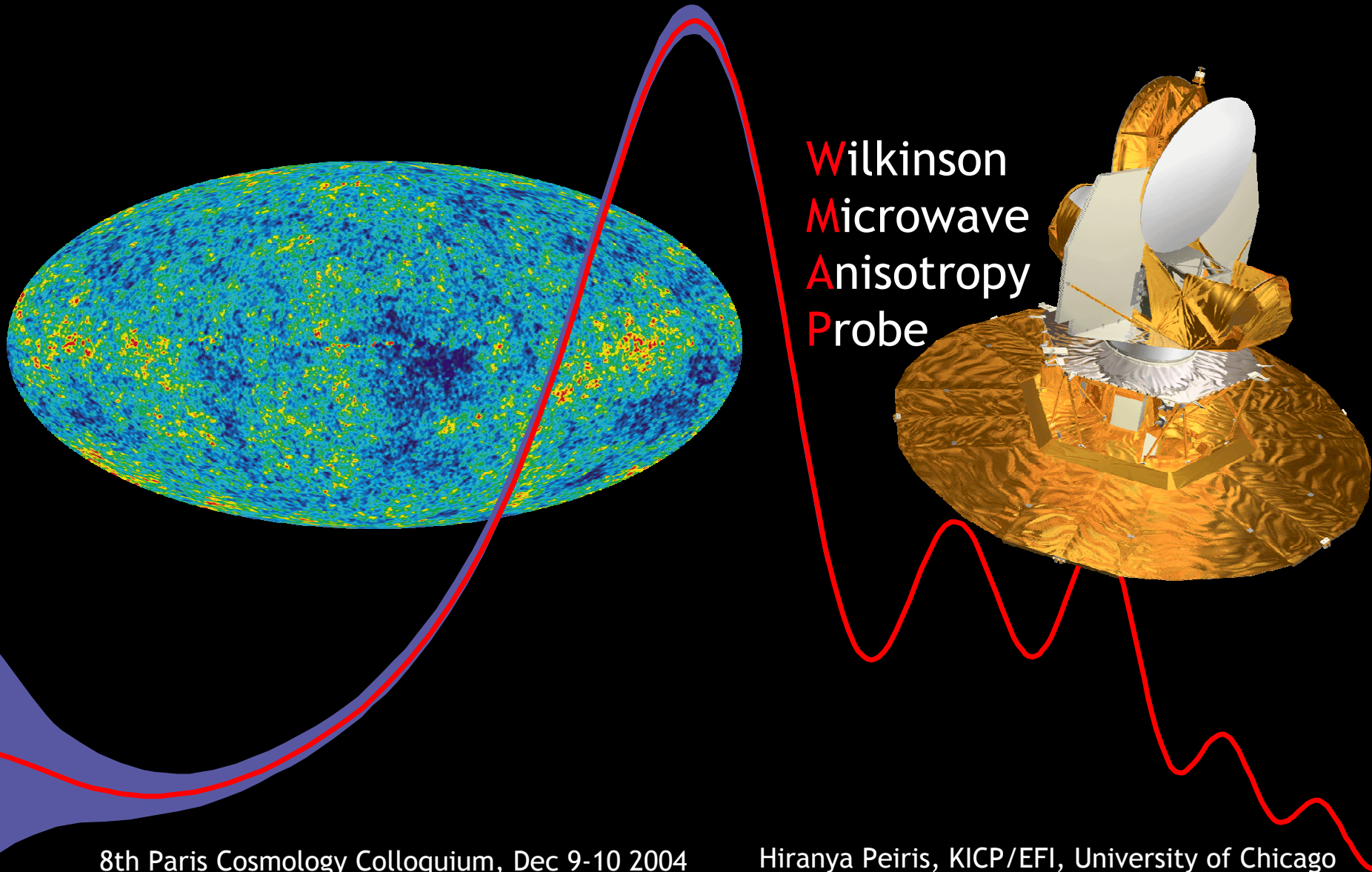


# WMAP and Early Universe Cosmology



# WMAP Science Team

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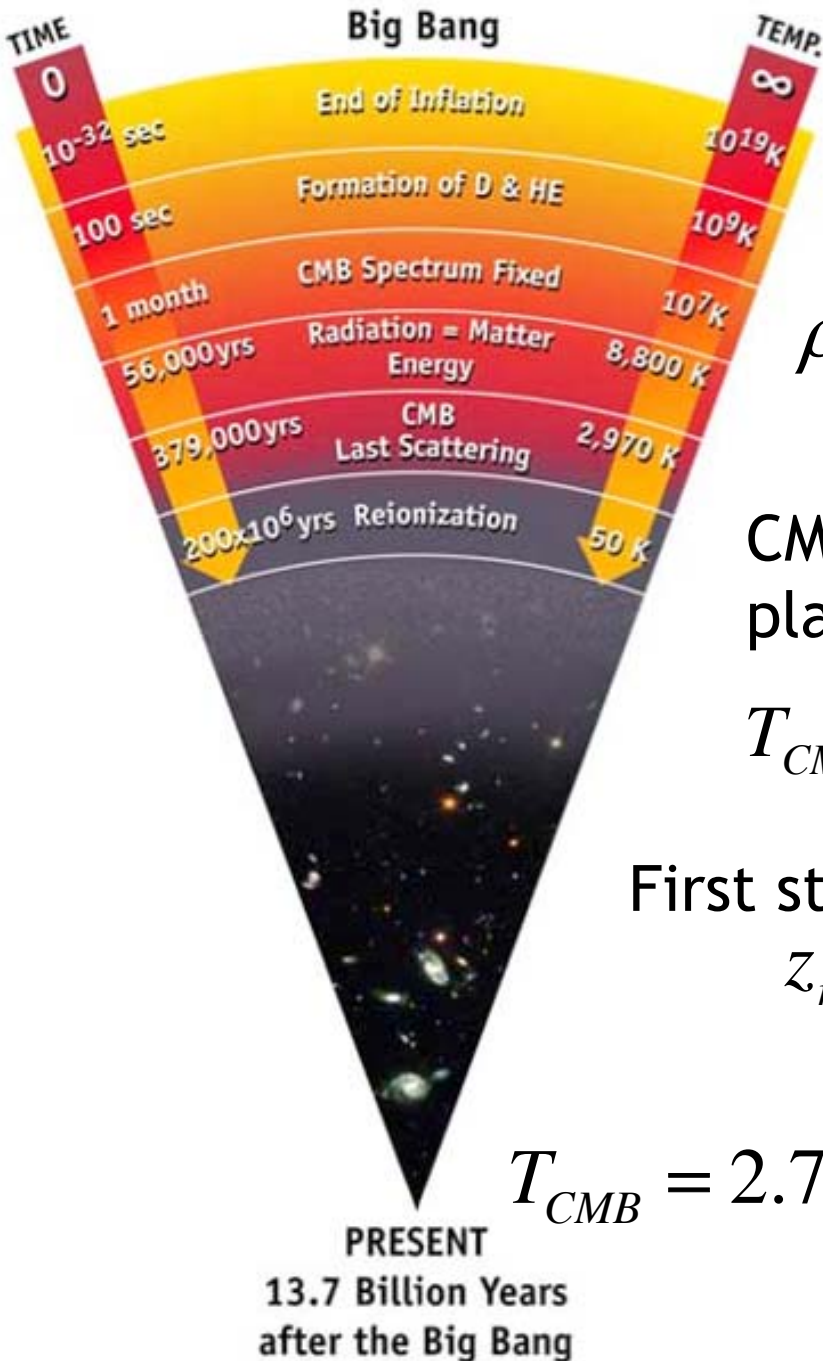
Michael Nolta

# WMAP Highlights

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- Microwave Maps
    - High fidelity full-sky maps in five bands (23,33,41,61,94 GHz).
    - High S/N CMB maps.
    - **Gaussianity of CMB.**
  - CMB Temperature
    - Power spectrum measured to  $l \sim 700$ .
    - **$\Lambda$ CDM fits.**
  - CMB Polarization
    - Temp-Pol cross-power spectrum measured to  $l \sim 400$ .
    - **Early reionization.**
- A lot more...
- Galaxy model
  - Point source catalogue
  - etc...
- Please visit  
<http://lambda.gsfc.nasa.gov>

# Cosmic History



Inflation-like epoch. **new**

$\rho_{\text{radiation}} = \rho_{\text{matter}}$  **new**

$z_{\text{eq}} = 3230$        $t_U = 56 \text{ kyr}$

CMB decouples from plasma **new**

$z_{\text{dec}} = 1089$        $t_U = 380 \text{ kyr}$

$T_{\text{CMB}} = 2970 \text{ K}$

First stars form **new**

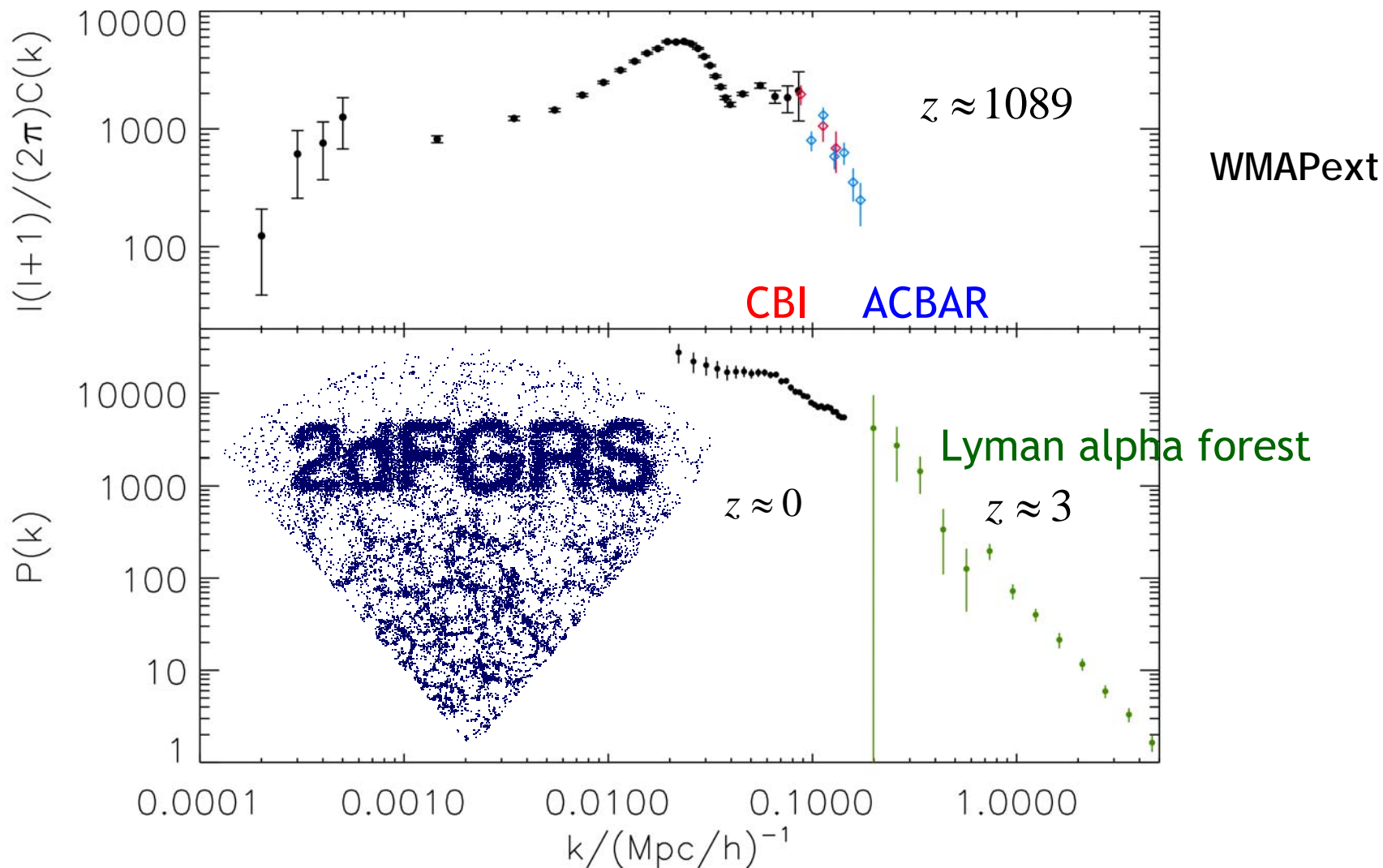
$z_r \sim 17$        $t_U = 200 \text{ Myr}$

$T_{\text{CMB}} = 2.725 \text{ K}$  **new**

$z = 0$        $t_U = 13.7 \text{ Gyr}$

PRESENT  
13.7 Billion Years  
after the Big Bang

# TEST MODEL CONSISTENCY and LIFT DEGENERACIES



Complementary in scales-redshift

# What Simple Inflationary Models predict

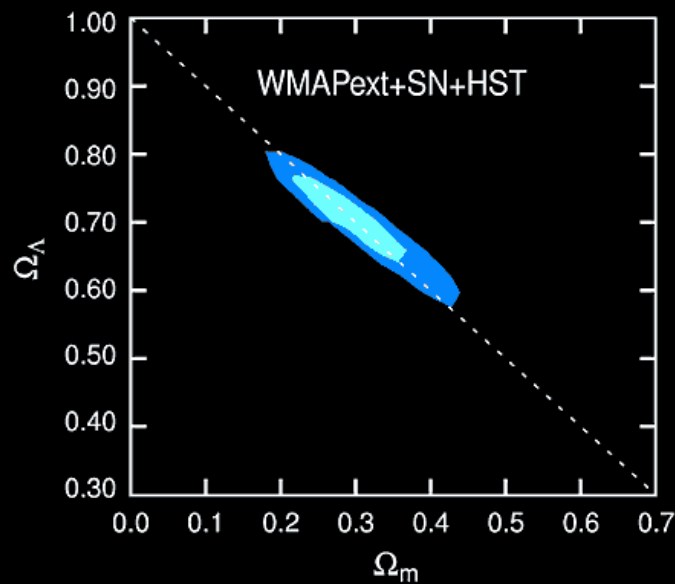
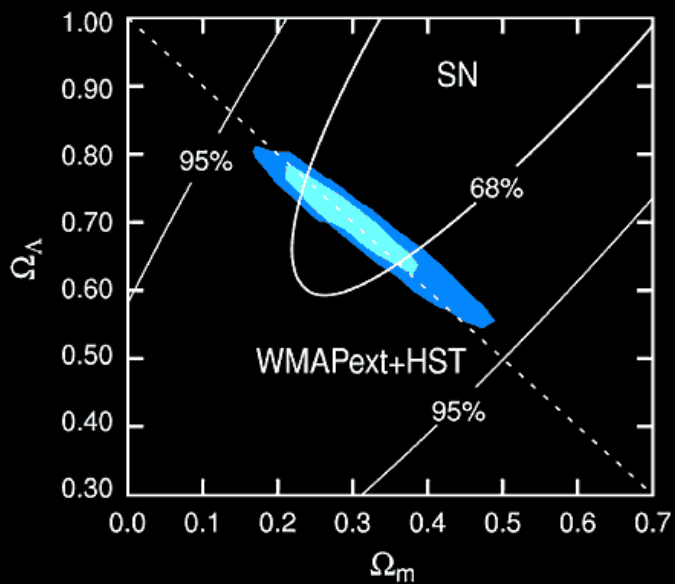
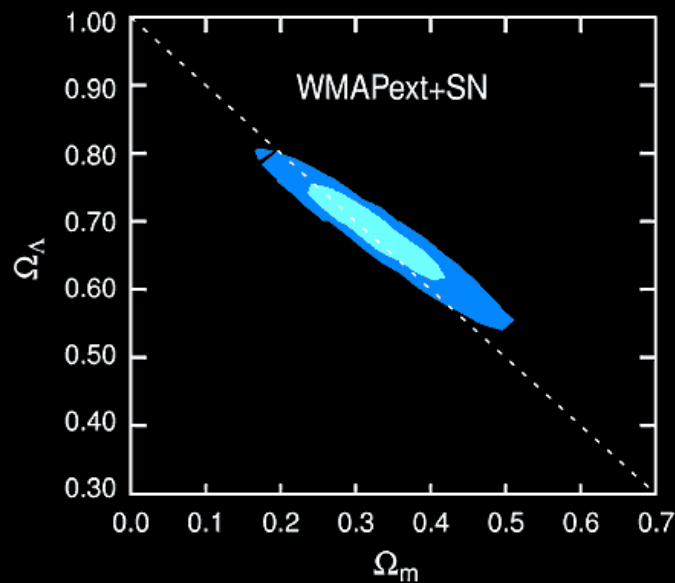
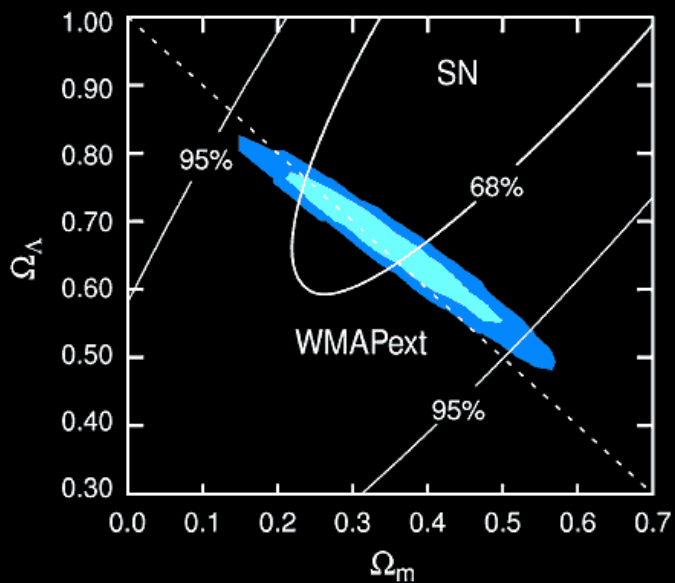
---

- Flat universe:  $\Omega_{\text{tot}} = 1$
- Gaussianity:  $f_{NL} \sim 1$   $\Phi(\vec{x}) = \Phi_{\text{gaus}}(\vec{x}) + f_{NL} \Phi_{\text{gaus}}^2(\vec{x})$   
See talk by S. Matarrese
- Power Spectrum spectral index nearly scale-invariant:  
 $n_s \sim 1$   $k^3 \langle \Phi \Phi \rangle \propto k^{n_s - 1}$
- Adiabatic superhorizon fluctuations

Guth (1981), Linde (1982), Albrecht & Steinhardt (1982), Sato (1981), Mukhanov & Chibisov (1981), Hawking (1982), Guth & Pi (1982), Starobinsky (1982), Bardeen et al. (1983), Mukhanov et al. (1992)

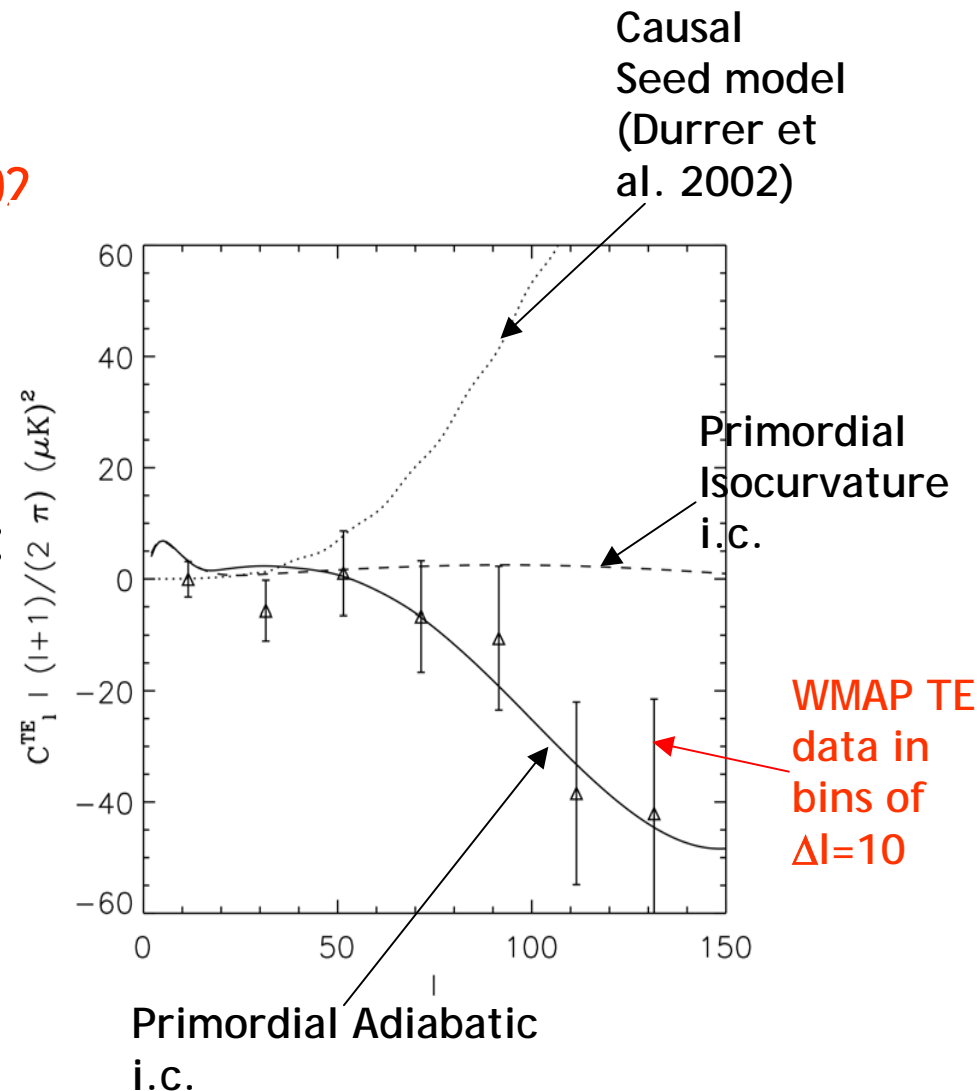
# $\Omega_\Lambda$ vs. $\Omega_m$

See Talk by O. Dore



# WMAP Supports Single Field Inflationary Models

- Flat universe:  $\Omega_{\text{tot}} = 1.02 \pm 0.07$
- Gaussianity:  $-58 < f_{NL} < 134$
- Power Spectrum spectral index nearly scale-invariant:  $n_s = 0.99 \pm 0.04$  (WMAP only)
- Adiabatic initial conditions
- Superhorizon fluctuations (TE anticorrelations)





# Detailed level of confrontation

---

- The primordial power spectrum is not a perfect power law.

$$n_s(k) - 1 = n_s(k_0) - 1 + \frac{dn_s}{d \ln k} \ln \left( \frac{k}{k_0} \right)$$

- There could be gravitational waves.

$$r \equiv \text{tensor-to-scalar ratio} = \frac{\langle h_{ij} h^{ij} \rangle(k_0)}{\langle \mathbf{R R} \rangle(k_0)}$$

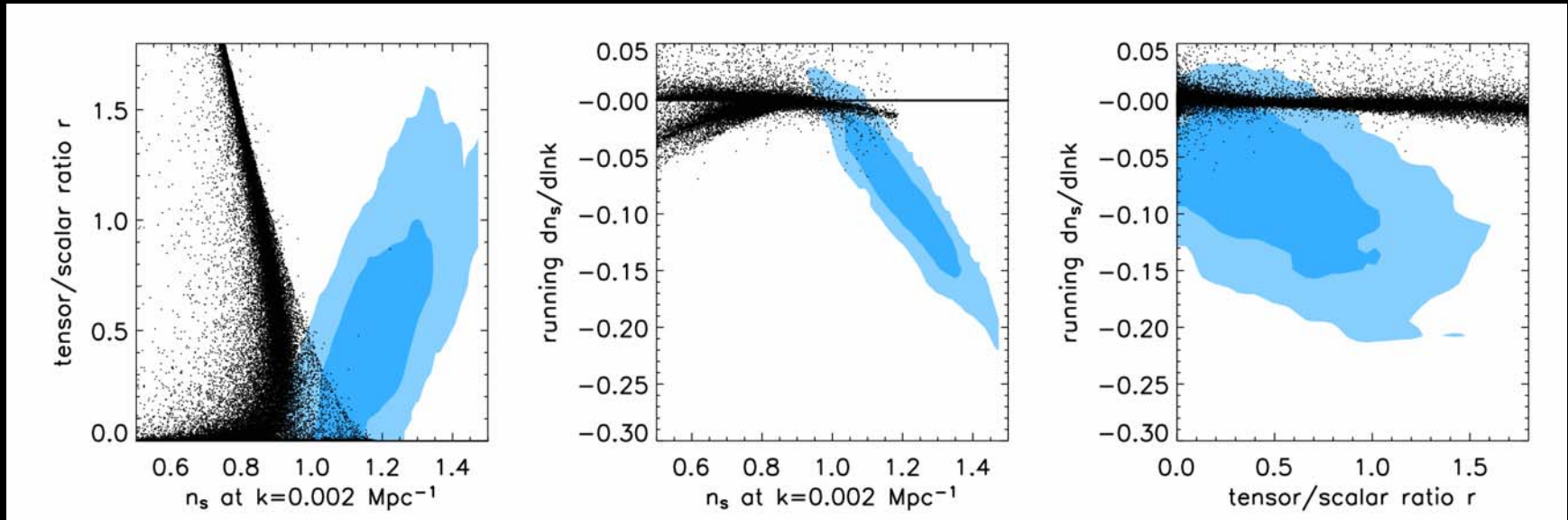
“running”

(The shape of the tensor power spectrum is determined by  $n_s$  and  $r$  using predictions of single field inflationary models.)

We use  $k_0 = 0.002 \text{ Mpc}^{-1}$  ( $l \sim 30$ )

# Generic predictions of single field slow roll models

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Each point is a “viable” slow roll model, able to sustain inflation for sufficient e-foldings to solve cosmological problems.

Monte Carlo simulations following Kinney (2002) and Easter and Kinney (2002)

# Classify Inflationary Models

---


- The shape of the scalar field potential,  $V(\phi)$ , determines the observables.
- We use three parameters to characterize the shape:
  - $\epsilon$ : “slope” of potential,  $(V'/V)^2$
  - $\eta$ : “curvature” of potential,  $V''/V$
  - $\xi$ : “jerk” of potential,  $(V'/V)(V'''/V)$

The curvature is the most important parameter in classifying the models.

# Classify Inflationary Models

---

$$\begin{cases} r = 16\varepsilon \\ n_s = 1 - 6\varepsilon + 2\eta \\ dn_s / d \ln k = -2\xi + 16\varepsilon\eta - 24\varepsilon^2 \end{cases}$$


$$\begin{cases} r = \frac{8}{3}(1 - n_s) + \frac{16}{3}\eta \\ dn_s / d \ln k + 2\xi = -\frac{2}{3} \left[ (1 - n_s)^2 - 4\eta^2 \right] \end{cases}$$

The curvature is the most important parameter in classifying the models.

# Negative curvature models

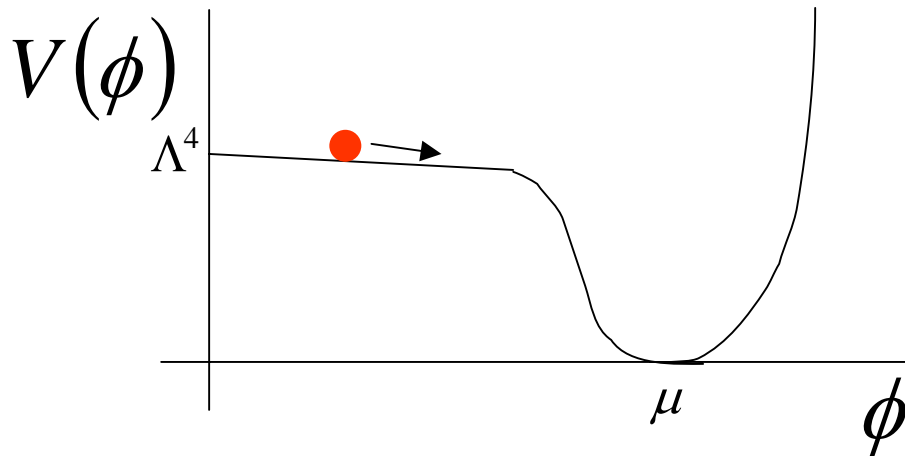
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$$\eta < 0$$

$$V(\phi) = \Lambda^4 \left[ 1 - (\phi/\mu)^p \right]$$

spontaneous  
symmetry  
breaking  
potential

e.g. new  
inflation



Albrecht & Steinhardt  
(1982), Linde (1982)

$n_s < 1$  (red tilt) & tiny  $r$  & tiny  $dn_s/d\ln k$

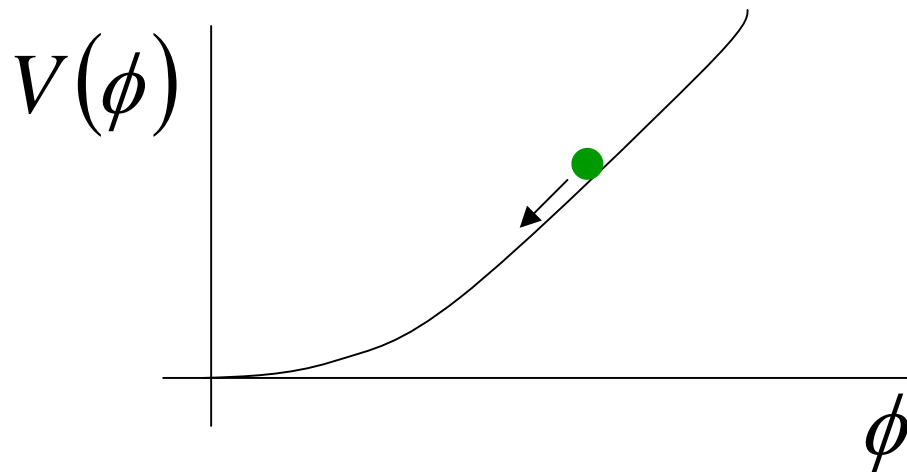
# Small positive curvature models

---

$$0 \leq \eta < 2\varepsilon$$

$$V(\phi) = \Lambda^4 (\phi/\mu)^p$$

e.g. chaotic inflation,  
extended inflation



Linde (1983), La &  
Steinhardt (1989)

$n_s < 1$  (red tilt) & large  $r$  & tiny  $dn_s/d\ln k$

# Large positive curvature models

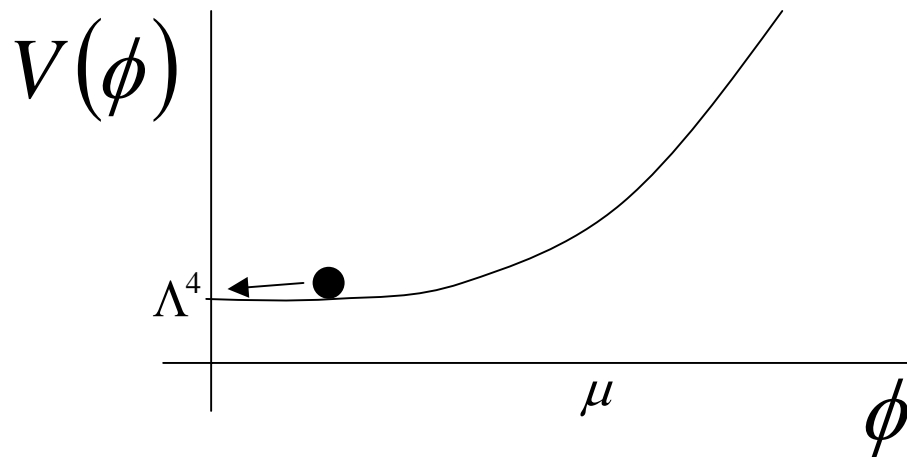
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$$\eta > 3\varepsilon$$

$$V(\phi) = \Lambda^4 \left[ 1 + (\phi/\mu)^p \right]$$

e.g. hybrid inflation

Linde (1994)



There are models in this class which give large  $dn_s/d\ln k$ .

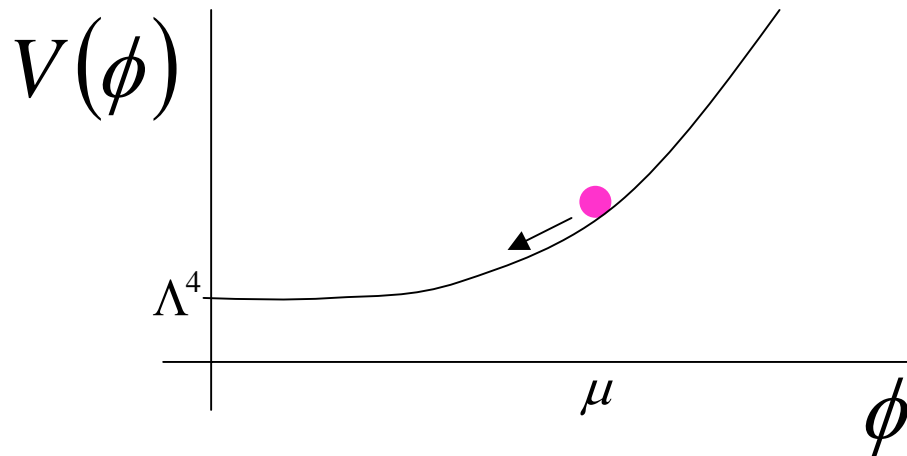
$n_s > 1$  (blue tilt) & tiny  $r$  & tiny  $dn_s/d\ln k$

# Intermediate positive curvature models

---

$$2\varepsilon \leq \eta \leq 3\varepsilon$$

$$V(\phi) = \Lambda^4 \left[ 1 + (\phi/\mu)^p \right]$$



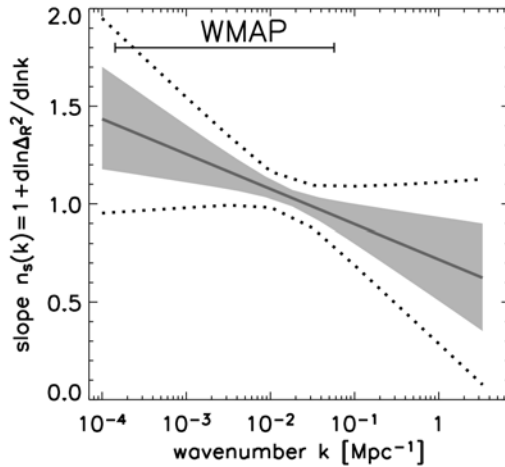
There are models in this class which give large  $dn_s/d\ln k$ .

$n_s < 1$  (red tilt) & large  $r$  & tiny  $dn_s/d\ln k$

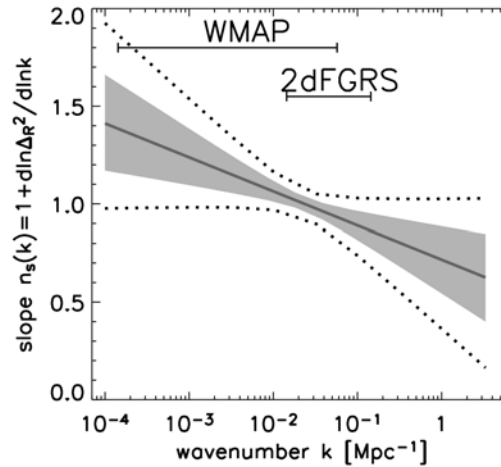


# A hint of something unusual

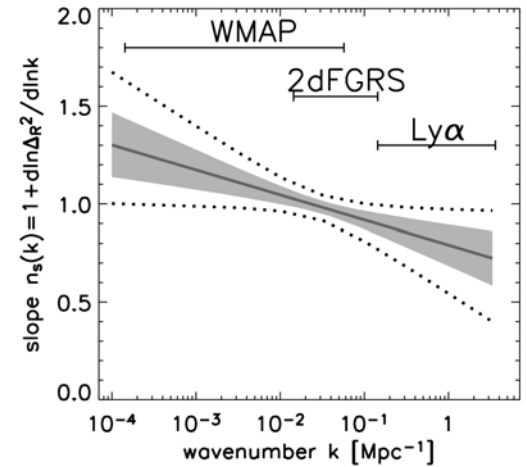
- Many inflationary models favour minimal “running”.



91% runs blue-to-red



95% runs blue-to-red



96% runs blue-to-red

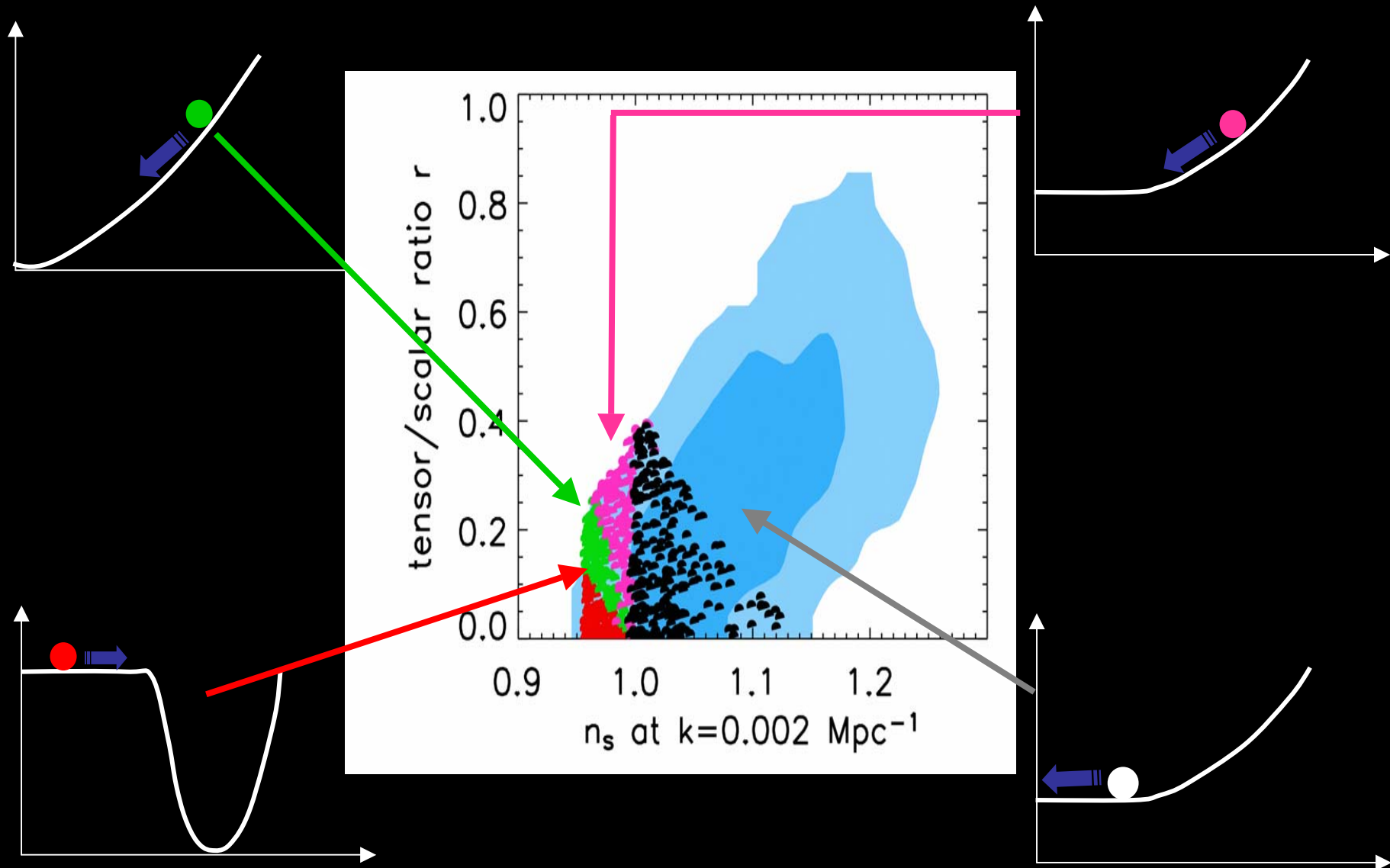
$$n_s = 1.13 \pm 0.08 \quad k = 0.002 \text{ Mpc}^{-1}$$

$$\frac{dn_s}{d \ln k} = -0.055 \pm 0.028$$

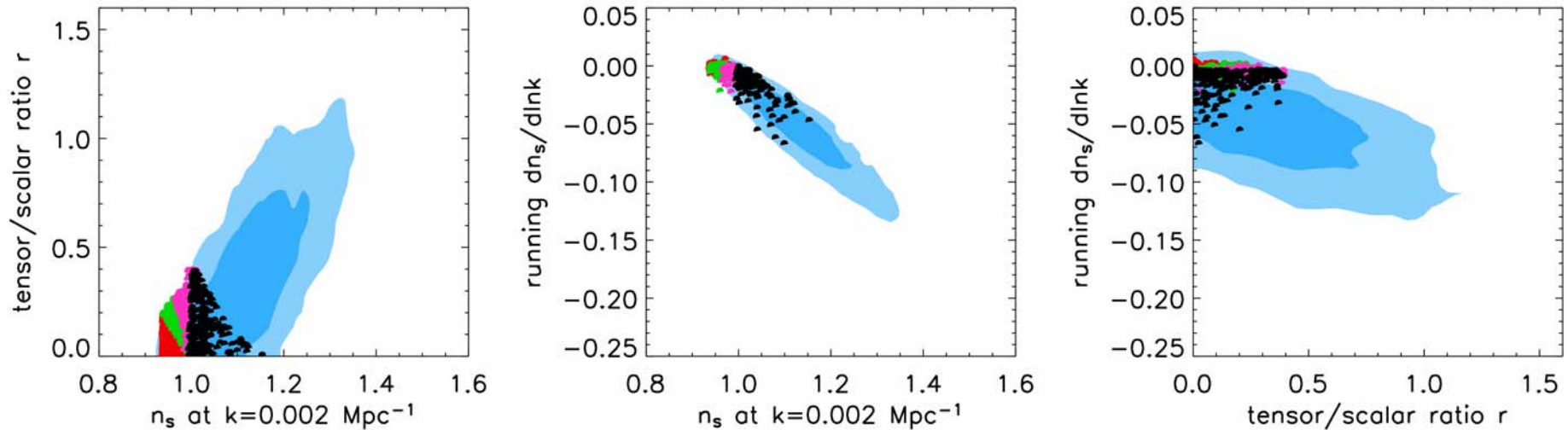
The data prefer, but do not require, a running spectral index.

If true, third derivative of potential is important.

# Testing inflation at a detailed level



# Categorizing single field slow roll models



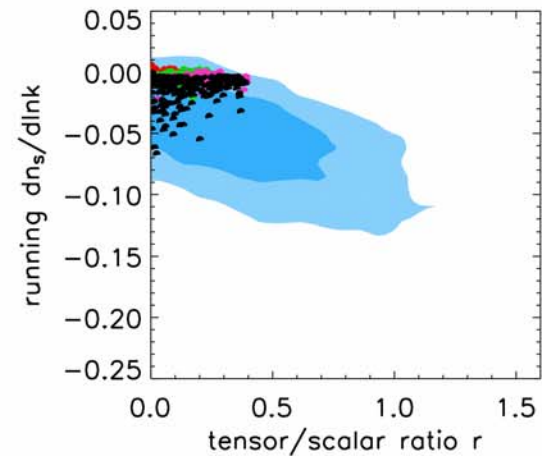
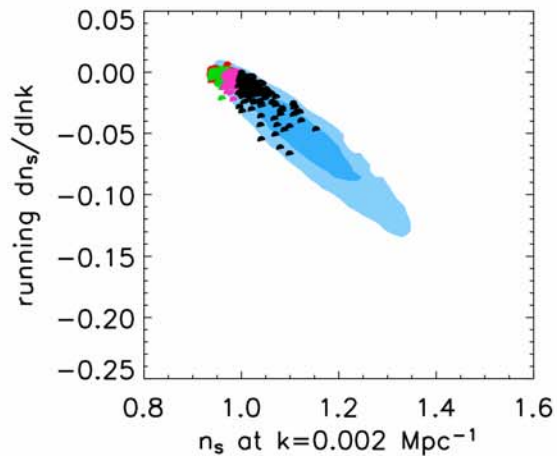
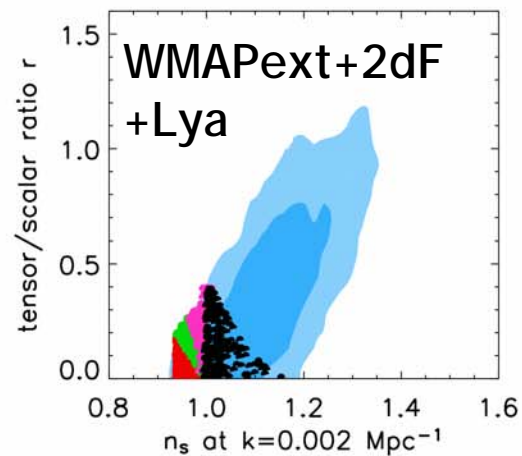
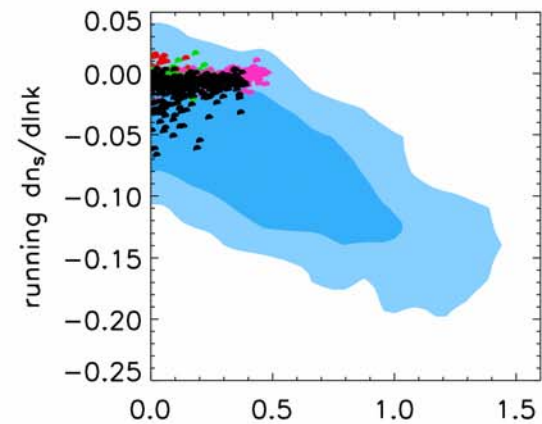
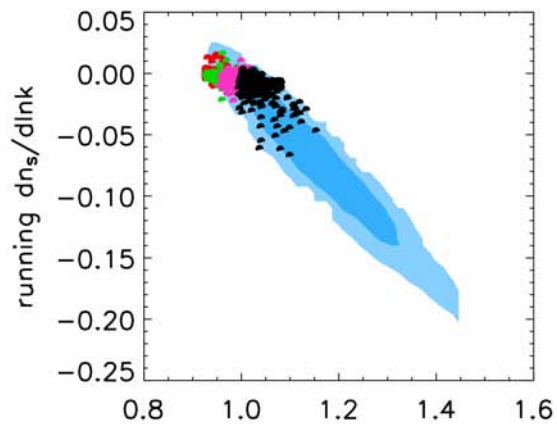
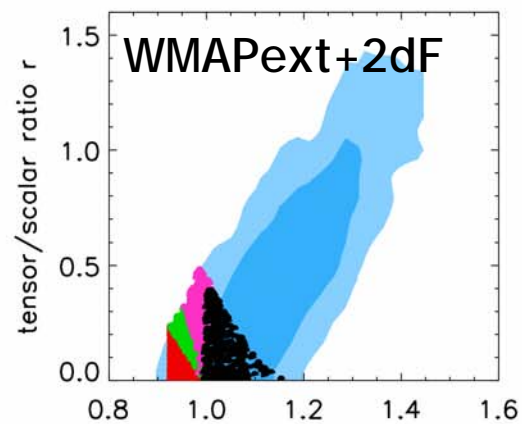
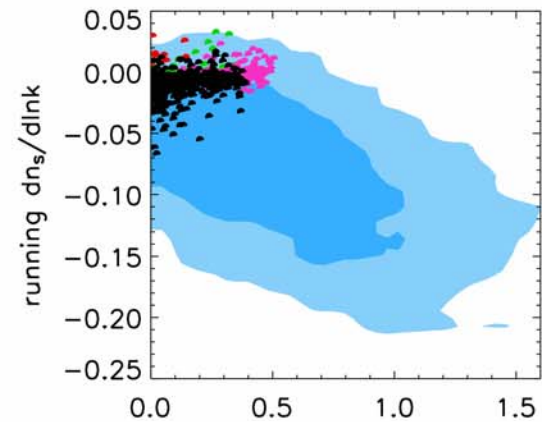
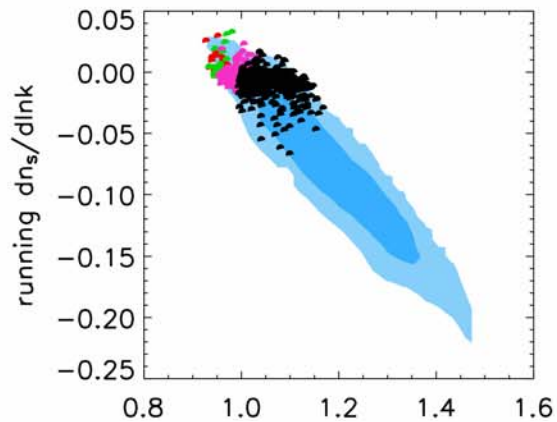
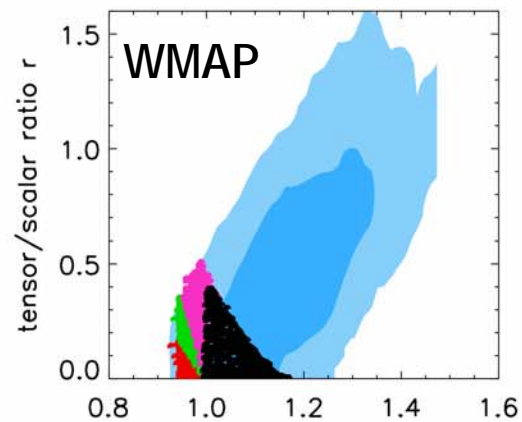
**Negative curvature** (e.g.: new inflation)

**Small positive curvature** (e.g.: chaotic inflation, extended inflation)

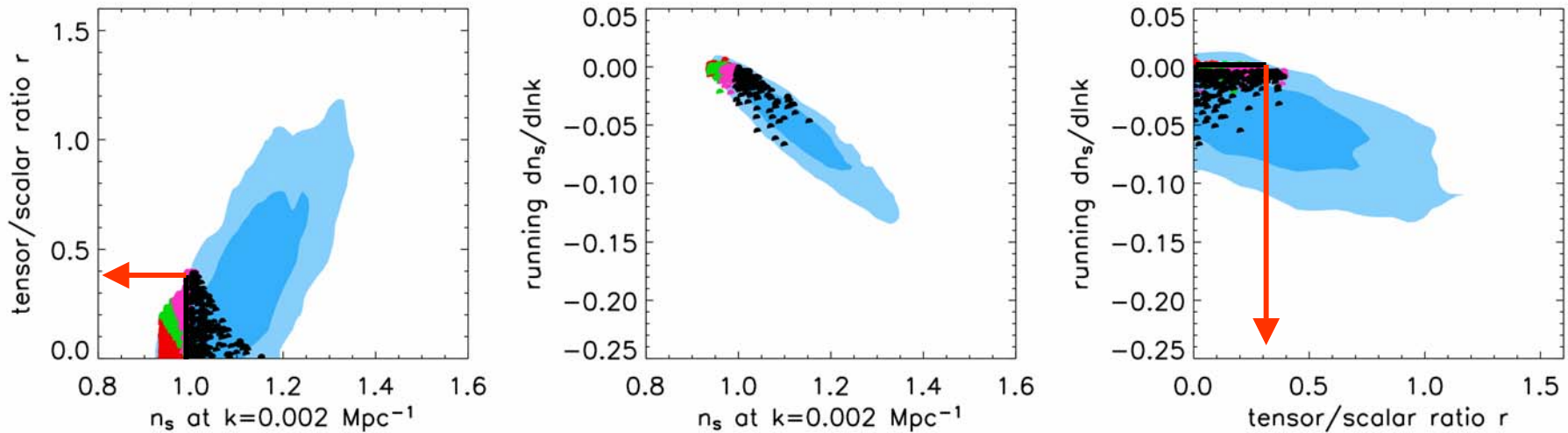
**Intermediate positive curvature**

**Large positive curvature** (e.g.: hybrid inflation)

Recommended: For given model, sit on that point and run likelihood analysis (may need to integrate mode equation directly).



# Constraining Inflation



## 95% Confidence Limits:

$r < 0.9$  (no priors)  $\Rightarrow$  Energy scale of inflation  
 $V^{1/4} < 3.3 \times 10^{16} \text{ GeV}$  (95% CL)

$r < 0.43$  (no running)

$r < 0.28$  (red tilt)

$\lambda\phi^4$  model,  
Excluded at  
more than  
3-sigma if  $N < 50$

# Double Field Models

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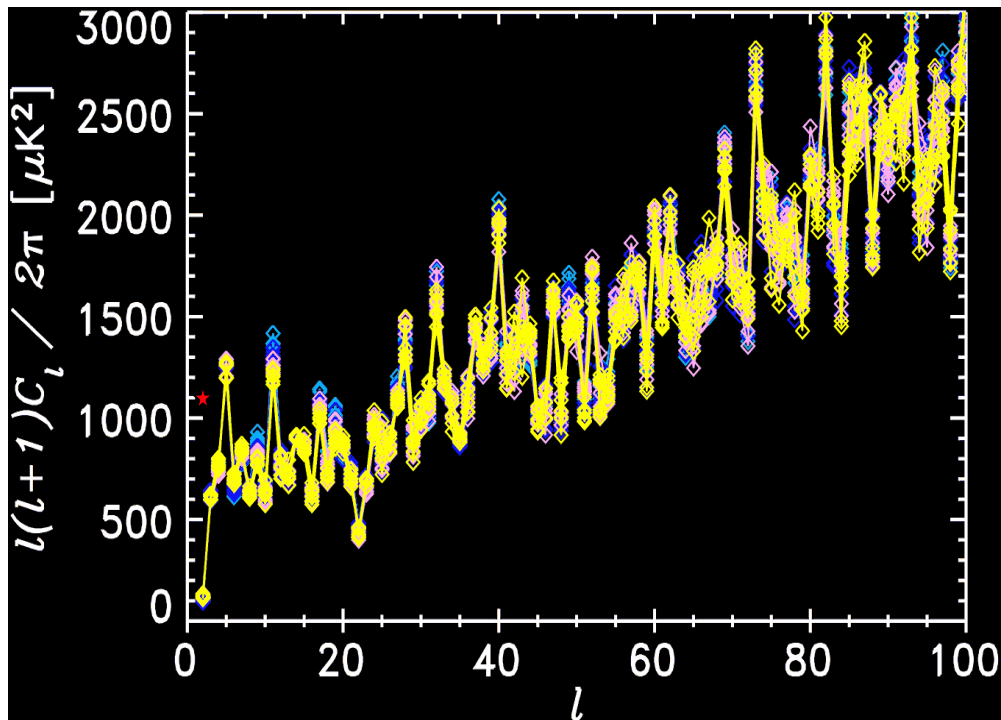
- If there is a scalar field, why not two?
- Distinctive signatures of double field models:
  - Fluctuations may not be purely adiabatic
  - Entropy (isocurvature) perturbations
- We perturb entropy between CDM and photons
  - Parameterize by two slopes ( $n_{\text{ad}}$ ,  $n_{\text{ent}}$ ), a fractional contribution of the entropy mode ( $f_{\text{ent}}$ ), and a correlation angle ( $\cos \Delta$ ).

See talk by J. Lesgourgues

# Double field models

- Motivation: Can we reduce low  $l$  anisotropy?

$$\frac{\Delta T}{T} = \frac{1}{5} \left( \hat{R}_{rad} - 2\hat{S}_{rad} \right)$$

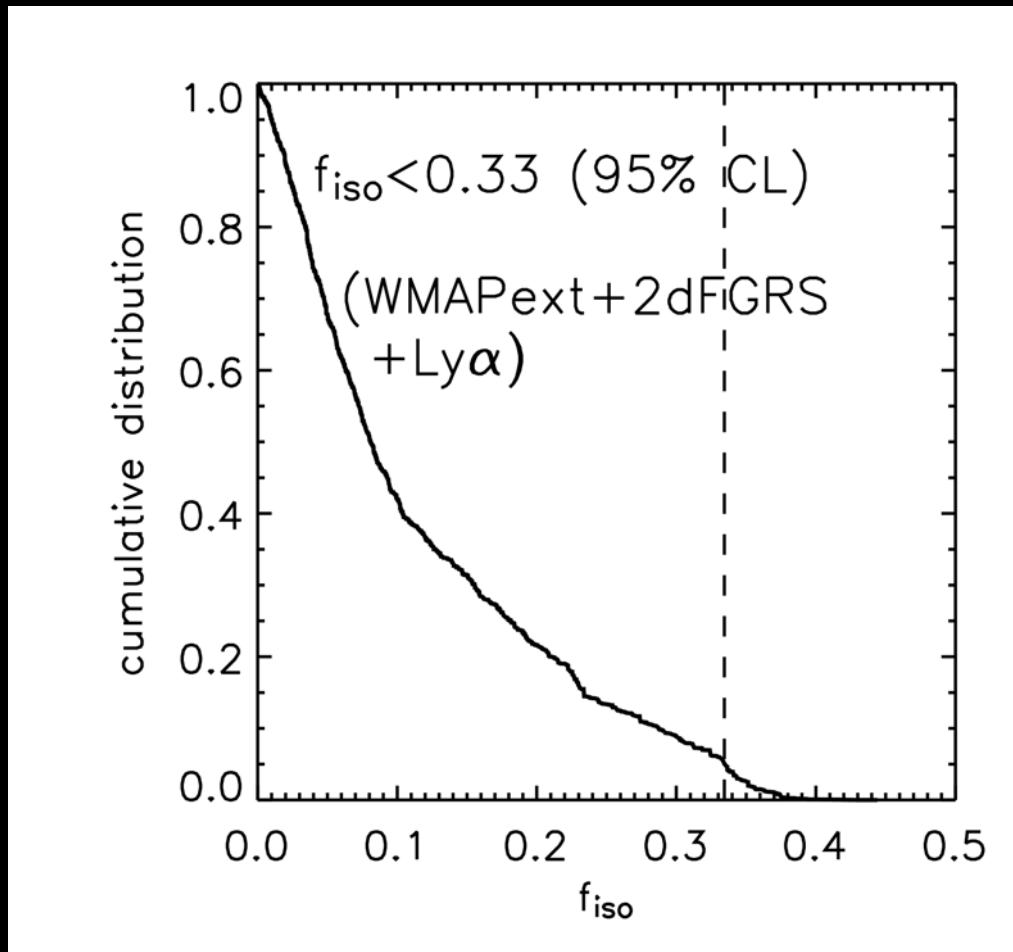


- CMB anisotropy can be **reduced**, when  $R$  and  $S$  are correlated.

See talk by J. Lesgourgues

# Constraint on correlated CDM isocurvature fraction

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Chi-square not improved by addition of three extra parameters

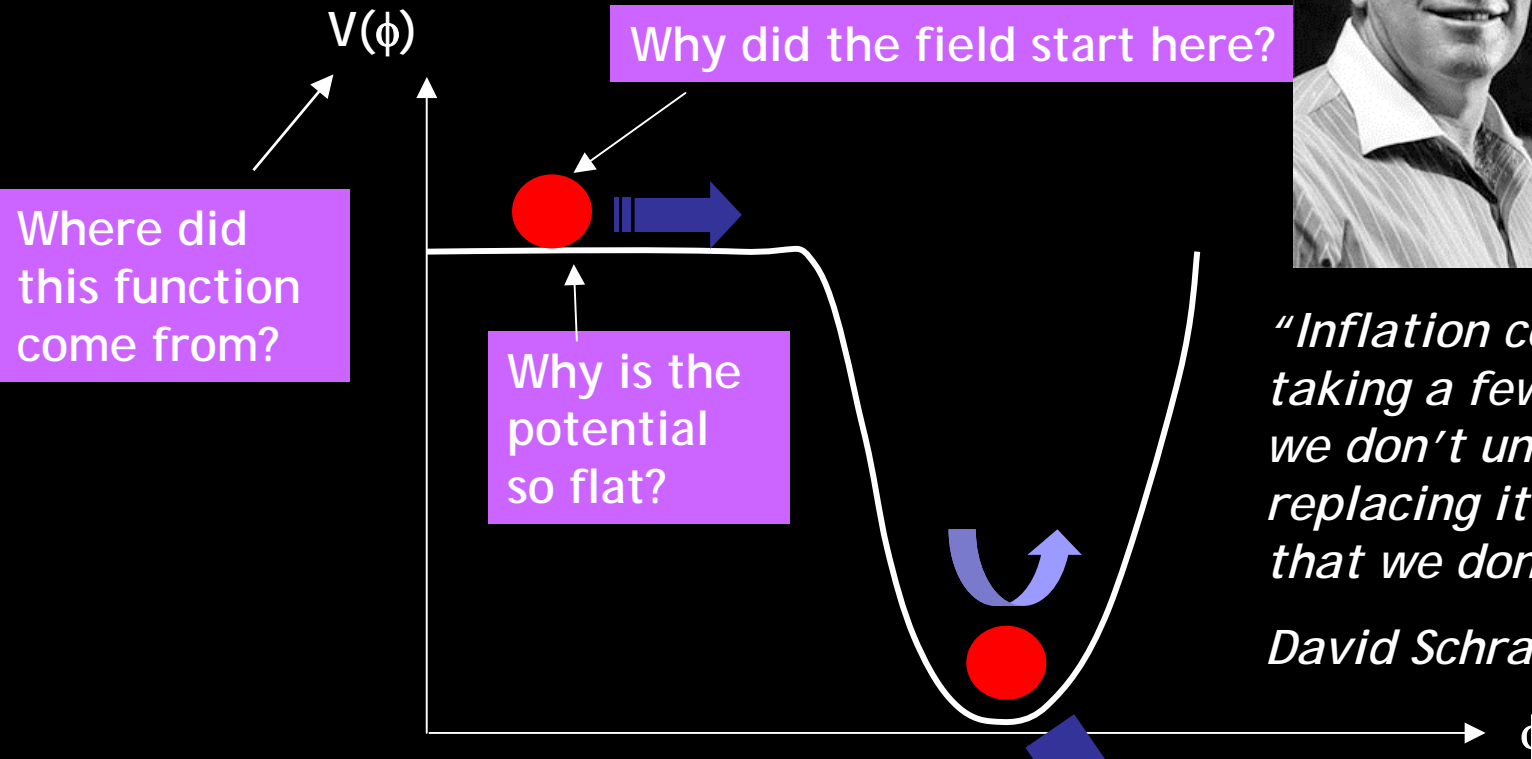


# What *WMAP* is telling us

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- No evidence for entropy perturbations between CDM and photons.
- Primordial fluctuations seem to be purely adiabatic.
- Supporting single field (simplest) inflationary models.
- No broad classes of inflation cannot currently be ruled out (because we cannot exclude  $n_s=1, r=0$ ), but specific models are starting to be able to be ruled out.

# Inflation: Problems



Where did this function come from?

Why did the field start here?

Why is the potential so flat?

How do we convert the field energy completely into particles?



*"Inflation consists of taking a few numbers that we don't understand and replacing it with a function that we don't understand"*

*David Schramm 1945 -1997*

# Is Inflation a Theory?

---

While the simplest versions of inflation have definite predictions (flat universe, Gaussian scale invariant spectrum of adiabatic fluctuations), inflationary models with more baroque forms of  $V(\phi)$  can produce non-flat universes, non-Gaussian fluctuations, non-adiabatic fluctuations, and deviations from scale invariance.

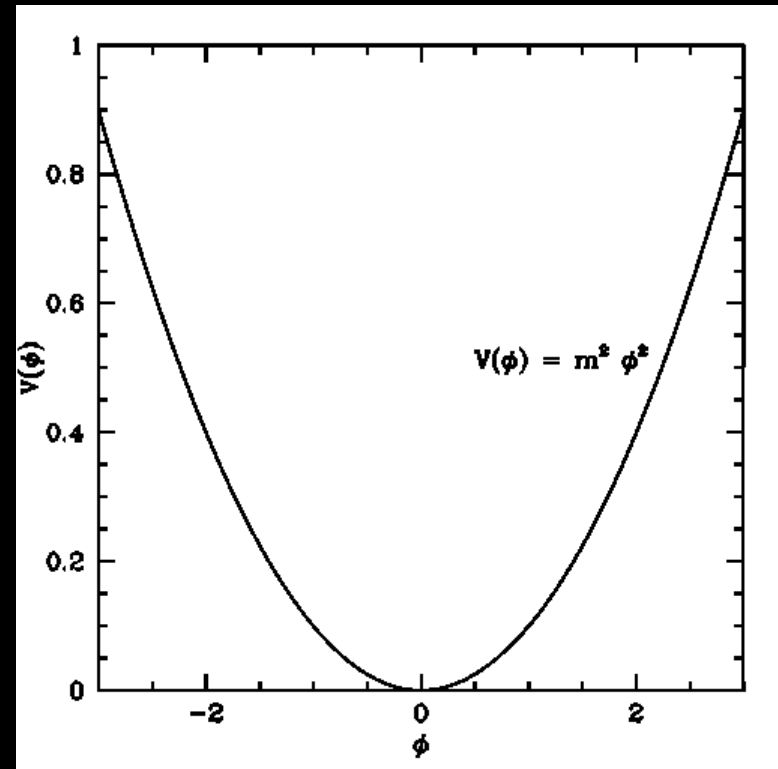
INFLATION IS A PARADIGM

# Testing Specific Inflationary Models

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- Cosmologists need to test simplest models
- Particle theorists need to motivate models
  - String theory
  - Multiple dimensional cosmologies

See talks by H. de Vega & N. Sanchez



# Testing the Simplest Models

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- Deviations from scale invariance
  - $m^2\phi^2$  inflation predicts  $n_s = 0.97$
  - WMAP + ACT (or Planck) should detect these deviation from scale invariance at greater than  $3 \sigma$
- Further tests of non-Gaussianity (See talk by S. Matarrese)
- Gravity waves from inflation
  - $m^2\phi^2$  inflation predicts a gravity wave background at a level detectable by WMAP (with  $\sim 8$  years of integration) and Planck

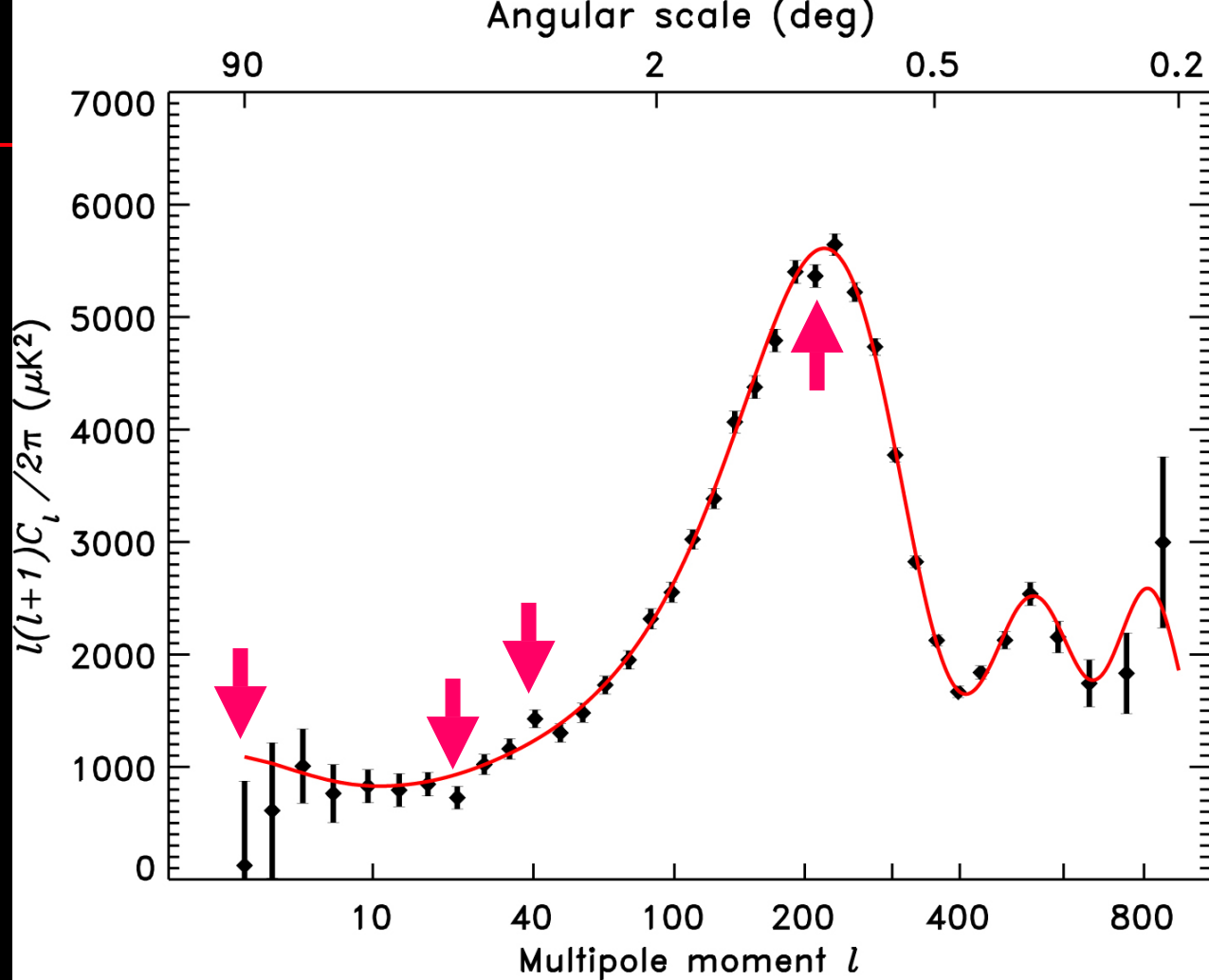
# Current Status of Inflation

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- Standard big bang models has a number of fundamental problems
- Inflationary scenario solves these problems by positing an inflaton potential
  - Inflaton potential form is currently ad-hoc

*Simplest Versions of Inflation are Testable!*

# Glitches?

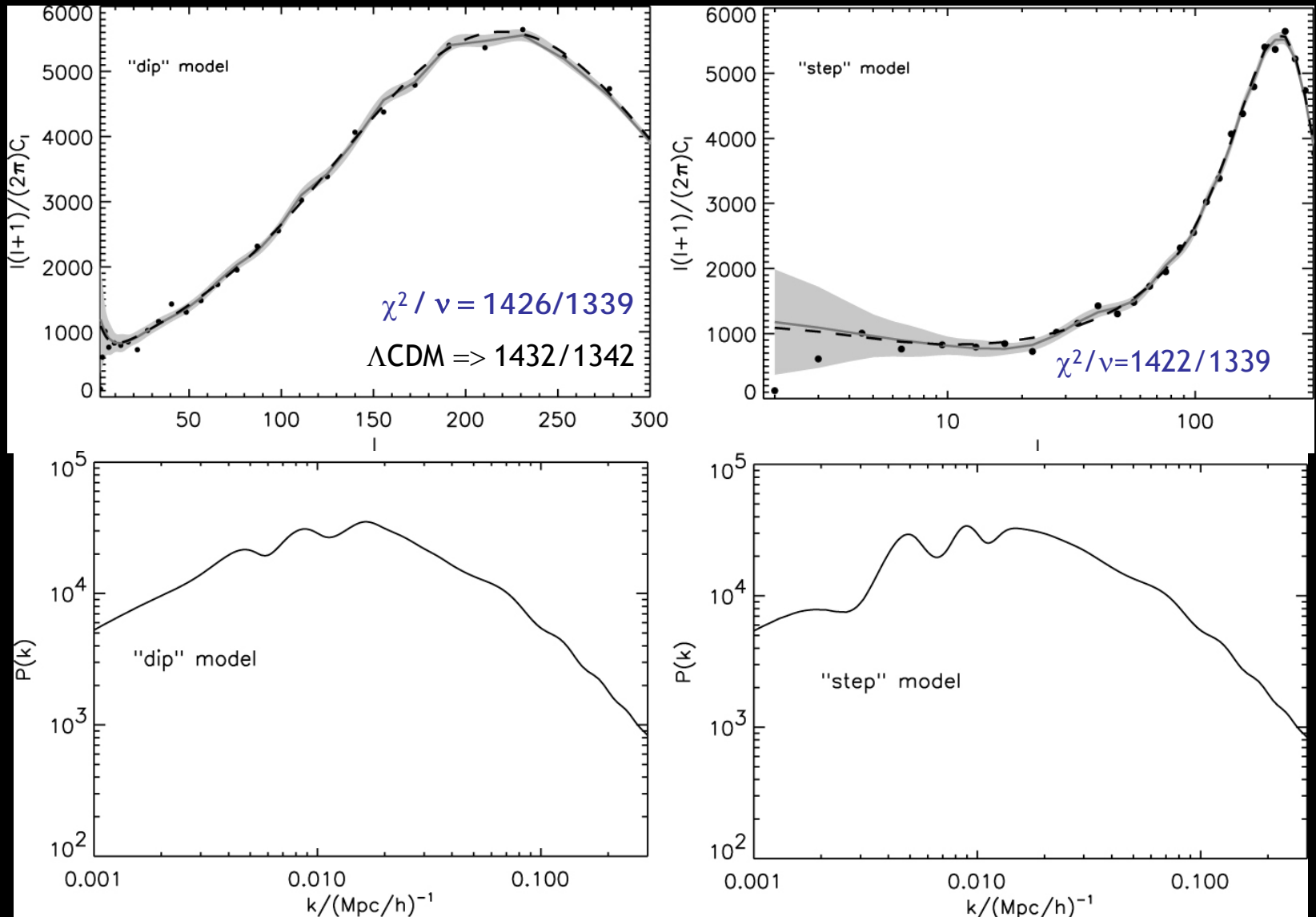


Reduced chisq for TT only 1.09  
Effects neglected on covariance matrix

0.5% to 1% error on the error!

- Weak lensing
- Beam asymmetries
- Non-gaussianity of the noise (striping)
- Features in inflation potential (Peiris et al. 2003)

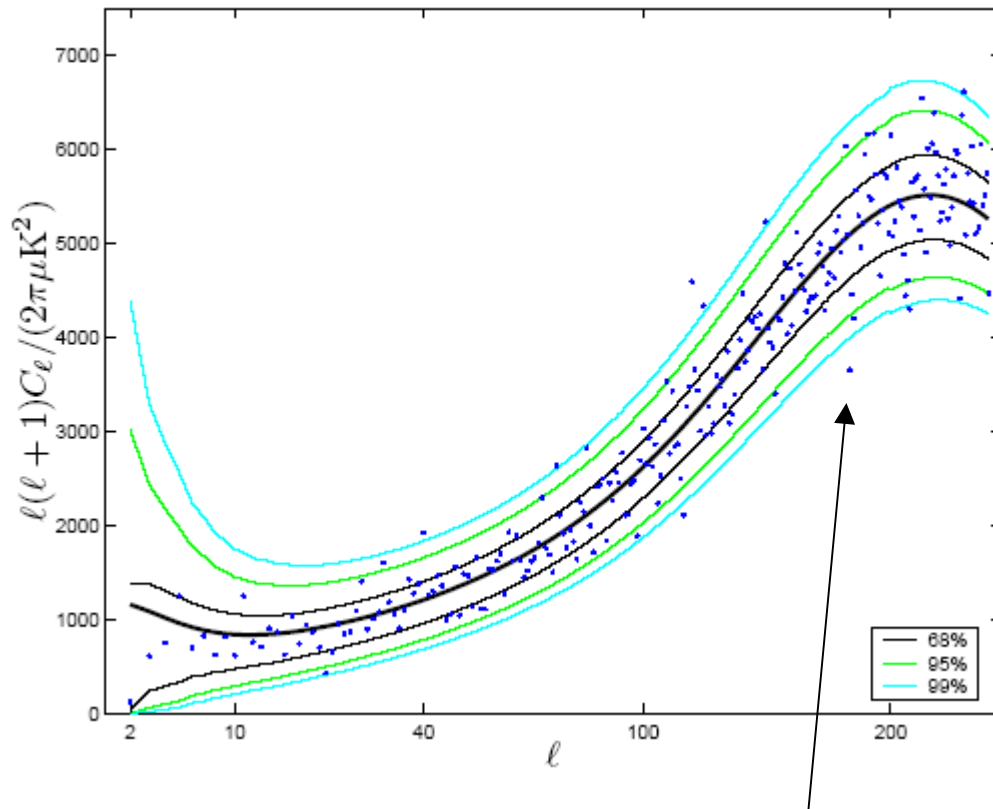
# Modifying the Inflaton Potential (0.1% change in amplitude)



See e.g. Adams, Cresswell & Easter (astro-ph/0102236), Hunt & Sarkar (astro-ph/0408138)

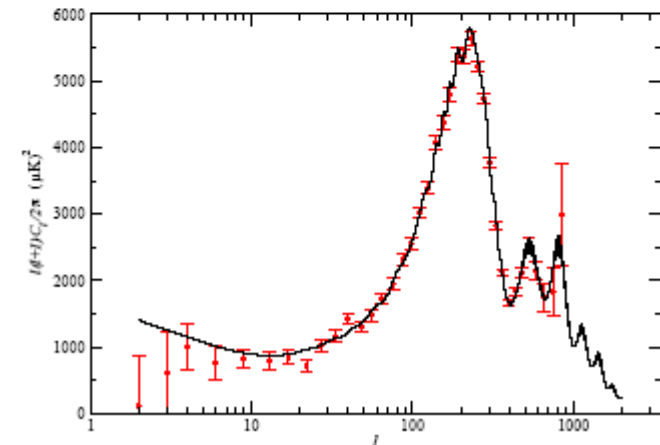
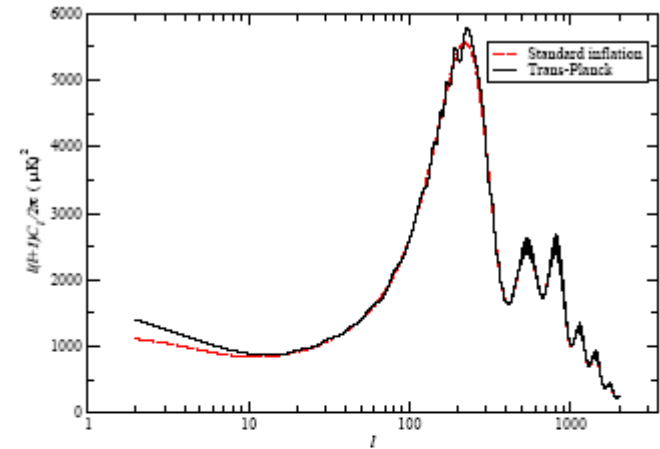


# More Power Spectrum Outliers



Lewis (astro-ph/0310186) observes that the number of  $3\sigma$  points (above) is high. Notes that only 3/16000 simulations have a lower value of  $C_{187}$ (arrow).

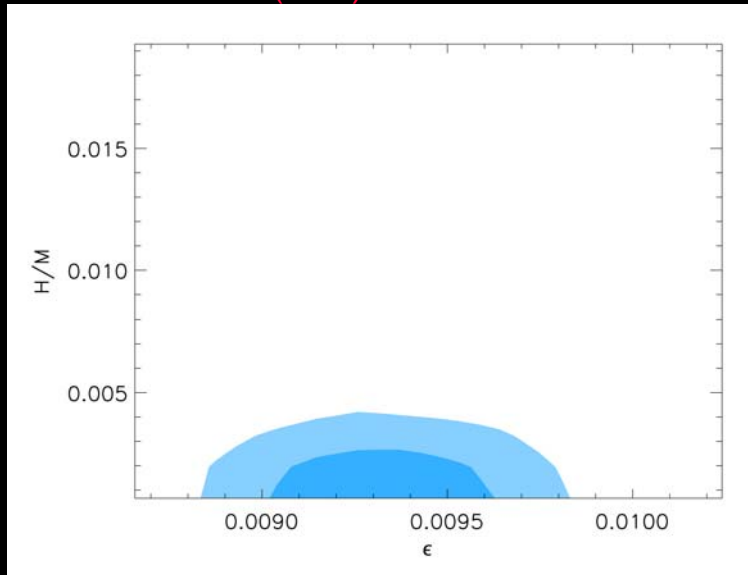
See Talk by O. Dore



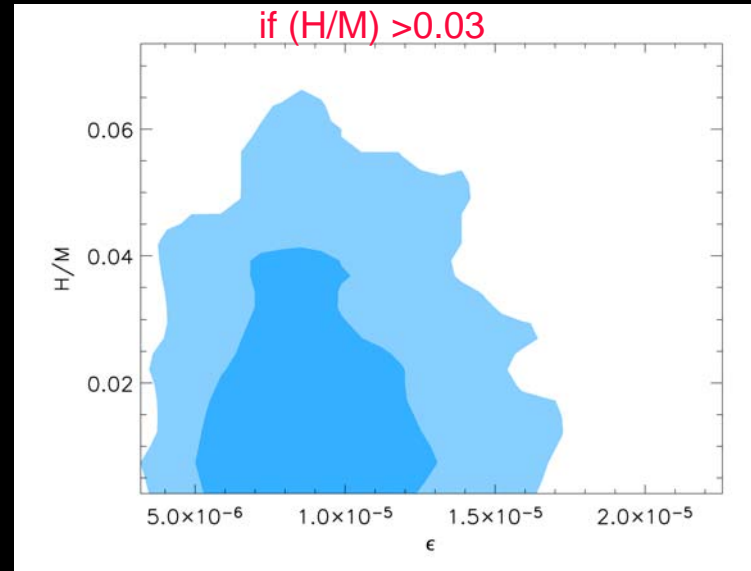
Martin & Ringeval (astro-ph/0310382) and Okamoto and Lim (astro-ph/0312284) fit toy trans-Planckian model to spectrum:  $\Delta\chi^2 = 16$  for 3(?) parameters.

# Are TP effects observable even in principle?

Fiducial model has  $r=16\epsilon=0.15$   
 $>3\sigma$  detection of TP oscillations  
if  $(H/M) > 0.004$



Fiducial model has  $r=16\epsilon=0.00013$   
 $>3\sigma$  detection of TP oscillations  
if  $(H/M) > 0.03$



- Feasibility study only
- Significant theoretical uncertainties (need a proper model to test!)
- Needs favorable piece of parameter space
- Tensor detection and  $(H/M)$  detection coupled

# Standard Model fits *WMAP* data

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- LCDM, Adiabatic, flat, composed of:
  - Baryonic Matter =  $4.7 \pm 0.6\%$
  - Dark Matter =  $24 \pm 7\%$
  - Dark Energy =  $71 \pm 7\%$
- Power Spectrum Slope =  $0.99 \pm 0.04$
- Hubble Constant =  $72 \pm 5$  km/s/Mpc [fits HST Key Project]
- $\Omega_b h^2 = 0.024$  [fits D/H]
- $\sigma_8 = 0.9$  (0.84 for running model) [fits lensing, clusters, etc]
- Age of the Universe  $13.7 \pm 0.2$  Gyr [fits stellar evolution]
- Fits LSS and Ly- $\alpha$  data

The numbers are consistent with a host of astronomical observations at different redshifts and scales.

# Future observational prospects

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- *Go to small scales!* Much better measurements of the primordial power spectrum shape.
  - Planck  $l \sim 3000$  ( $k \sim 0.2/\text{Mpc}$ )
  - ACT  $l \sim 10000$  ( $k \sim 0.7/\text{Mpc}$ ) [secondary effects]
  - Galaxies  $k \sim 1/\text{Mpc}$  [non-linearity & bias]
  - Lyman alpha  $k \sim 5/\text{Mpc}$  [gas phys. & radiation feedback]
  - Reionization  $k \sim 50/\text{Mpc}$  [much is unknown]
- Detecting gravitational waves
  - QUEST, QuAD, BICEP, PolarBear, EBEX, CLOVER, QUIET, Planck, CMBpol, Inflation Probe etc... [See G. Smoot Talk]
- Detecting non-Gaussianity from 2<sup>nd</sup> order gravity
  - Can we detect  $f_{NL} \sim 1$ ?

# Theoretical directions

---

- More detailed predictions from a specific, physically motivated model
  - More accurate reheating scenario is necessary!!
  - No slow-roll approximations
  - Test (constrain) models one by one

THE END

