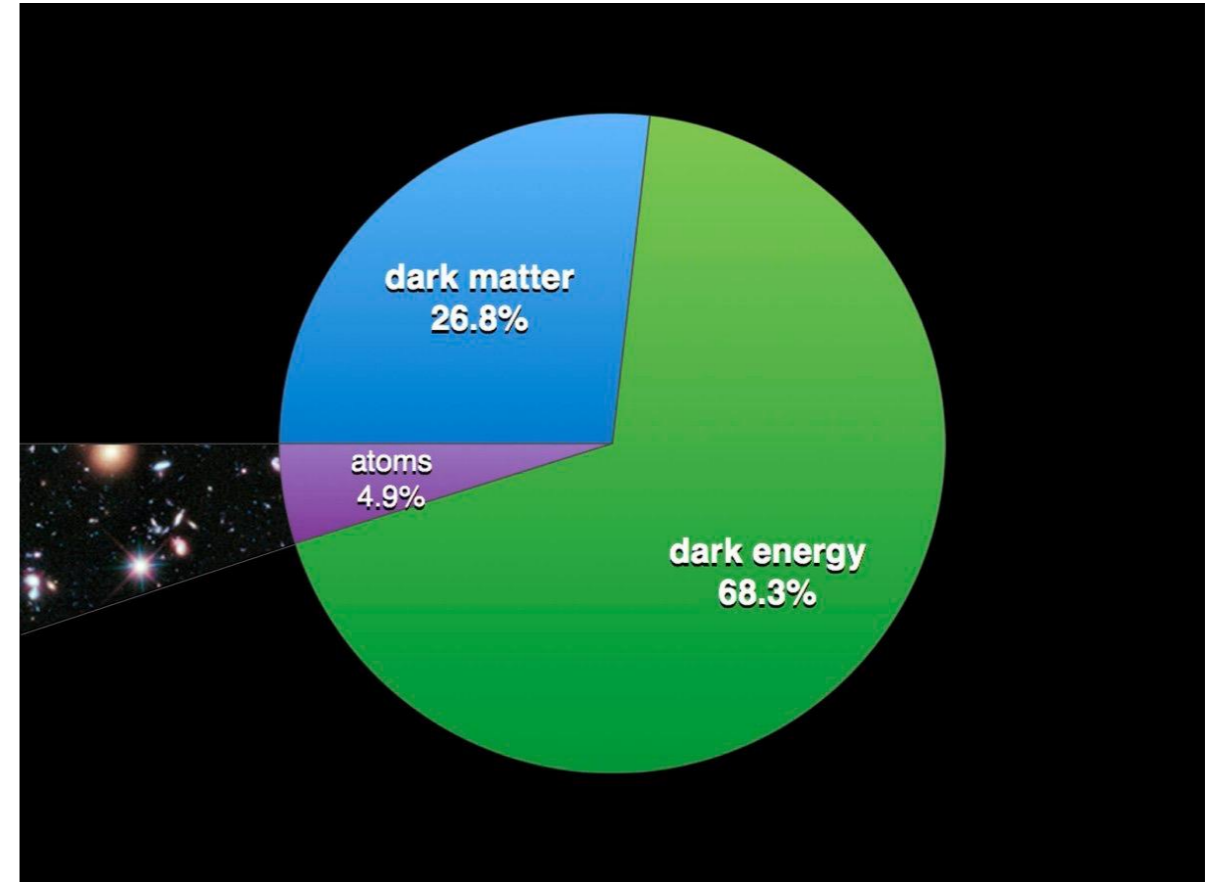


Fundamental properties of the dark and the luminous matter from Low Surface Brightness discs

Paolo Salucci



**Universe Webinar | keV Warm Dark Matter in Agreement
with Observations in Tribute to Héctor J. De Vega**

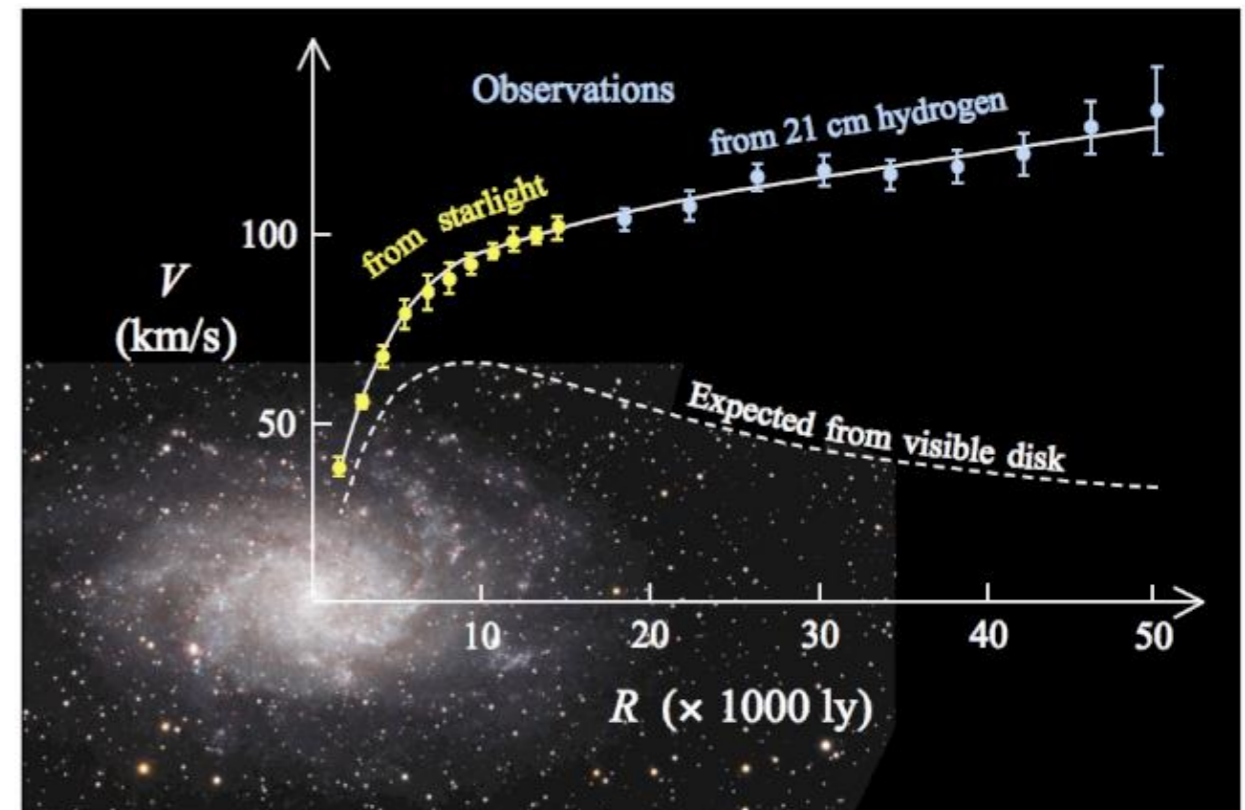
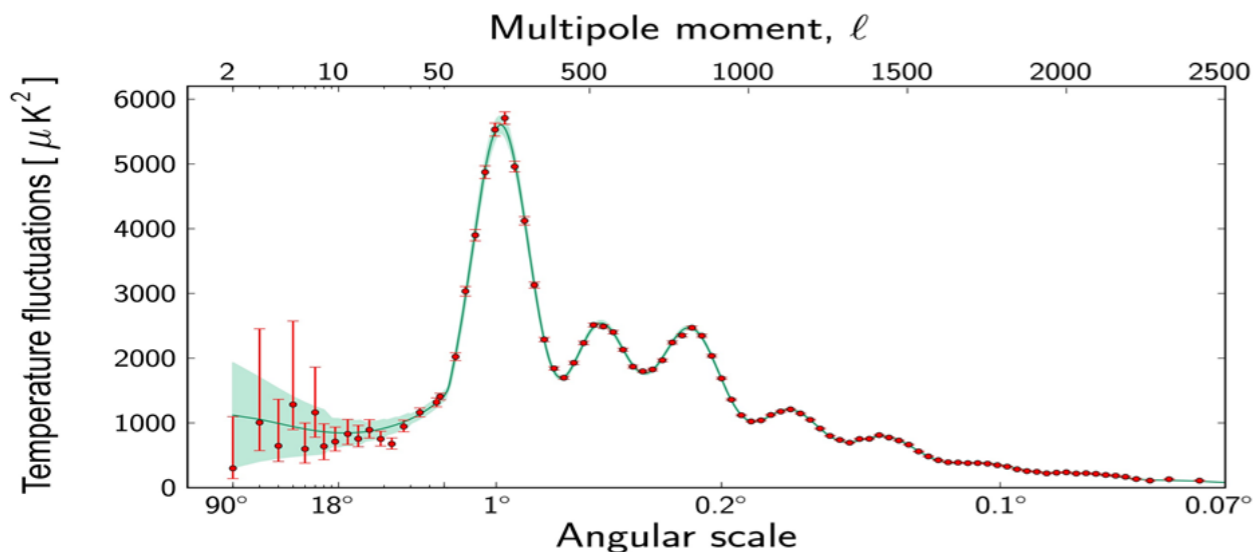
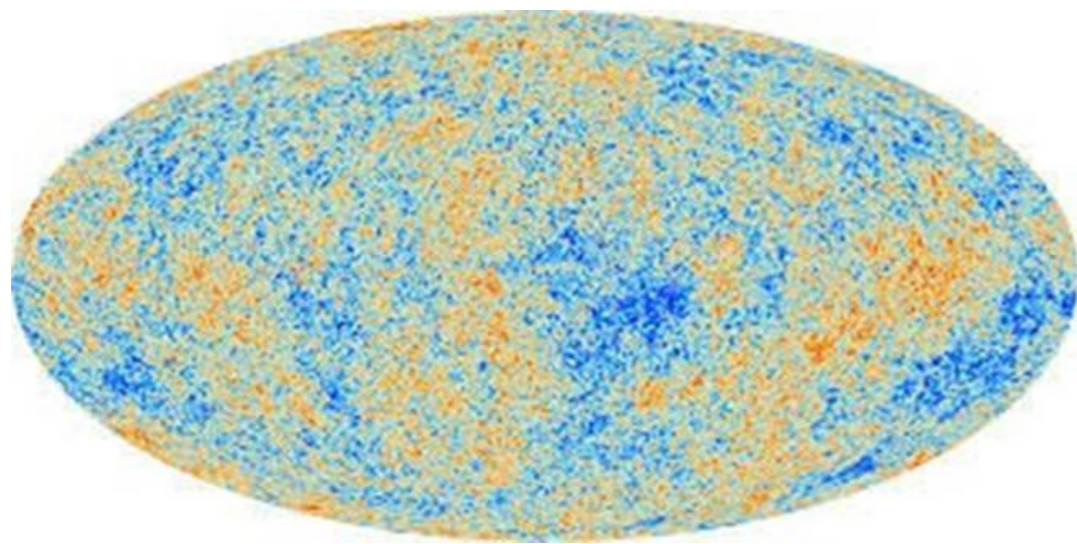


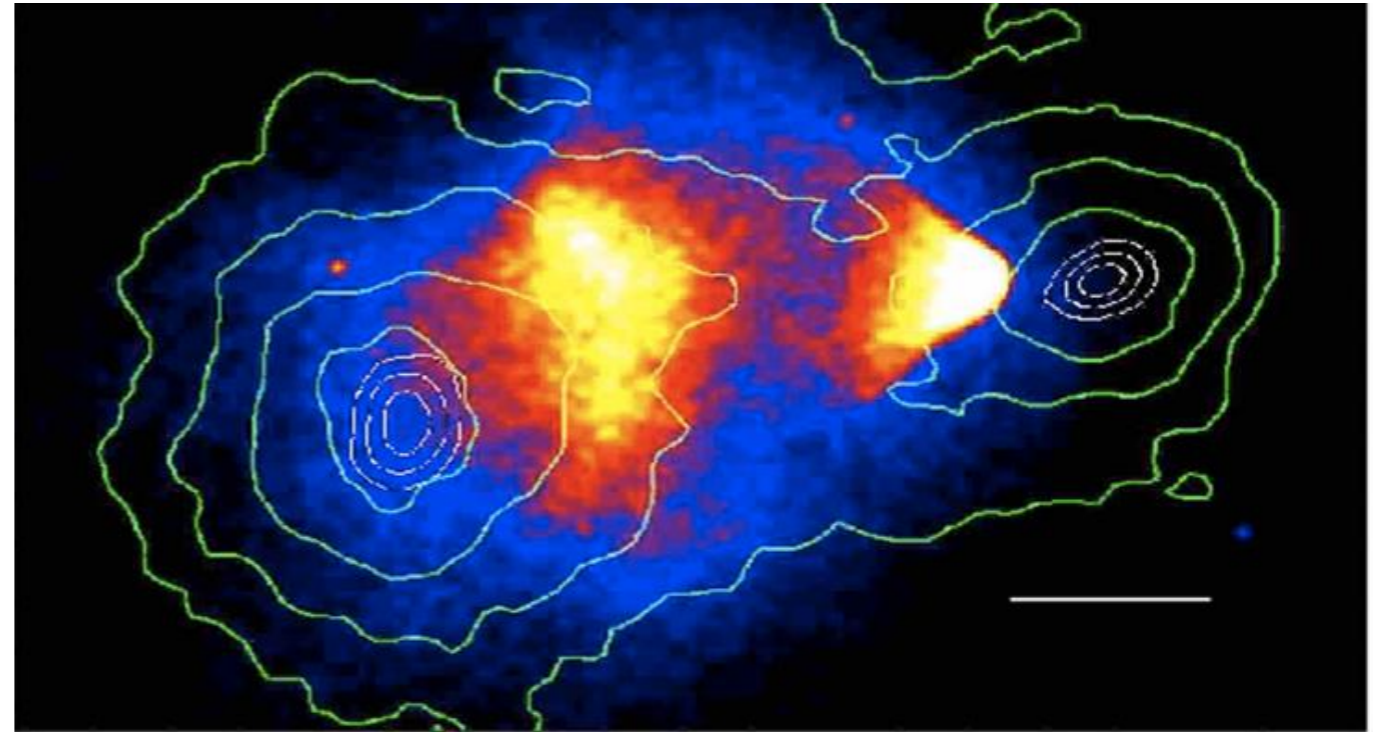
What we can observe by means of telescopes is the light emitted by stars, dust and gas, but they are only the tip of an iceberg.

Dark matter (DM) is a type of matter hypothesized to account for effects that appear to be the result of invisible mass.

The existence and properties of dark matter can be inferred from its gravitational effects on visible matter.

Many OBSERVATIONS:





What is the **NATURE** of dark matter (**DM**)?

It must barely interact with ordinary baryonic matter and radiation, except through gravity

It remains unknown

DM particle CANDIDATES

WIMP (Weakly interacting massive particles)

Relic DM particle from the early Universe

- DM abundance in the Universe
 - Interaction via the electroweak force
- 
- 100 GeV
mass range

with non-relativistic velocities
since its decoupling time

Clump in small structures (galaxies) and
aggregation to form larger structures
(bottom-up theory)

Supersymmetric extensions of the standard model of particle physics predict a new particle with these properties

Cold Dark Matter (CDM)

Collisionless CDM

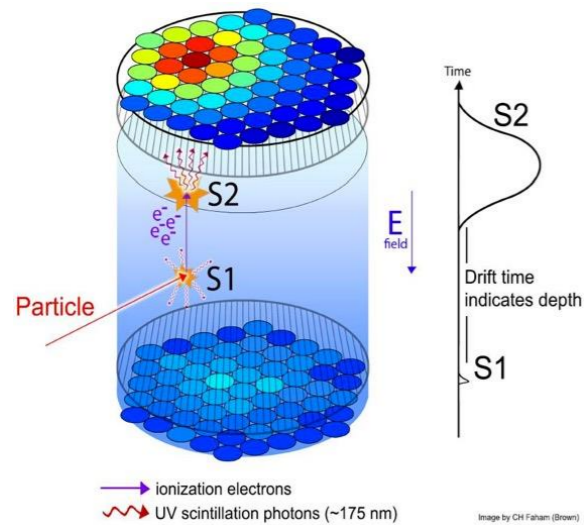
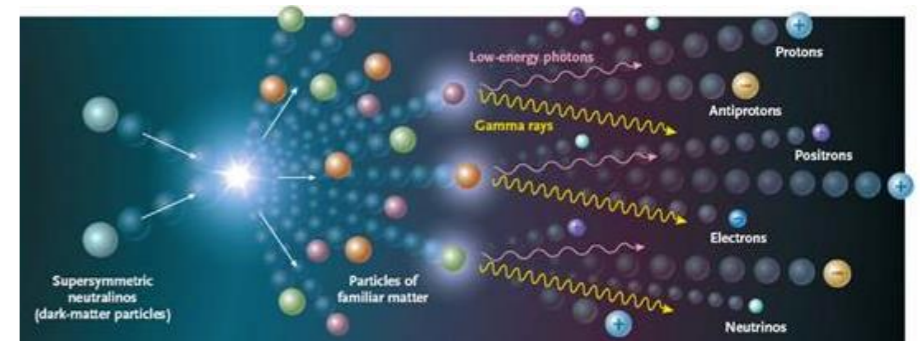
Λ CDM N-body simulations

CUSPY profile

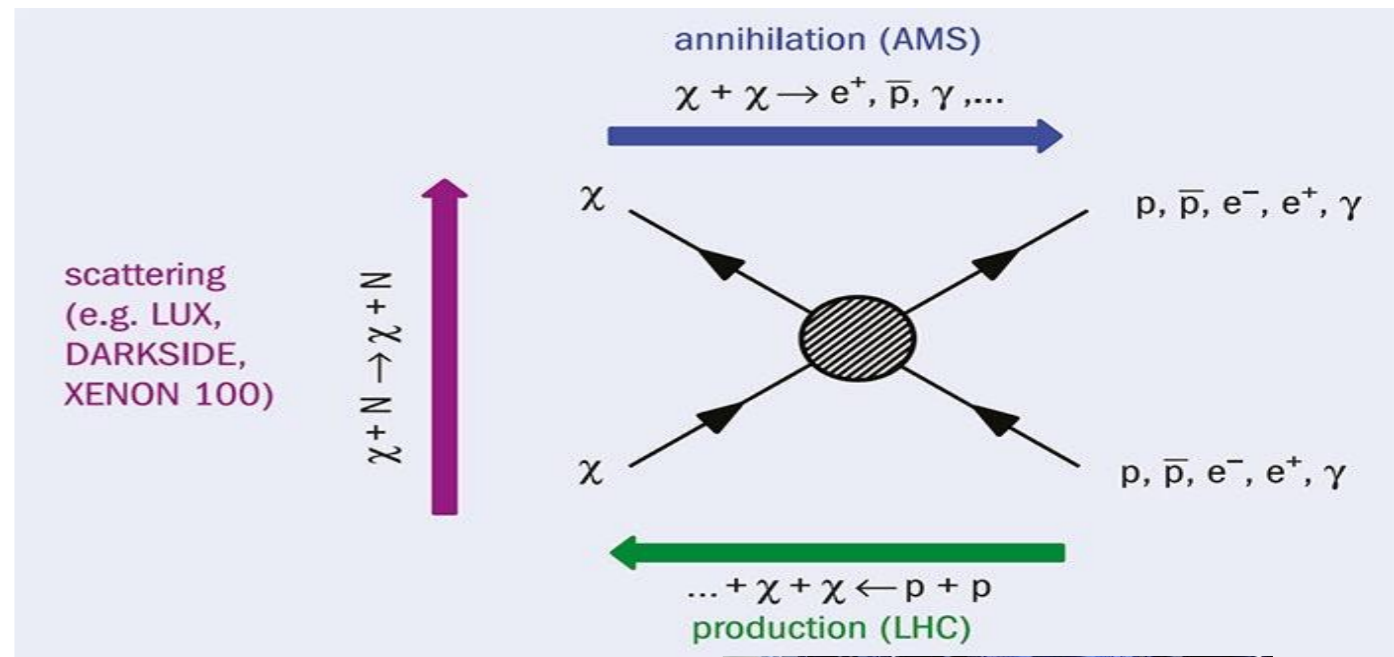
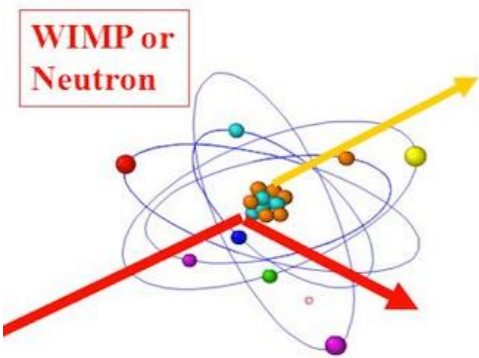
$$\rho_{NFW}(r) = \frac{\rho_s}{(r/r_s)(1 + r/r_s)^2}$$

WIMP particle has not yet been observed directly

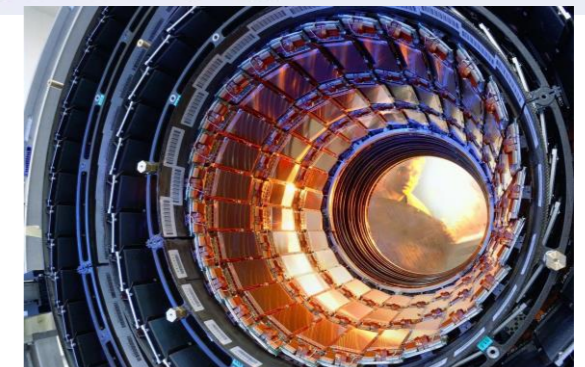
indirect
research



direct
research



colliders
production



LHC

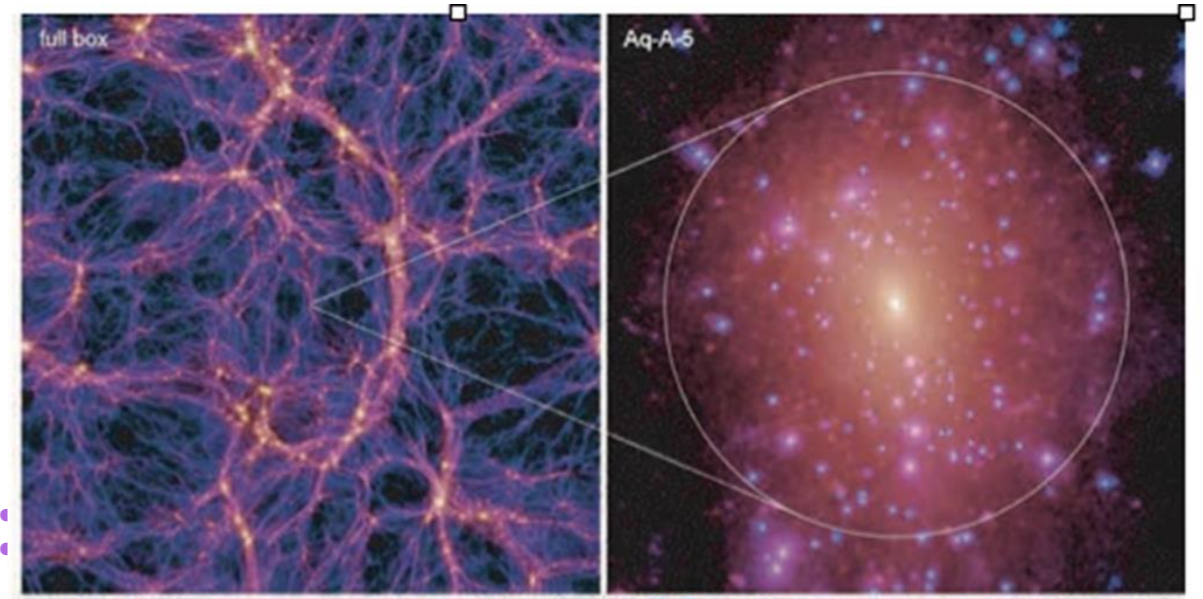
Issues with the main DM scenario

N-body simulations Λ CDM

Success on large scale structure $\gtrsim 50 \text{ kpc}$

missing satellite problem
too big to fail problem

Core - cusp problem in any galaxy:



- CDM



cusp

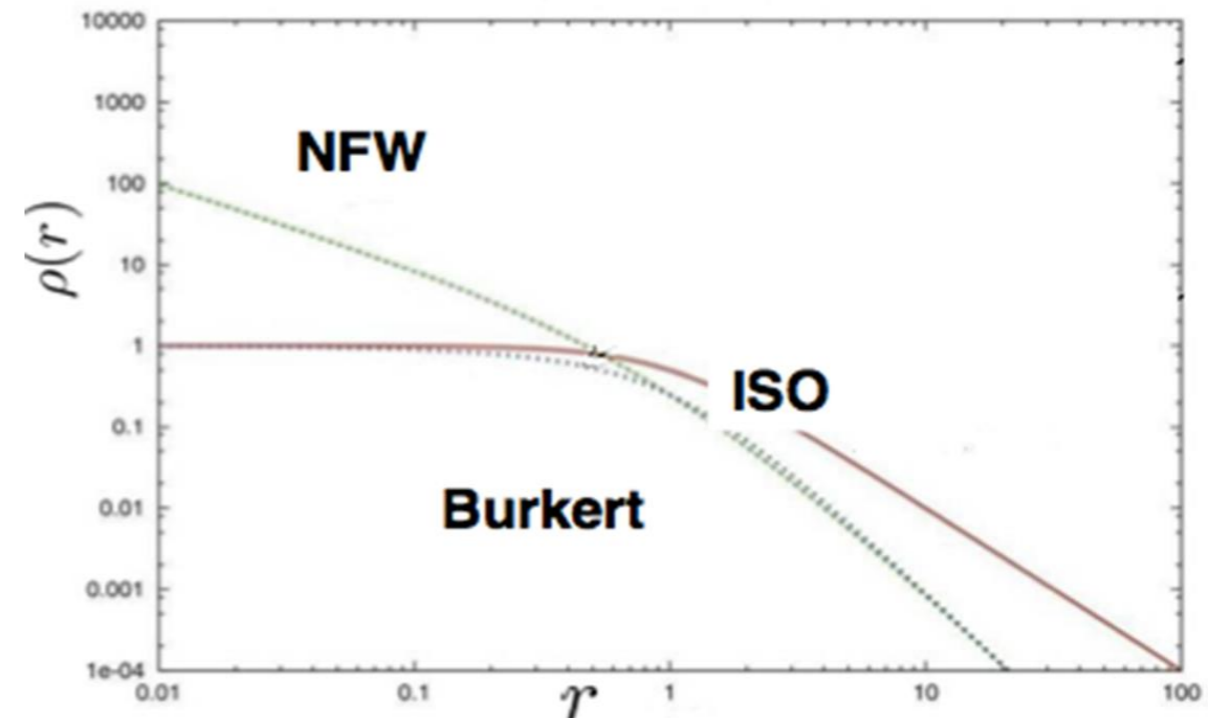
$$\rho(r) \sim r^{-1}$$

- Observations



core

$$\rho(r) = \text{const.}$$



Different DM particle CANDIDATE

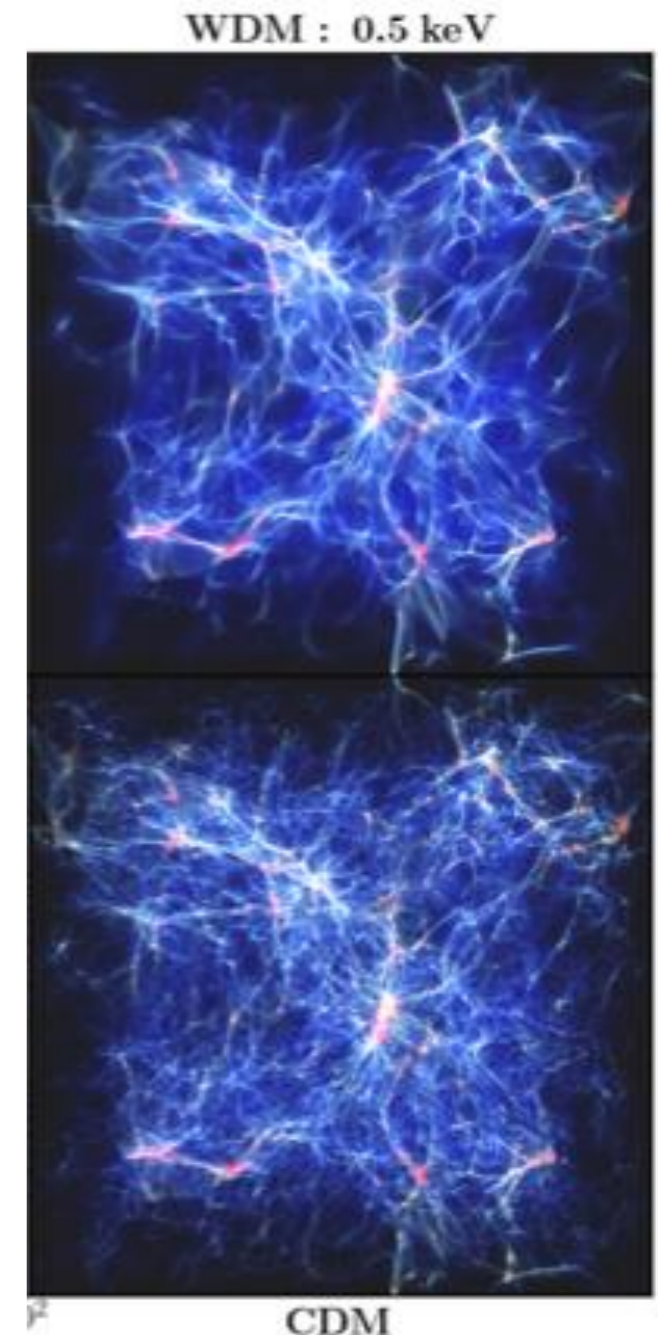
Warm dark matter particle (**WDM**)
keV mass range

- DM particles decouples from the cosmological plasma when it is still mildly relativistic.

It overcomes the problem on small scales

- WDM is fermionic a quantum pressure emerges

possibility of forming cored distribution



CDM Proposed solution:

baryonic matter feedback from supernovae explosions, stellar winds, gas outflow

However:

- challenges in dwarf and large spirals
- challenges in fine-tuned process



in Low Surface Brightness galaxies?

DM distribution in galaxies

Ellipticals

- dominated by random motions
- kinematics is very uncertain
- nuisance anisotropy parameter, which is font of degeneracy



Disc galaxies

- rotational supported systems
- rather simple kinematics

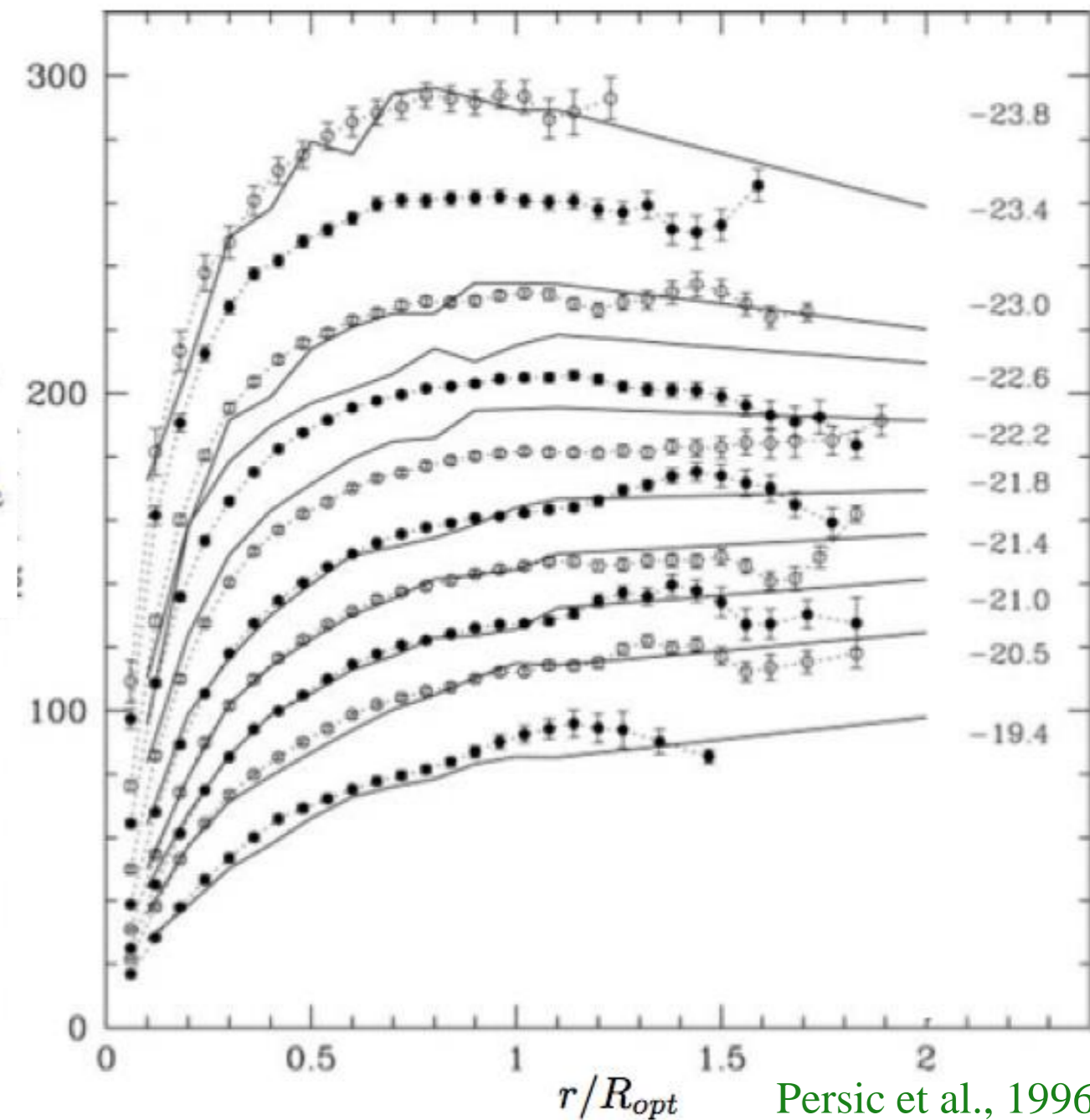
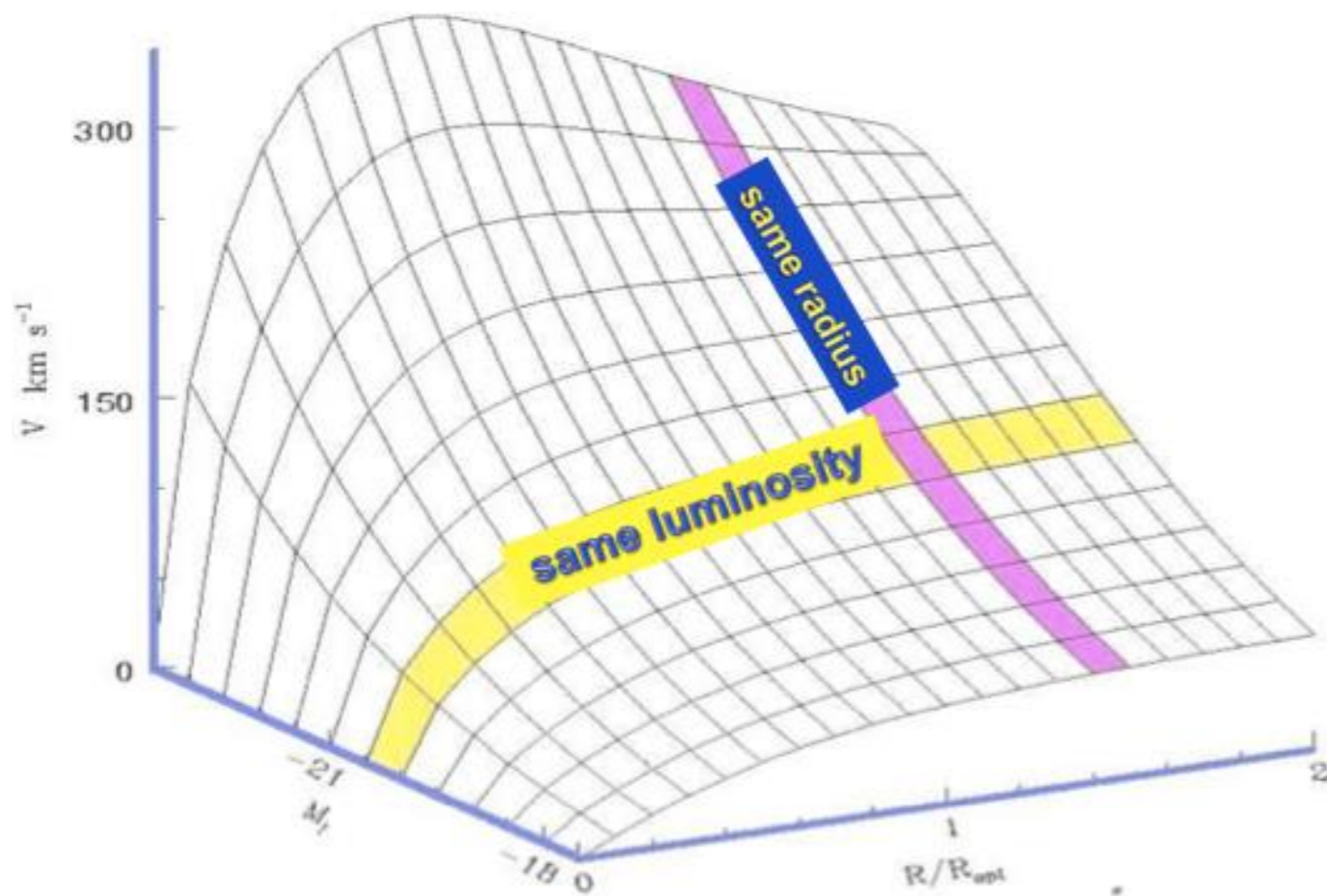
$$V_{tot}^2(r) = r \frac{d}{dr} \phi_{tot}(r) = V_d^2(r) + V_{HI}^2(r) + V_{bu}^2(r) + V_h^2(r)$$



Universal Rotation Curve

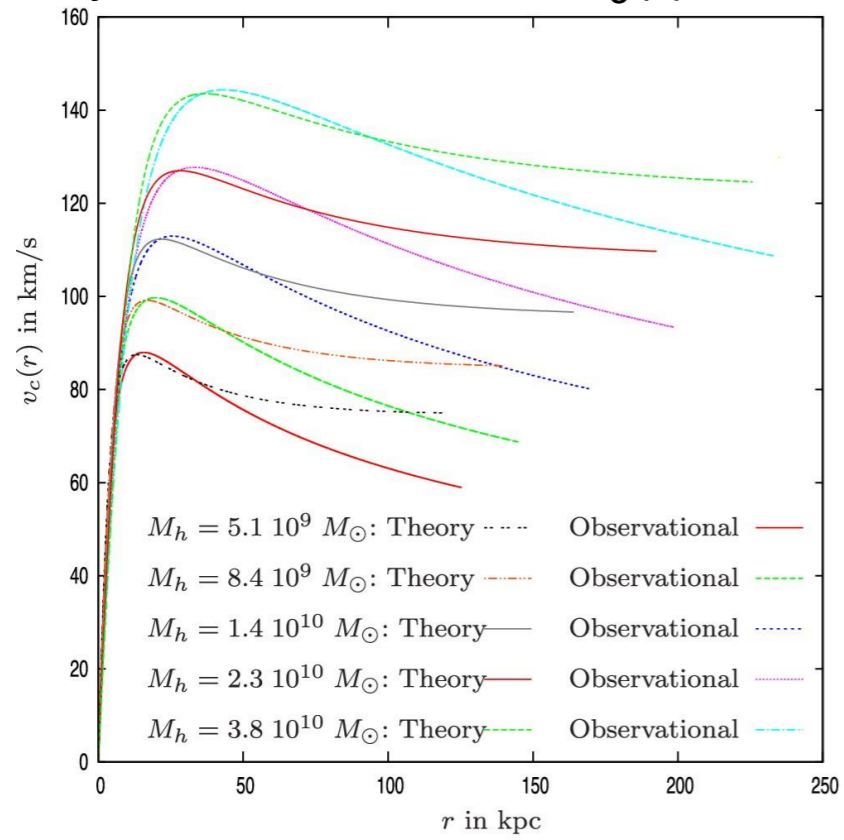
(URC)

stacked analysis



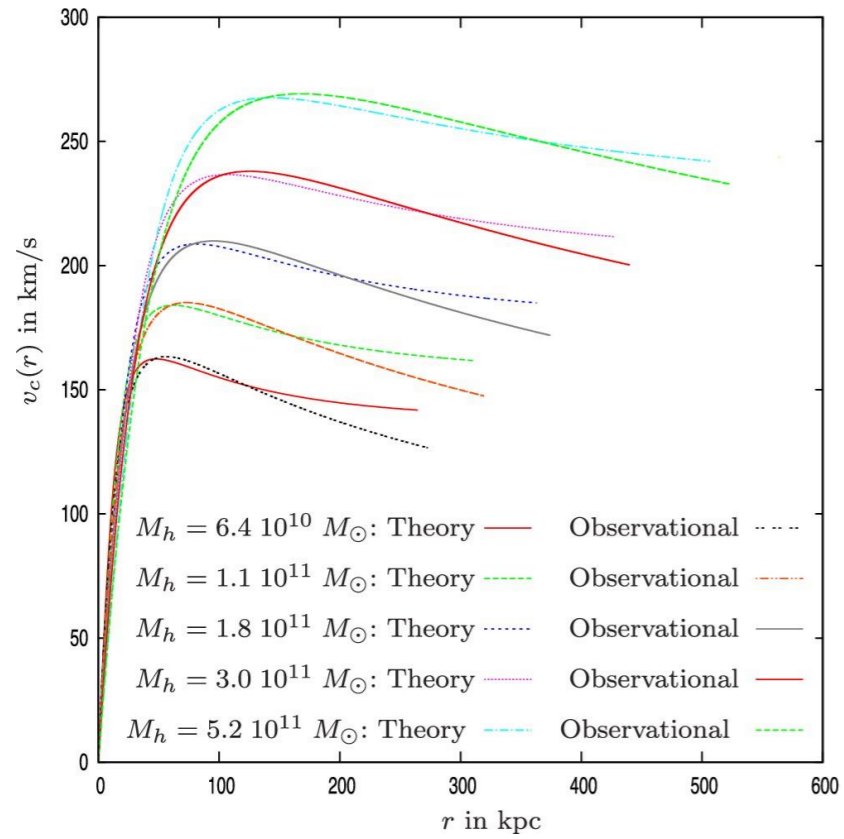
RCs gradual change
from small to large galaxies

The velocity rotation curves $v_c(r)$ in km s^{-1} .

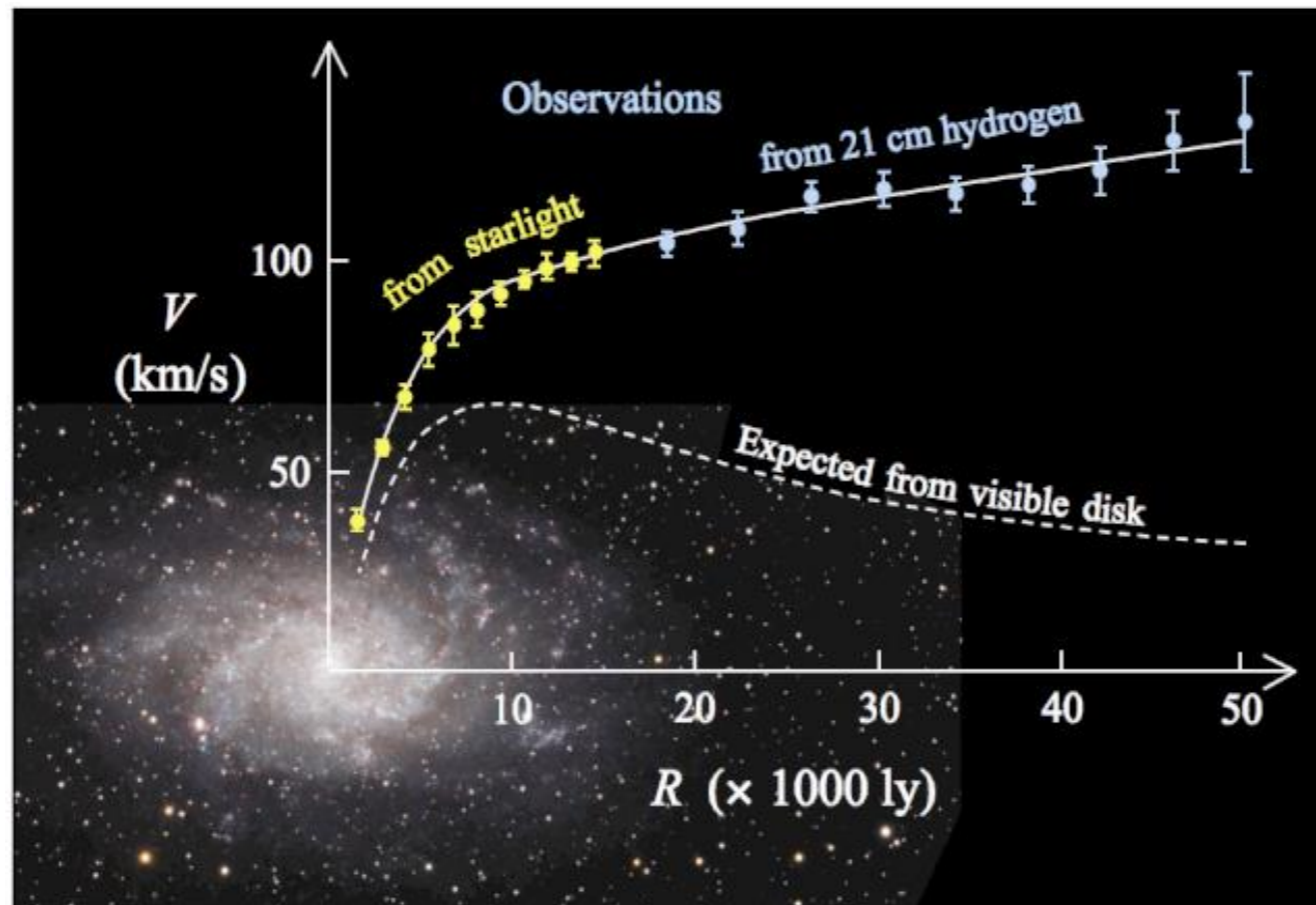


$$\frac{d^2 \mu}{dr^2} + \frac{2}{r} \frac{d\mu}{dr} = -4\pi G m \rho(r)$$

$$= -\frac{4 G m^2}{\pi \hbar^3} \int_0^\infty dp p^2 f\left(\frac{p^2}{2m} - \mu(r)\right),$$



DM distribution in disc galaxies



circular velocity
rotation curves
(RC)

Corbelli and Salucci, 2000

$$V_{tot}^2(r) = r \frac{d}{dr} \phi_{tot}(r) = V_d^2(r) + V_{HI}^2(r) + V_{bu}^2(r) + V_h^2(r)$$

Exponential
(Freeman)
stellar disc

$$\begin{cases} \mu(R) = \mu_0 e^{-R/R_d} \\ V_d^2(r) = \frac{1}{2} \frac{G M_d}{R_d} (3.2 r/R_{opt})^2 (I_0 K_0 - I_1 K_1) \end{cases}$$

$$V_{tot}^2(r) = r \frac{d}{dr} \phi_{tot}(r) = V_d^2(r) + V_{HI}^2(r) + V_{bu}^2(r) + V_h^2(r)$$

Exponential
(Freeman)
stellar disc



$$R_{opt} = 3.2 R_d$$

optical radius
size of the stellar disk
(83 % of total luminosity)

$$V_{tot}^2(r) = r \frac{d}{dr} \phi_{tot}(r) = V_d^2(r) + V_{HI}^2(r) + V_{bu}^2(r) + V_h^2(r)$$

gaseous
HI disc



$$\Sigma_{HI}(R) = \Sigma_{HI,0} e^{-R/3 R_d}$$

$$V_{HI}^2(R) = \frac{1}{2} \frac{GM_{HI}}{3R_D} (1.1 R/R_{opt})^2 (I_0 K_0 - I_1 K_1)$$

$$V_{tot}^2(r) = r \frac{d}{dr} \phi_{tot}(r) = V_d^2(r) + V_{HI}^2(r) + V_{bu}^2(r) + V_h^2(r)$$

stellar
bulge



$$V_{bu}^2(r) = \alpha_{bu} V_{in}^2 \left(\frac{r}{r_{in}} \right)^{-1}$$

$$V_{tot}^2(r) = r \frac{d}{dr} \phi_{tot}(r) = V_d^2(r) + V_{HI}^2(r) + V_{bu}^2(r) + V_h^2(r)$$

spherical
dark matter
halo



$$\rho_{DM}(r) = \frac{\rho_0 r_0^3}{(r + r_0)(r^2 + r_0^2)}$$

Burkert cored profile

$$M_{DM}(r) = \int_0^r 4\pi \tilde{r}^2 \rho_{DM}(\tilde{r}) d\tilde{r}$$

$$V_h^2(r) = G \frac{M_{DM}(r)}{r}$$

RC fit



best fit parameters



DM properties

Low Surface Brightness (LSB) disc galaxies

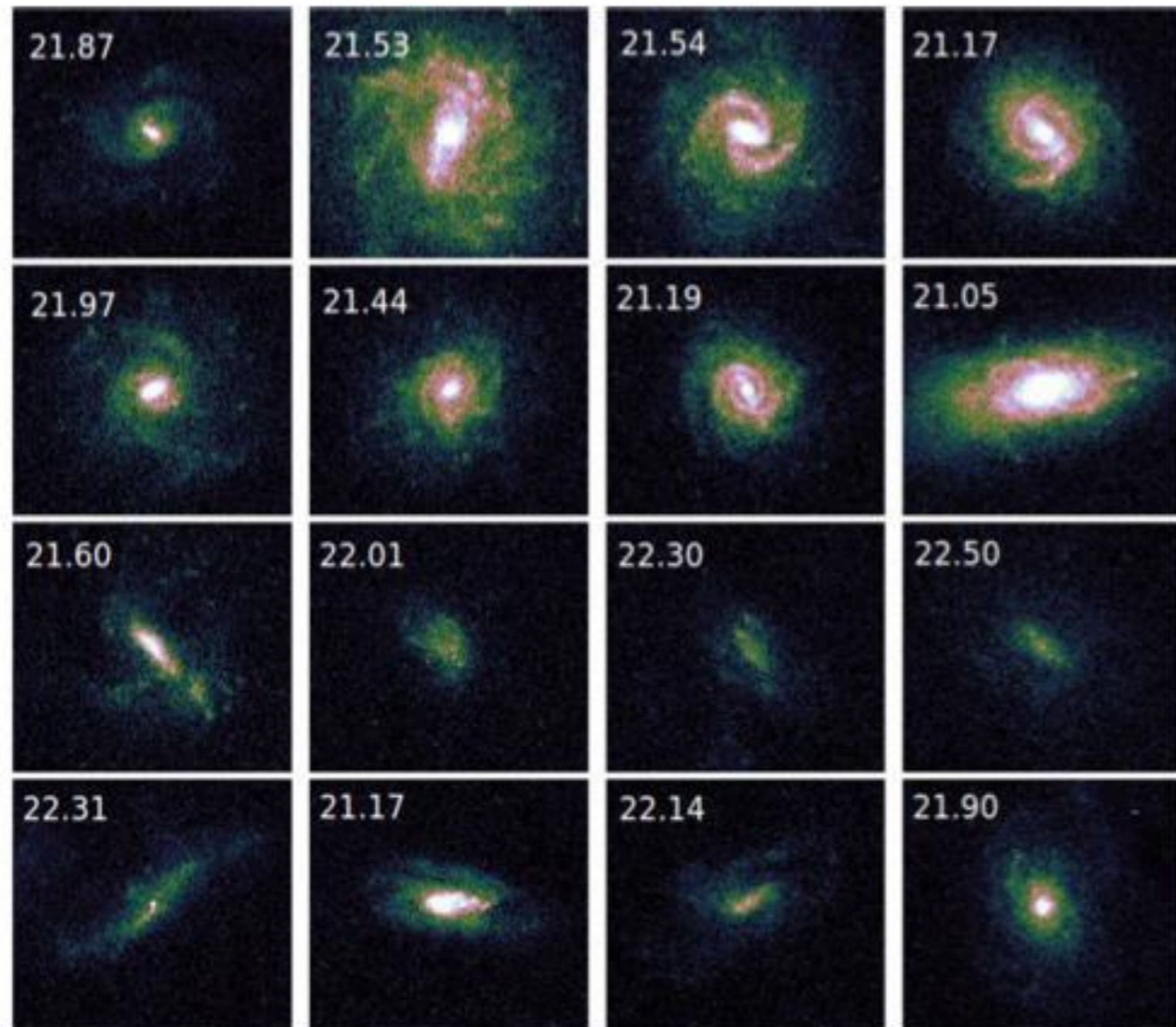
rotating disc systems
which emit an amount of
light per area smaller
than normal spirals
(4-10 times fainter)

$$\mu_{0,B} \gtrsim 23 \text{ mag arcsec}^{-2}$$

$$\mu_{0,R} \gtrsim 21 \text{ mag arcsec}^{-2}$$

diffuse, low-density
exponential stellar discs

$$\Sigma_* \simeq 12.3 M_{\odot}/pc^2$$



Low Surface Brightness (LSB) disc galaxies

full population of galaxies, ranging from small ($10^7 M_\odot$) to very large (more than $10^{10} M_\odot$) stellar disc mass M_d

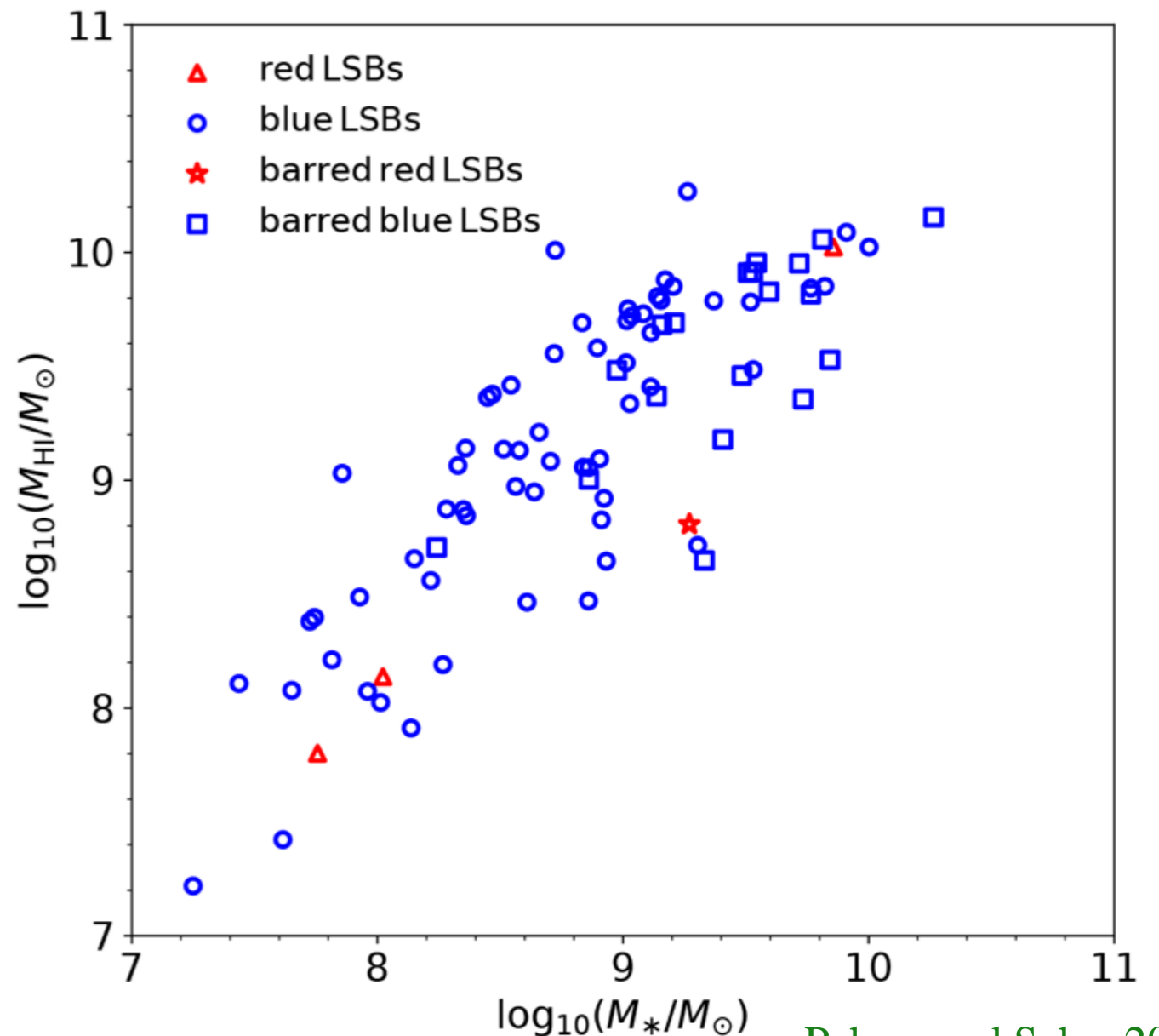
Typical size:

fraction of kpc $< R_d <$ tens of kpc

Typical magnitude:

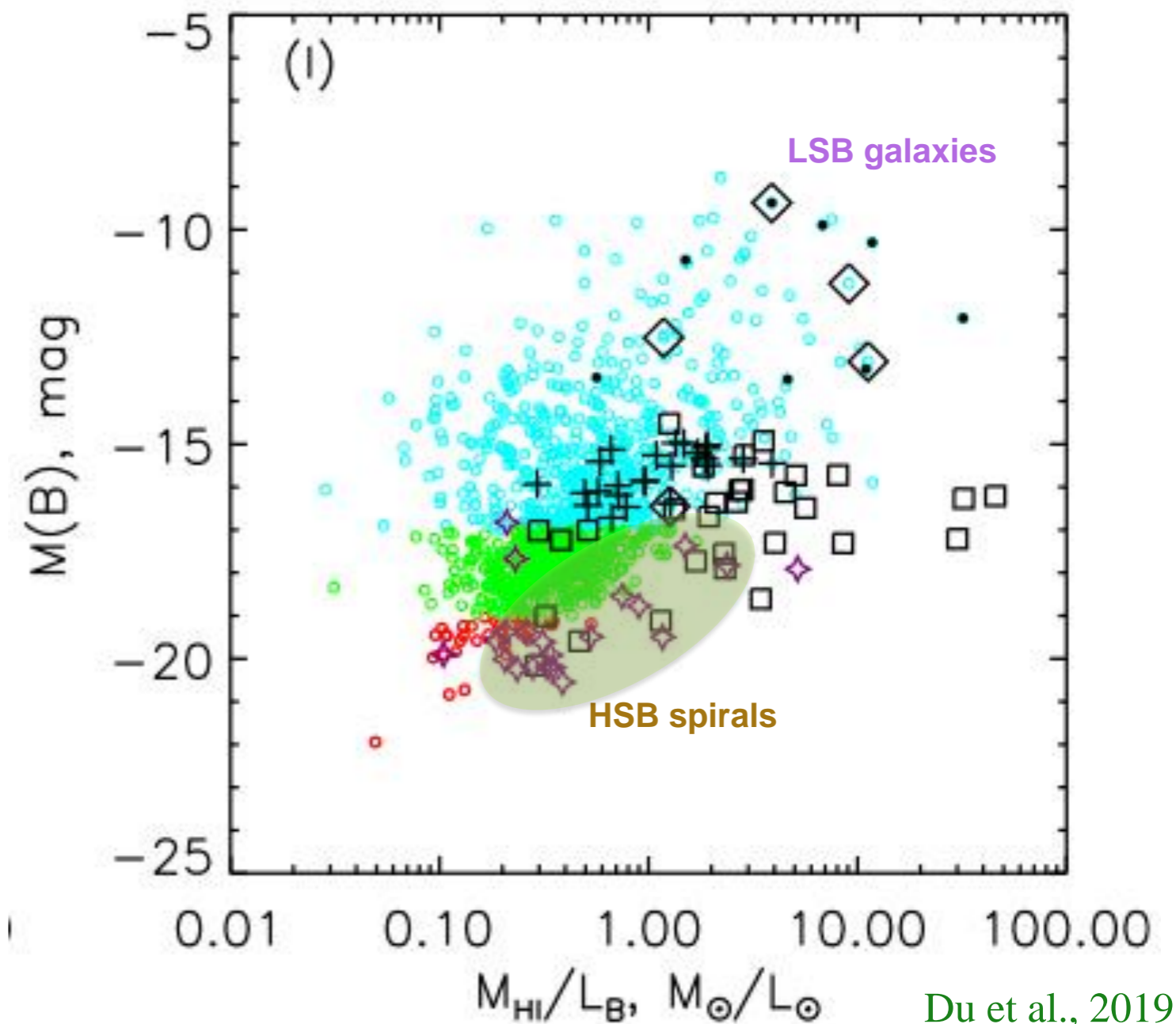
$$-20 \lesssim M_B \lesssim -10$$

$$-23 \lesssim M_R \lesssim -14$$



Low Surface Brightness (LSB) disc galaxies

Values of M_{HI}/L_B
in LSBs range
from $\simeq 0.1$ to $\simeq 10$



Typical LSBs gaseous disc

$$M_{HI} \simeq M_d$$

Characteristic low density

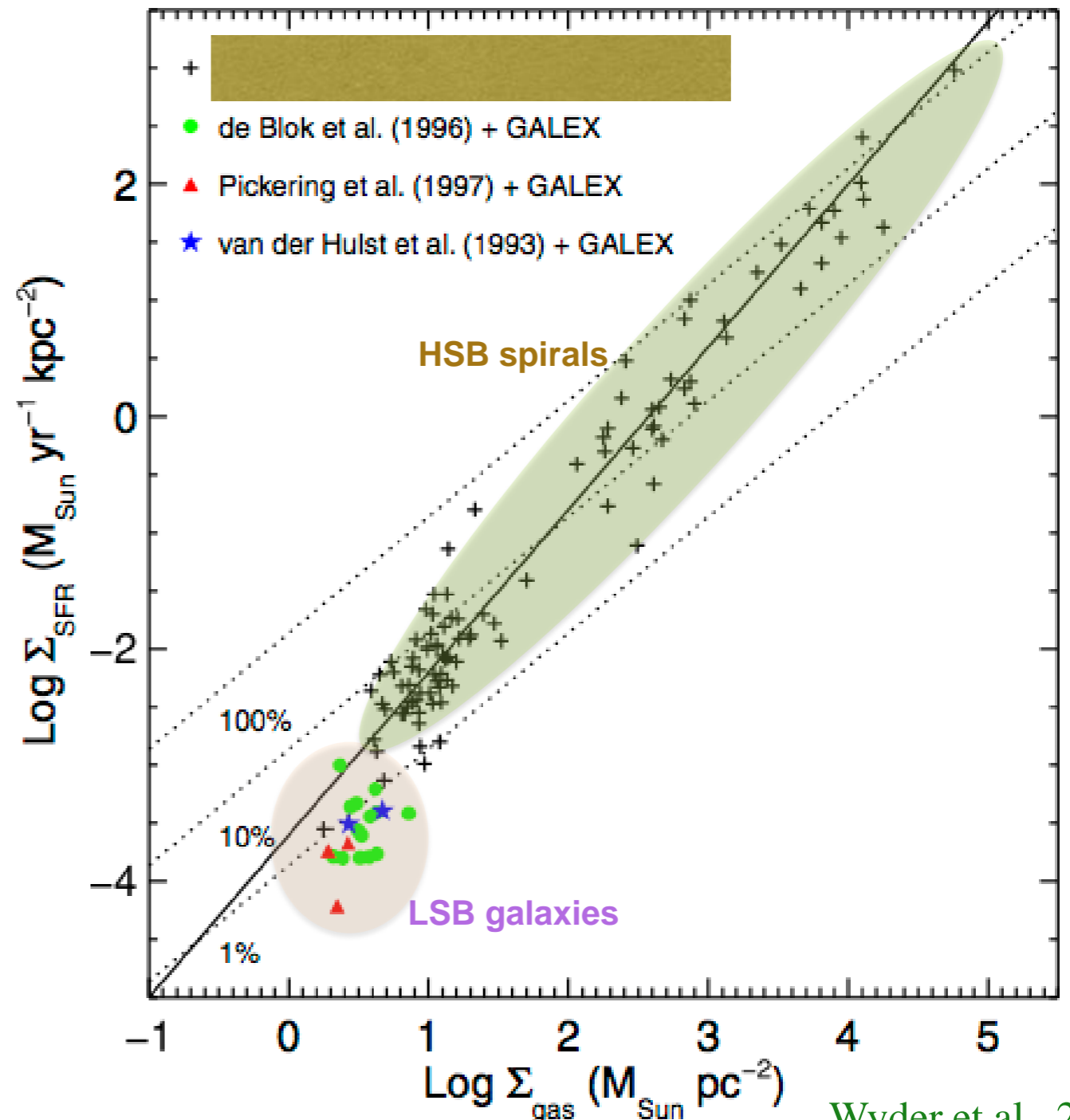
$$\Sigma_{HI} \simeq 5 M_{\odot} pc^{-2}$$



likely prevents
an efficient star formation

$$SFR \lesssim 0.1 M_{\odot} yr^{-1}$$

$$\Sigma_{SFR} \lesssim 10^{-3} M_{\odot} yr^{-1} kpc^{-2}$$



Low Surface Brightness (**LSB**) disc galaxies

very low content in metal and dust

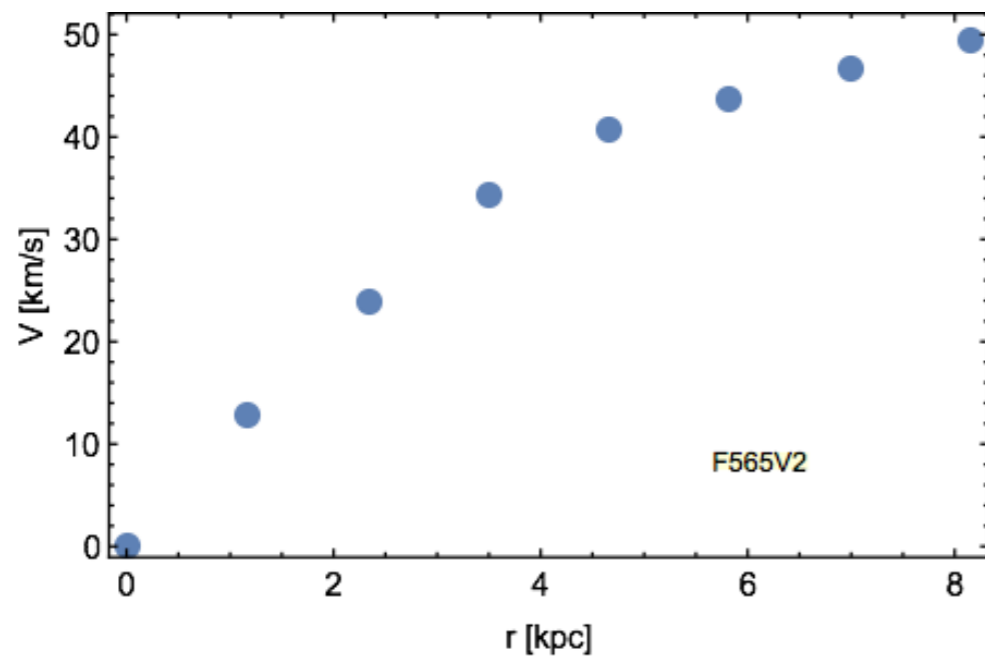
stellar population appears to be uniformly distributed in the stellar disc

generally isolated systems, located on the edges of large-scale structure

- LSBs are not the faded remnants of HSBs that have ceased star formation,
LSBs are likely **slowly evolving galaxies**
-

URC method applied to the LSB galaxies

LSB rotation curve



72 Low Surface Brightness galaxies
(Di Paolo, Salucci, Erkurt (2019))
1601 circular velocity measurements

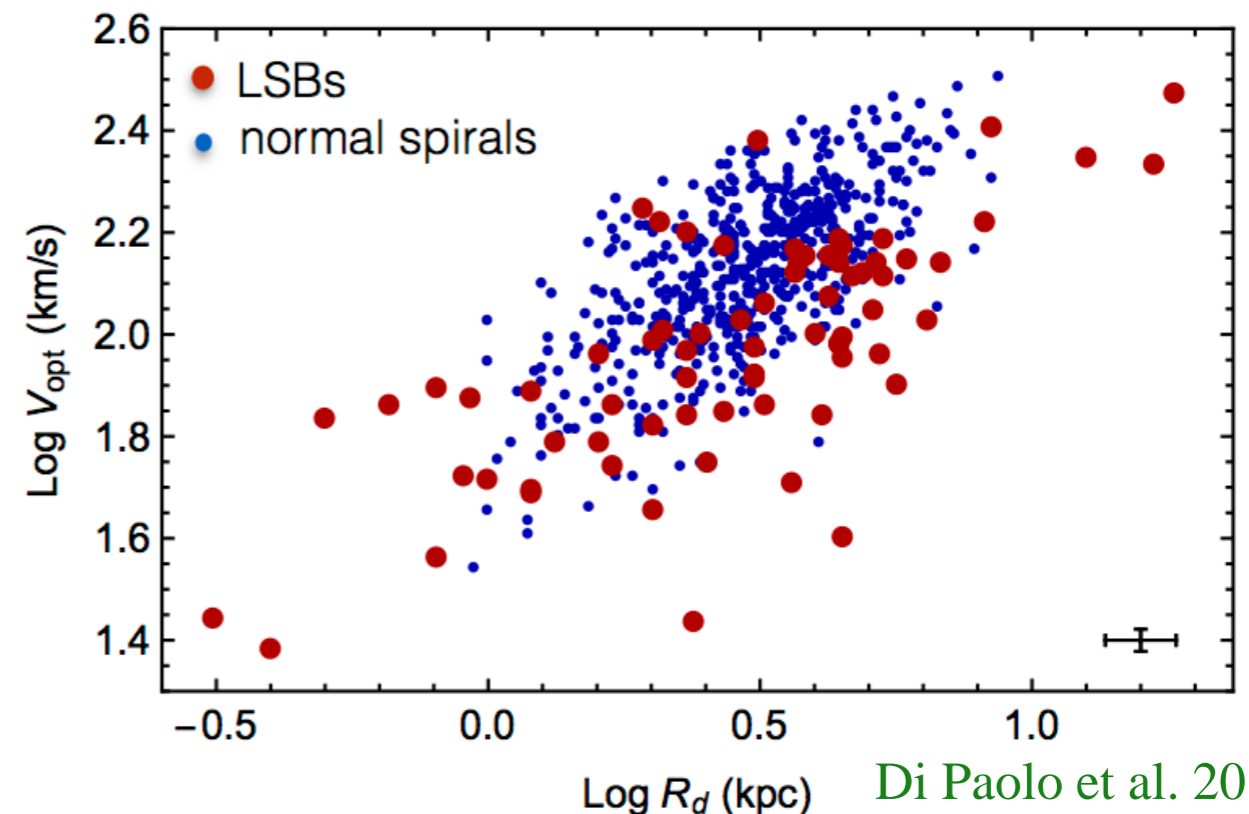
R_d = disk scale length
exponential stellar disk

$R_{opt} = 3.2R_d$ → 83 % total
luminosity

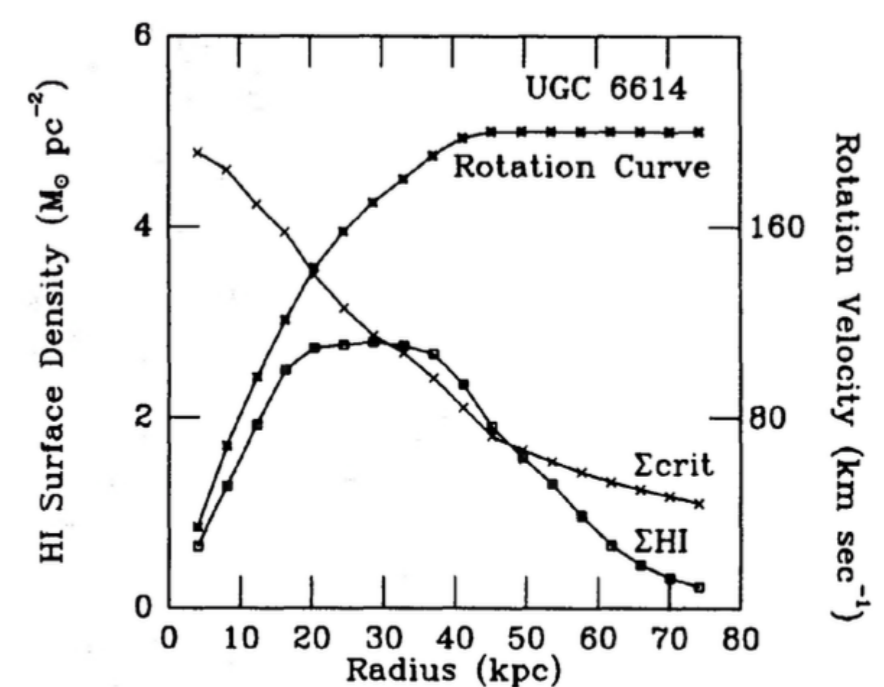
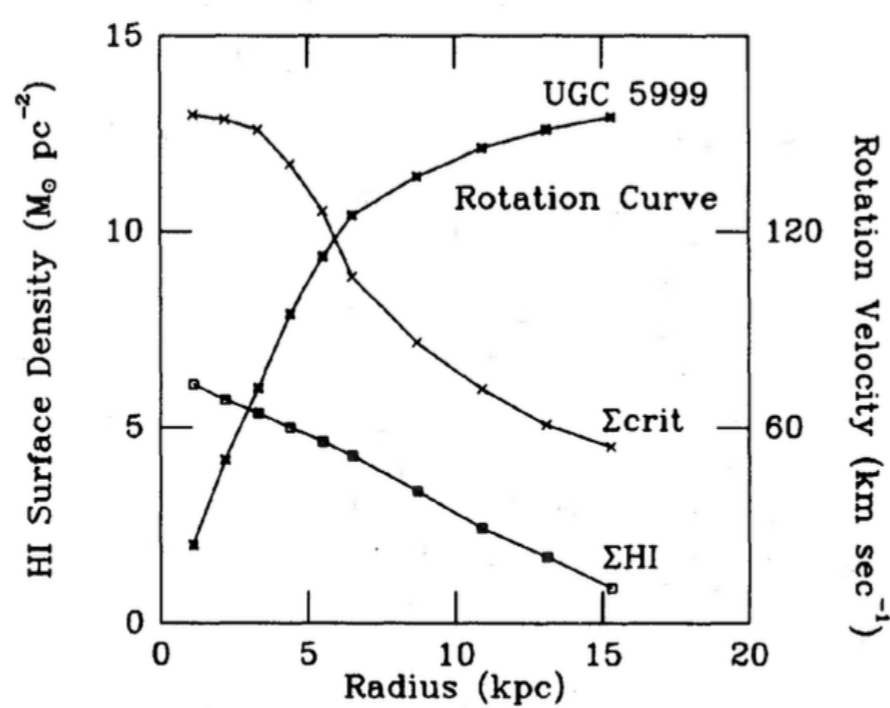
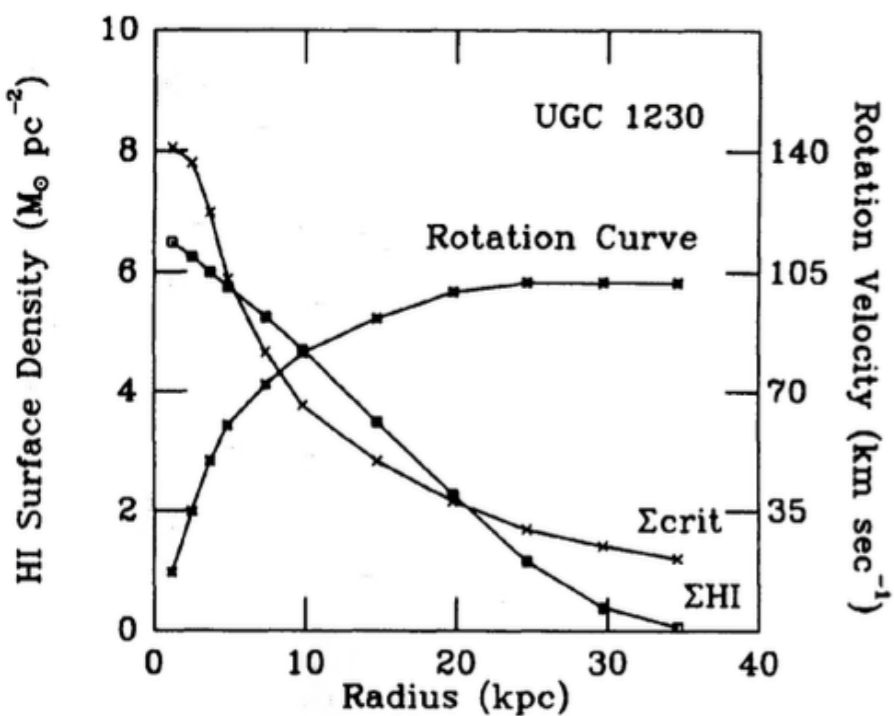
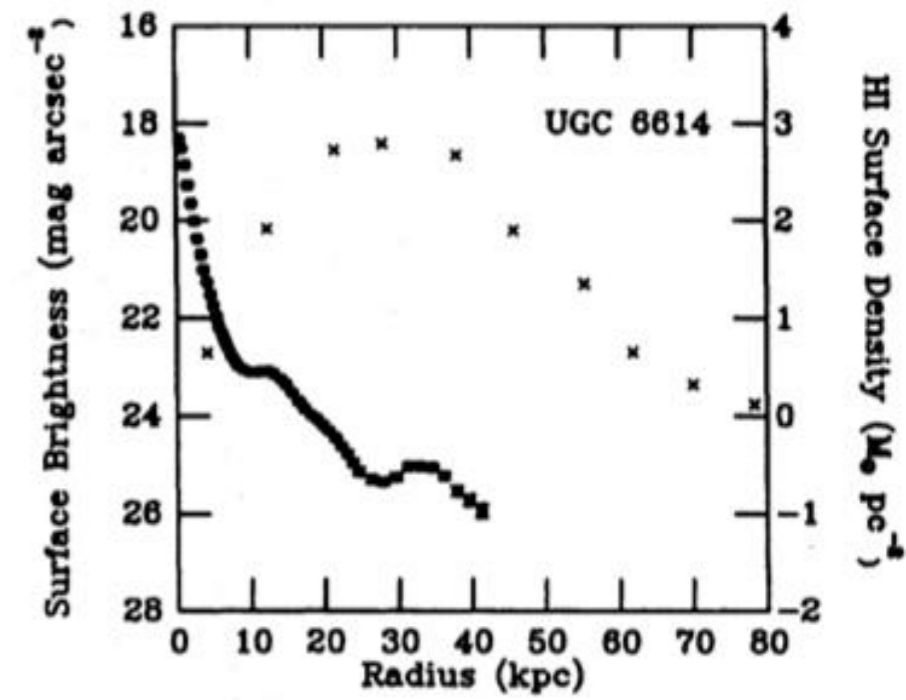
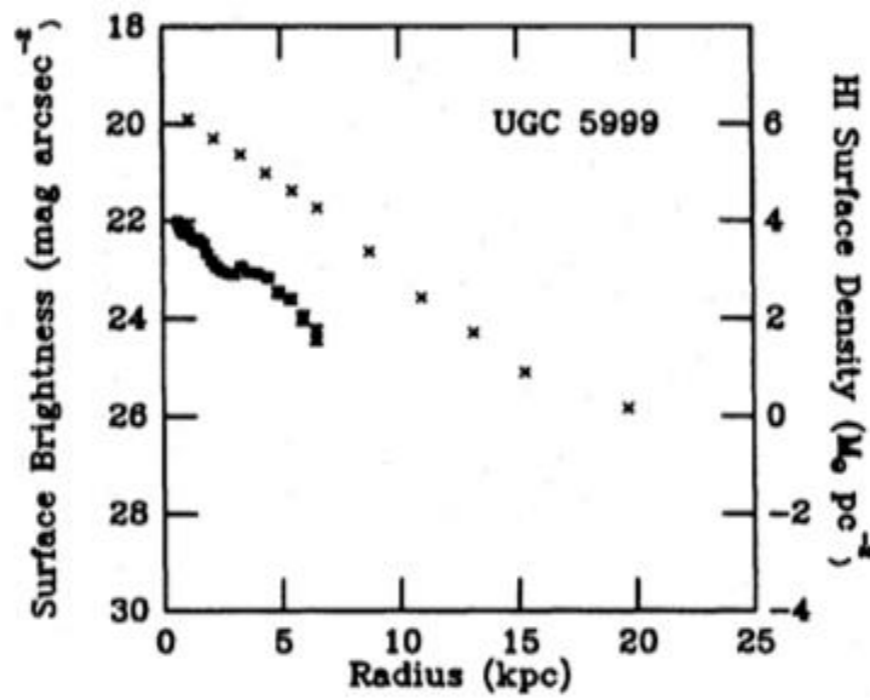
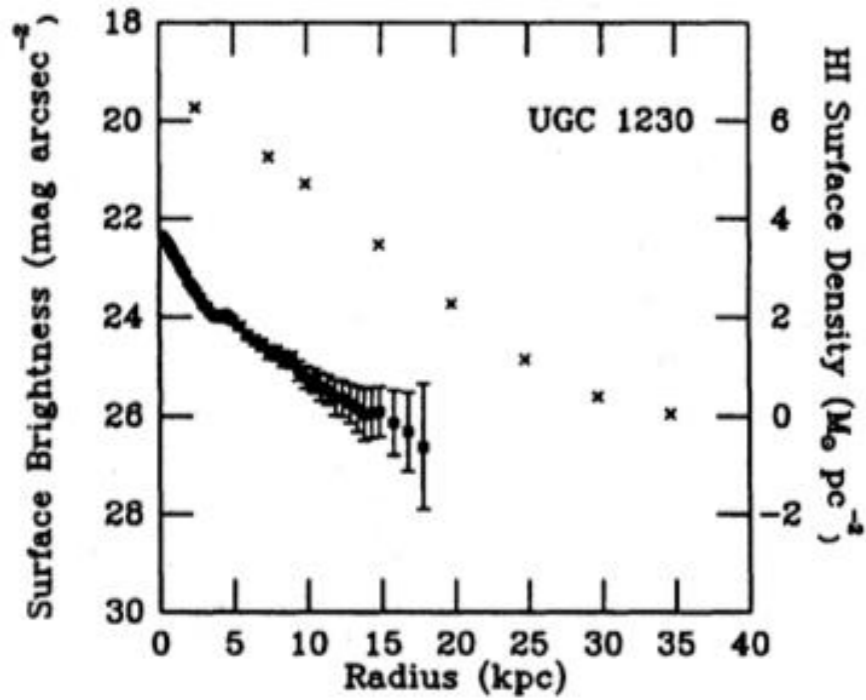
$V_{opt} = V(R_{opt})$

optical
radius

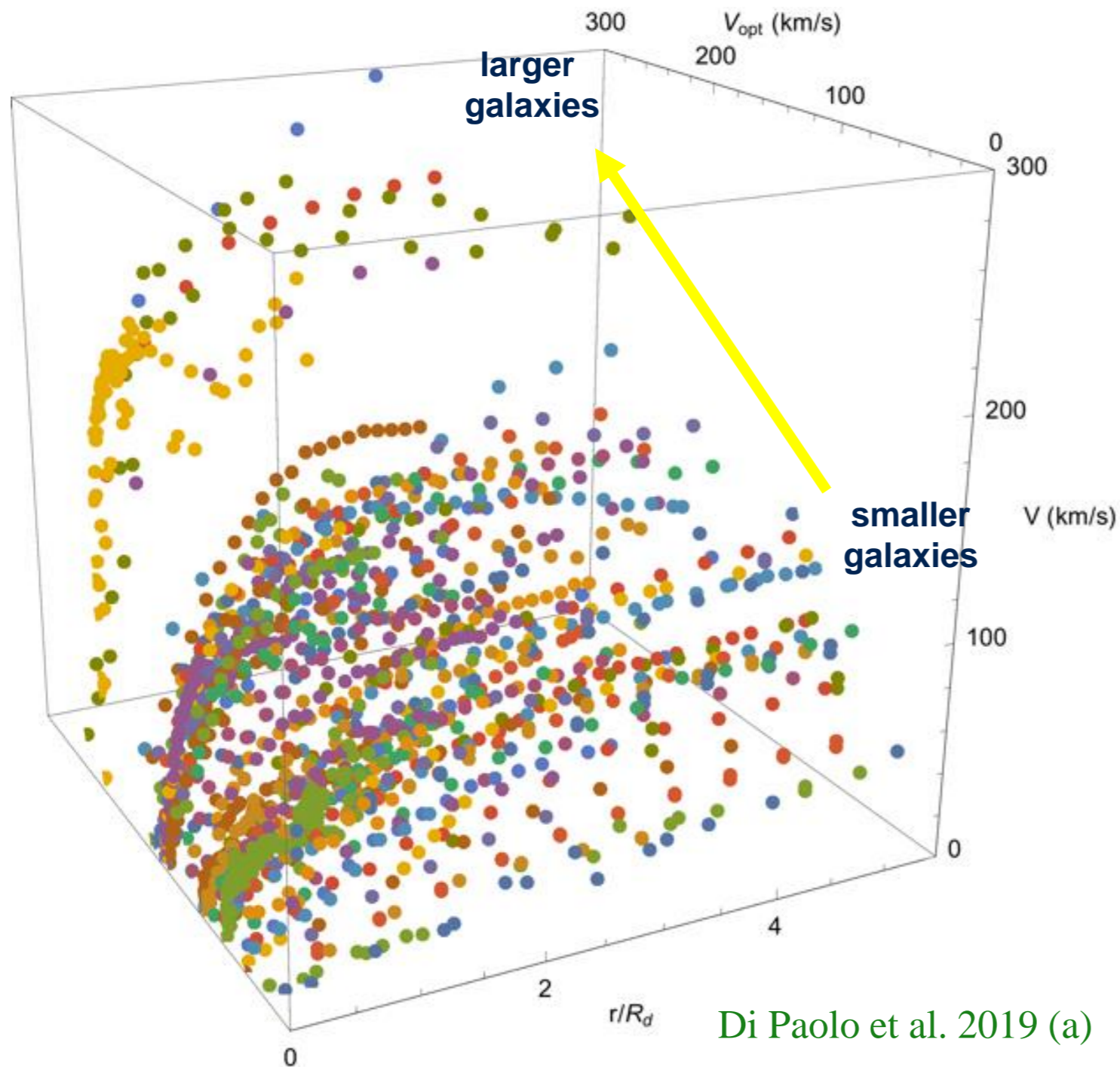
optical
velocity



Low Surface Brightness galaxies (LSBs)

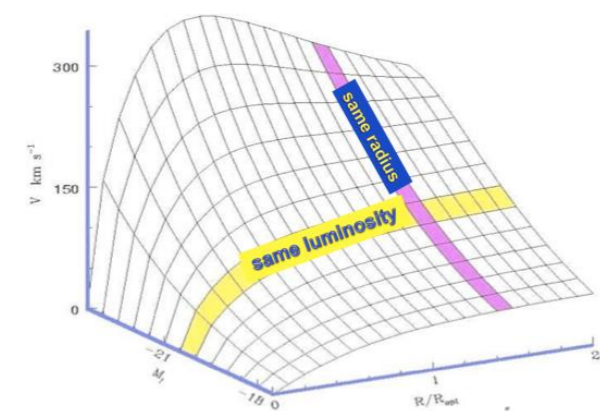


URC method applied to the LSB galaxies



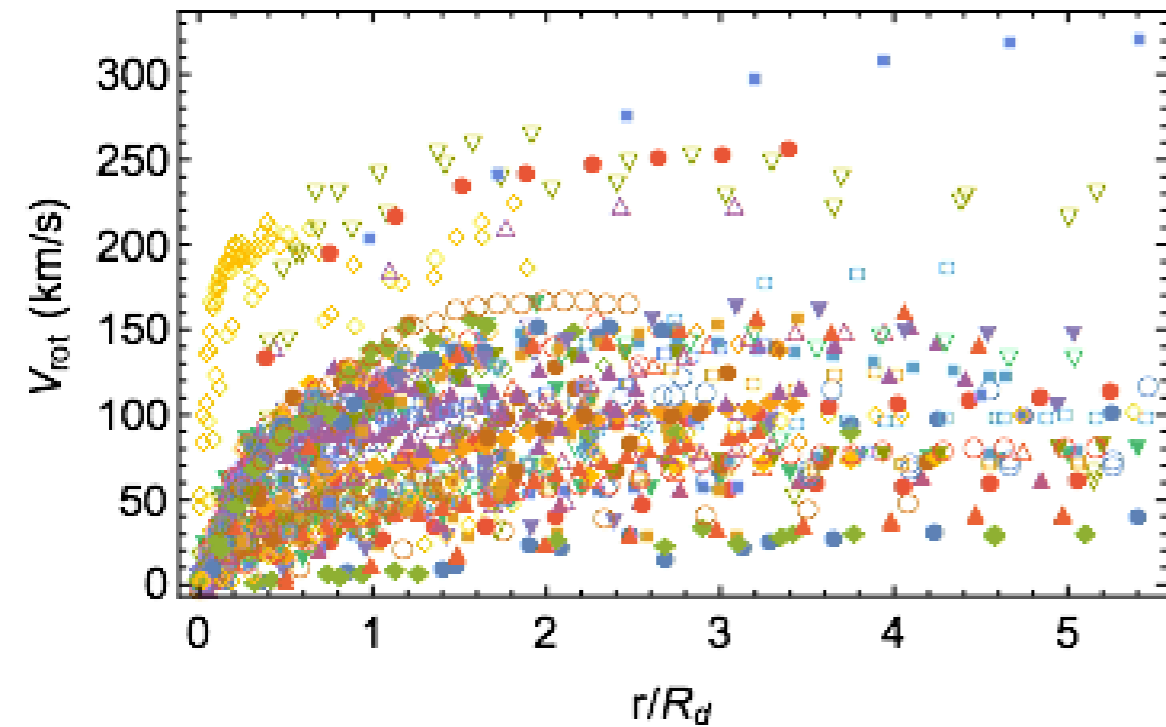
Di Paolo et al. 2019 (a)

LSBs rotation curves
Show a
universal trend



Persic, Salucci, Stel, 2007

Rotation curves BINNING

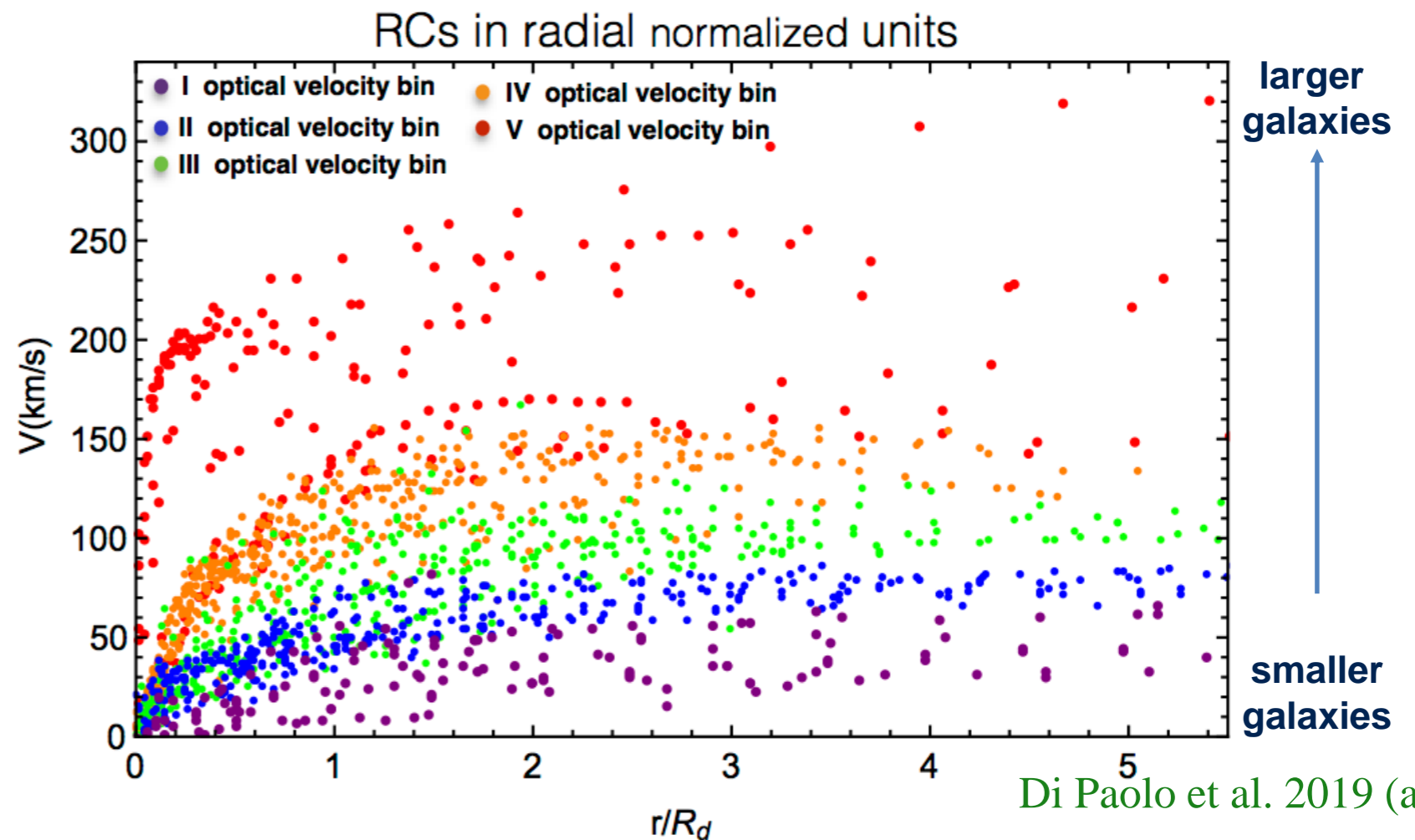


Vel.Bin	Vel.Range km/s	N.galaxies	$\langle V_{opt} \rangle$ km/s	$\langle R_D \rangle$ kpc	N.data
(1)	(2)	(3)	(4)	(5)	(6)
1	24-60	13	43.5	1.7	151
2	60-85	17	73.3	2.2	393
3	85-120	17	100.6	3.7	419
4	120-154	15	140.6	4.5	441
5	154-300	10	205.6	7.9	210

Division in
5 velocity bins

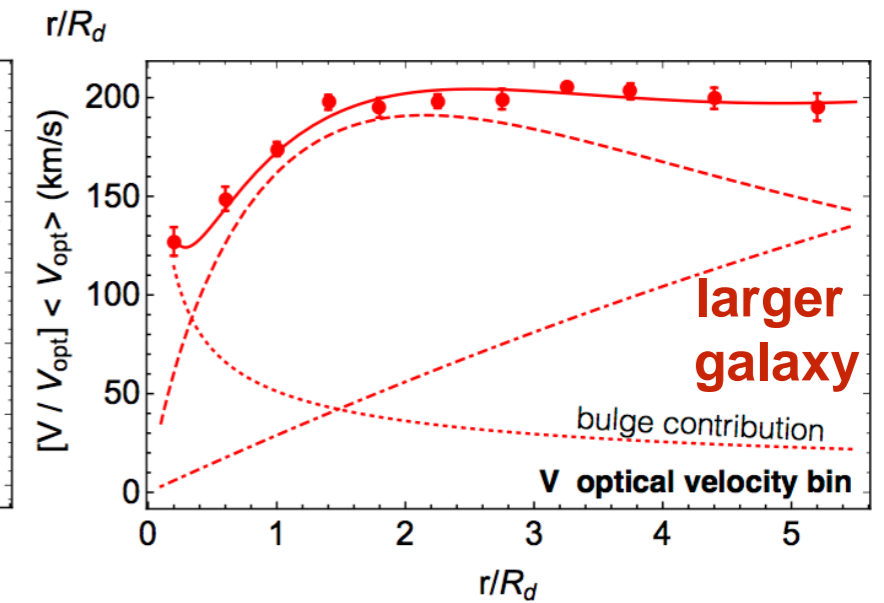
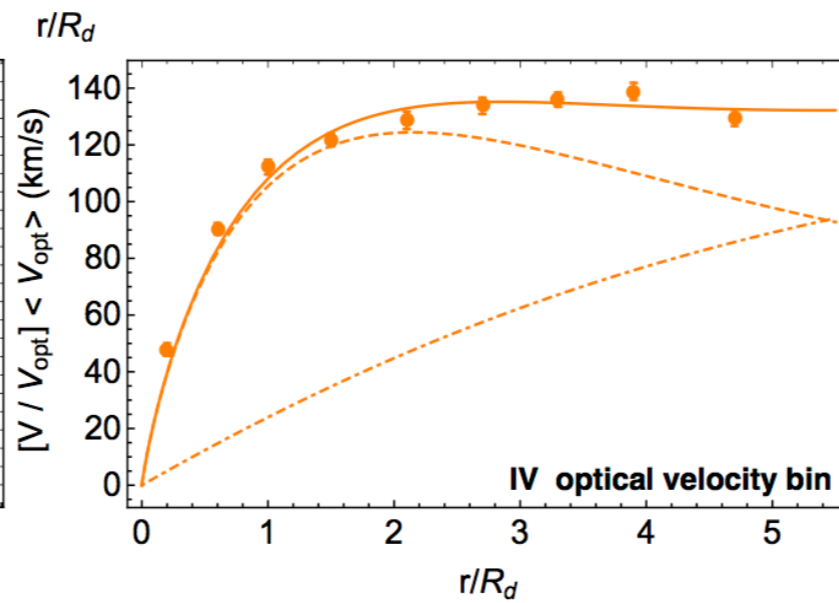
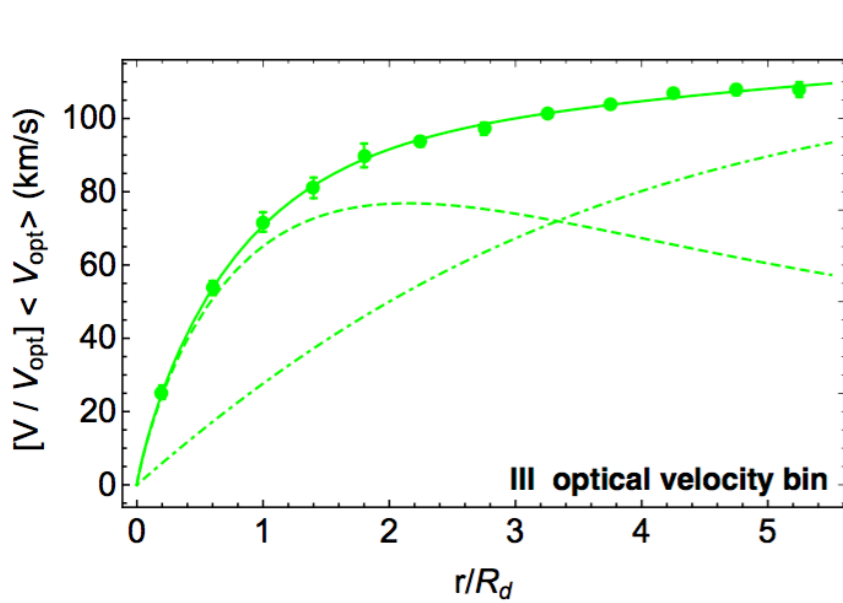
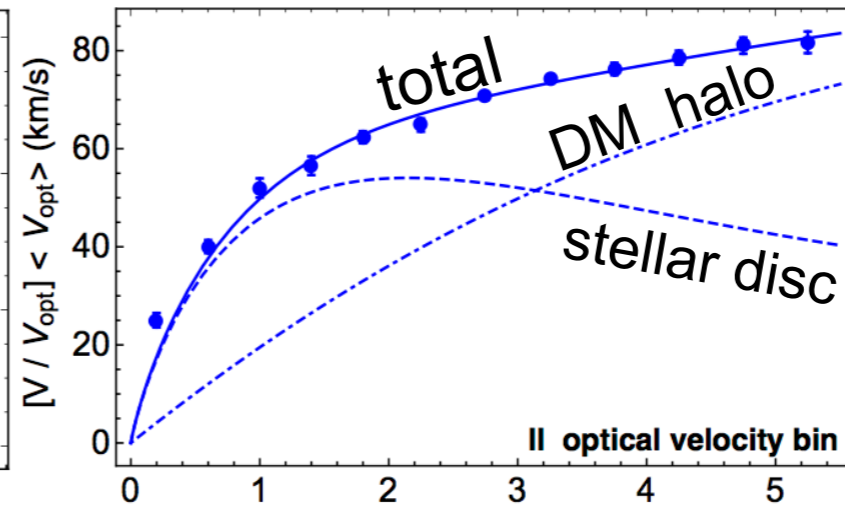
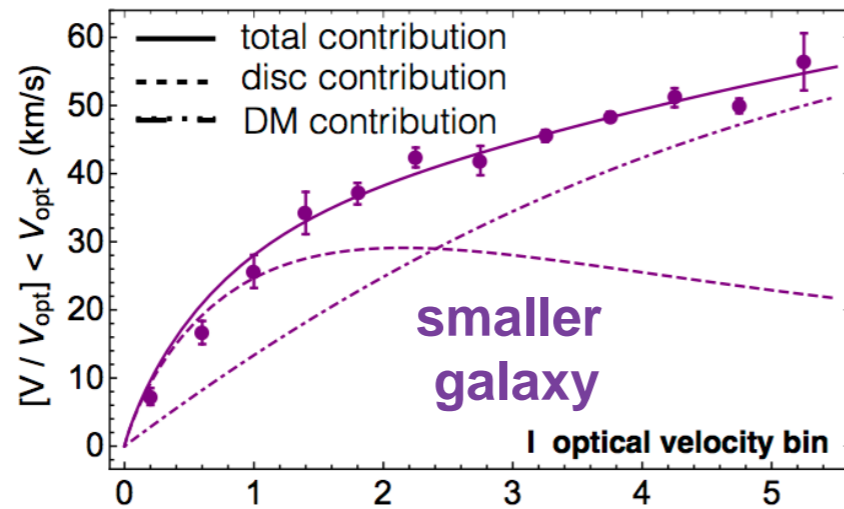


inner curvature
more similar



Mass Modeling 5 coadded RCs

coadded RCs in radial normalized units



$$V^2(r) = V_d^2(r) + V_{DM}^2(r)$$

exponential stellar disc
DM spherical cored halo (Burkert)

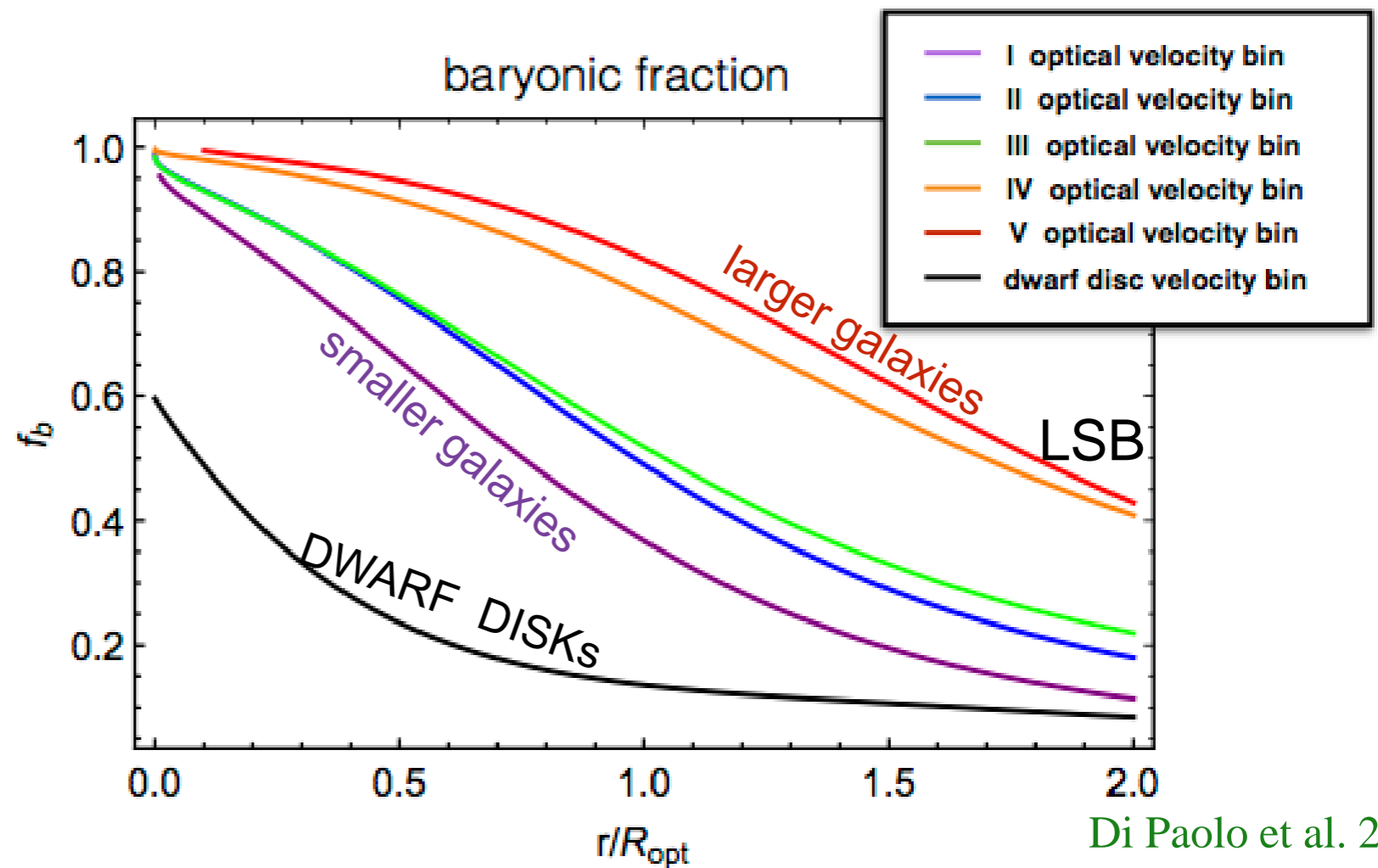
3 free parameters :

- M_d = stellar disc mass
- ρ_0 = DM halo's central mass density
- r_0 = DM halo's core radius

Mass Modeling 5 coadded RCs

Baryonic fraction :

$$f_b(r) = \frac{V_b^2(r)}{V^2(r)}$$



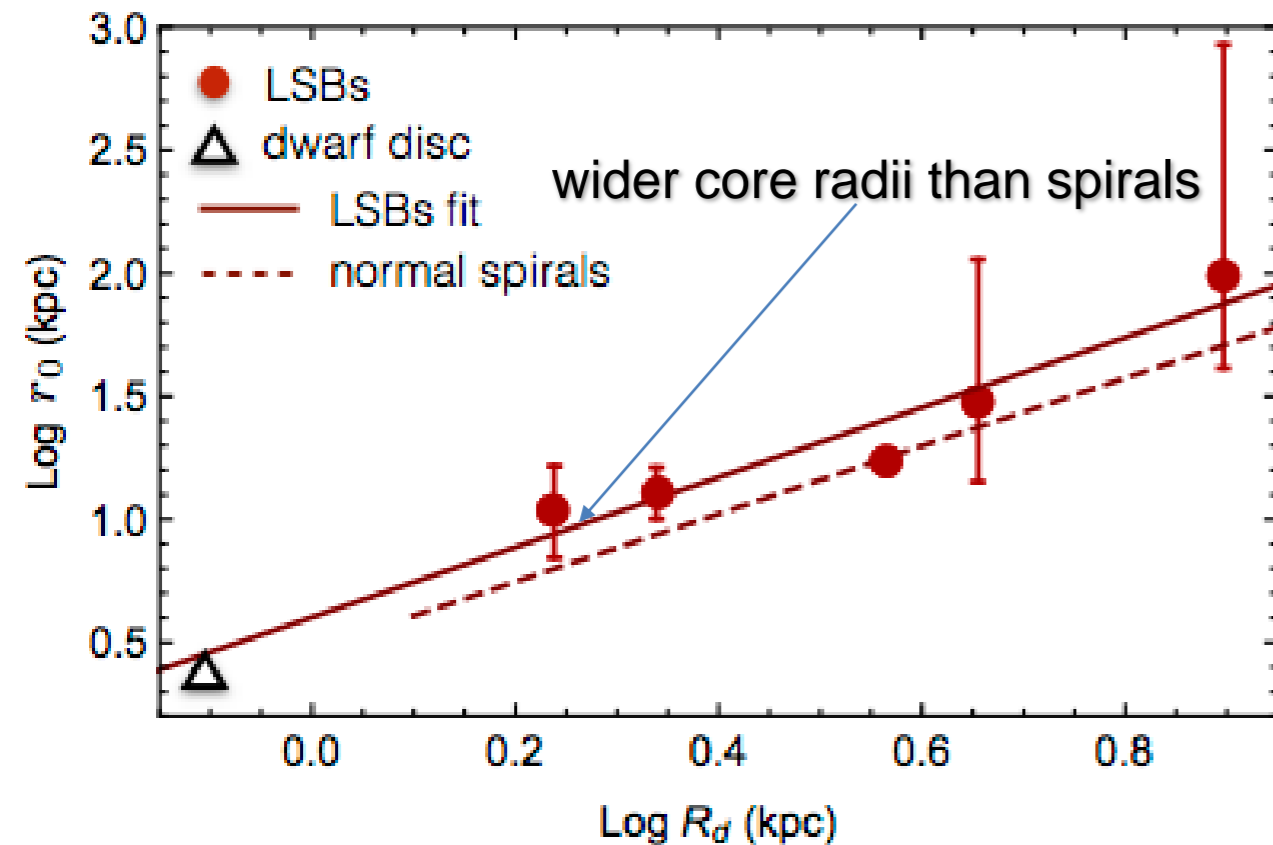
NOTE:

- radial dependence of f_b
- different $f_b(r)$ in galaxies of different size
- different $f_b(r)$ in galaxies of different morphology

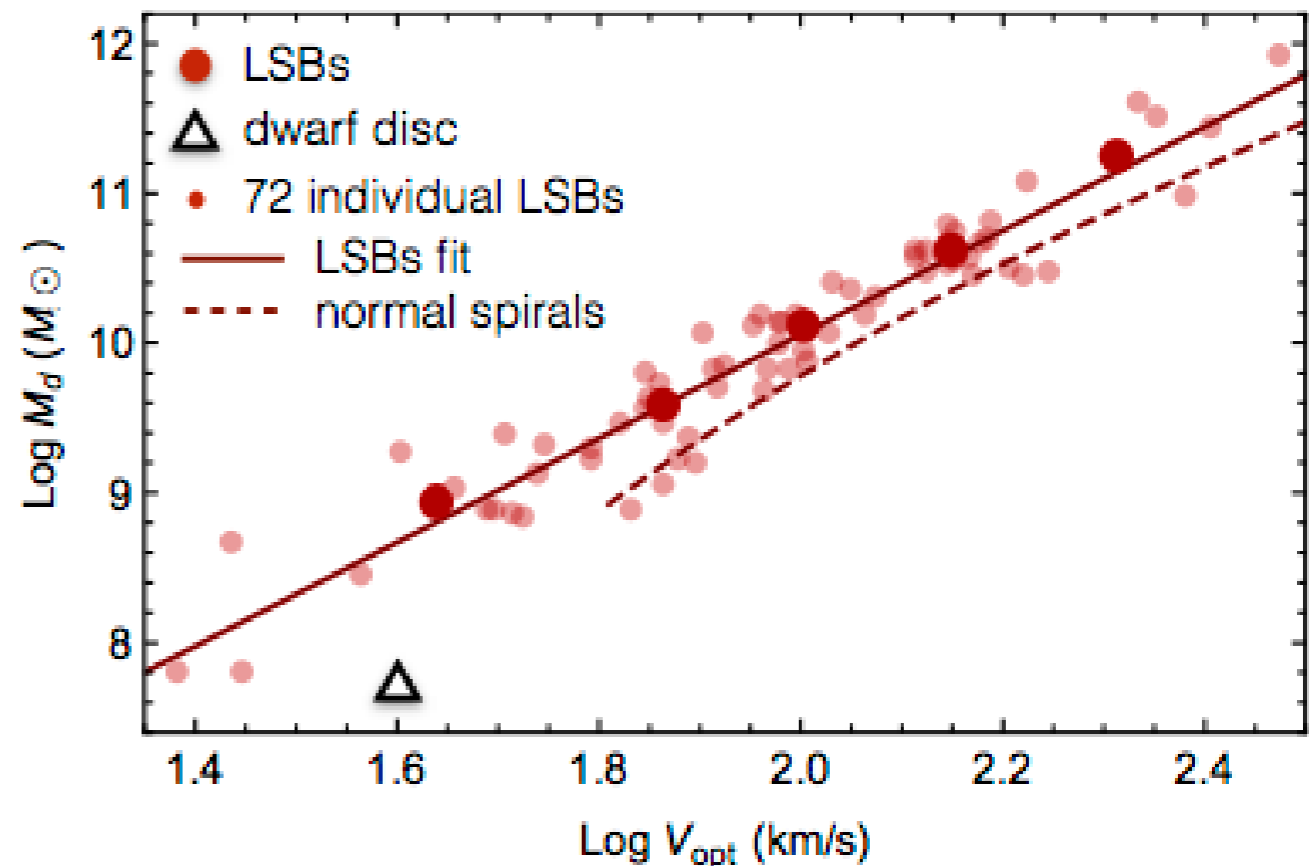
Scaling laws for LSBs

Similar to the HSB spirals scaling laws

$$\text{Log } r_0 = 0.60 + 1.42 \text{ Log } R_d$$

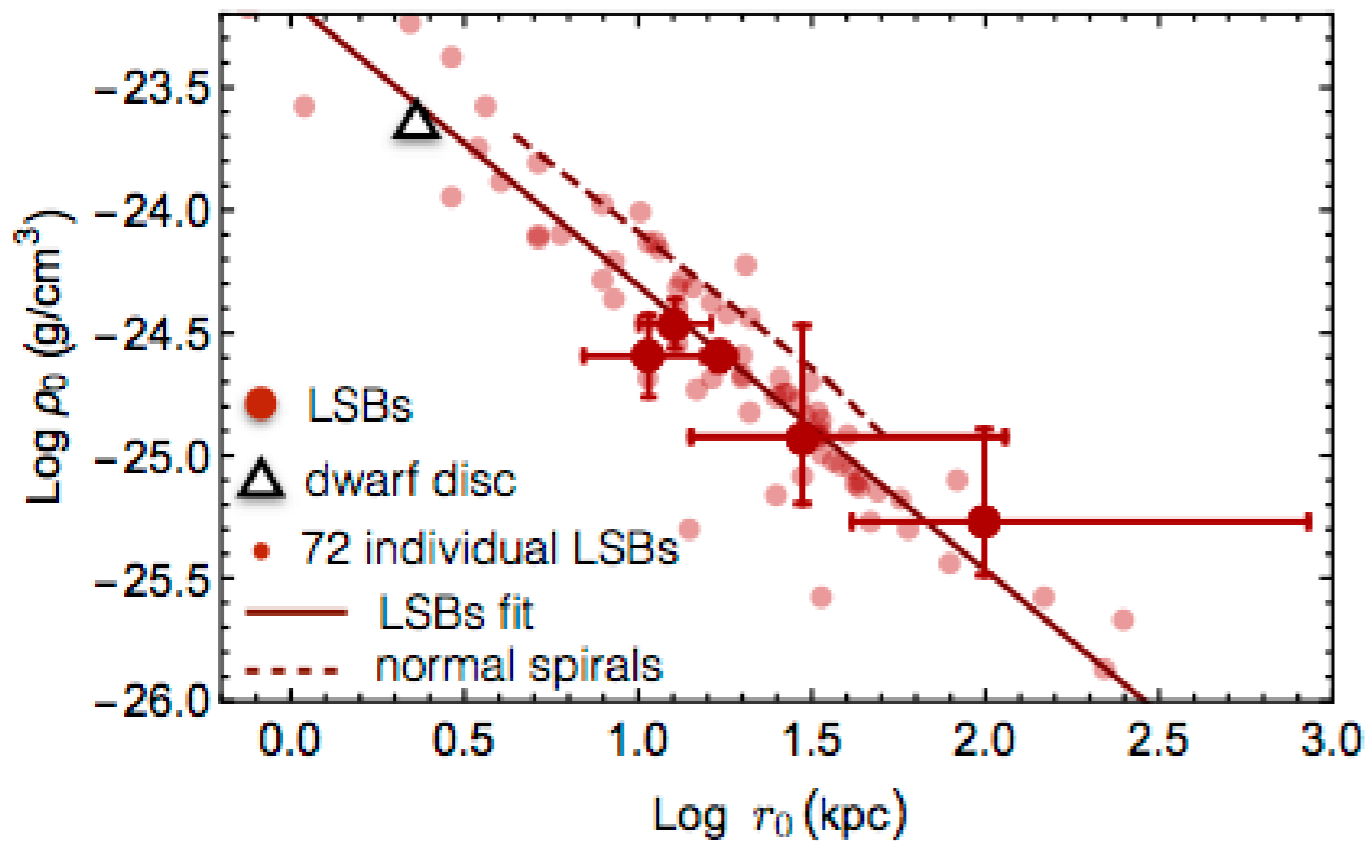


$$\text{Log } M_d = 3.12 + 3.47 \text{ Log } V_{opt}$$

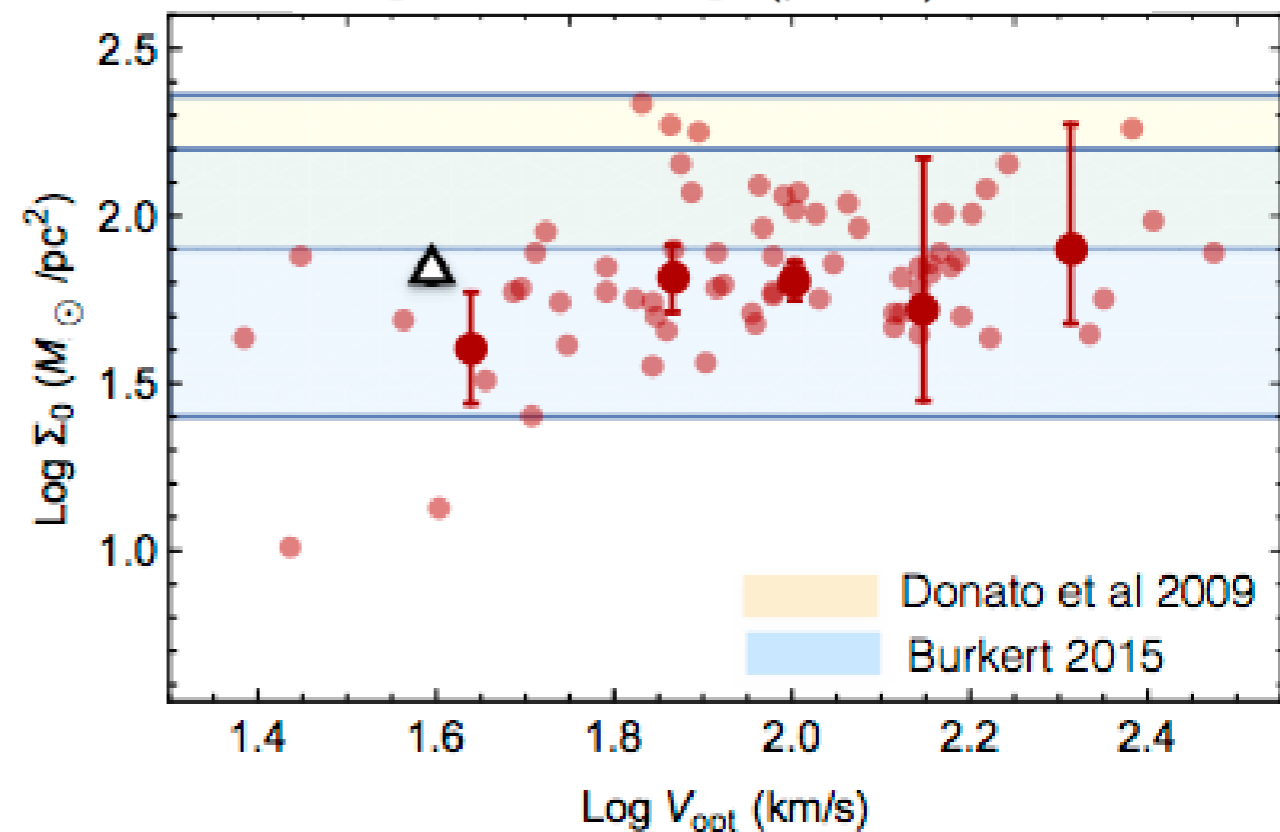


Scaling laws for LSBs

$$\text{Log } \rho_0 = -23.15 - 1.16 \text{Log } r_0$$

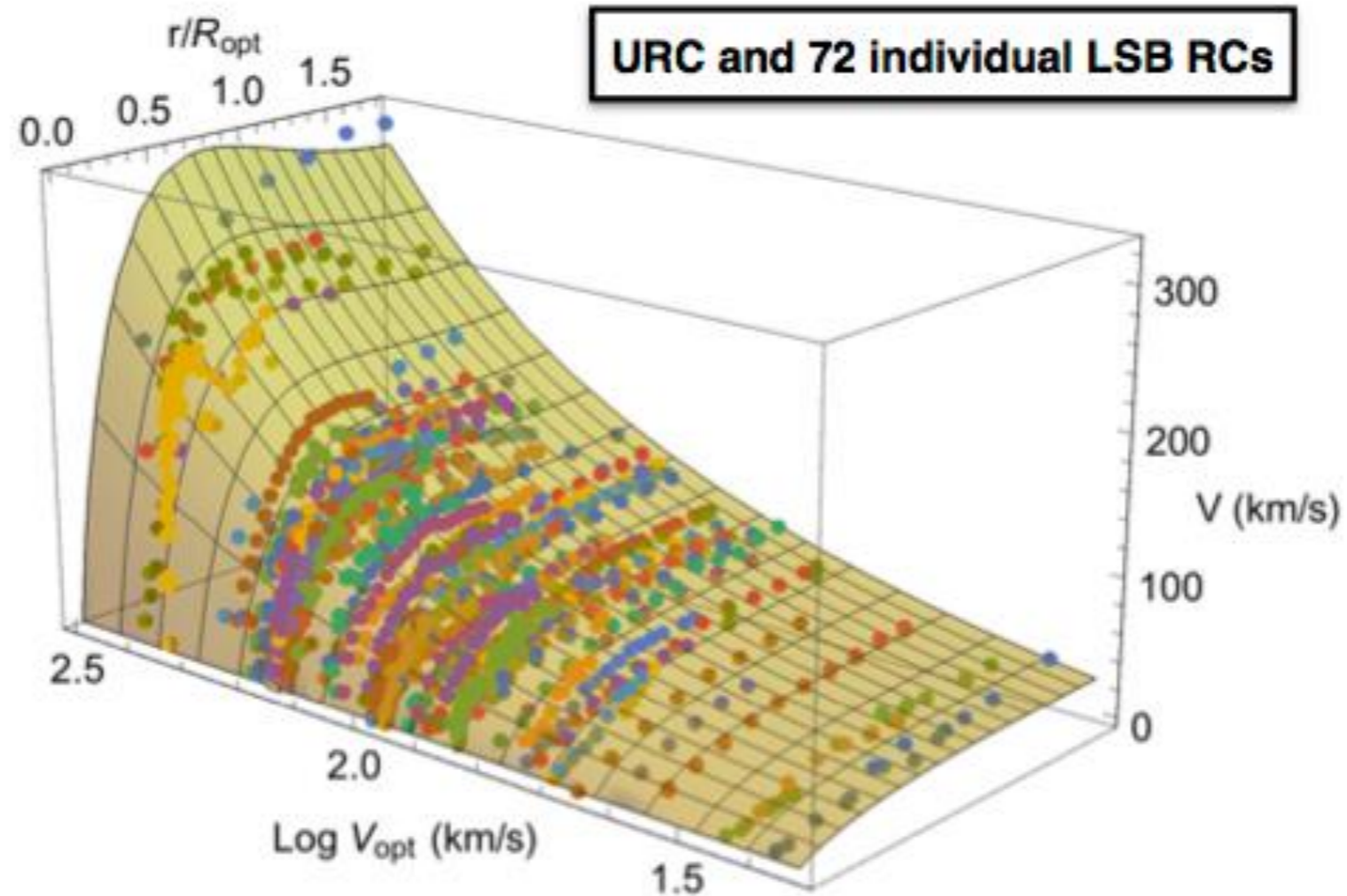


$$\text{Log } \Sigma_0 = \text{Log } (\rho_0 r_0) \simeq 1.9$$



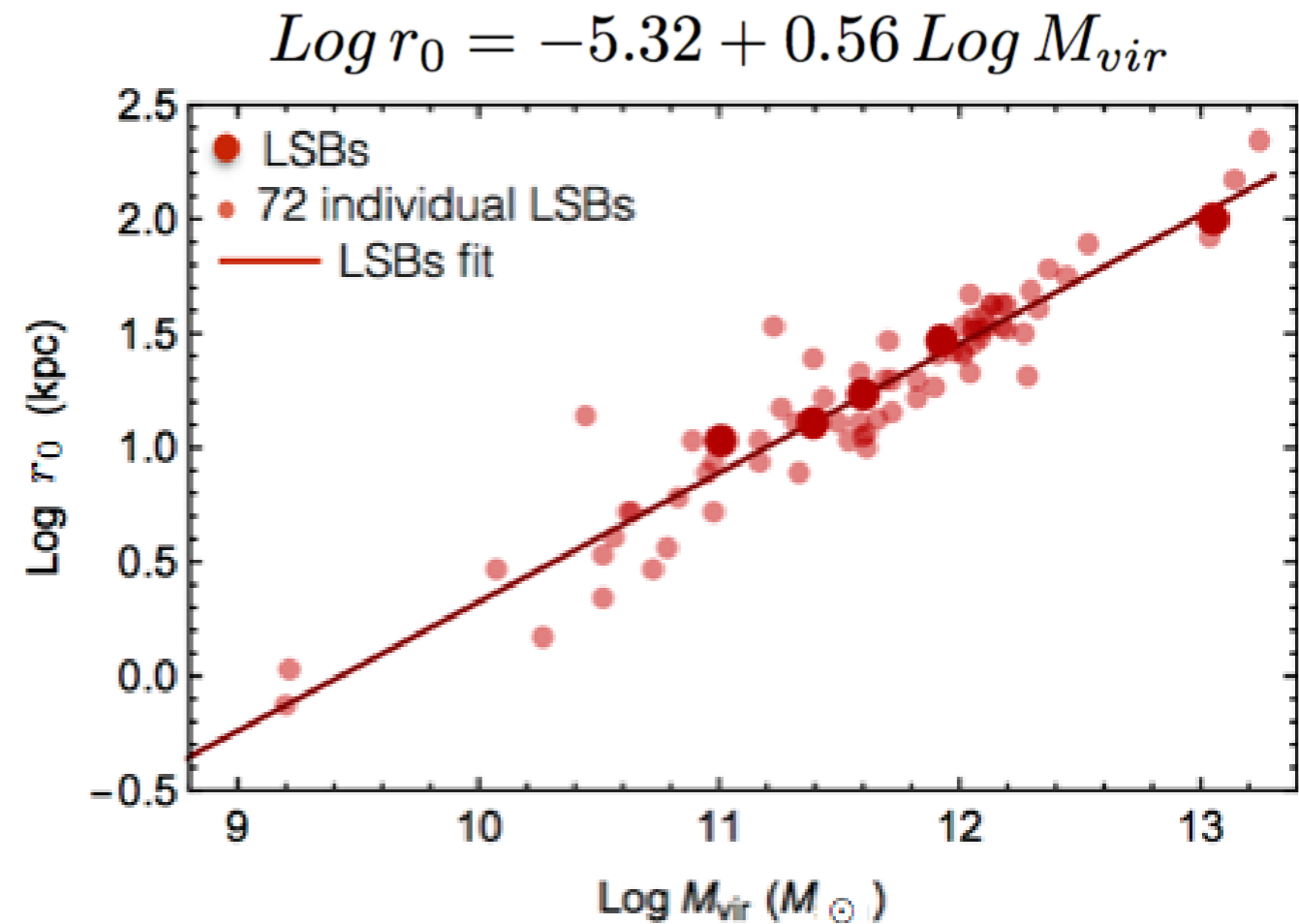
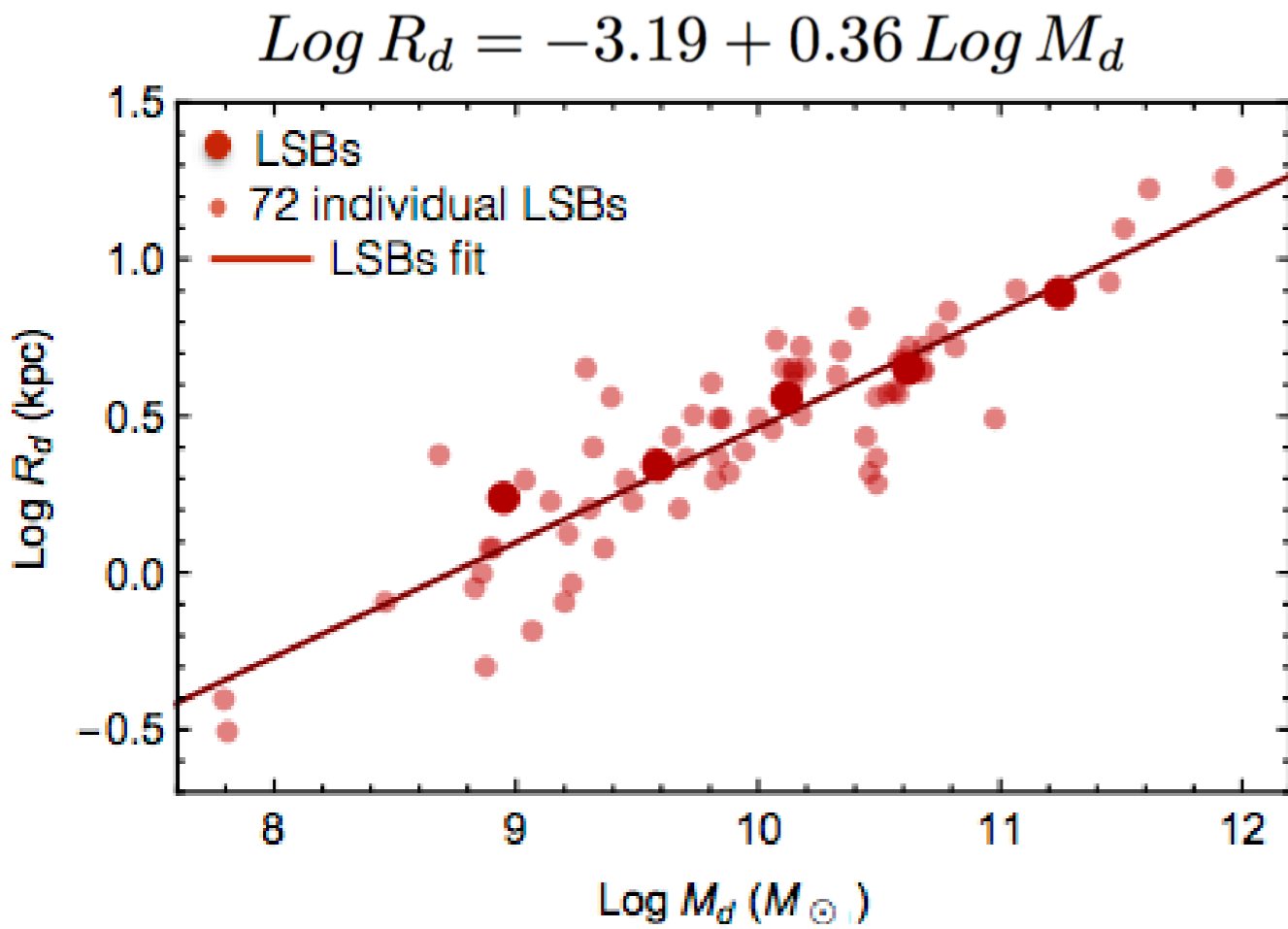
Universal Rotation Curve

$$\Delta V / V \simeq 8\%$$



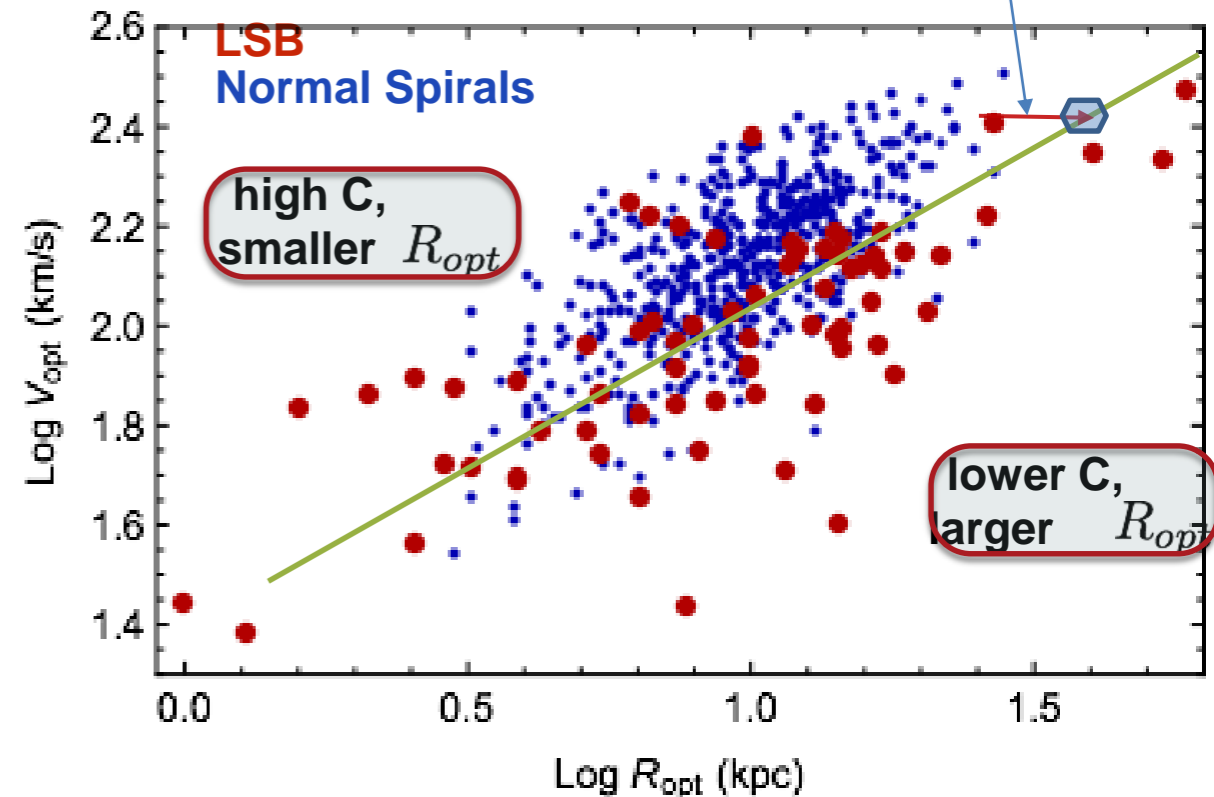
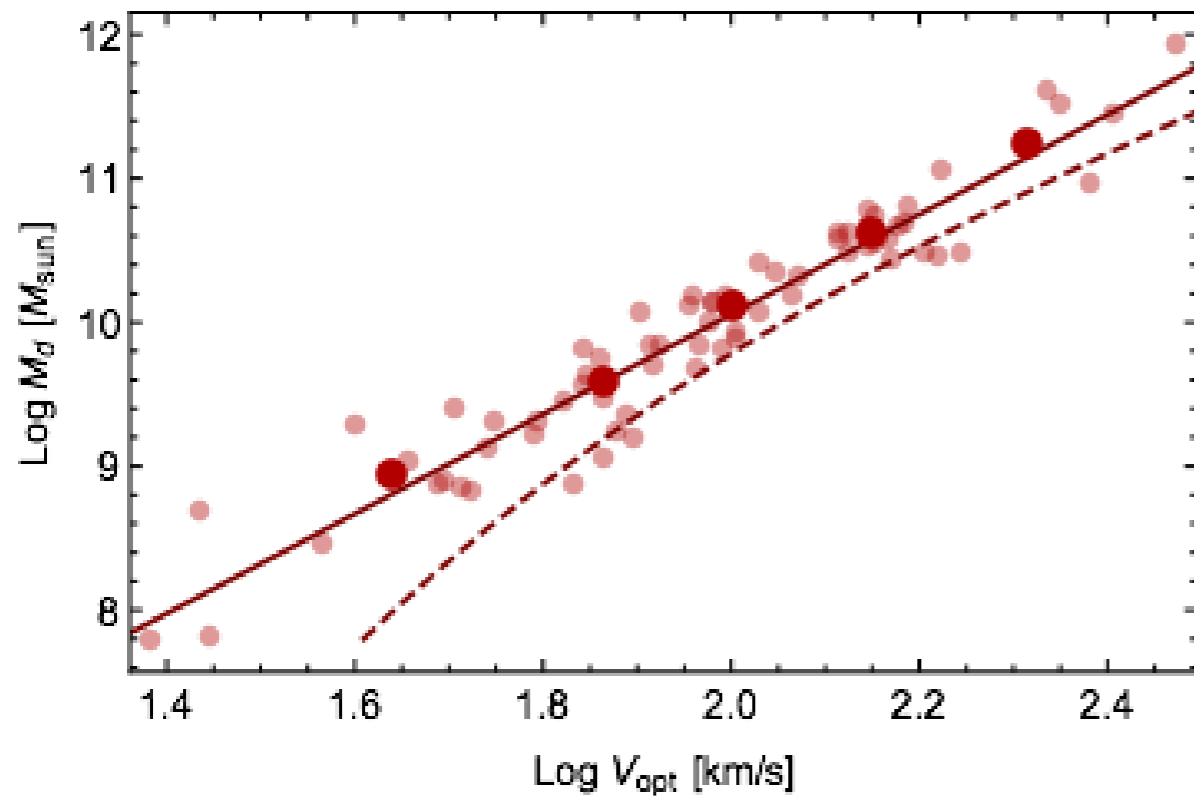
$$V^2(r/R_{\text{opt}}) = V_d^2(r/R_{\text{opt}}) + V_{DM}^2(r/R_{\text{opt}}) \longrightarrow \text{function of } V_{\text{opt}}$$

Scaling laws for LSBs



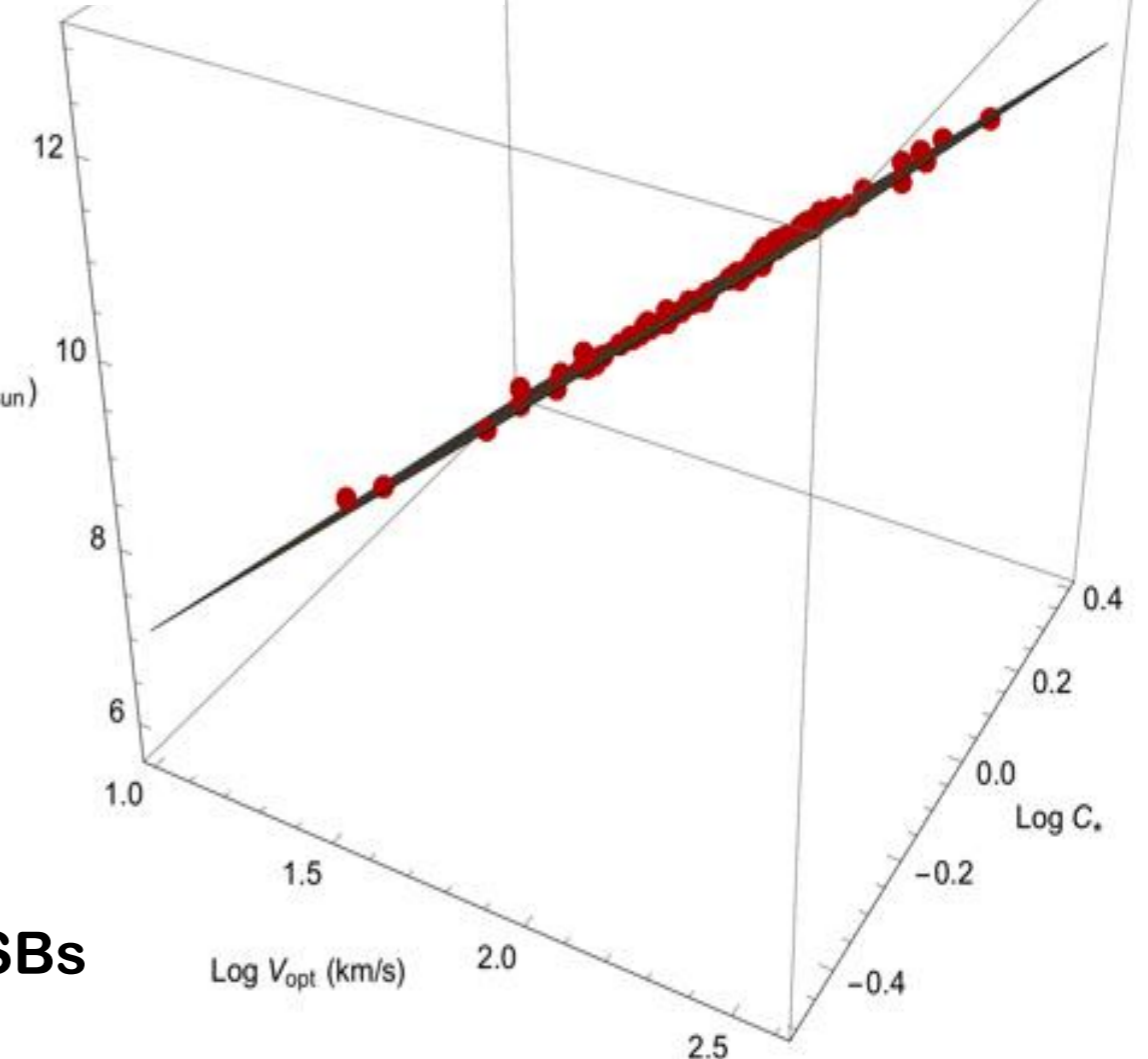
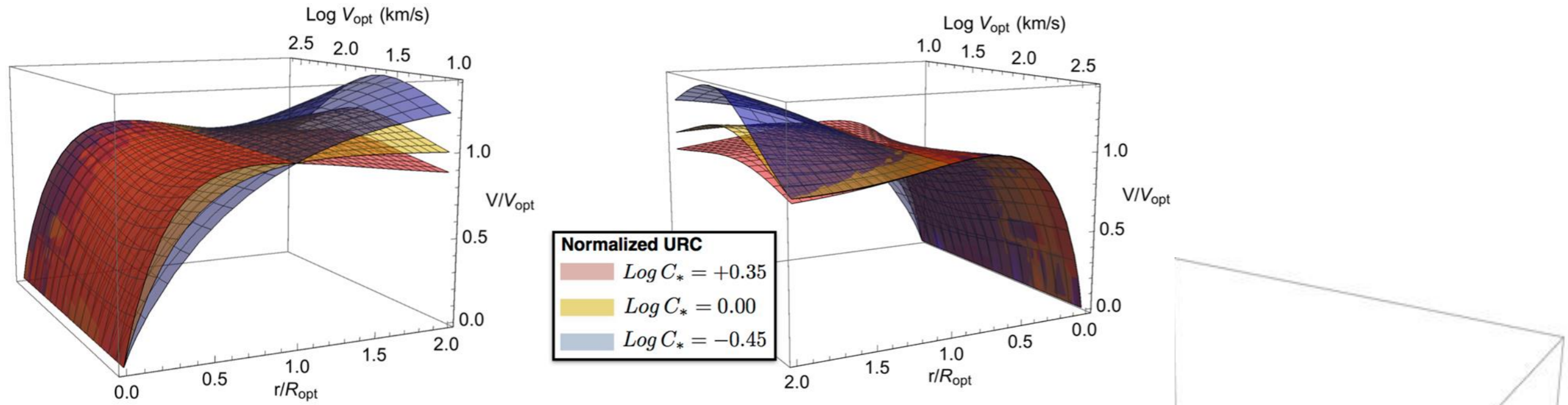
Compactness

Log C



COMPACTNESS (C):
discrepancy between the measured R_{opt}
and the expected value \bar{R}_{opt}

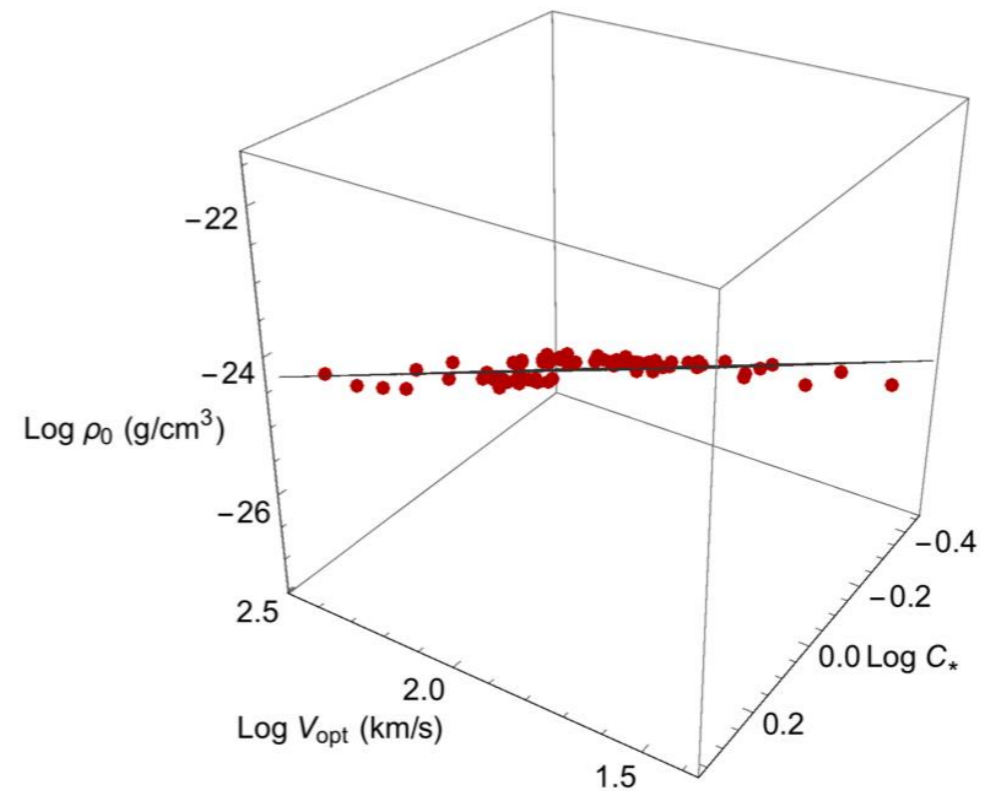
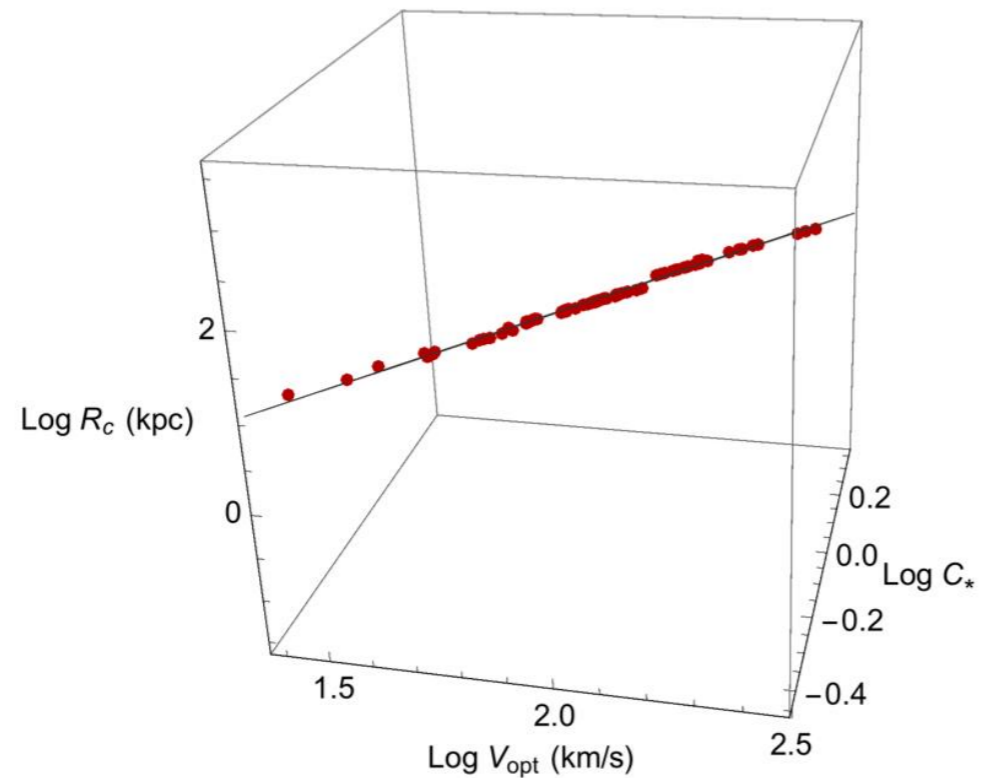
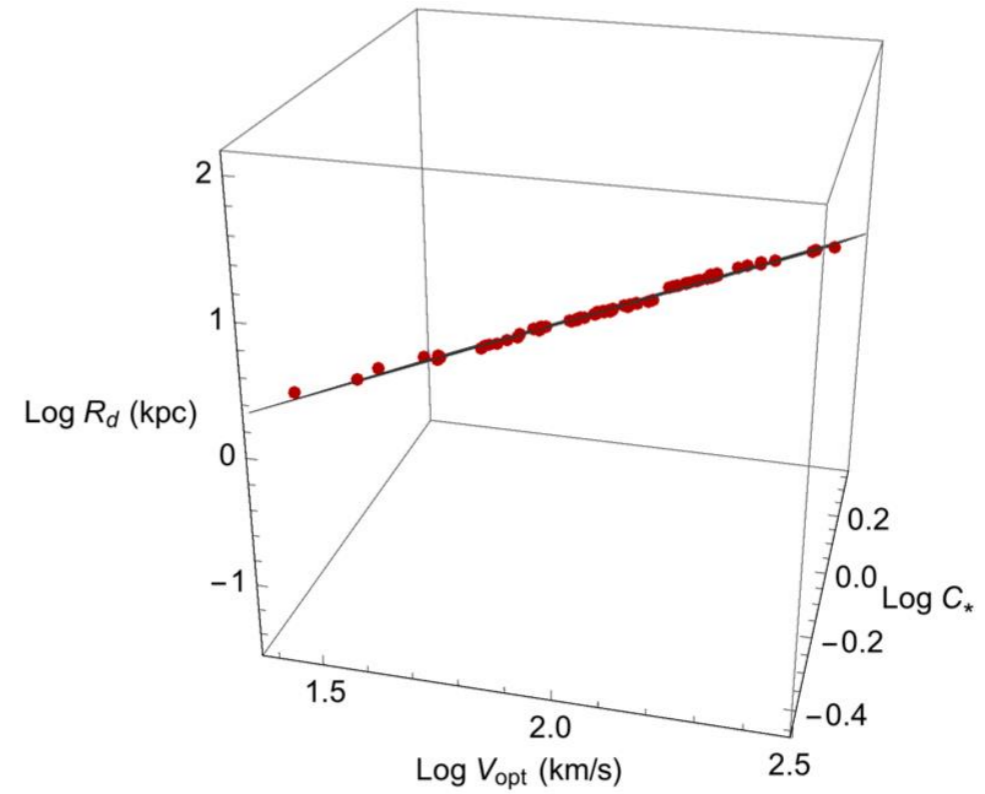
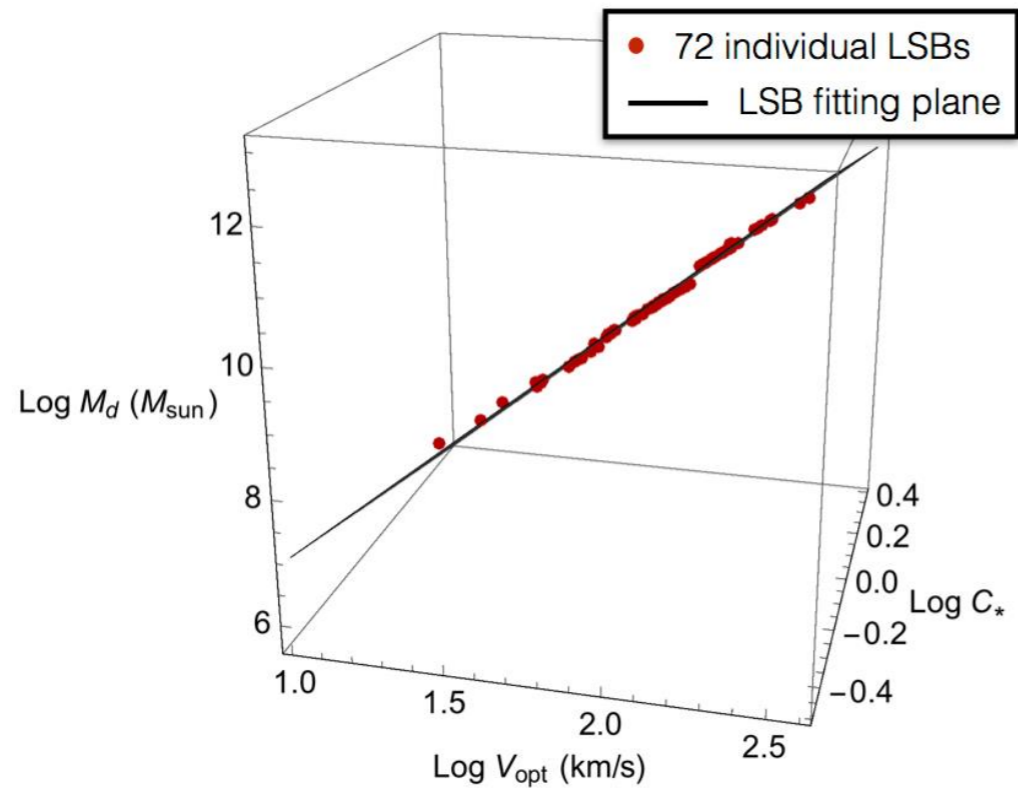
$$V_{\text{urc}}(r, R_D, V_{\text{opt}}, C)^2 = V_d(r, R_D(V_{\text{opt}}, C), M_D(V_{\text{opt}}, C))^2 + V_h(r, r_0(V_{\text{opt}}, C), \rho_0(V_{\text{opt}}, C))^2$$



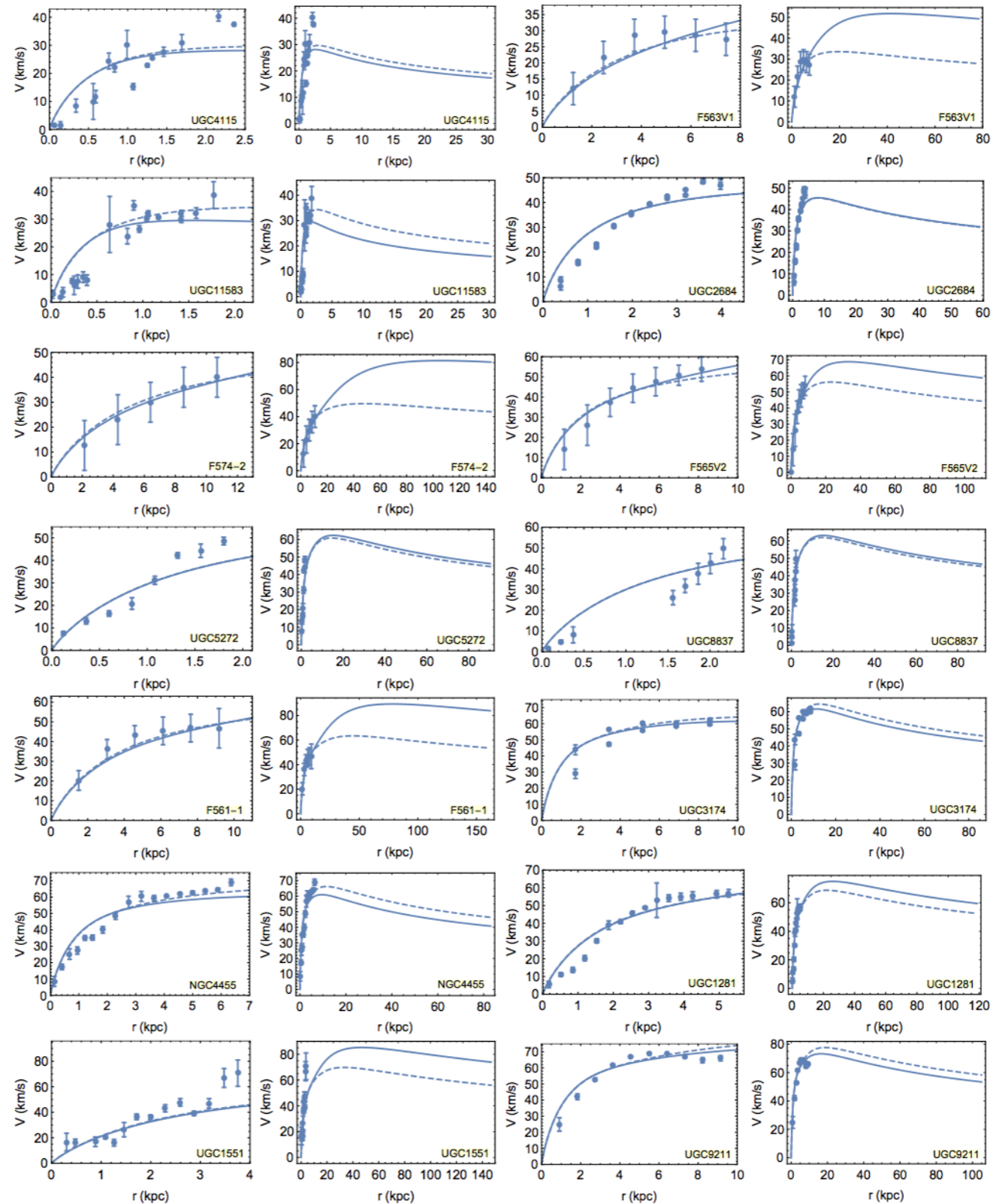
Reduced scatter after the **COMPACTNESS** introduction

important parameter with V_{opt} R_{opt}
 In order to have a “better universality” in LSBs

Low Surface Brightness galaxies (LSBs)



Universal Rotation Curve



$$\Delta V / V \simeq 8\%$$

WHAT IS DARK MATTER?

Constant density central regions in all LSBs

- R_d vs r_0

Relation & Co.

Portal btween Dark
and Luminous matter

in the Λ CDM scenario + baryonic feedback :

- challenge in dwarf and large spirals
- challenge in fine-tuning parameters
- Challenge at higher levels in LSB galaxies




Specific power spectrum of dark matter AND DM pressure
plus some direct interaction between dark matter (DM) and luminous matter (LM).

Requirements possible for keV WDM

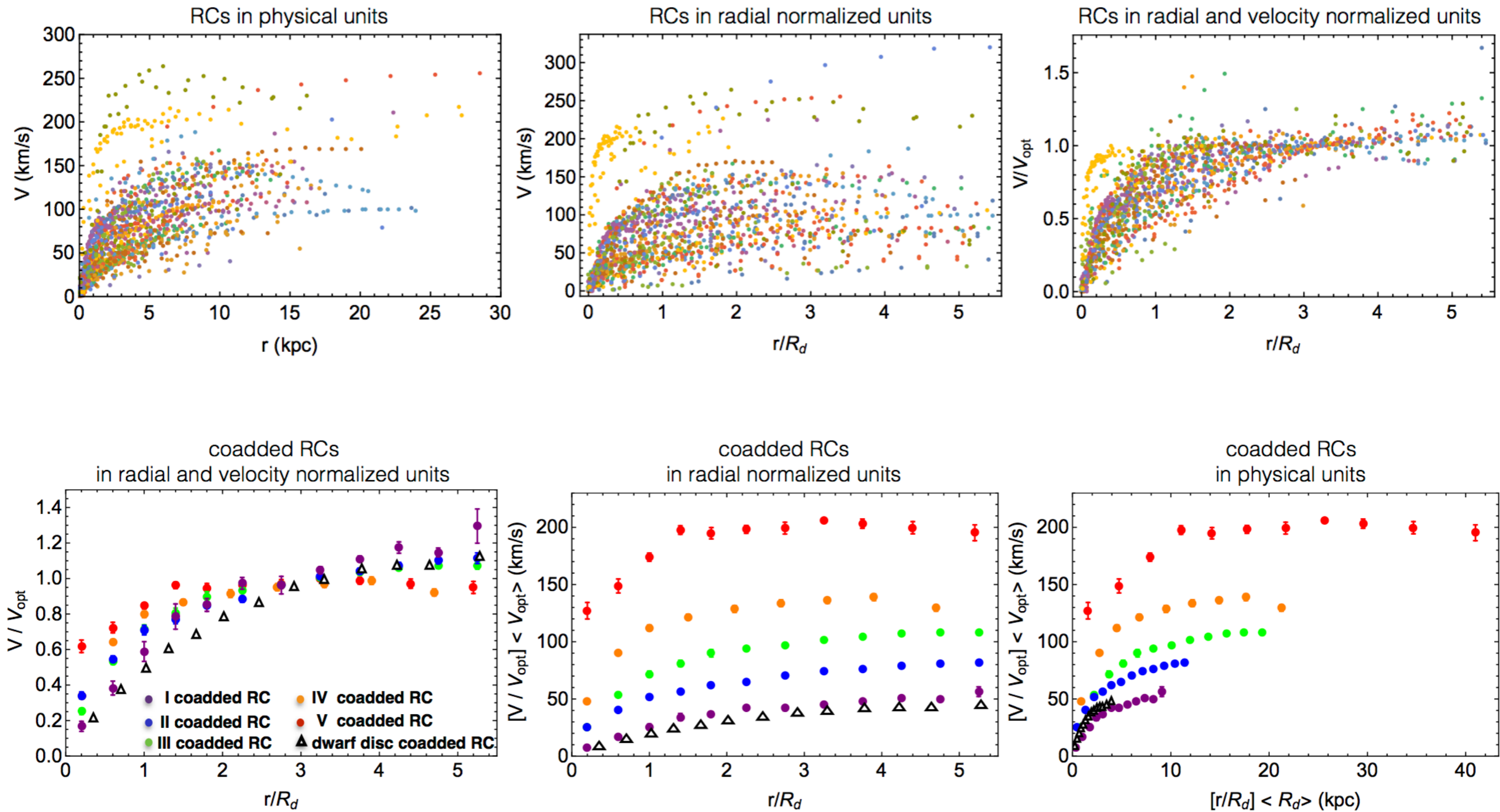
Start from halos made by collisionless particle inconsistent with present data

Conclusion

- The LM-DM interaction might be a necessary key to understand the DM phenomenon: WDM passes
The test

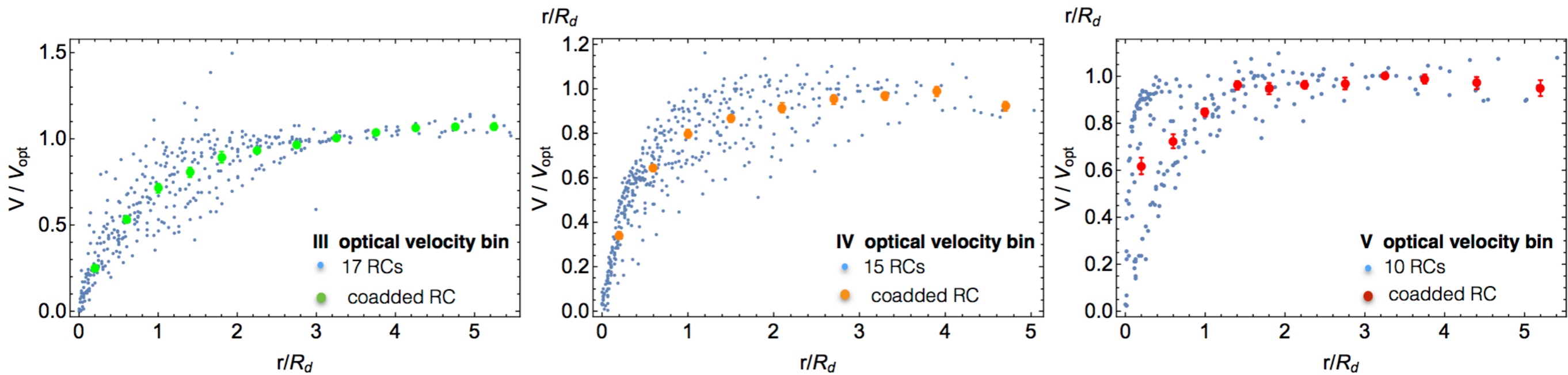
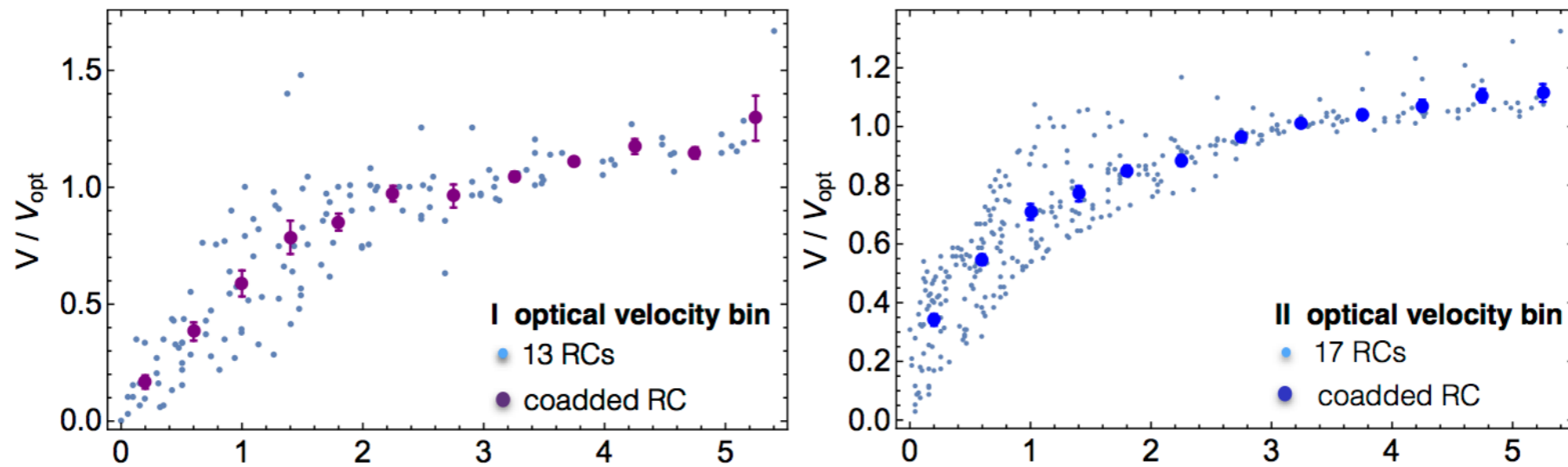
- 
- to reproduce the observed core in the galactic DM halo
 - the empirical relationships between the galactic properties

Low Surface Brightness galaxies (LSBs)



Low Surface Brightness galaxies (LSBs)

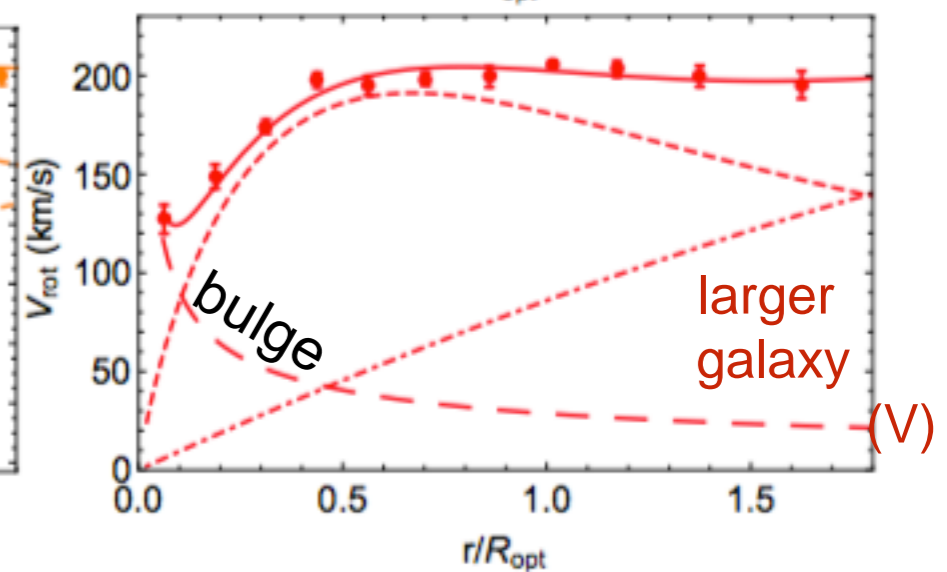
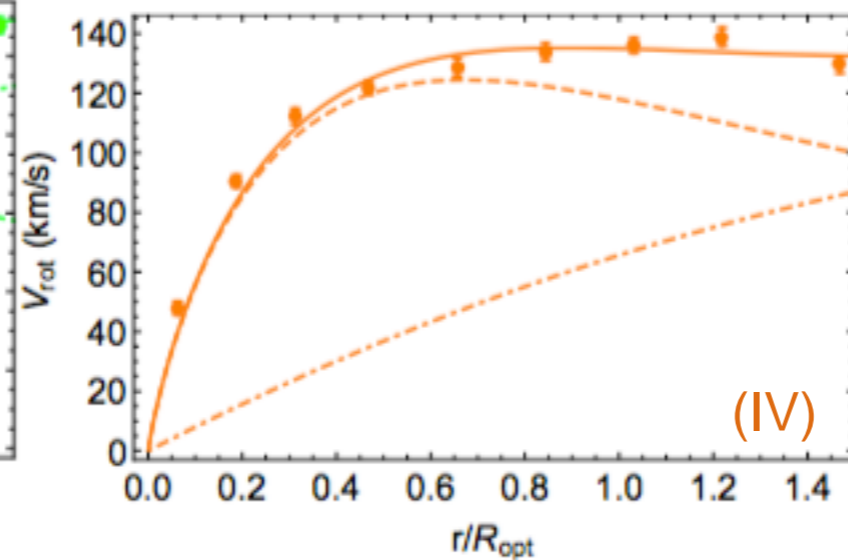
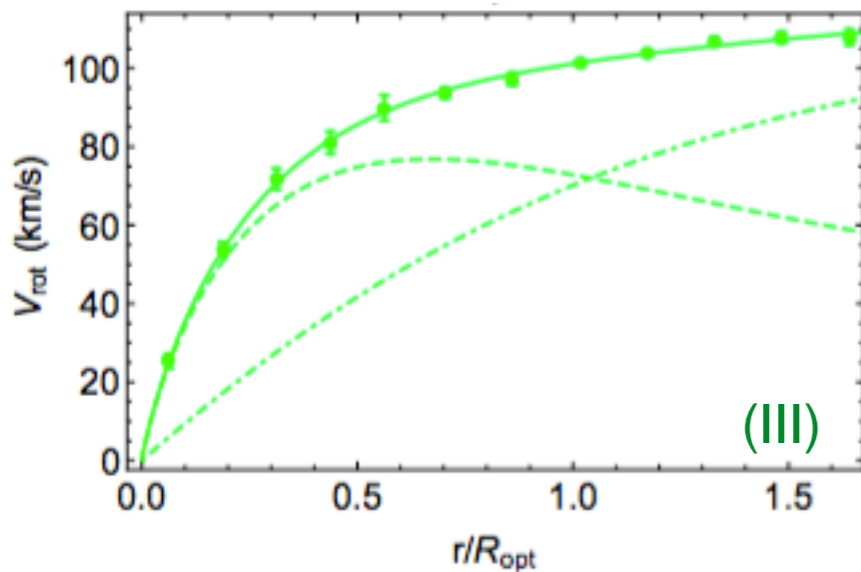
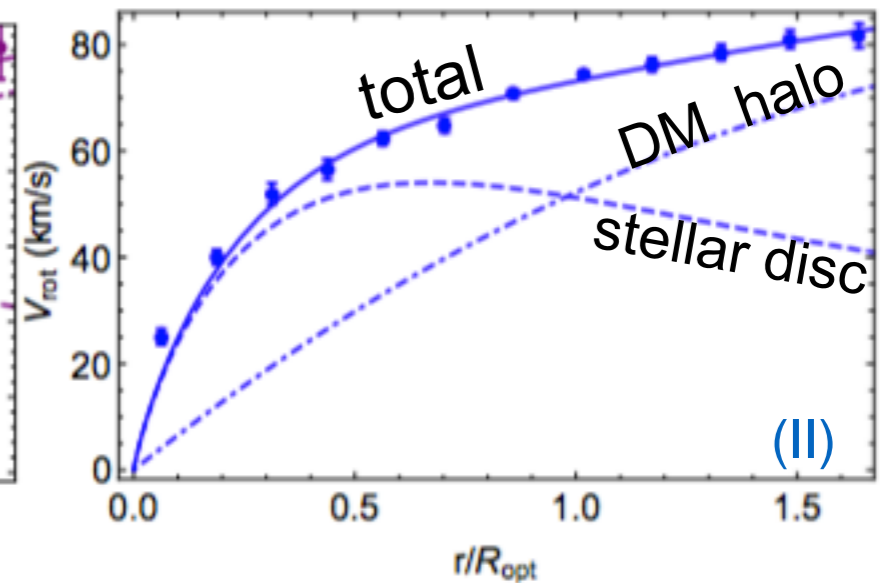
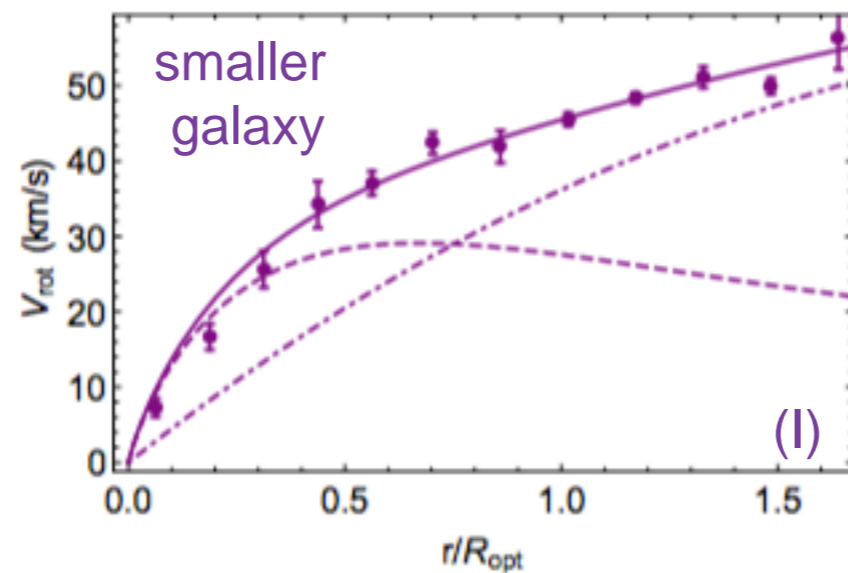
RCs and their coadded RCs in radial and velocity normalized units



Low Surface Brightness galaxies (LSBs)

Mass Modelling

5 co-added RCs



Vel. Bin	$\langle \rho_0 \rangle$ $10^{-3} M_{\odot}/pc^3$	$\langle R_c \rangle$ kpc	$\langle M_D \rangle$ $10^8 M_{\odot}$	$\langle M_{vir} \rangle$ $10^{11} M_{\odot}$	$\alpha(R_{opt})$
(1)	(2)	(3)	(4)	(5)	(6)
1	3.7 ± 1.4	10.7 ± 4.3	8.8 ± 1.8	1.0 ± 0.4	0.37
2	5.1 ± 1.1	12.8 ± 3.0	38 ± 3	2.4 ± 0.9	0.49
3	3.7 ± 0.5	17.1 ± 1.9	130 ± 5	4.0 ± 1.3	0.52
4	$1.7^{+3.2}_{-1.1}$	$29.7^{+84.1}_{-22.0}$	421 ± 40	8.4 ± 3.5	0.76
5	$0.8^{+1.1}_{-0.4}$	$99.1^{+750.5}_{-87.5}$	1730 ± 117	112 ± 55	0.82

$$V_i^2(r) = G \frac{M_i(r)}{r}$$

$$\alpha = \frac{\langle V_D^2(R_{opt}) \rangle}{\langle V_{tot}^2(R_{opt}) \rangle}$$

↓
baryonic fraction

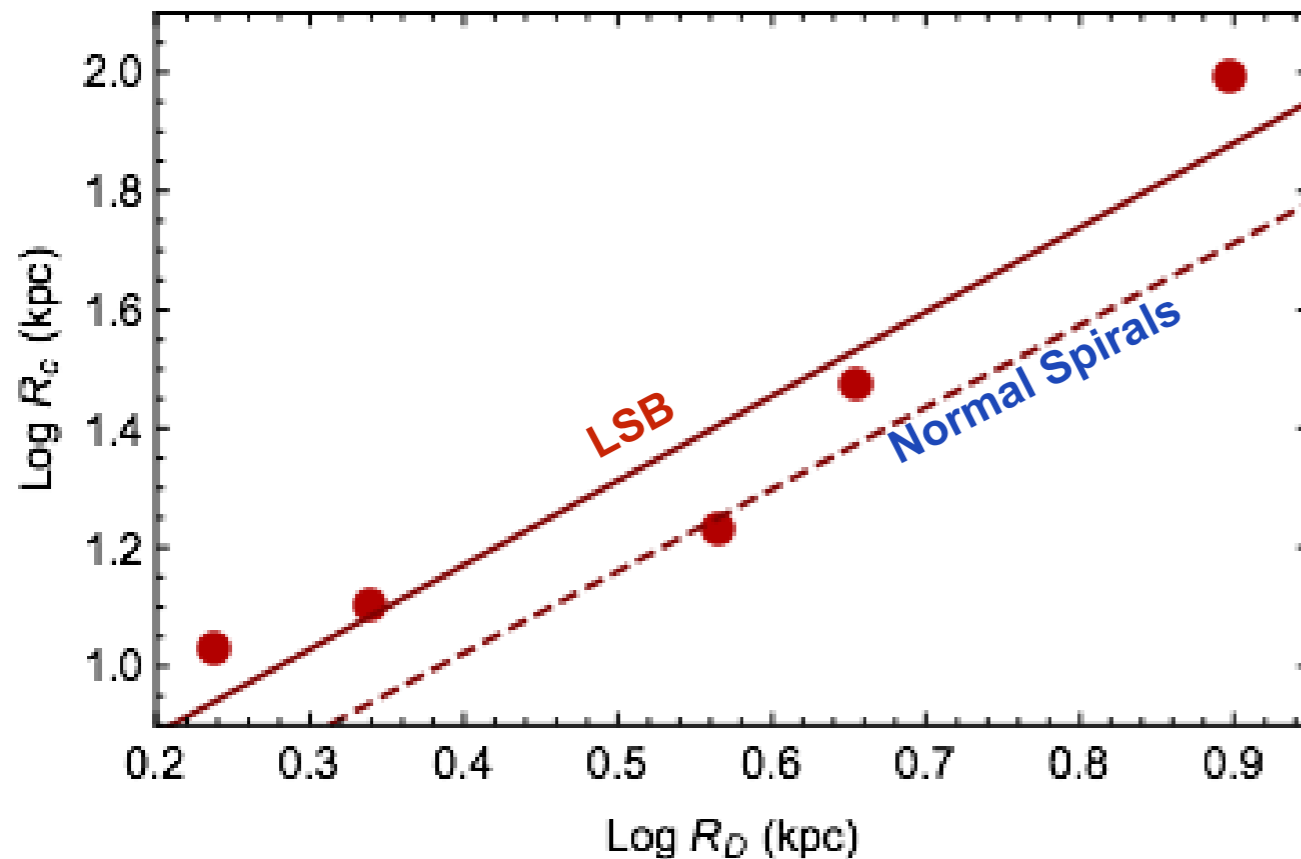
Low Surface Brightness galaxies (LSBs)

DENORMALIZATION

takes into account that all the double normalised RCs are similar to their co-added double normalised RC in **each single velocity bin**

↓
good approximation :

the relations obtained for the co-added RCs are assumed to be true also for the single galaxies



$R_c/R_d^{1.42} = const.$
one relation
in all velocity bins

▶ R_c

$\frac{M_d}{V_{opt}^2 R_{opt}} = const.$
one different value
in each velocity bin

▶ M_d

$$M_{DM}(R_{opt}) = G^{-1}(1 - \alpha)V_{opt}^2 R_{opt}$$

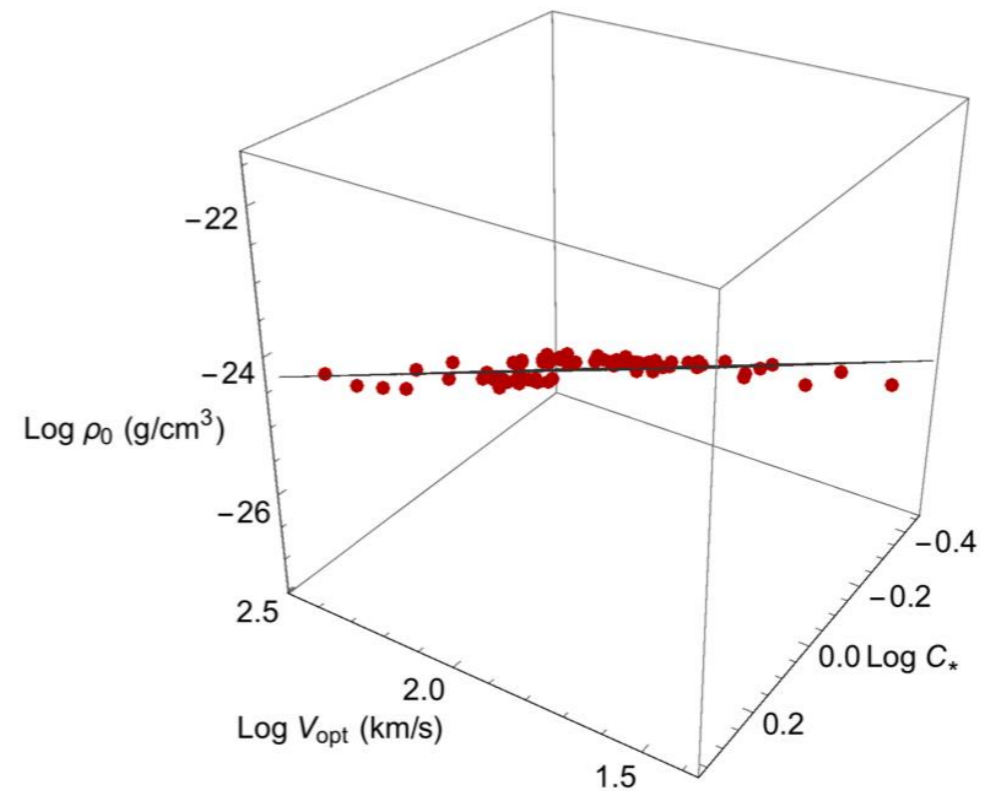
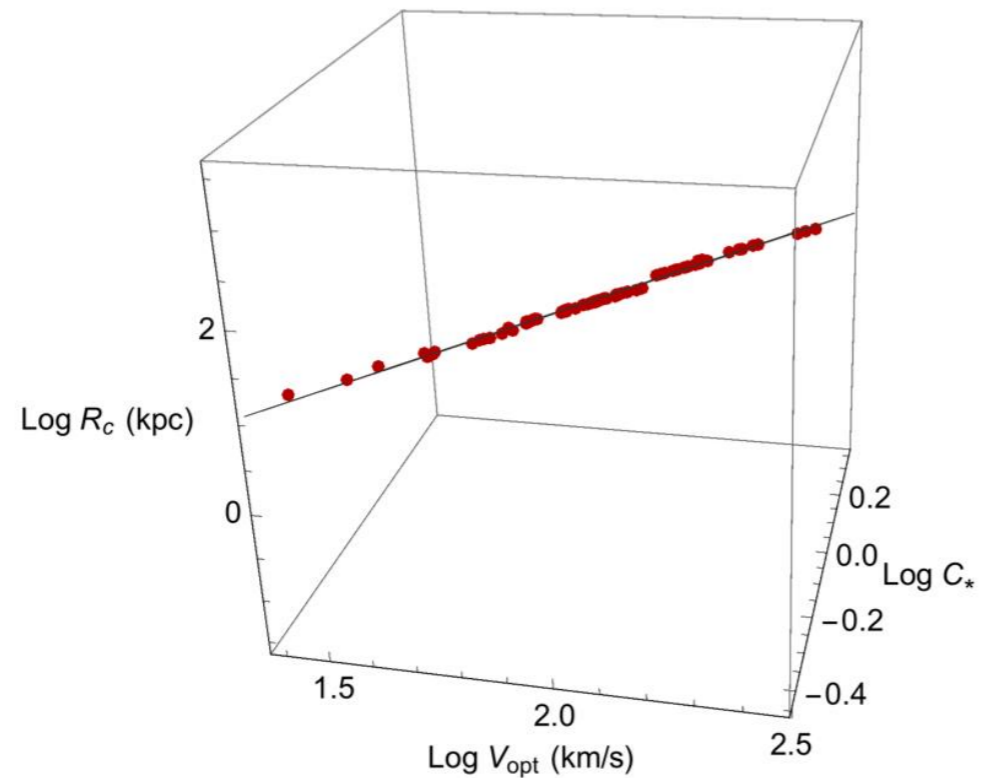
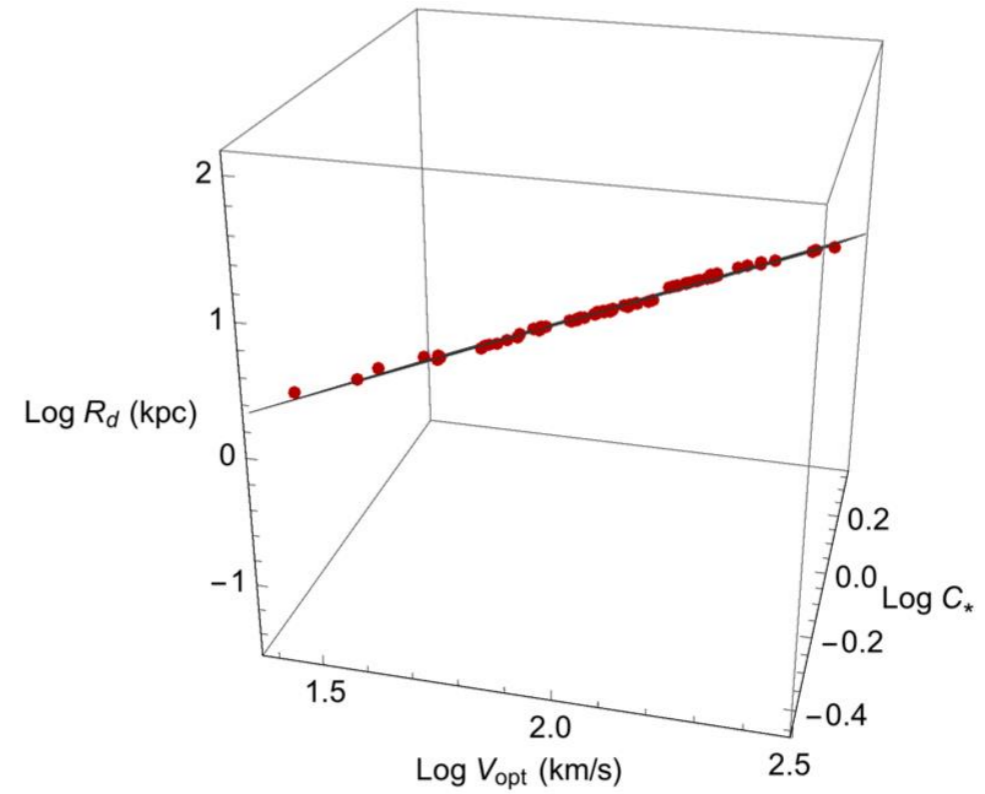
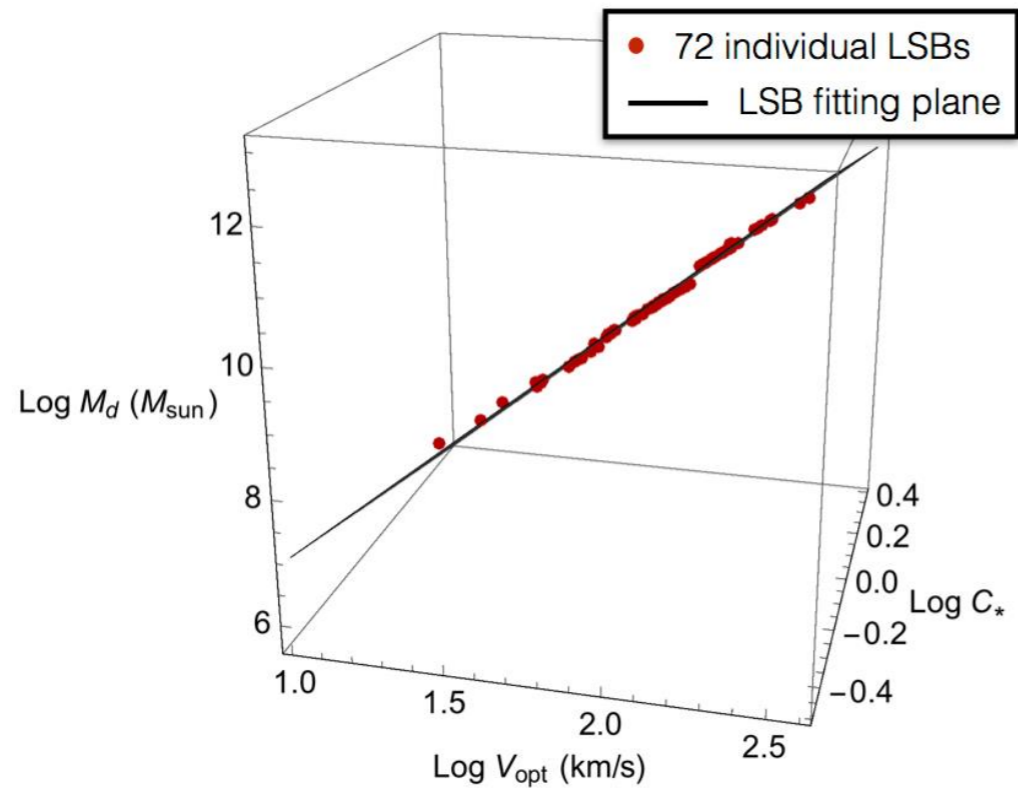
$\alpha = \frac{V_d^2(R_{opt})}{V^2(R_{opt})}$ = baryonic fraction at optical radius,
one different value in each velocity bin

▶ ρ_0

$$M_{DM}(r) = 2\pi\rho_0 R_c^3 [\ln(1 + r/R_c) - tg^{-1}(r/R_c) + 0.5\ln(1 + (r/R_c)^2)]$$

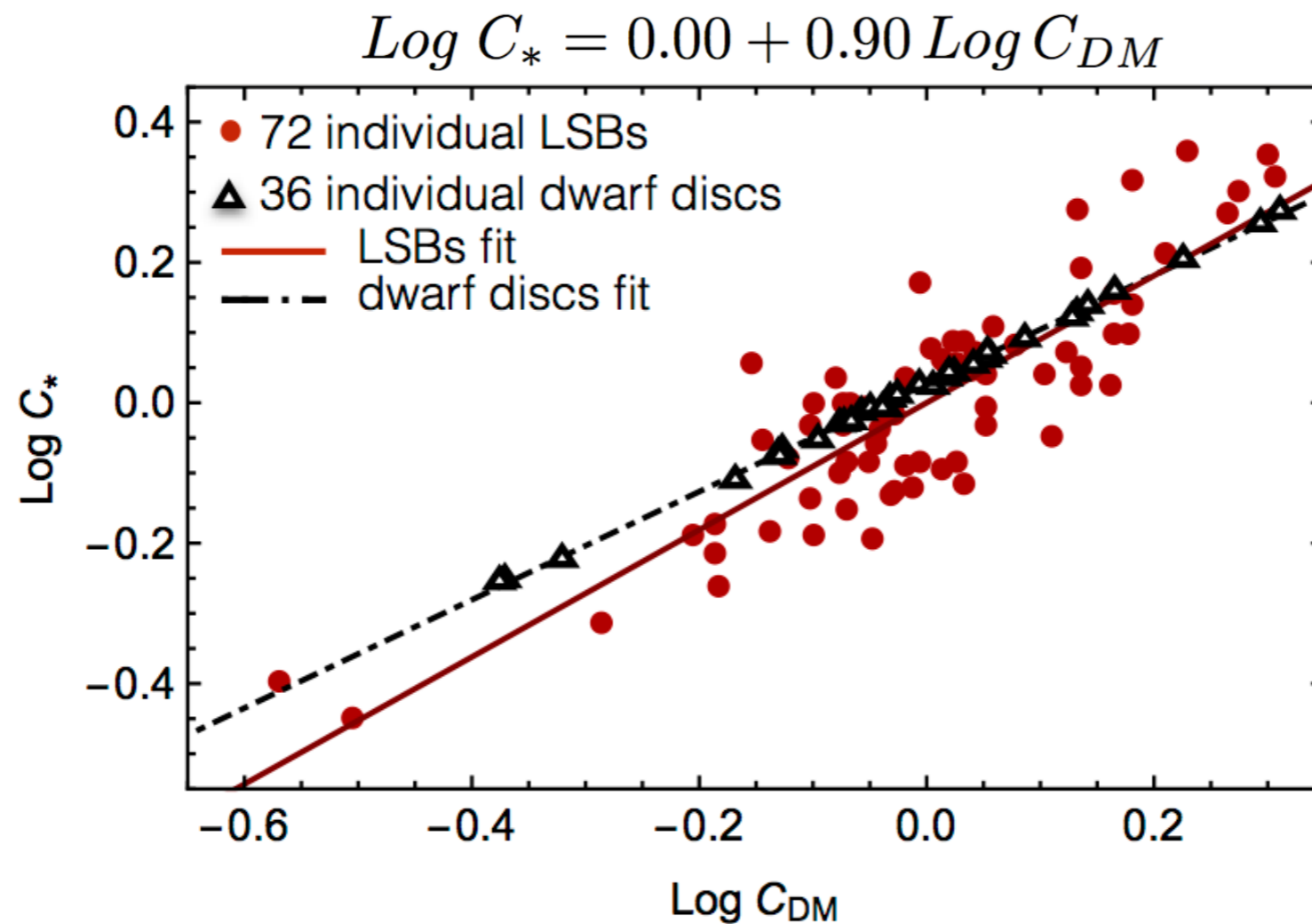
for a DM cored
Burkert profile

Low Surface Brightness galaxies (LSBs)



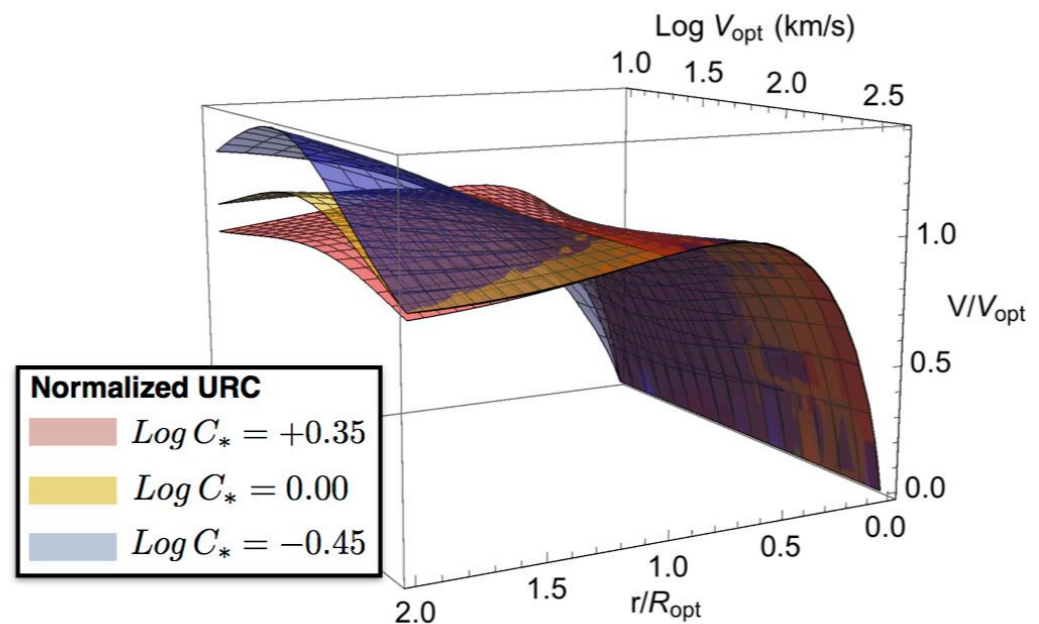
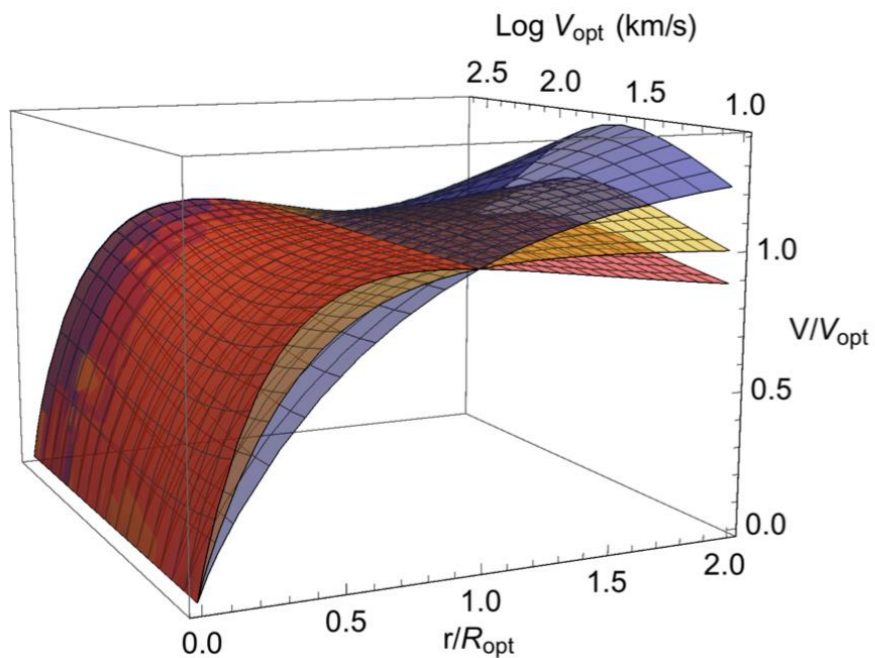
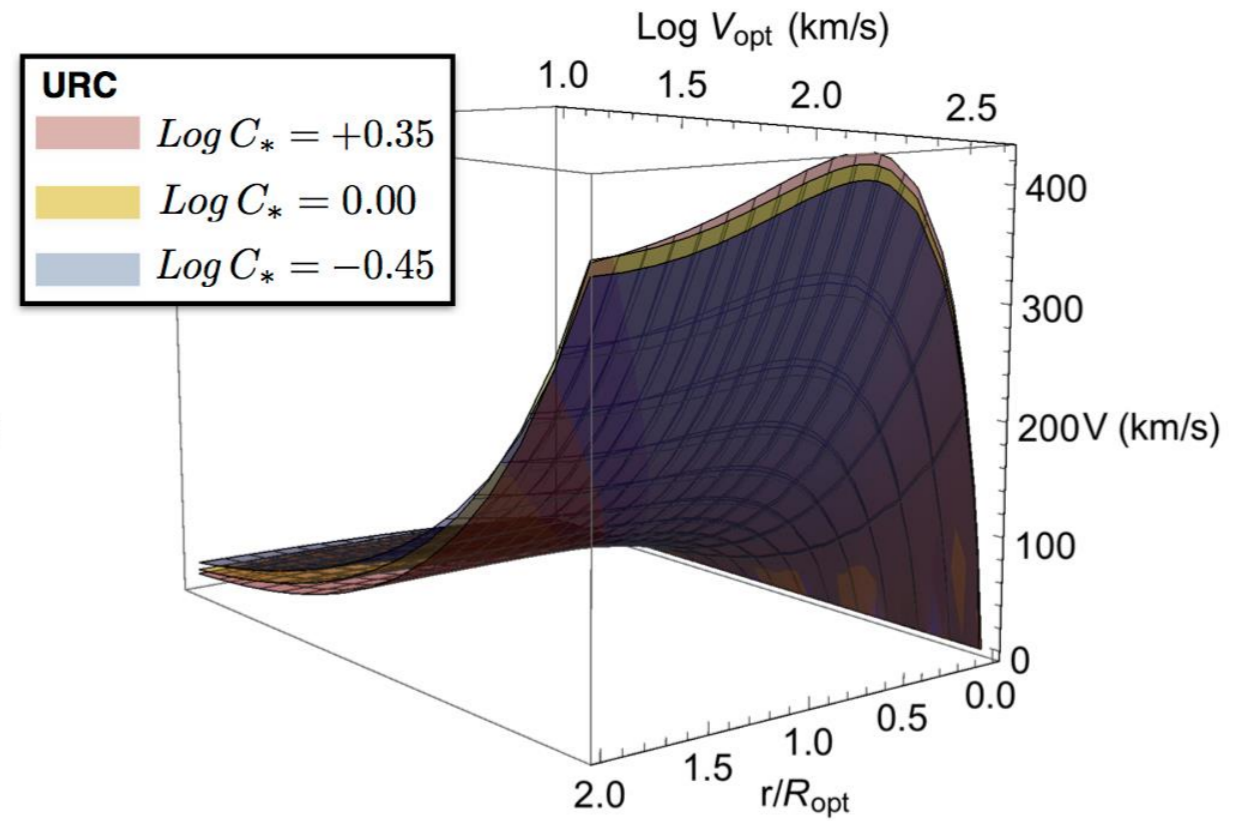
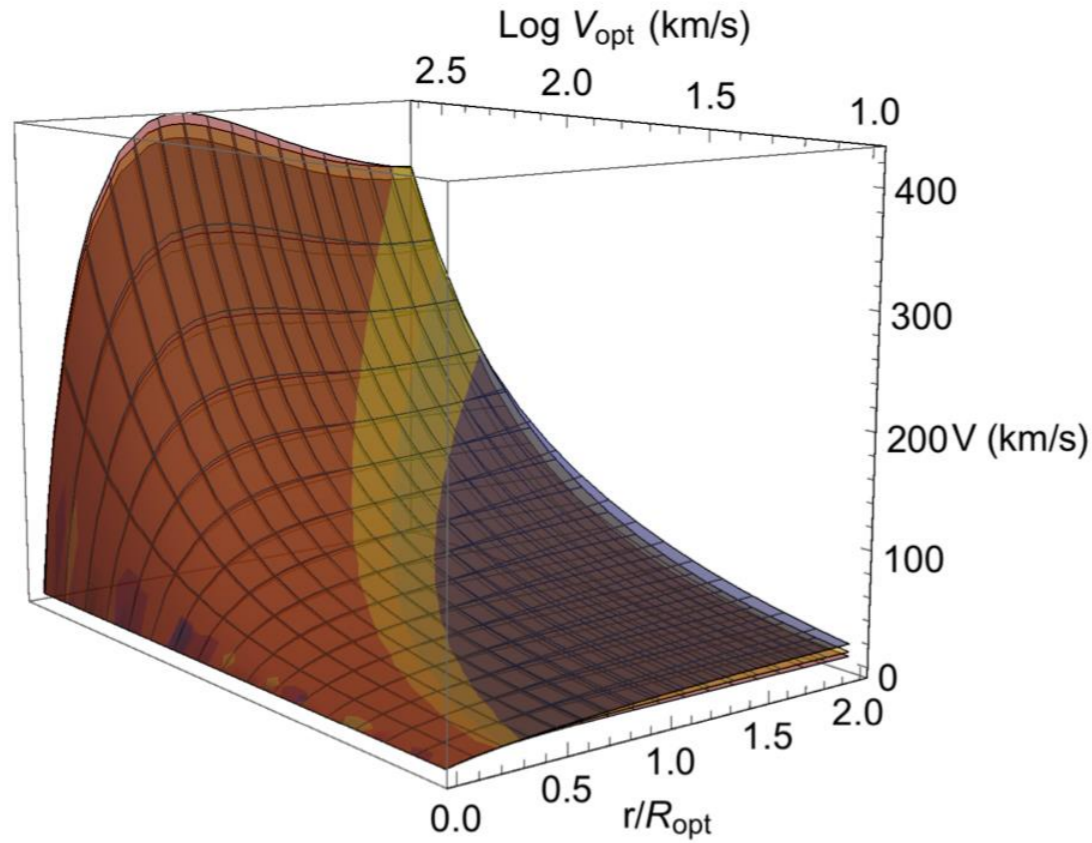
Compactness

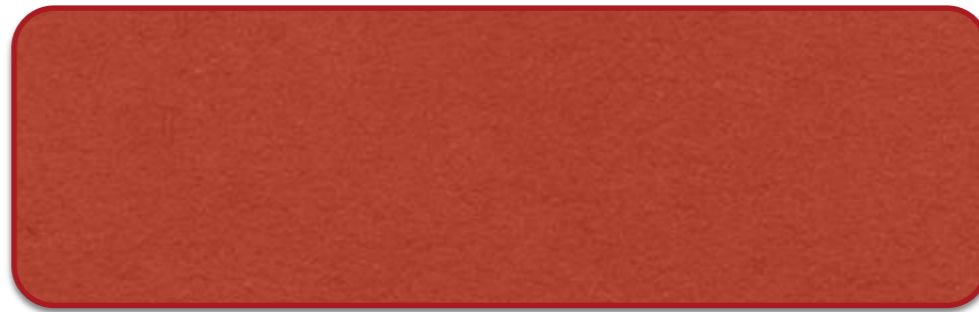
Similar to the dwarf discs relation



Strong correlation between the stellar and the DM compactness

Low Surface Brightness galaxies (LSBs)





the lower bound on the mass of a **fermionic** DM candidate from known small **Dwarf Spheroidal galaxies**

ADVANTAGES:

- model independent \longrightarrow does not require any assumption on initial distribution or evolution of DM particles
- relaxed the relation $R_h = R_{\frac{1}{2}}$

THOMAS-FERMI MODEL

galactic halo made of a fermionic self-gravitating gas

$\Phi(r)$: mean field gravitational potential, with spherical symmetry at a finite temperature T_0

Poisson equation

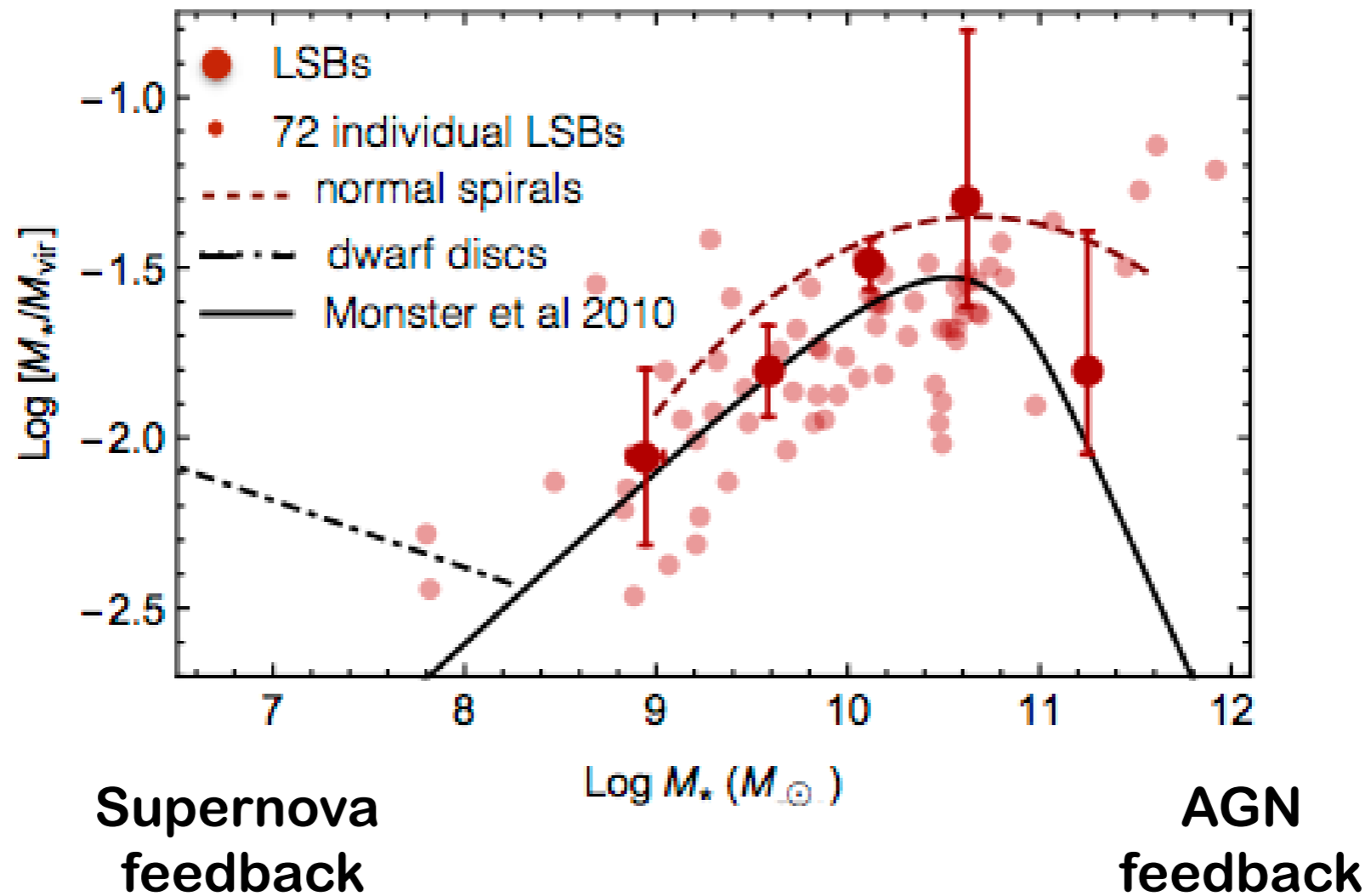
$$\begin{cases} \frac{d\phi(r)}{dr} = G \frac{M(r)}{r^2} \\ \frac{dM(r)}{dr} = 4\pi r^2 \rho(r) \end{cases}$$

$$\rho(r) = mn(r) = \frac{gm}{2\pi^2 \hbar^3} \int_0^\infty p^2 dp f \left[\frac{p^2}{2m} - \mu(r) \right] \quad g = 2$$

$$f(E) \quad E = p^2 / (2m) - \mu(r) \quad \mu(r) = \tilde{\mu}_0 - m\Phi(r)$$

$$f_{FD}(E) = \frac{1}{1 + \exp(E/T_0)}$$

Fraction of baryonic matter in LSBs versus their mass in stars



 masses for DM particle

Not so stringent

but **very robust** result because:

- ~ generic fermionic DM particle
- ~ model independent \longrightarrow based only on present phase-space density of DM particles in galaxies
- ~ relaxed the relation $R_h = R_{\frac{1}{2}}$