

Environmental high-energy astrophysics in the context of space missions as LISA, Solar Orbiter and JWST, and its implications for space weather science

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Raphael (Raffaello Sanzio)
1483-1520



1506
UNIVERSITÀ
DEGLI STUDI
DI URBINO
CARLO BO



Istituto Nazionale di Fisica Nucleare

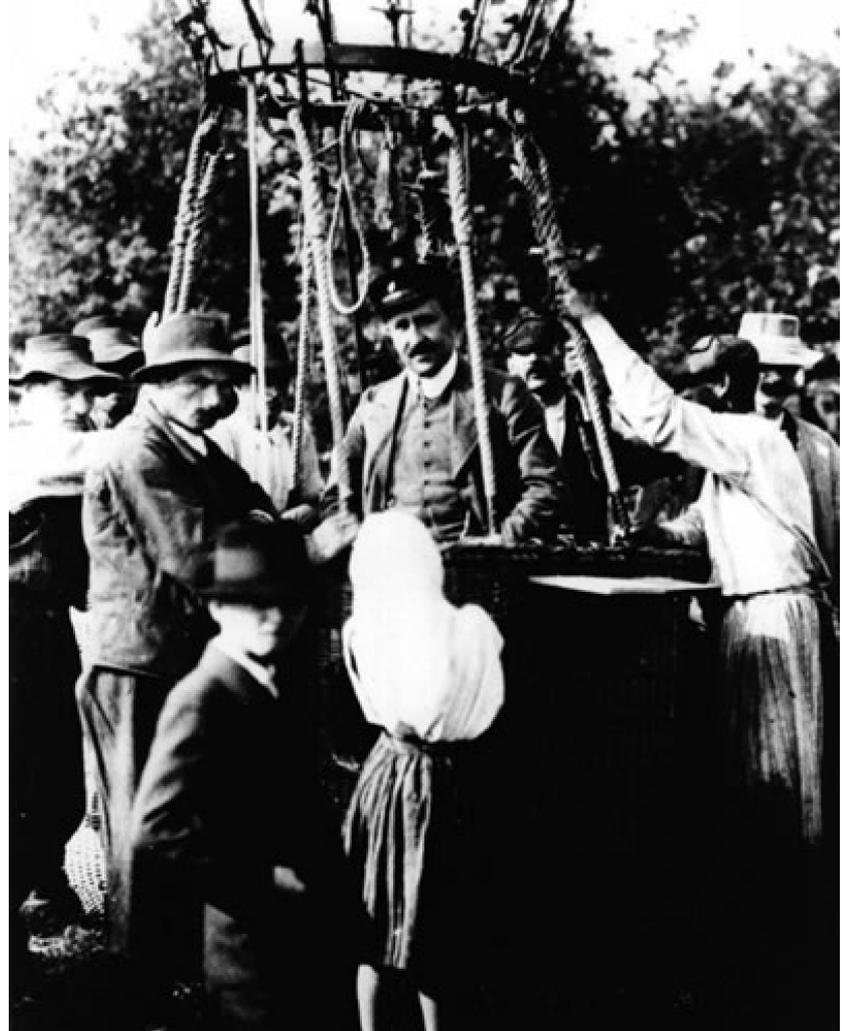


Contents

- Antimatter search in cosmic rays: pulsars as most plausible astrophysical sources of high-energy positrons in cosmic rays
- Cosmic-ray role in limiting the performance of space experiments (LISA Pathfinder, LISA, Solar Orbiter): instrument diagnostics and space weather science
- Circumpulsar (magnetar) disc formation and positron measurements near Earth: clues from JWST observations and gravitational wave detection
- Conclusions

The discovery of cosmic rays

- 1912 cosmic-ray discovery by Victor F. Hess
- What are cosmic rays made of?
- Photons? No
- Energetic positively charged particles (protons and ions)!
Latitude effect and east-west asymmetry
- 1000 particles $\text{m}^{-2} \text{s}^{-1}$ at the top of the atmosphere



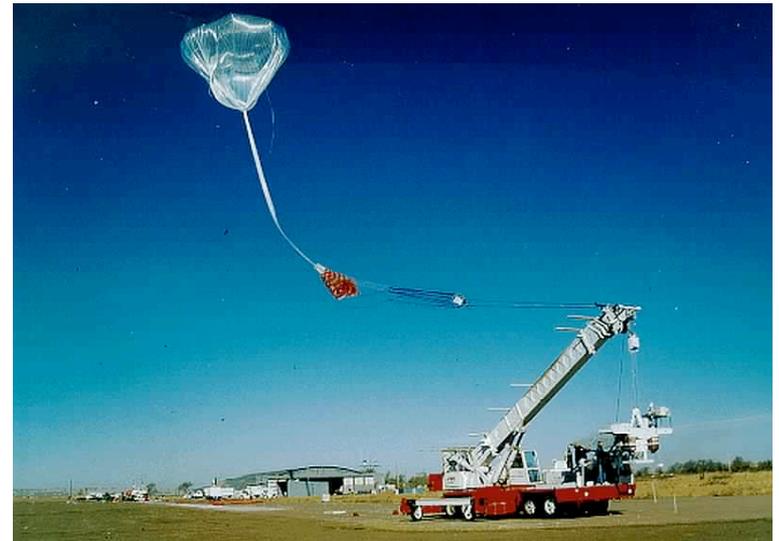
Cosmic-ray research@ NMSU 1990-1992



NSBF, NASA

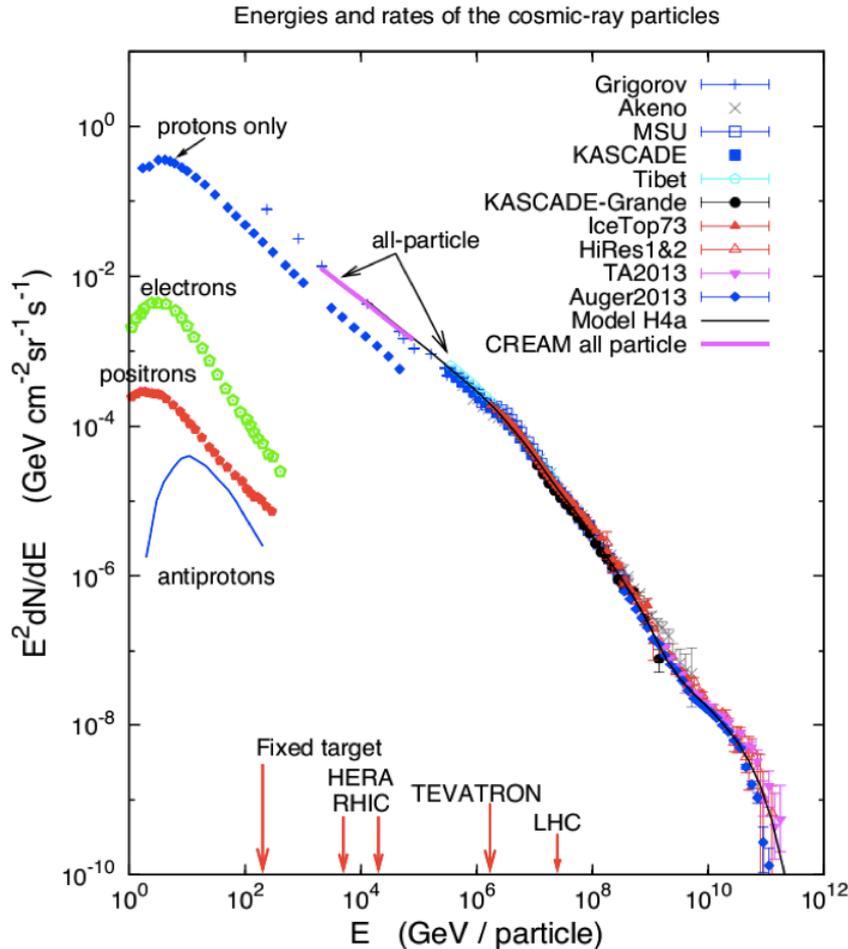


Fort Sumner
NM, USA



Cosmic-ray energy spectra

Blasi, 2012



The cosmic-ray differential flux varies by 32 orders of magnitude when the energy varies by 12 orders of magnitude

$$F(E) = A (E+b)^{-\alpha} E^{\beta}$$

particles/($\text{m}^2 \text{sr s GeV}$)
 $E < 100 \text{ GeV} (/n)$

$$F(E) = A E^{-\gamma}$$

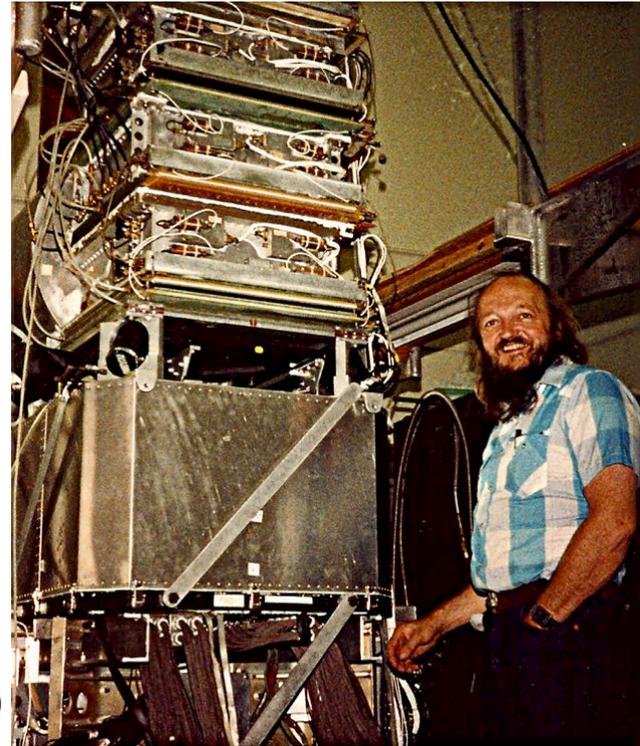
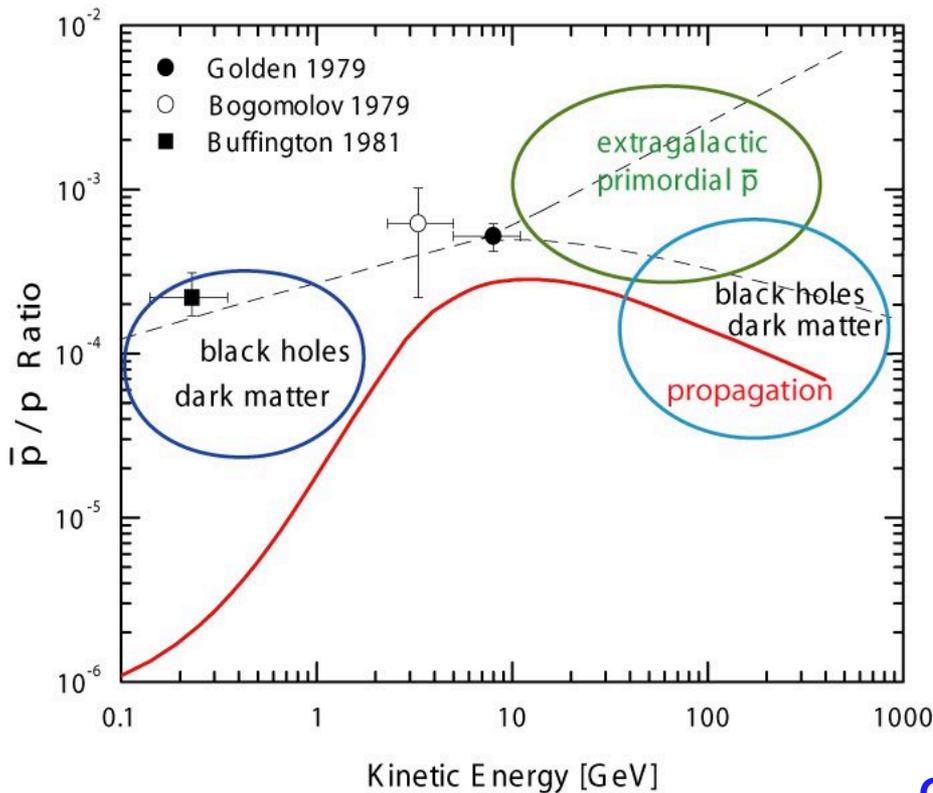
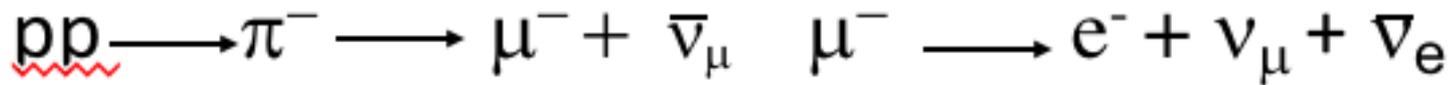
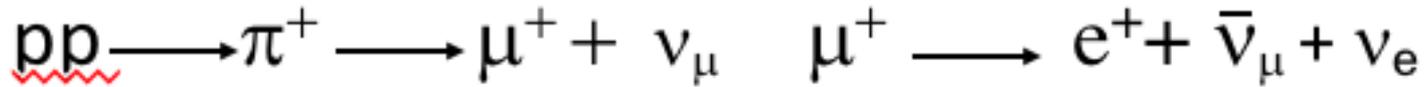
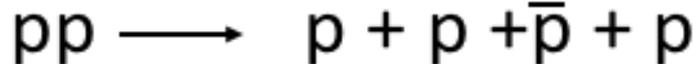
particles/($\text{m}^2 \text{sr s GeV}$)
 $E > 100 \text{ GeV} (/n)$
 with changing γ

$\gamma = 2.7$ up to the knee ($3 \times 10^{15} \text{ eV}$)
 $\gamma = 3-3.15$ up to 10^{18} eV
 $\gamma = 2.7$ above the ankle ($3 \times 10^{18} \text{ eV}$)
 then there is a possible cut-off

Search for antimatter in cosmic rays

- Positron discovery in cosmic rays in 1964 (De Shong et al. 1964)
- Antiproton discovery in cosmic rays in 1979 (Bogomolov et al.; 1979 and Golden et al.; 1979)
- No antinucleus detection up to present time
- Which is the origin of antimatter in cosmic rays?
- Both positrons and antiprotons are produced in primary proton interactions in the interstellar medium (1 atom hydrogen cm^{-3} , for comparison on Earth surface there are 10^{19} atoms cm^{-3}). Positrons are in excess with respect to the secondary component, while the same does not apply to antiprotons.

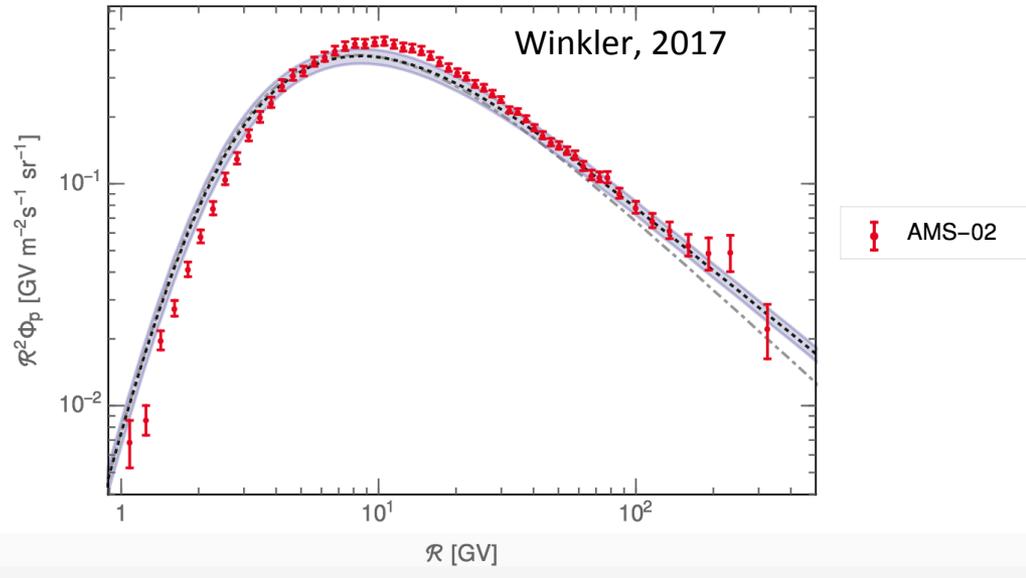
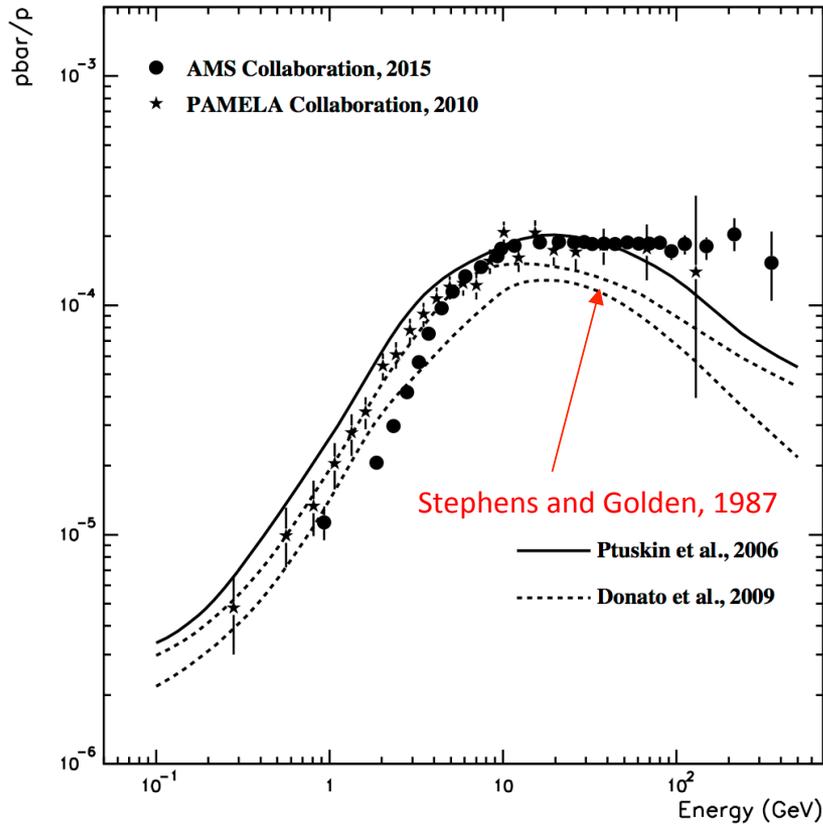
Antiprotons and positrons as secondary particles



Antiproton discovery in cosmic rays in 1979

Golden et al.; 1979 and Bogomolov et al.; 1979

Antiproton observations in cosmic rays



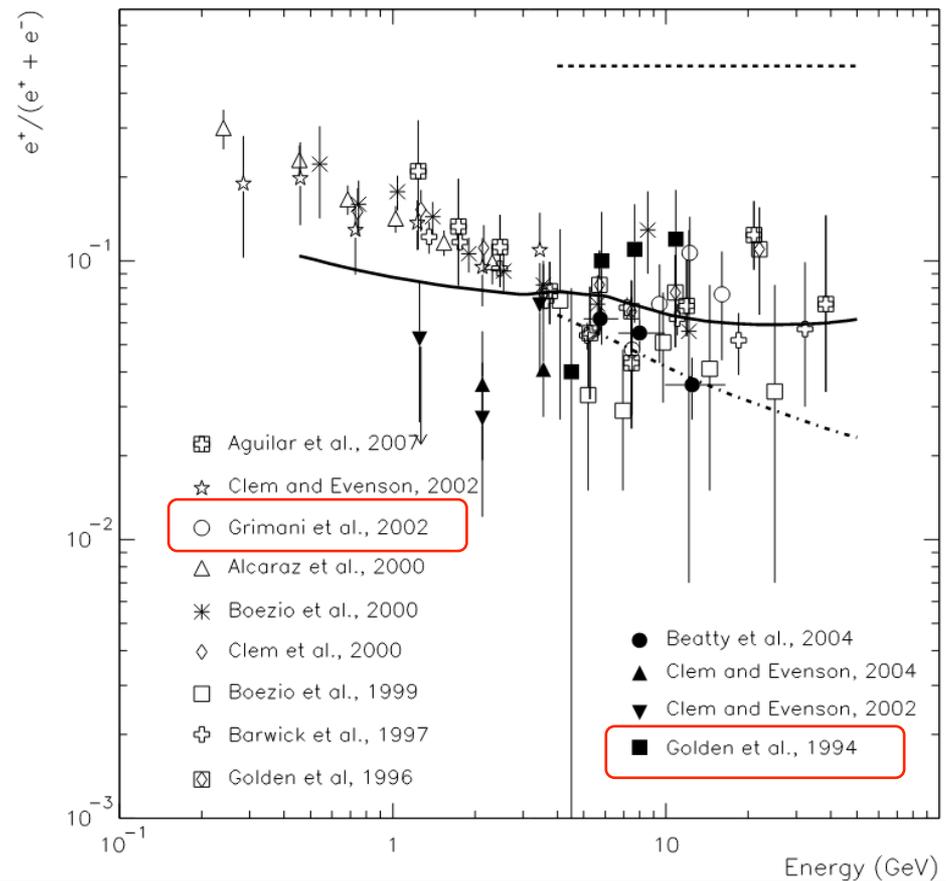
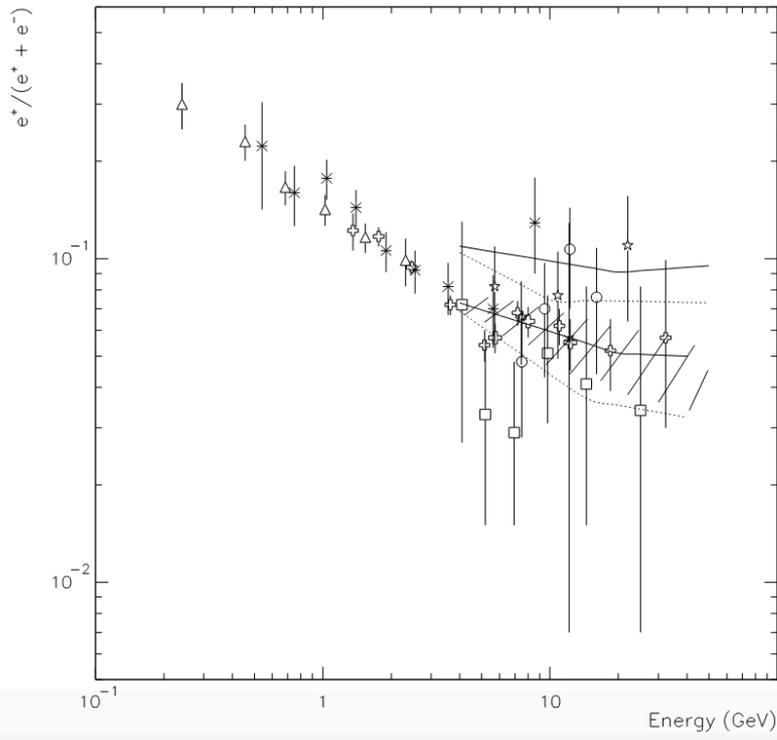
CG, Astrophysics and Space Science, 268, 465, 1999

Early positron measurements

Pulsars as astrophysical sources of high-energy positrons
 Pulsar birthrate from pulsar characteristics and vice versa

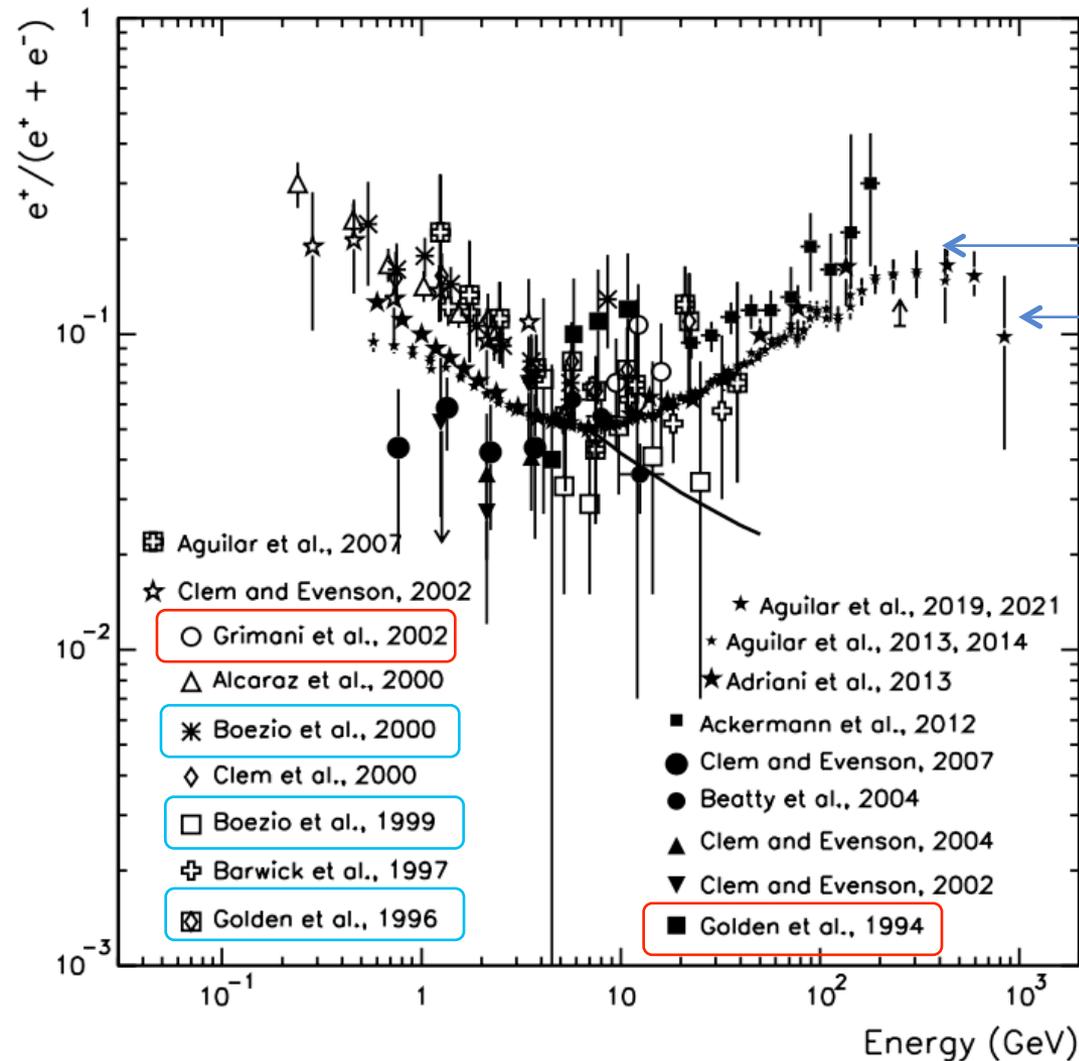
Pulsar Birthrate⁻¹=200+/-100 years

Pulsar Birthrate⁻¹= 33 years



CG, A&A, 418, 649-653, 2004
CG, A&A, 474, 339-343, 2007

Positron observations in cosmic rays



Average $7 \text{ GeV} < E < 50 \text{ GeV}$
 0.064 ± 0.003
 Before PAMELA, Fermi
 and AMS-01 AMS-02

$284^{+91}_{-64} \text{ GeV}$

Cut-off $810^{+310}_{-180} \text{ GeV}$
 Aguilar et al.; 2021
 AMS-02

Positron observations are compatible with sources of TeV $e^- - e^+$ located between 100 and 800 pc from the Solar System ([Attallah 2016](#)).

Possible sources of positron excess

- 1) Primordial Black Hole Evaporation
- 2) ^{56}Co decay in Supernova Remnants
- 3) Supersymmetric particle annihilation
- 4) $\gamma\gamma$ interaction

5) Pulsar magnetosphere

(Polar Cap Models- Outer Gap Models)

Bushing et al, 2007

Hooper, Blasi &

Serpico, 2009

Di Mauro et al., 2014

Feng & Zhang, 2016

CG, A&A, 2004, 2007

CG, ICRC 2015

Pulsar polar cap models

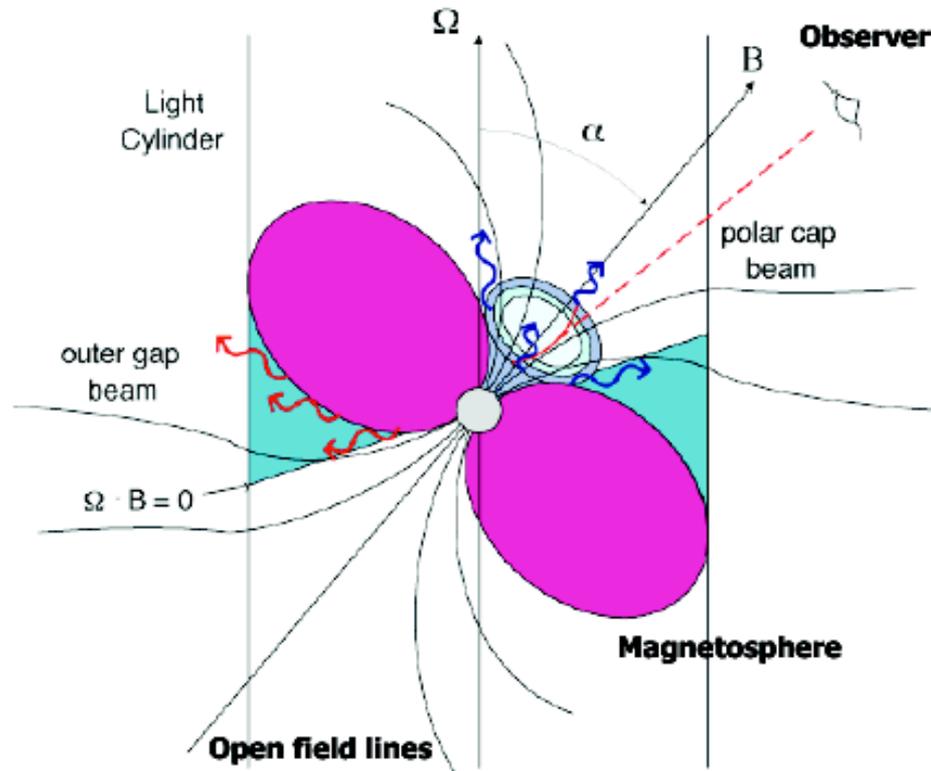
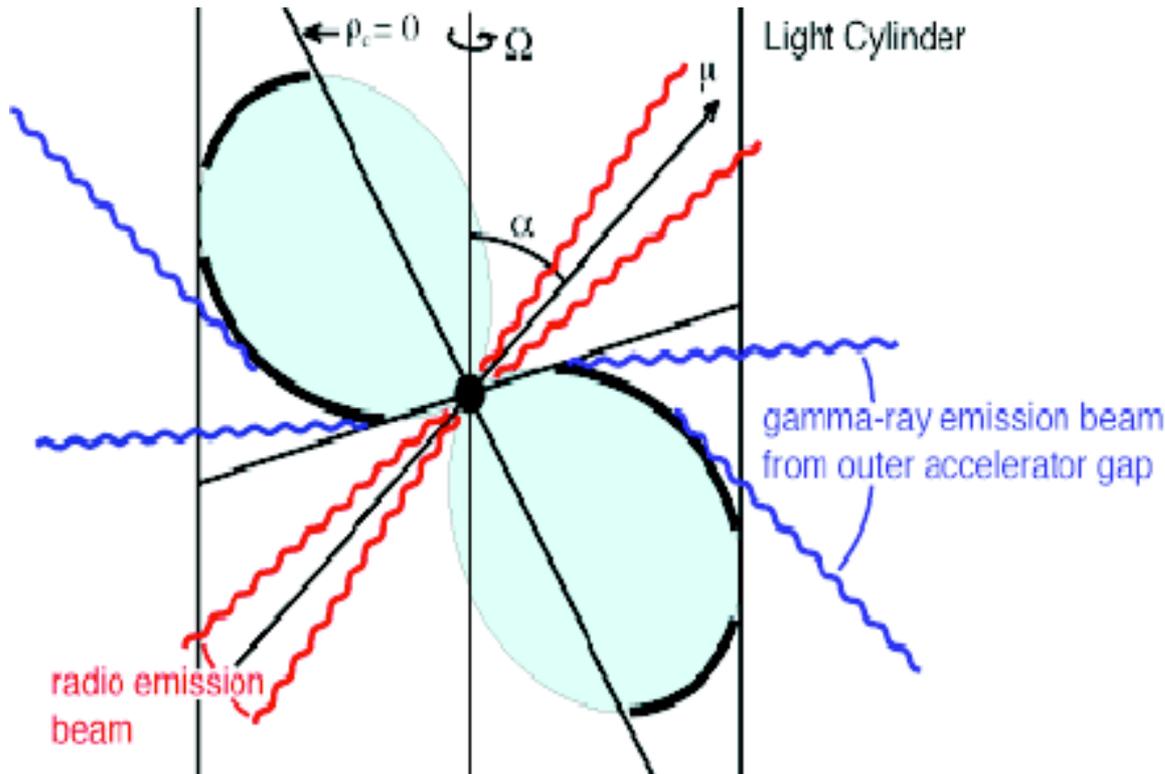


Figure from <http://cossc.gsfc.nasa.gov/images/epo/gallery/pulsars/>

Goldreich & Julian, 1969
Harding & Ramaty, 1987

- Strong electric fields are induced by the rotating neutron star
- Electrons are extracted from the pulsar outer layer and accelerated
- Open field lines originate at polar caps ($r_{pc} = 8 \times 10^2$ m)
- Electron and positron synchrotron energy losses limit the particles to GeV energies

Pulsar outer gap models



Cheng, Ho & Ruderman, 1986

- *Electrons are accelerated in the outer magnetosphere in vacuum gaps within a charge separated plasma
- *Electrons interact through synchrotron radiation or inverse Compton scattering
- * e^+e^- pairs are produced by $\gamma\gamma$ interaction

<http://coss.c.gsfc.nasa.gov/images/epo/gallery/pulsars/>

Different cut-off energies are predicted by polar cap and outer gap models for the pulsed gamma-ray spectra!

C. Grimani ECRS Florence August 31st - September 3rd 2004

Pulsar TeV pulsed gamma-ray emission

- From Crab up to 2 TeV (Ansoldi et al.; 2016)
- From Vela up to 7 TeV (Djannati-Atai et al. 2017)
- From Geminga above 20 TeV (Lopez et al. 2019)

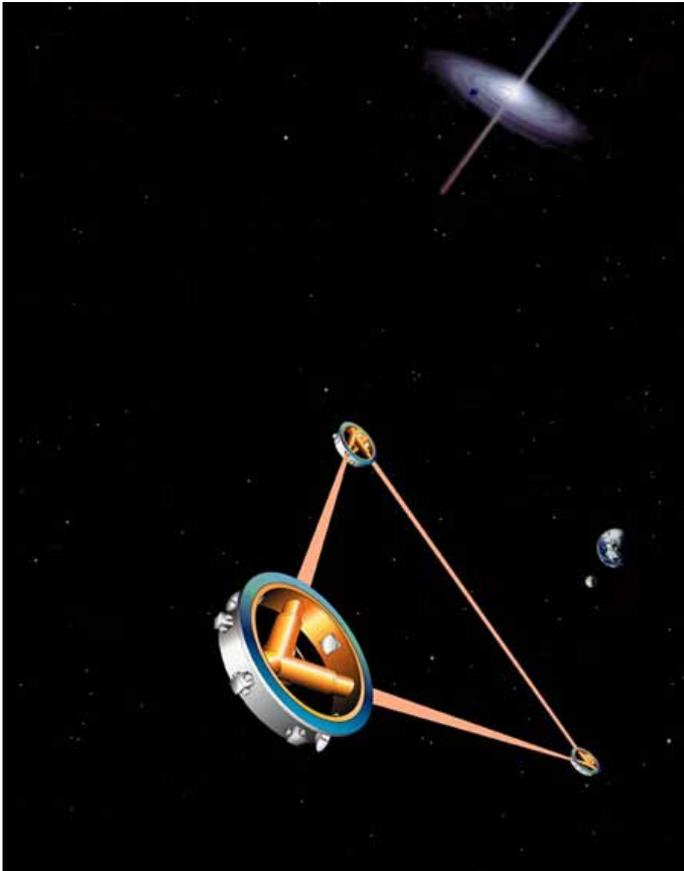
Electron and positron Lorentz factors $> 5 \times 10^6$

These evidences reveal that outer gap models are favoured over polar cap models and that pulsars are CERTAIN sites of positron acceleration up to tens of TeV. May the pulsar environment limit the electron and positron propagation in the interstellar medium?

PHYSICS OF COSMIC RAYS FOR SPACE INSTRUMENT DIAGNOSTICS AND SPACE WEATHER SCIENCE

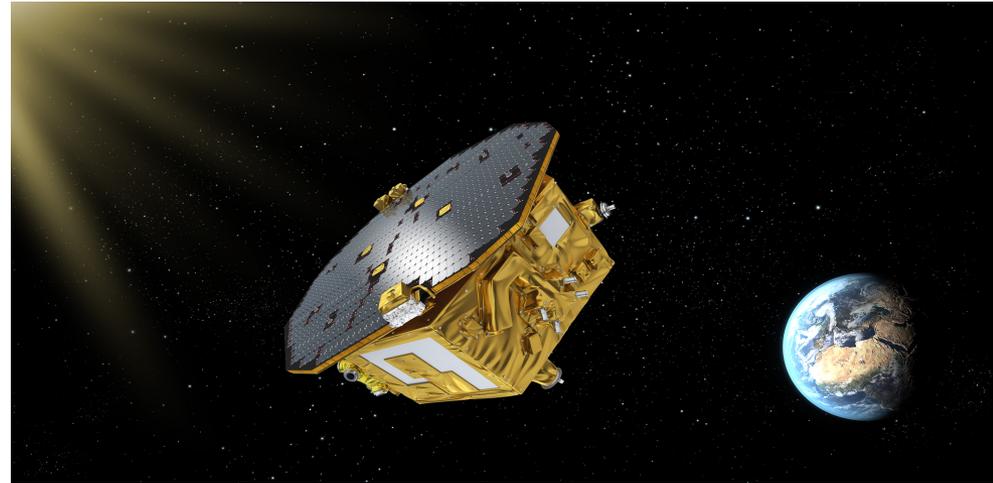


The LISA mission



Lunch: 2035 **Arm:** $2.5 \cdot 10^6$ km
Orbit: Heliocentric - 3 S/C trailing Earth at 50 million km

The LISA Pathfinder mission



Lunch: December 3, 2015 4:04 GMT
from Kourou (French Guiana)

Orbit: L1

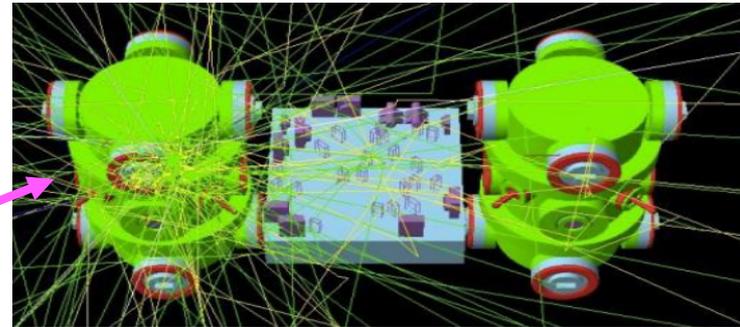
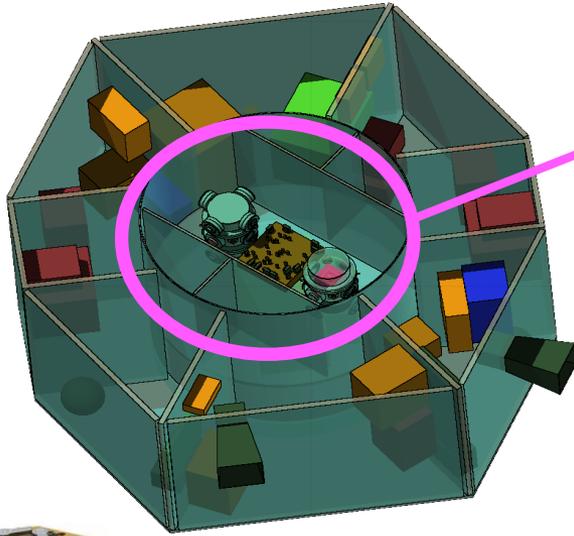
Mission end: July 18, 2017

<http://sci.esa.int/lisa-pathfinder/>

<https://sci.esa.int/web/lisa>

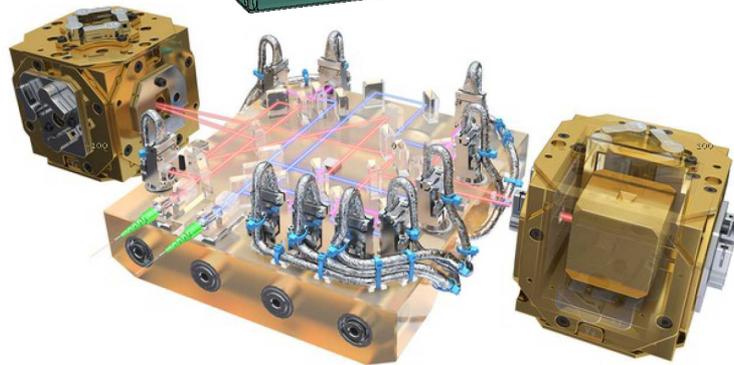
<https://www.elisascience.org/>

LISA Pathfinder test masses



Approximately 13 g cm^{-2} of material surrounded the test masses

Test masses:
cubes of 70% Au-30% Pt
4.6 cm side
2 kg mass
38 cm distance



Armano et al., PRL, 116, 231101, 2016
Armano et al., PRL, 120, 061101, 2018

Test-mass net and effective charging

$$\lambda_{NET} = \sum_{j=-\infty}^{+\infty} j \lambda_j$$

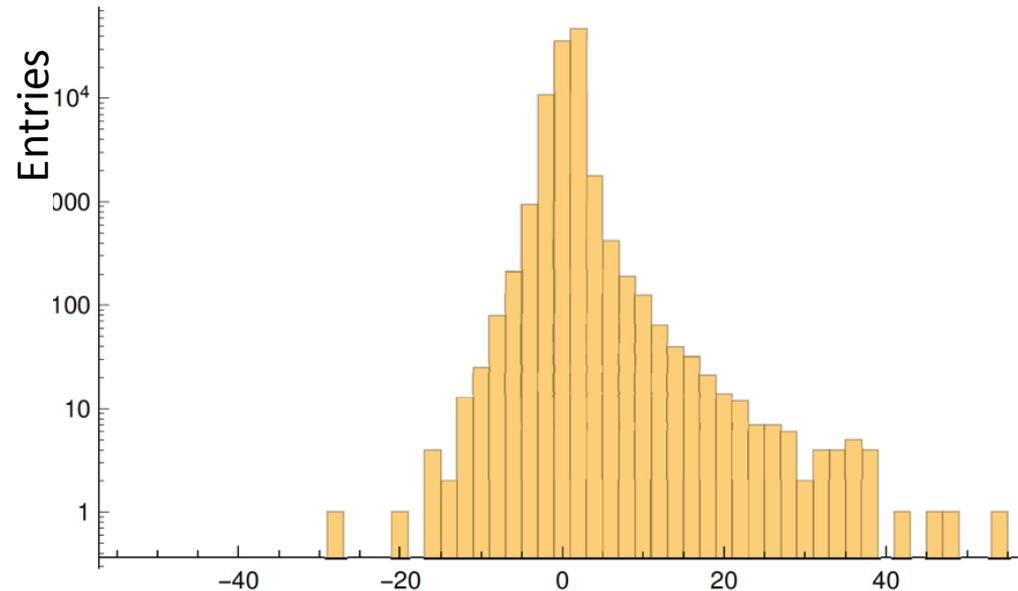
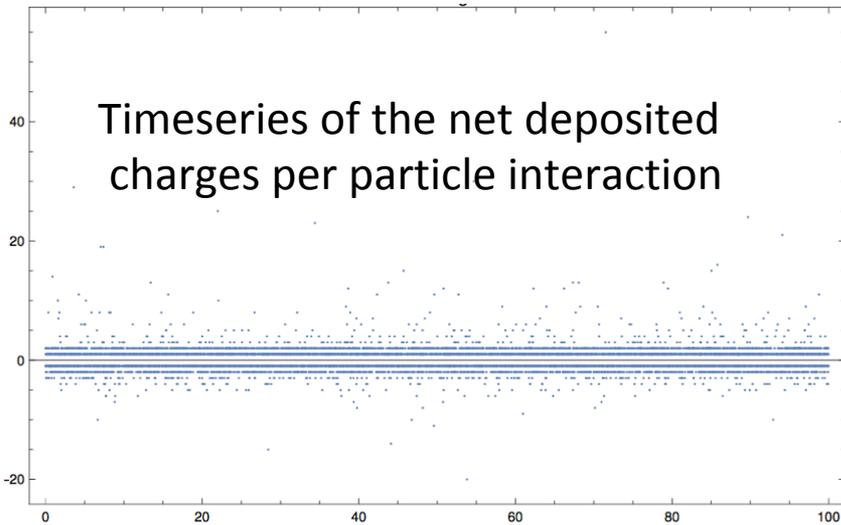
$$\lambda_{EFF} = \sum_{j=-\infty}^{+\infty} j^2 \lambda_j$$

Where: j is the amplitude of the charge released by each particle and λ_j is the rate of occurrence of that event.

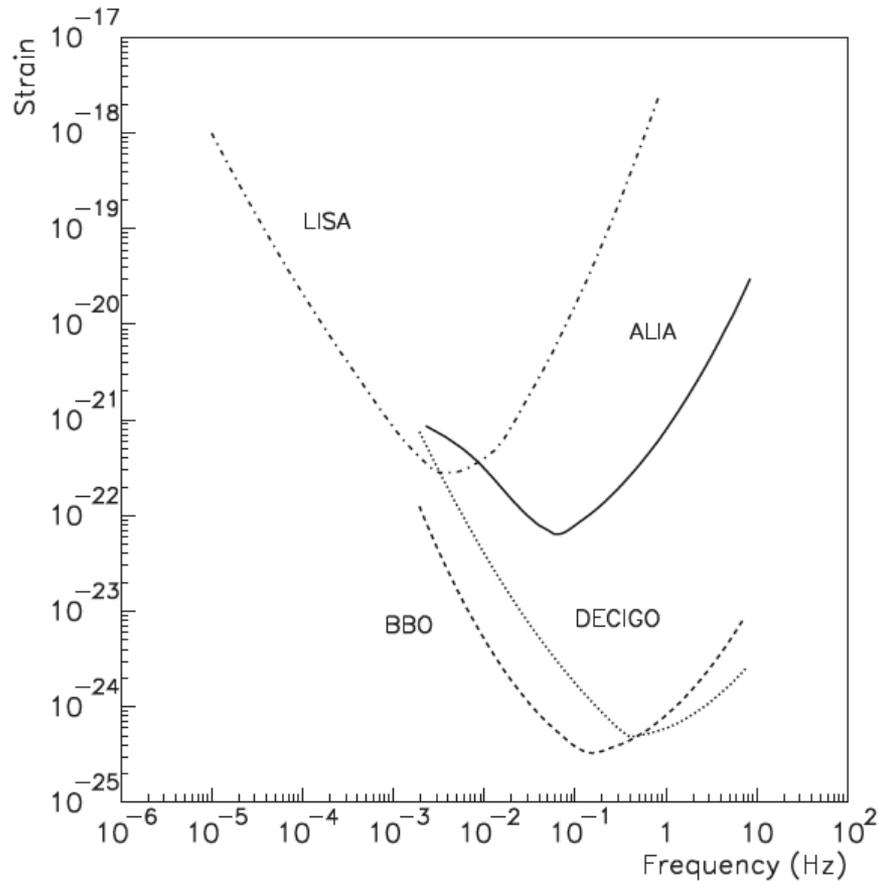
Ferrari et al., 2005-10 (Geneva:CERN)
 Battistoni et al.,
 AIP Conf. Proc. 896,
 M. Albrow and R. Raja, 31-49

$$S_Q(\omega) = \frac{S}{\omega} = \frac{\sqrt{2e^2 \lambda_{eff}}}{\omega}$$

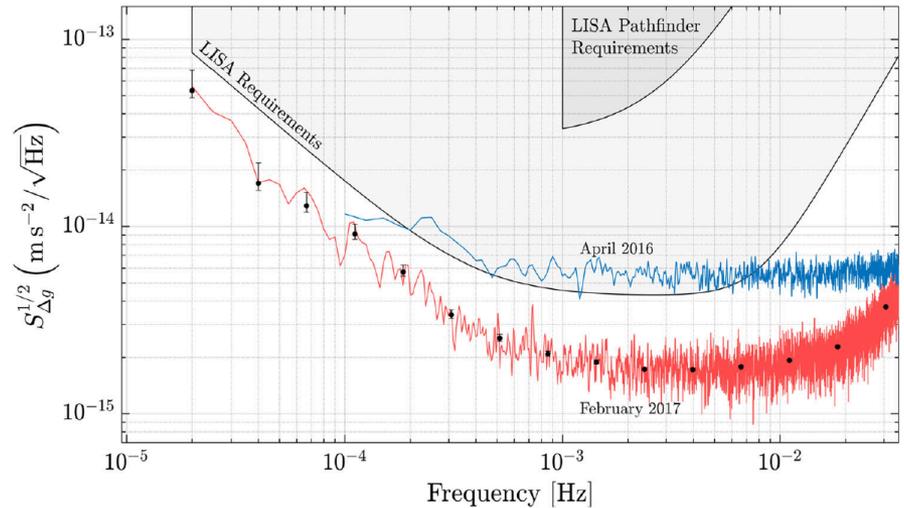
$$S = \sqrt{2e^2 \lambda_{eff}} \quad e \text{ s}^{-1} \text{ Hz}^{-1/2}$$



Sensitivities of future interferometers for gravitational wave detection in space



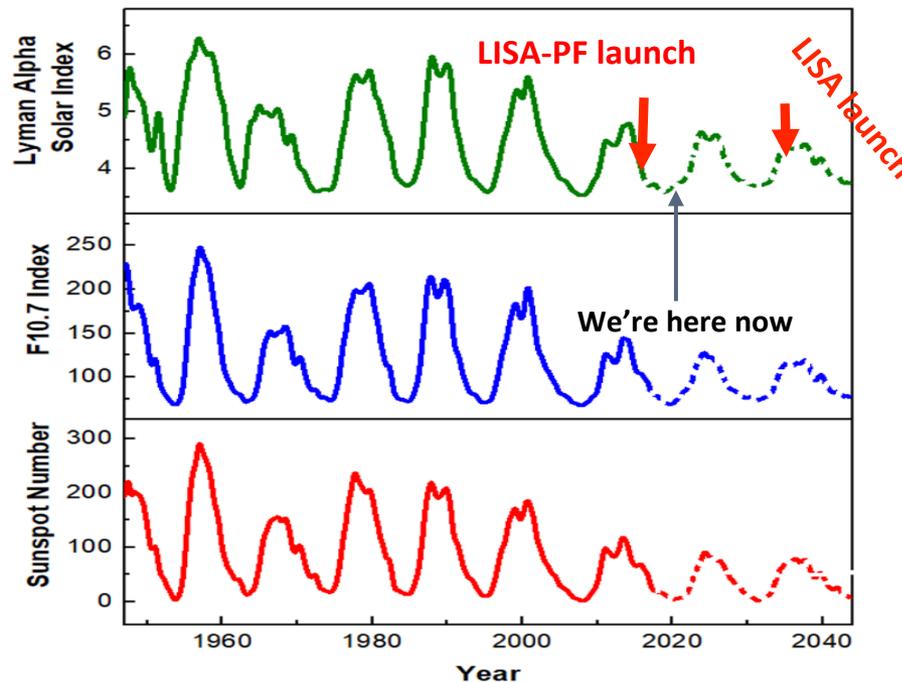
CG, MNRAS, 507 (1), 261-266, 2021



Armano et al., PRL, 120, 061101, 2018

Amplitude spectral density:
 at $1.74 \pm 0.01 \text{ fm s}^{-2}/\text{Hz}^{0.5}$ above 2 mHz and
 $(6 \pm 1) \times 10 \text{ fm s}^{-2}/\text{Hz}^{0.5}$ at 20 μHz

Predictions of the solar cycles 25-26



From the analysis of:

Ly- α index
F10.7 cm
Sunspot number
recorded in the last 70 years

In agreement with:
Cameron, Jiang and
Schussler
Apj, 823, L22, 2016
for the solar cycle 25

Fig. 5 Monthly variations of the observed and predicted (dotted) values of the sunspot numbers, F10.7 cm index and Lyman alpha index

Singh & Bhargawa *Astrophys. Space Sci.* (2019) 364:12

Cosmic-ray flux prediction and interpolation

NO DRIFT INCLUDED

Gleeson and Axford, Ap. J., 154, 1011, 1968

$$\frac{J(r, E, t)}{E^2 - E_0^2} = \frac{J(\infty, E + \Phi)}{(E + \Phi)^2 - E_0^2}$$

J: particle flux

r: distance from Sun

E: particle total energy

t: time

E_0 = particle mass

$\Phi = Ze\phi$ particle energy loss from ISM
(different for each particle species)

$$\begin{aligned}\Phi_{m_LISA} &= 200 \text{ MV/c} \\ \Phi_{M_LISA} &= 1200 \text{ MV/c}\end{aligned}$$

$$\begin{aligned}\Phi_{m_LPF} &= 320 \text{ MV/c} \\ \Phi_{M_LPF} &= 550 \text{ MV/c}\end{aligned}$$

$$F(E) = A (E + b)^{-\alpha} E^\beta \quad \text{particles (m}^2 \text{ sr s GeV n}^{-1}\text{)}^{-1},$$

Papini, Grimani and Stephens, Nuovo Cimento, 1996

Φ from

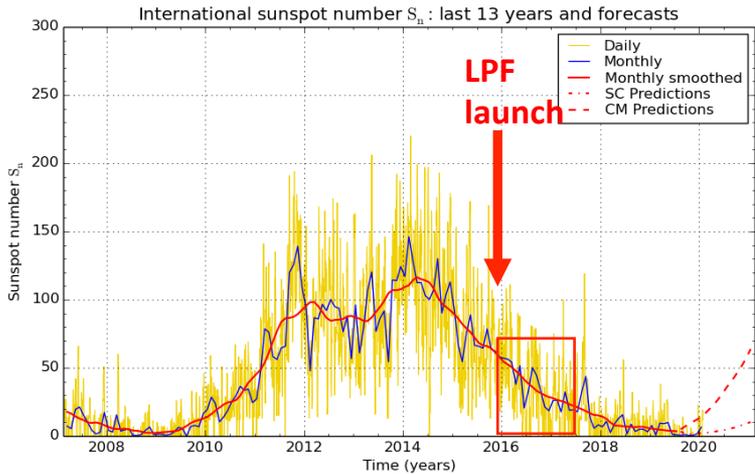
http://cosmicrays oulu.fi/phi/Phi_mon.txt

$$J_p(\infty, E + \Phi)$$

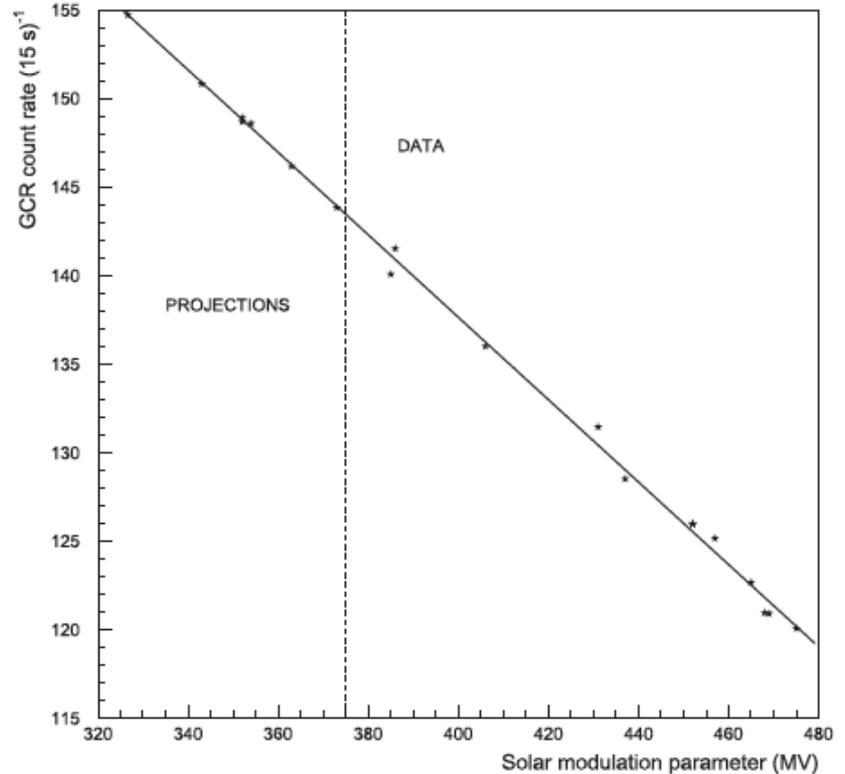
Burger, R. A., Potgieter, M. S. and Heber, B., JGR, 105, 27447, 2000

LISA Pathfinder monthly cosmic-ray observations

The LISA PATHFINDER mission
Orbit: L1 in 2016-2017



SILSO graphics (<http://sidc.be/silso>) Royal Observatory of Belgium 2020 February 1



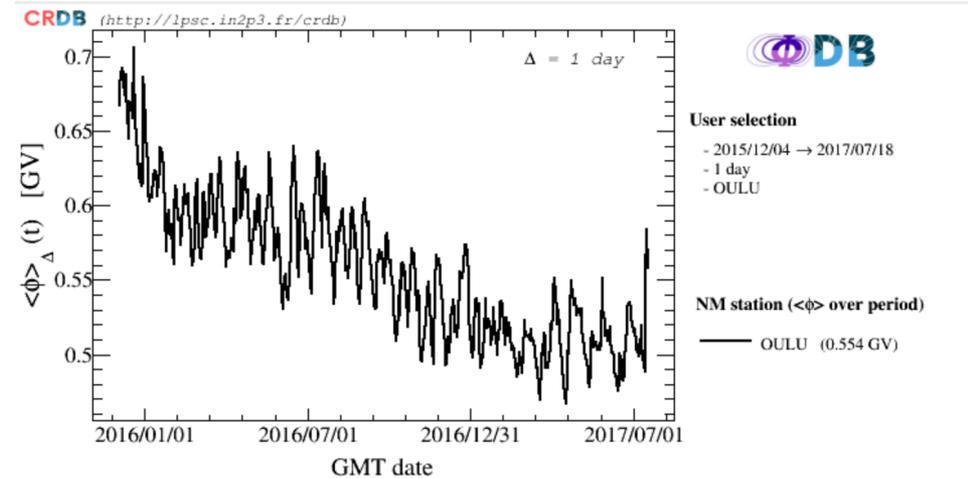
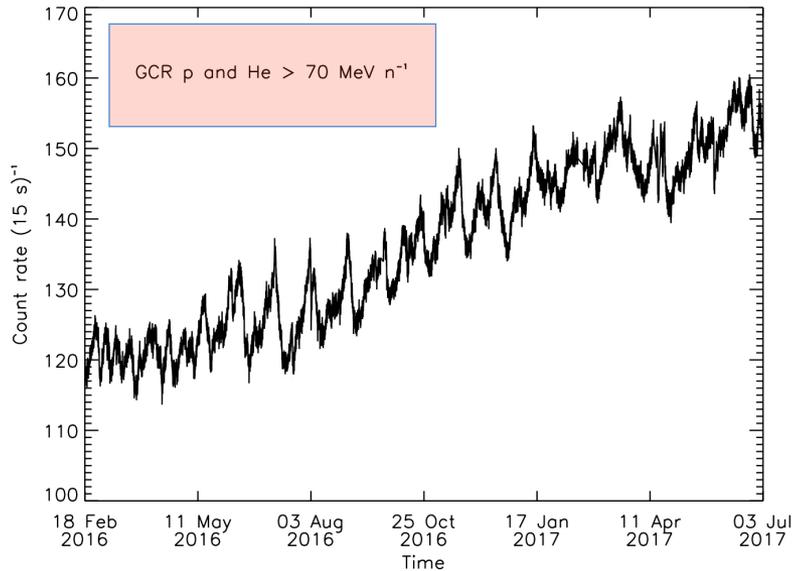
2015	695	686	743	705	655	666	617	612	611	601	585	561	644
2016	500	468	475	468	464	447	464	438	436	407	385	386	444
2017	366	357	348	367	359	350	383	420	444	405	367	360	377

http://cosmicrays oulu.fi/phi/Phi_mon.txt

Monthly and daily solar modulation in 2016-2017: lessons learned with LISA Pathfinder

Armano et al.; ApJ,854, 113, 2018

Armano et al.; ApJ,874, 167, 2019



Increasing trend (decreasing solar modulation)

45 galactic cosmic-ray short-term recurrent depressions > 2 days

3 Forbush decreases

23 galactic cosmic-ray short-term non-recurrent variations < 2 days

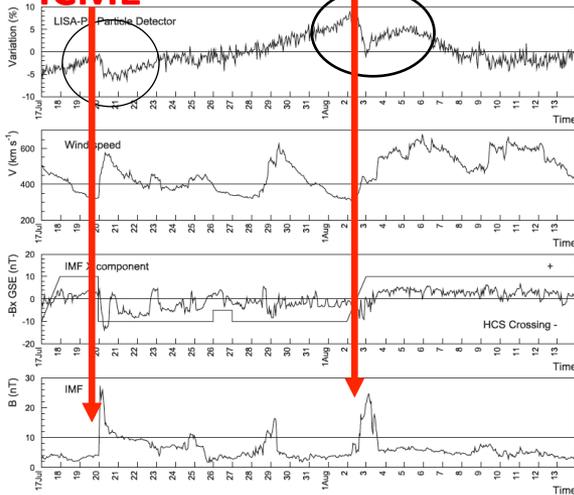
Cosmic-ray short-term variations with LISA Pathfinder

Forbush Decreases

ICME

Bartels rotation 2496

ICME



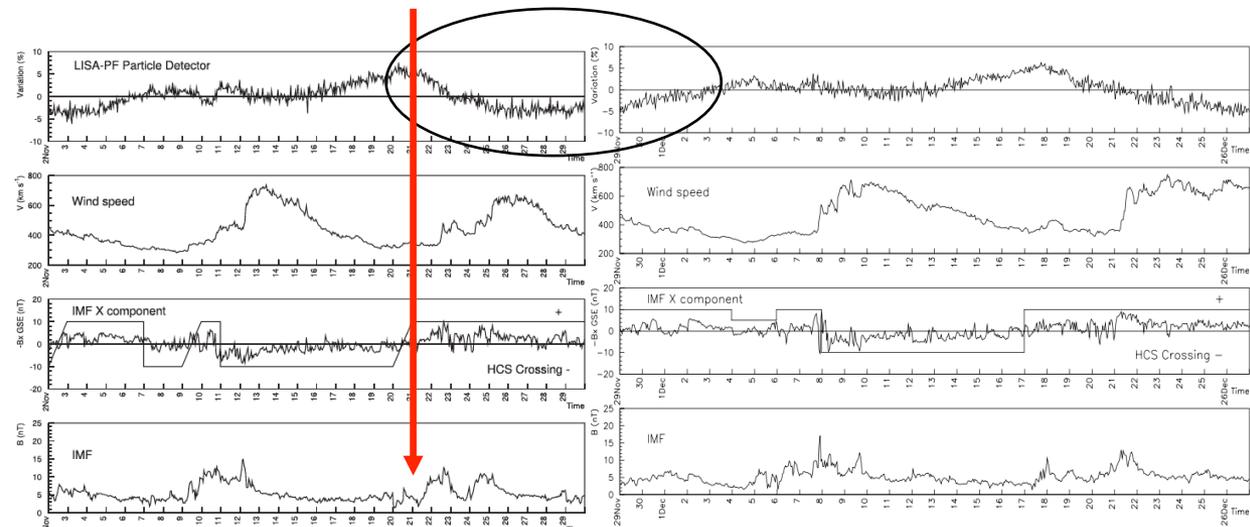
Armano et al.; ApJ, 854, 113, 2018
 Armano et al.; ApJ, 874, 167, 2019

Recurrent galactic cosmic-ray flux variations

HSS

Bartels rotation 2500

Bartels rotation 2501

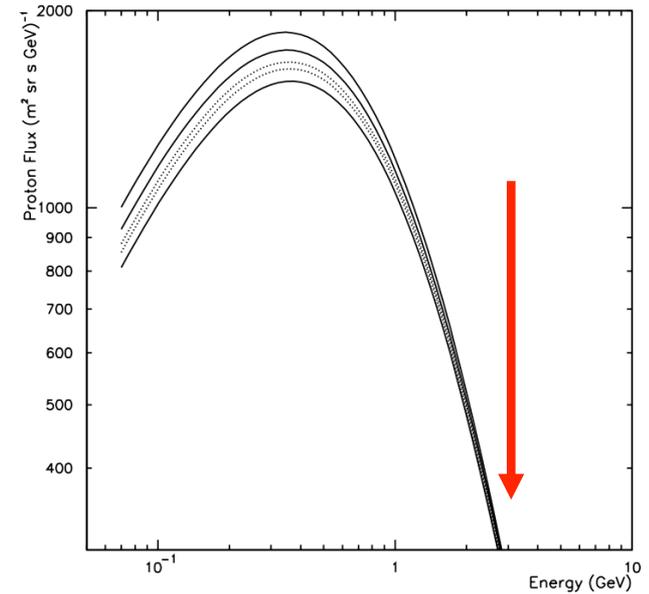
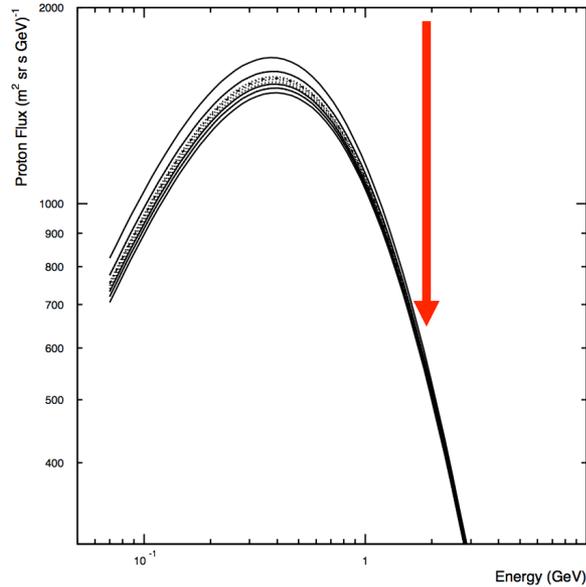
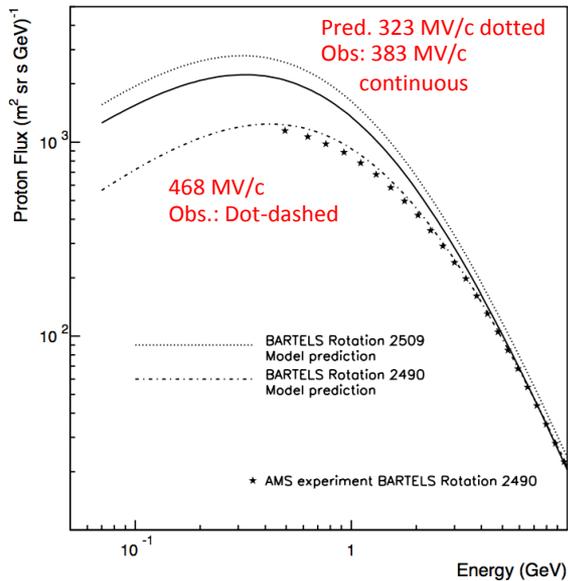


CG et al.; ApJ, 904, 64, 2020

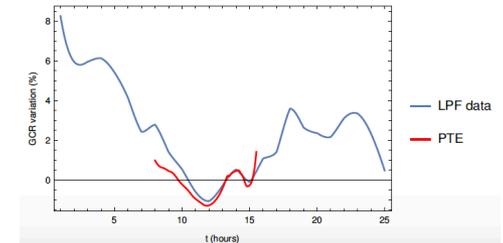
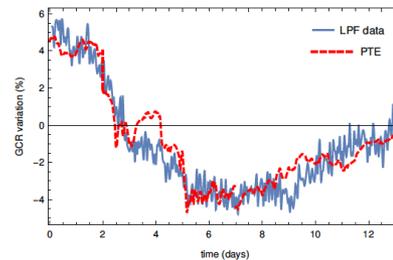
Energy-dependence of galactic cosmic-ray long and short-term variations

Recurrent variations

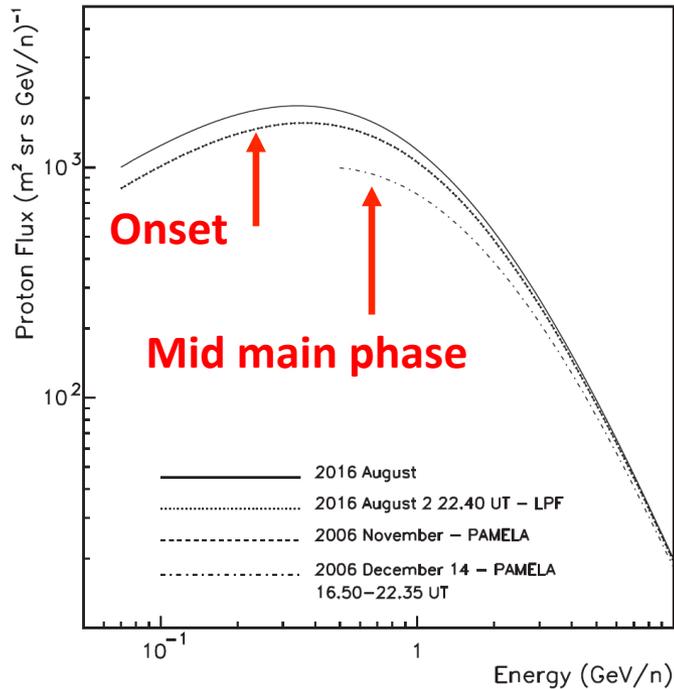
Forbush decrease



Fokker-Planck approach to reproduce galactic cosmic-ray short-term variations



The PAMELA December 14, 2006 Forbush decrease



Onset November 2006

Net charging: $26.0 +e s^{-1}$

Effective charging: $1169 e s^{-1}$

Mid main phase December 14, 2006 17.20 UT

Net charging: $10.2 +e s^{-1}$

Effective charging: $536.7 e s^{-1}$

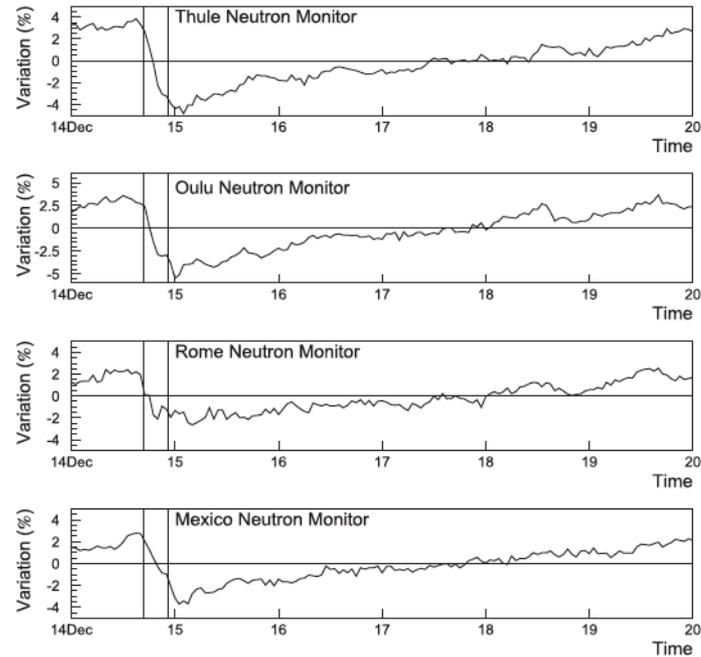
Dip December 15, 2006 00.00 UT

Net charging: $6.7 +e s^{-1}$

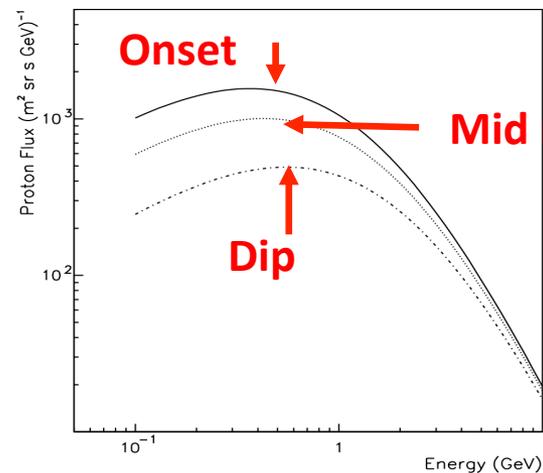
Effective charging: $483.4 e s^{-1}$

ESA ITT 10081

48%
depression
In space



7%
depression
on polar
NMs



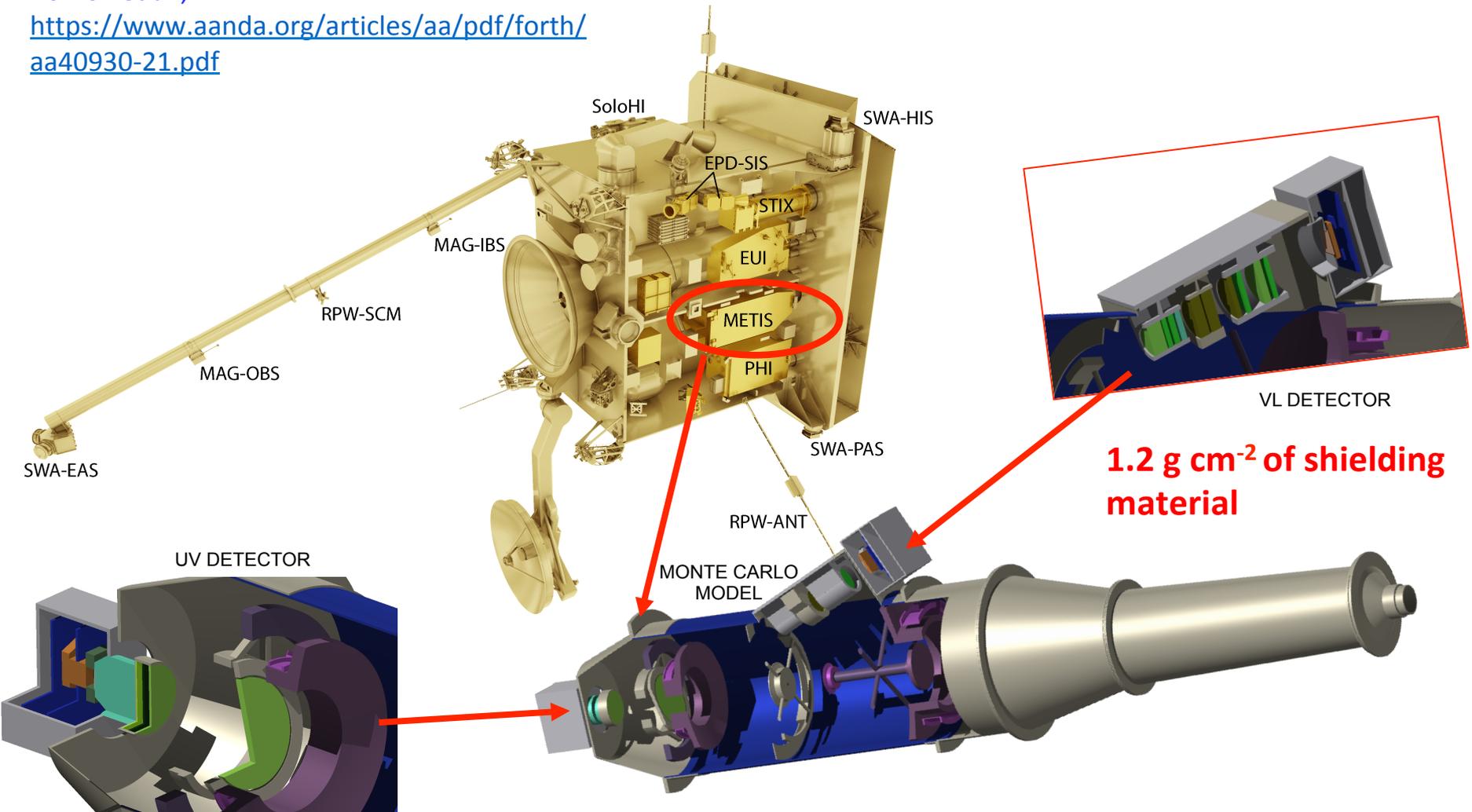
Adriani et al. 2011,
ApJ, 742, 102
CG et al.; ApJ, 904,
64, 2020

Solar Orbiter and Metis

Antonucci et al., A&A 642 (A10), 2020

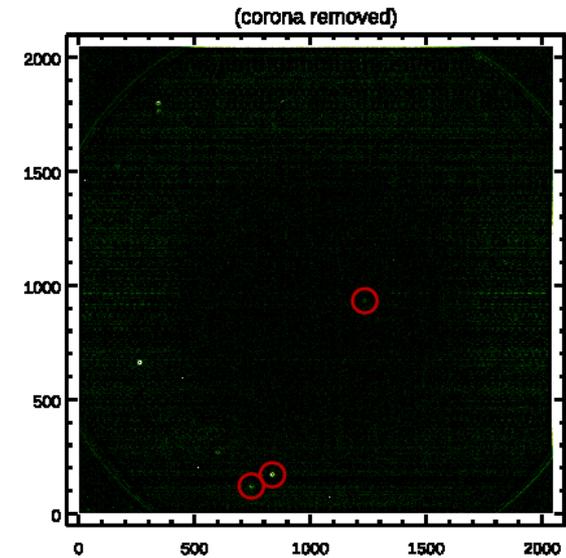
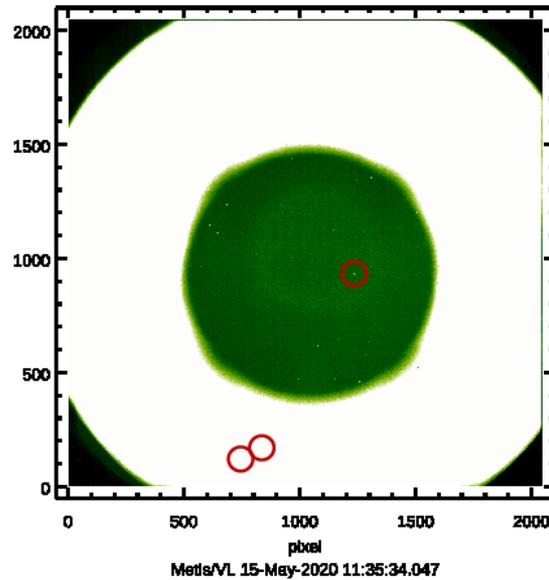
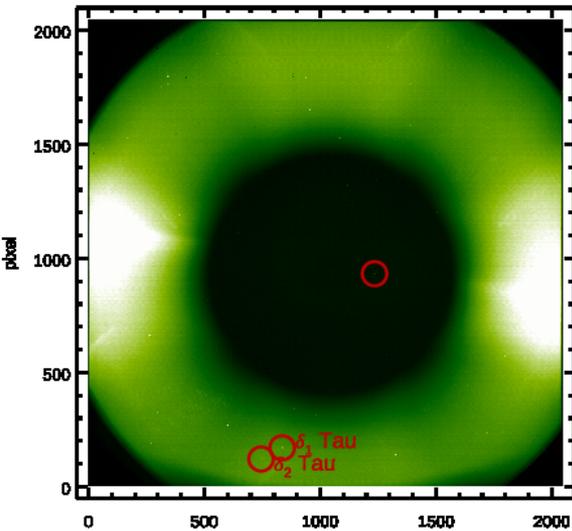
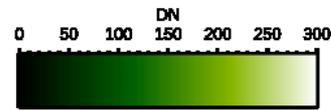
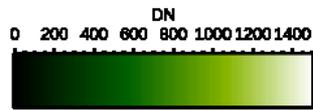
Romoli et al.;

<https://www.aanda.org/articles/aa/pdf/forth/aa40930-21.pdf>



Cosmic-ray tracks in the Metis VL detector

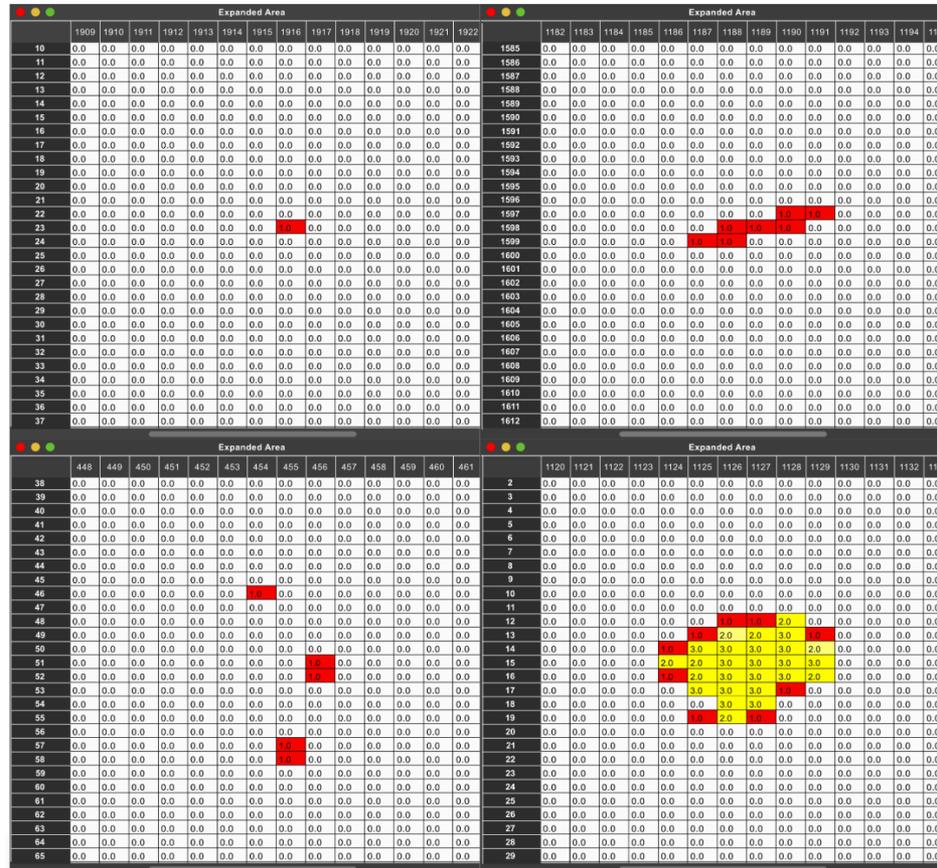
$\lambda = 580 - 640$ nm



Courtesy of V. Andretta

The APviewer for cosmic-ray analysis in the Metis VL images

Single track



Single slant track

Multiple tracks

No cosmic ray

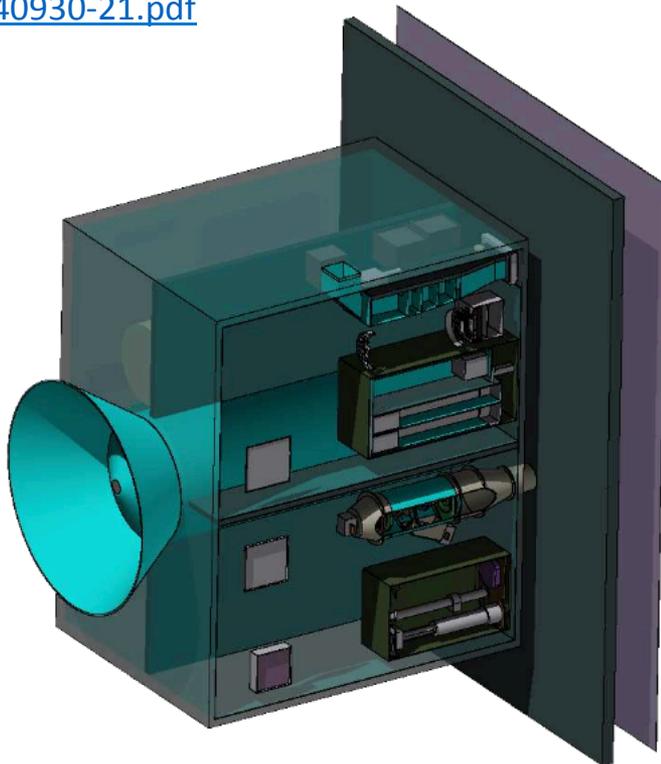
METIS MONTE CARLO SIMULATIONS WITH FLUKA

FLUKA FLAIR SoIo SIMULATION GEOMETRY

CG et al.;

<https://arxiv.org/abs/2104.13700>

<https://www.aanda.org/articles/aa/pdf/forth/aa40930-21.pdf>



← EUI

← METIS

← PHI

FLUKA SIMULATION OUTCOMES:

- Approximately 50% of incident protons are stopped in the spacecraft material surrounding Metis
- The spurious pixels fired by cosmic rays are 10^{-4} of the total sample in the VL detector
- The multiple tracks are not secondaries
- The observed number of tracks is compatible with primary and secondary particles generated by protons only (IMP!!!)
- Metis may play the role of a proton monitor

Thanks to the PHI and EUI Collaborations

Circumpulsar/magnetar disc formation

- Coplanar planets surround the millisecond pulsar PSR 1257+12 (Wolszczan & Frail 1992)
- A fallback disc was detected around the anomalous X-ray pulsar (AXP) 4U 0142+61 (Wang, Chakrabarty & Kaplan 2006)
- The presence of a disc around the RRAT B0656+14 was reported in Perna, Hernquist & Narayan (2000). In Grimani (2013) it was pointed out that a precessing disc would have explained the transient emission from this pulsar
- The hypothesis of a disc surrounding Crab just beyond the LC was discussed in Menou, Perna & Hernquist (2001)
- The presence of a disc may be responsible of pulsar spin down
- The presence of a disc may also affect high-energy electron and positron propagation in the pulsar and magnetar near environment

The case of the AXP 4U 0142+61

- In addition to the disc around the AXP 4U 0142+61 detected by Spitzer (Wang, Chakrabarty & Kaplan 2006; Ertan et al. 2007)...
- ...the **Nuclear Spectroscopic Telescope Array (NuSTAR)** experiment detected an X-ray emission modulated at **55 ks, or half this value,** from the same AXP (Makishima et al. 2018).
- It is shown here that this modulation is compatible with the presence of a Keplerian precessing disc with the characteristics observed by Wang et al. (2006) hiding periodically the emission region.

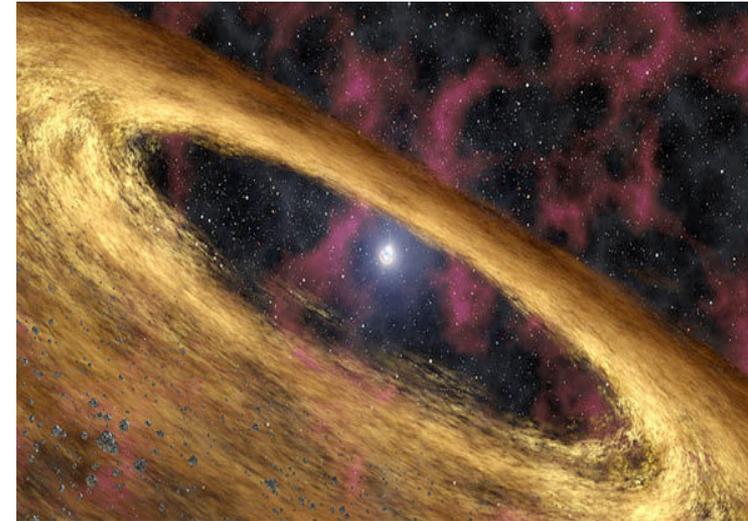
Characteristics of AXP 4U 0142+61

- Mass: 2.8×10^{30} kg
- Radius: 10 km
- Period: 8.68917 s
- Magnetic field: $(4.75 \pm 0.02) \times 10^{14}$ G
- Blackbody temperature: 3.59×10^6 K
- Blackbody radiation: 0.309 ± 0.001 keV
- Distance from Earth: 3.6 kpc
- Momentum of inertia: 1.12×10^{38} kg m²
- Angular momentum: 8.10×10^{37} kg m² s⁻¹

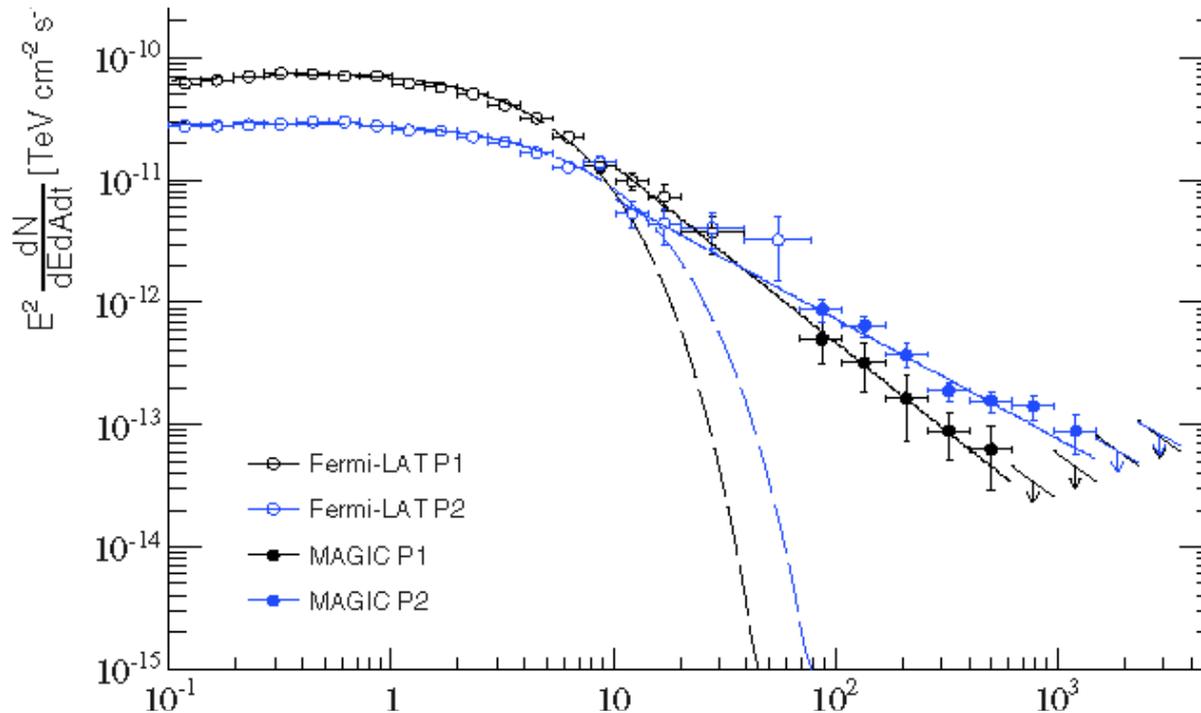
Characteristics of the disc around AXP 4U 0142+61

<https://www.spitzer.caltech.edu/video/ssc2006-10v1-birth-of-phoenix-planets>

- Inner radius (R_i): 2.02×10^9 m
- Outer radius (R_o): 6.75×10^9 m
- Disc height: $0.035 R_i$
- Mass: 5.97×10^{25} kg
- Blackbody Temperature: 920 K
- Blackbody $\lambda_{\max} = 3.15 \times 10^{-6}$ m (mid-infrared)
- Disc temperature compatible with 5% absorption of the magnetar illuminating energy
- Momentum of inertia: $I_D = 1.48 \times 10^{45}$ kg m²
- Angular momentum: $L_D = 2.23 \times 10^{41}$ kg m² s⁻¹



Pulsed photon observations from Crab with MAGIC



Ansoldi et al., A &A, 585, A133 (2016)

IC scattering dominates above 50 GeV. Aharonian et al. suggested this takes place at 50-60 light cylinder radius distance (10^7 - 10^8 m; when measurements were available up to 400 GeV): may infrared emission from a disc possibly contribute in limiting the propagation of TeV electrons and positrons?

James Webb Space Telescope

- The James Webb Space Telescope (JWST; Gardner et al. 2006), scheduled to launch in December 2021, will definitely open the infrared observational window on disc formation around magnetars and pulsars.
- Gravitational wave emission from precessing discs may also help in detecting these discs

4U 0142+61: disc precession or magnetar precession forced by the disc?

- The NuSTAR observations are compatible with a precessing Keplerian disc
- Keplerian frequency: $V = R_l \Omega_K$ $V = [(GM_{\text{mag}})/R_l]^{1/2}$
- $V = 3.0 \times 10^5 \text{ m s}^{-1}$ and $\Omega_K = 1.505 \times 10^{-4} \text{ s}^{-1}$
- $\omega_{\text{mag}} = 0.723 \text{ s}^{-1}$
- $P = 2\pi/\phi$ with $\phi = (45 M_D \cos \theta' \omega_{\text{mag}})/(32 \pi \rho R_l^3)$

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where $\theta' = 15^\circ$, $\rho = 10^{17} \text{ kg m}^{-3}$ for the NS density with θ' representing the angle formed by the rotation axis of the magnetar and the perpendicular to the disc plane

$P = 2.77 \times 10^{20} \text{ s}$ different from NuSTAR observations: this scenario would work only for heavy discs formed near young pulsars

May be a free magnetar precession?

The magnetar free precession would imply a magnetic field of the magnetar of 10^{16} G (Makishima et al.; 2018) that however, is not in agreement with observations: $(4.75 \pm 0.02) \times 10^{14}$ G (Guver et al., 2008)

Gravitational wave emission from precessing discs

Table 1. Gravitational wave amplitude generated by precessing discs and detection distances for $\theta = 10^\circ$. The signal-to-noise ratio (S/N) is inferred from the *LISA* (below 10^{-3} Hz) and *BBO* (above 10^{-3} Hz) sensitivities reported in Fig. 2. Maximum distances for precessing disc detection, when feasible, are set for S/N = 1. The detection probability would increase with the square root of the observational time.

	Disc mass (kg)	$\omega_1(\omega_2)$ (Hz)	$h_{\omega_1}(h_{\omega_2})$	S/N $_{\omega_1}$ (S/N $_{\omega_2}$)	Distance (kpc)
4U 0142+61	5.97×10^{25}	$3.1 \times 10^{-4}(6.2 \times 10^{-4})$	$1.1 \times 10^{-26}(4.5 \times 10^{-26})$	$\ll 1(\ll 1)$	0.05(0.05)
Crab	5.97×10^{25}	$3.9 \times 10^{-2}(7.8 \times 10^{-2})$	$6.7 \times 10^{-26}(2.7 \times 10^{-25})$	$\ll 1(< 1)$	0.05(0.05)
	10^{27}	$3.9 \times 10^{-2}(7.8 \times 10^{-2})$	$1.1 \times 10^{-24}(4.5 \times 10^{-25})$	$< 1(1)$	0.05(0.5)
	10^{28}	$3.9 \times 10^{-2}(7.8 \times 10^{-2})$	$9.3 \times 10^{-25}(4.8 \times 10^{-25})$	1(1)	0.6(4.65)
	10^{29}	$3.9 \times 10^{-2}(7.8 \times 10^{-2})$	$8.7 \times 10^{-25}(4.5 \times 10^{-25})$	1(1)	6.8(50.0)

Table 2. Same as Table 1, but for $\theta = 30^\circ$. For masses of 10^{29} kg, in principle, disc precession detection may extend to the whole Galaxy (WG).

	Disc mass (kg)	$\omega_1(\omega_2)$ (Hz)	$h_{\omega_1}(h_{\omega_2})$	S/N $_{\omega_1}$ (S/N $_{\omega_2}$)	Distance (kpc)
4U 0142+61	5.97×10^{25}	$3.5 \times 10^{-4}(7.0 \times 10^{-4})$	$1.9 \times 10^{-25}(4.8 \times 10^{-25})$	$\ll 1(\ll 1)$	0.05(0.05)
Crab	5.97×10^{25}	$4.4 \times 10^{-2}(8.8 \times 10^{-2})$	$7.2 \times 10^{-25}(4.8 \times 10^{-25})$	1(1)	0.05(0.3)
	10^{27}	$4.4 \times 10^{-2}(8.8 \times 10^{-2})$	$8.0 \times 10^{-25}(4.2 \times 10^{-25})$	1(1)	0.75(5.7)
	10^{28}	$4.4 \times 10^{-2}(8.8 \times 10^{-2})$	$8.1 \times 10^{-25}(4.2 \times 10^{-25})$	1(1)	7.5(57.1)
	10^{29}	$4.4 \times 10^{-2}(8.8 \times 10^{-2})$	$7.1 \times 10^{-25}(4.1 \times 10^{-25})$	1(1)	84.5(WG)

CG, CQG, 2009; CG Adv. Sp. Res. 2013; CG MNRAS, 2016; CG MNRAS, 2021

Electron and positron propagation in the region of the disc formation

- Blackbody photon density near a disc:
- $N/V = 8 \pi (kT/hc) \times 2.405$
with $V = 1 \text{ cm}^3$ $T = 920 \text{ K}$ $N = 1.57 \times 10^{10} \text{ photons cm}^{-3}$
and $T = 1800 \text{ K}$ $N = 1.18 \times 10^{11} \text{ photons cm}^{-3}$
- Average energy is 0.21 eV and 0.43 eV, respectively, in the near infrared region
- Pulsar blackbody radiation: $N = 1.15 \times 10^9 \text{ photons cm}^{-3}$
- The infrared radiation emission is approximately one order of magnitude smaller than that of the disc (Penny 1982)

Discs that may limit the propagation of TeV electrons and positrons

Klein-Nishina total cross section

$$\sigma_{\text{KN}} = \sigma_{\text{T}} \frac{3}{4} \left[\frac{1+x}{x^3} \left(\frac{2x(1+x)}{1+2x} - \ln(1+2x) \right) + \frac{1}{2x} \ln(1+2x) + \frac{1+3x}{(1+2x)^2} \right], \quad (19)$$

where $x = hv_0/(m_e c^2)$. Since $hv_0/(m_e c^2) > 1$ the above equation can be reduced as follows:

$$\sigma_{\text{KN}} = \sigma_{\text{T}} \frac{3}{8} \frac{1}{x} \left(\ln 2x + \frac{1}{2} \right).$$

Where hv_0 is the maximum photon energy in the electron rest frame (20)

The total cross section and associated collision length is $9.9 \times 10^{-30} \text{ m}^2$ and $8.6 \times 10^{11} \text{ m}$. Only discs of a.u. dimensions and hundreds of degrees temperature may limit the propagation of TeV electrons and positrons.

This is the case of the XTIN J0806.44123 for instance.

Conclusions

- Any astrophysical source being possibly at the origin of the positron excess in cosmic rays should be properly considered before evaluating the role of exotic sources
- Interplanetary physics of cosmic rays may give a major contribution to space weather science and mission diagnostics
- Fallback matter and disc formation near pulsars and magnetars may limit the maximum energy of electrons and positrons reaching the inner heliosphere
- Direct observations of discs through infrared emission are expected to be allowed by JWST. Future gravitational wave detection from precessing discs may help in solving the puzzle in the whole Galaxy