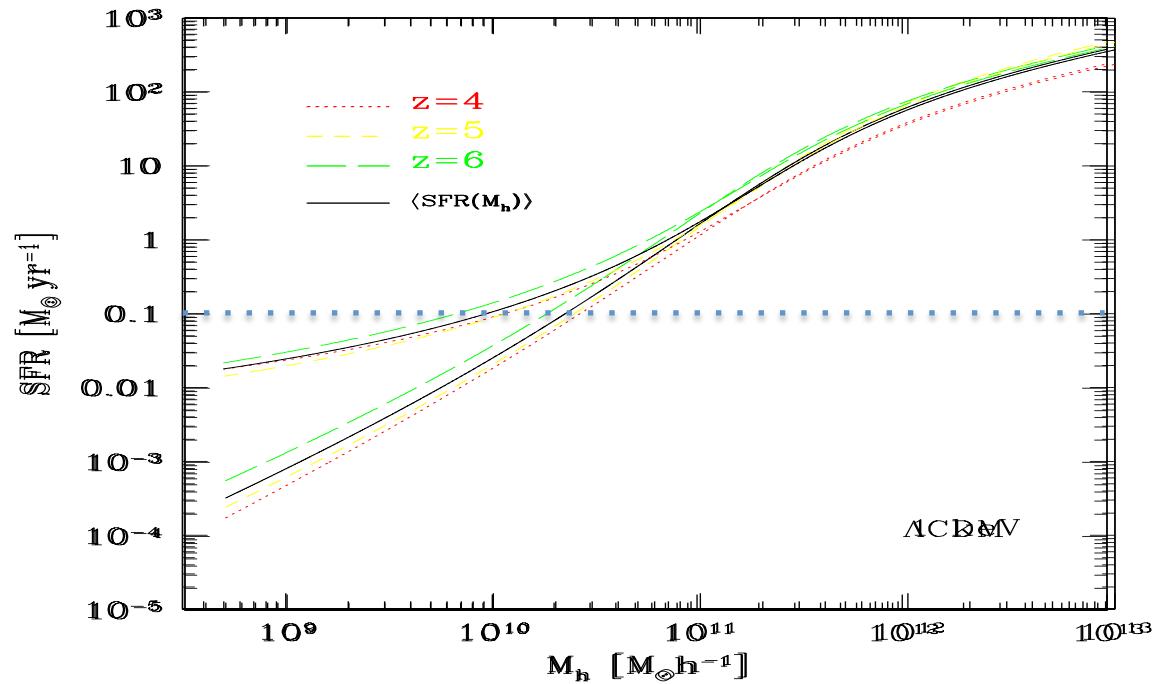
# SFR density function

## $\Phi(\text{SFR})$ at $4 < z < 6$



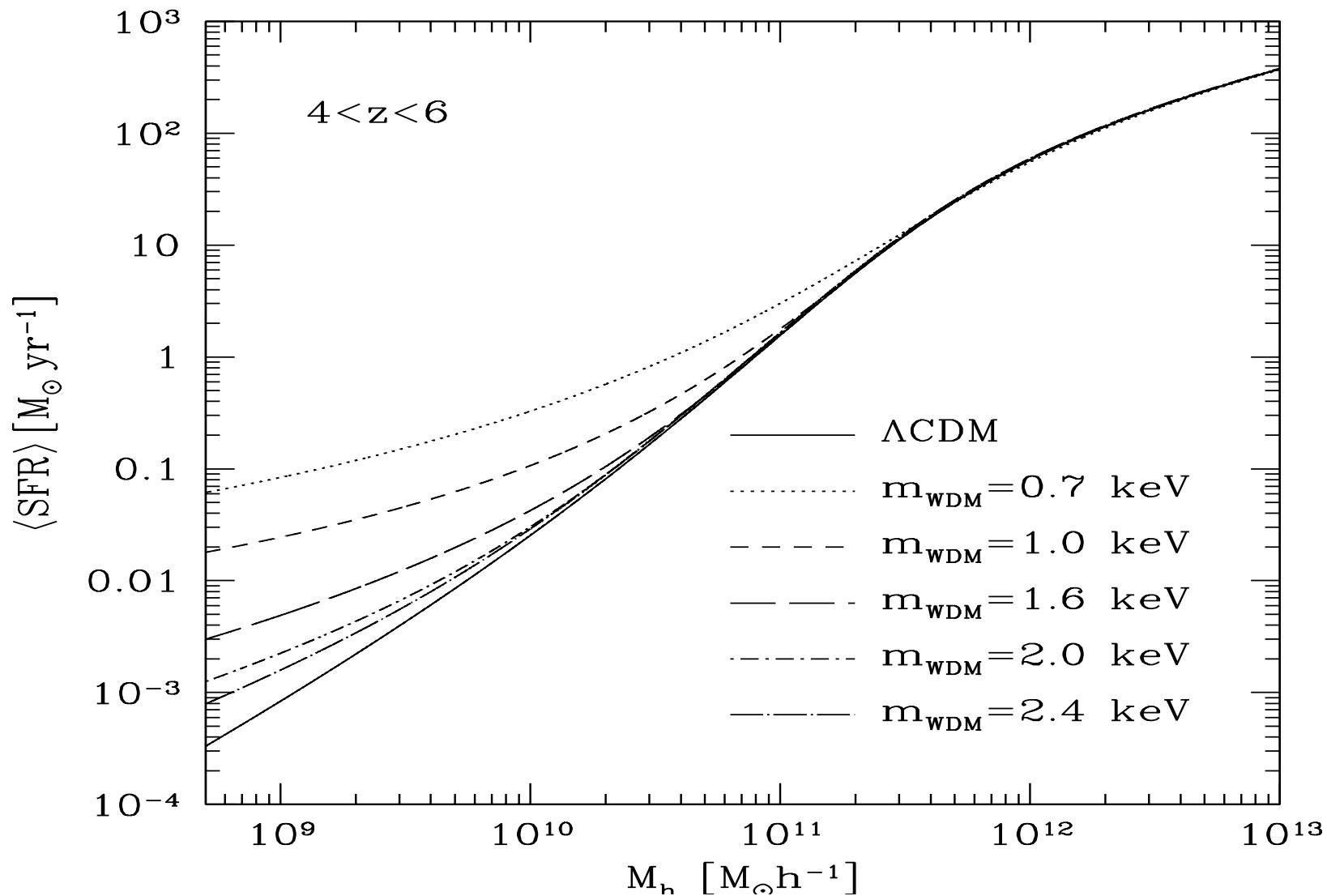
- Fit  $\Phi(\text{SFR})$  with Press-Schechter function
- HAM

$$\int_M d\tilde{M}_h \left( \frac{dn}{d\tilde{M}_h} \right)_{NDM} = \int_{SFR} dSFR' \phi(SFR')$$



- SFR- $M_h$  evolves in amplitude
- Redshift average

# Average SFR-M<sub>h</sub>



# Modeling Luminosity Function at $z > 6$

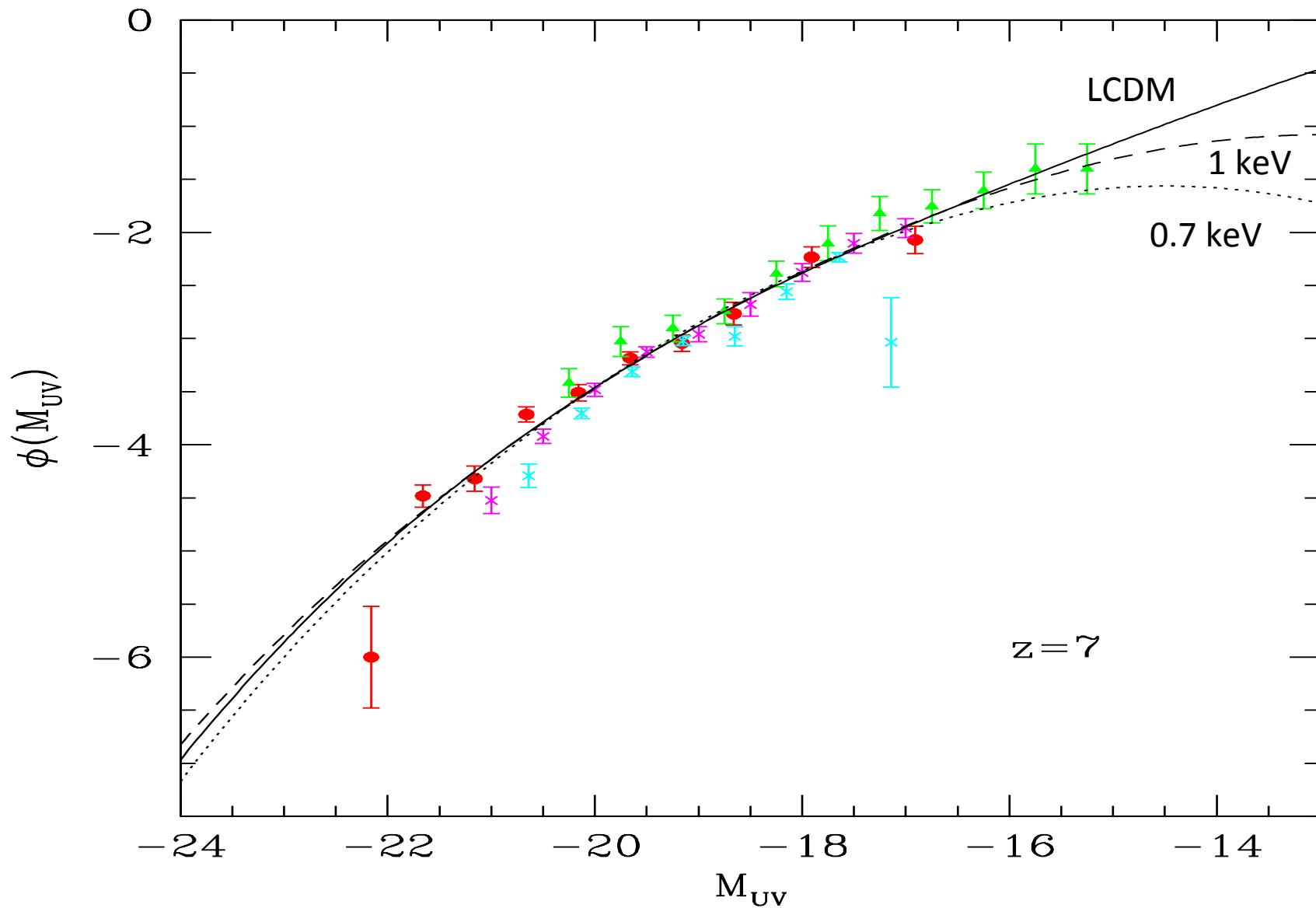
## Intrinsic Scatter SFR- $M_h$ relation

- Compute

$$\phi(SFR, z) = \frac{1}{\sigma_{\text{int}}^2 \sqrt{2\pi} SFR} \int dM_h \frac{dn}{dM_h}(M_h, z) e^{-\frac{\log_2^2[SFR/(\varepsilon \cdot \langle SFR(M_h) \rangle)]}{2\sigma_{\text{int}}^2}}$$

- Convert to UV luminosities
- Add extinction effect
- Estimate  $\Phi(M_{\text{UV}})$
- Fit against the data  $\varepsilon, \sigma_{\text{int}}$

# Fitting Galaxy Luminosity Function at $z > 6$



# Conclusions

- Small scale clustering of matter as probe of non-standard DM models (*free your mind*)
- Artificial fragmentation in simulation of models with suppressed spectrum at small scales needs to be accounted for (*be inquisitive*)
- Galaxy formation cannot occur in the same way in NCDM models (*question your beliefs*)
- Testing SFR halo mass relation can provide key insights (*truth arises more from error than confusion*)